

## 45. Monte Carlo Particle Numbering Scheme

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The Monte Carlo particle numbering scheme presented here is intended to facilitate interfacing between matrix-element generators, event generators, detector simulators, and analysis packages used in particle physics, and is widely accepted as the “industry standard”. The numbering scheme was introduced in 1988 [1], and has since been revised and expanded in the light of new information on hadronic resonances, and expanded to encompass as yet undiscovered and hypothetical particles.

The general form is a 7–digit number:

$$\pm n \ n_r \ n_L \ n_{q_1} \ n_{q_2} \ n_{q_3} \ n_J .$$

This encodes information about the particle’s spin, flavor content, and internal quantum numbers. A 9–digit extension is used specifically for tetra- and penta-quark states, and a distinct 10–digit scheme is used for encoding (hyper)nuclear states. The details are as follows:

1. Particles are given positive numbers, antiparticles negative numbers. The PDG convention for mesons is used, so that  $K^+$  and  $B^+$  are particles.
2. Quarks and leptons are numbered consecutively starting from 1 and 11 respectively; to do this they are first ordered by family and within families by weak isospin.
3. In composite quark systems (diquarks, mesons, and baryons)  $n_{q_{1-3}}$  are quark numbers used to specify the quark content, while the rightmost digit  $n_J = 2J + 1$  gives the system’s spin (except for the  $K_S^0$  and  $K_L^0$ ). The scheme does not cover particles of spin  $J > 4$ .
4. Diquarks have 4–digit numbers with  $n_{q_1} \geq n_{q_2}$  and  $n_{q_3} = 0$ .
5. The numbering of mesons is guided by the nonrelativistic ( $L$ – $S$  decoupled) quark model, as listed in Tables 15.2, 15.3, and 15.4.
  - (a) The numbers specifying the meson’s quark content conform to the convention  $n_{q_1} = 0$  and  $n_{q_2} \geq n_{q_3}$ . The special case  $K_L^0$  is the sole exception to this rule.
  - (b) The quark numbers of flavorless, light ( $u, d, s$ ) mesons are: 11 for the member of the isotriplet ( $\pi^0, \rho^0, \dots$ ), 22 for the lighter isosinglet ( $\eta, \omega, \dots$ ), and 33 for the heavier isosinglet ( $\eta', \phi, \dots$ ). Since isosinglet mesons are often large mixtures of  $u\bar{u} + d\bar{d}$  and  $s\bar{s}$  states, 22 and 33 are assigned by mass and do not necessarily specify the dominant quark composition.
  - (c) The special numbers 310 and 130 are given to the  $K_S^0$  and  $K_L^0$  respectively.
  - (d) The fifth digit  $n_L$  is reserved to distinguish mesons of the same total ( $J$ ) but different spin ( $S$ ) and orbital ( $L$ ) angular momentum quantum numbers. For  $J > 0$  the numbers are:  $(L, S) = (J - 1, 1)$   $n_L = 0$ ,  $(J, 0)$   $n_L = 1$ ,  $(J, 1)$   $n_L = 2$  and  $(J + 1, 1)$   $n_L = 3$ . For the exceptional case  $J = 0$  the numbers are  $(0, 0)$   $n_L = 0$  and  $(1, 1)$   $n_L = 1$  (*i.e.*  $n_L = L$ ). See Table 45.1.
  - (e) If a set of physical mesons correspond to a (non-negligible) mixture of basis states, differing in their internal quantum numbers, then the lightest physical state gets the smallest basis state number. For example the  $K_1(1270)$  is numbered 10313 ( $1^1P_1 K_{1B}$ ) and the  $K_1(1400)$  is numbered 20313 ( $1^3P_1 K_{1A}$ ).
  - (f) The sixth digit  $n_r$  is used to label mesons radially excited above the ground state.
  - (g) Numbers have been assigned for complete  $n_r = 0$   $S$ - and  $P$ -wave multiplets, even where states remain to be identified.

**Table 45.1:** Meson numbering logic. Here  $qq$  stands for  $n_{q_2} n_{q_3}$ .

$J$	$L = J - 1, S = 1$			$L = J, S = 0$			$L = J, S = 1$			$L = J + 1, S = 1$		
	code	$J^{PC}$	$L$	code	$J^{PC}$	$L$	code	$J^{PC}$	$L$	code	$J^{PC}$	$L$
0	—	—	—	00qq1	0 <sup>-+</sup>	0	—	—	—	10qq1	0 <sup>++</sup>	1
1	00qq3	1 <sup>--</sup>	0	10qq3	1 <sup>+-</sup>	1	20qq3	1 <sup>++</sup>	1	30qq3	1 <sup>--</sup>	2
2	00qq5	2 <sup>++</sup>	1	10qq5	2 <sup>-+</sup>	2	20qq5	2 <sup>--</sup>	2	30qq5	2 <sup>++</sup>	3
3	00qq7	3 <sup>--</sup>	2	10qq7	3 <sup>+-</sup>	3	20qq7	3 <sup>++</sup>	3	30qq7	3 <sup>--</sup>	4
4	00qq9	4 <sup>++</sup>	3	10qq9	4 <sup>-+</sup>	4	20qq9	4 <sup>--</sup>	4	30qq9	4 <sup>++</sup>	5

- (h) In some instances assignments within the  $q\bar{q}$  meson model are only tentative; here best guess assignments are made.
- (i) Many states appearing in the Meson Listings are not yet assigned within the  $q\bar{q}$  model. Here  $n_{q_{2-3}}$  and  $n_J$  are assigned according to the state's likely flavors and spin; all such unassigned light isoscalar states are given the flavor code 22. Within these groups  $n_L = 0, 1, 2, \dots$  is used to distinguish states of increasing mass. These states are flagged using  $n = 9$ . It is to be expected that these numbers will evolve as the nature of the states are elucidated. Codes are assigned to all mesons which are listed in the one-page table at the end of the Meson Summary Table as long as they have a preferred or established spin. Additional heavy meson states expected from heavy quark spectroscopy are also assigned codes.
6. The numbering of baryons is again guided by the nonrelativistic quark model, see Table 15.6. This numbering scheme is illustrated through a few examples in Table 45.2.
- (a) The numbers specifying a baryon's quark content are such that in general  $n_{q_1} \geq n_{q_2} \geq n_{q_3}$ .
- (b) Two states exist for  $J = 1/2$  baryons containing 3 different types of quarks. In the lighter baryon ( $\Lambda, \Xi, \Omega, \dots$ ) the light quarks are in an antisymmetric ( $J = 0$ ) state while for the heavier baryon ( $\Sigma^0, \Xi', \Omega', \dots$ ) they are in a symmetric ( $J = 1$ ) state. In this situation  $n_{q_2}$  and  $n_{q_3}$  are reversed for the lighter state, so that the smaller number corresponds to the lighter baryon.
- (c) For excited baryons a scheme is adopted, where the  $n_r$  label is used to denote the excitation bands in the harmonic oscillator model, see Sec. 15.5. Using the notation employed there,  $n_r$  is given by the  $N$ -index of the  $D_N$  band identifier.
- (d) Further degeneracies of excited hadron multiplets with the same excitation number  $n_r$  and spin  $J$  are lifted by labelling such multiplets with the  $n_L$  index according to their mass, as given by its  $N$  or  $\Delta$ -equivalent.
- (e) In such excited multiplets extra singlets may occur, the  $\Lambda(1520)$  being a prominent example. In such cases the ordering is reversed such that the heaviest quark label is pushed to the last position:  $n_{q_3} > n_{q_1} > n_{q_2}$ .
7. The gluon, when considered as a gauge boson, has official number 21. In codes for glueballs, however, 9 is used to allow a notation in close analogy with that of hadrons.
8. The pomeron and odderon trajectories and a generic reggeon trajectory of states in QCD are assigned codes 990, 9990, and 110 respectively, where the final 0 indicates the indeterminate

**Table 45.2:** Some examples of octet (top) and decuplet (bottom) members for the numbering scheme for excited baryons. Here  $qqq$  stands for  $n_{q_1}n_{q_2}n_{q_3}$ . See the text for the definition of the notation. The numbers in parenthesis correspond to the mass of the baryons. The states marked as (?) are not experimentally confirmed.

$J^P$	$(D, L_N^P)$	$n_r n_L n_{q_1} n_{q_2} n_{q_3} n_J$	$N$	$\Lambda_8$	$\Sigma$	$\Xi$	$\Lambda_1$
Octet							
$1/2^+$	<b>(56, <math>0_0^+</math>)</b>	00qqq2	(939)	(1116)	(1193)	(1318)	—
$1/2^+$	<b>(56, <math>0_2^+</math>)</b>	20qqq2	(1440)	(1600)	(1660)	(1690)	—
$1/2^+$	<b>(70, <math>0_2^+</math>)</b>	21qqq2	(1710)	(1810)	(1880)	(?)	(?)
$1/2^-$	<b>(70, <math>1_1^-</math>)</b>	10qqq2	(1535)	(1670)	(1620)	(1750)	(1405)
$J^P$	$(D, L_N^P)$	$n_r n_L n_{q_1} n_{q_2} n_{q_3} n_J$	$\Delta$	$\Sigma$	$\Xi$	$\Omega$	
Decuplet							
$3/2^+$	<b>(56, <math>0_0^+</math>)</b>	00qqq4	(1232)	(1385)	(1530)	(1672)	
$3/2^+$	<b>(56, <math>0_2^+</math>)</b>	20qqq4	(1600)	(1690)	(?)	(?)	
$1/2^-$	<b>(70, <math>1_1^-</math>)</b>	11qqq2	(1620)	(1750)	(?)	(?)	
$3/2^-$	<b>(70, <math>1_1^-</math>)</b>	12qqq4	(1700)	(?)	(?)	(?)	

nature of the spin, and the other digits reflect the expected “valence” flavor content. We do not attempt a complete classification of all reggeon trajectories, since there is currently no need to distinguish a specific such trajectory from its lowest-lying member.

9. Two-digit numbers in the range 21–30 are provided for the Standard Model gauge and Higgs bosons.
10. Codes 81–100 are reserved for generator-specific pseudoparticles and concepts. Codes 901–930, 1901–1930, 2901–2930, and 3901–3930 are for additional components of Standard Model parton distribution functions, where the latter three ranges are intended to distinguish left/right/ longitudinal components. Codes 998 and 999 are reserved for GEANT tracking purposes.
11. The search for physics beyond the Standard Model is an active area, so these codes are also standardized as far as possible.
  - (a) A standard fourth generation of fermions is included by analogy with the first three.
  - (b) The graviton and the boson content of a two-Higgs-doublet scenario and of additional  $SU(2) \times U(1)$  groups are found in the range 31–40.
  - (c) “One-of-a-kind” exotic particles are assigned numbers in the range 41–80. The subrange 61–80 can be used for new heavier fermions in generic models, where partners to the SM fermions would have codes offset by 60. If required, however, other assignments could be made.
  - (d) Fundamental supersymmetric particles are identified by adding a nonzero  $n$  to the particle number. The superpartner of a boson or a left-handed fermion has  $n = 1$  while the superpartner of a right-handed fermion has  $n = 2$ . When mixing occurs, such as between the winos and charged Higgsinos to give charginos, or between left and right sfermions, the lighter physical state is given the smaller basis state number.
  - (e) Technicolor states have  $n = 3$ , with technifermions treated like ordinary fermions. States which are ordinary color singlets have  $n_r = 0$ . Color octets have  $n_r = 1$ . If a state has non-trivial quantum numbers under the topcolor groups  $SU(3)_1 \times SU(3)_2$ , the quantum numbers are specified by tech, $ij$ , where  $i$  and  $j$  are 1 or 2.  $n_L$  is then  $2i + j$ . The coloron,

$V_8$ , is a heavy gluon color octet and thus is 3100021.

- (f) Excited (composite) quarks and leptons are identified by setting  $n = 4$  and  $n_r = 0$ .
  - (g) Within several scenarios of new physics, it is possible to have colored particles sufficiently long-lived for color-singlet hadronic states to form around them. In the context of supersymmetric scenarios, these states are called  $R$ -hadrons, since they carry odd  $R$ -parity.  $R$ -hadron codes, defined here, should be viewed as templates for corresponding codes also in other scenarios, for any long-lived particle that is either an unflavored color octet or a flavored color triplet. The  $R$ -hadron code is obtained by combining the SUSY particle code with a code for the light degrees of freedom, with as many intermediate zeros removed from the former as required to make place for the latter at the end. (To exemplify, a sparticle  $n00000n_{\tilde{q}}$  combined with quarks  $q_1$  and  $q_2$  obtains code  $n00n_{\tilde{q}}n_{q_1}n_{q_2}n_J$ .) Specifically, the new-particle spin decouples in the limit of large masses, so that the final  $n_J$  digit is defined by the spin state of the light-quark system alone. An appropriate number of  $n_q$  digits is used to define the ordinary-quark content. As usual, 9 rather than 21 is used to denote a gluon/gluino in composite states. The sign of the hadron agrees with that of the constituent new particle (a color triplet) where there is a distinct new antiparticle, and else is defined as for normal hadrons. Particle names are  $R$  with the flavor content as lower index.
  - (h) A black hole in models with extra dimensions has code 5000040. Kaluza-Klein excitations in models with extra dimensions have  $n = 5$  or  $n = 6$ , to distinguish excitations of left- or right-handed fermions or, in case of mixing, the lighter or heavier state (cf. 11d). The nonzero  $n_r$  digit gives the radial excitation number, in scenarios where the level spacing allow these to be distinguished. Should the model also contain supersymmetry, excited SUSY states would be denoted by an  $n_r > 0$ , with  $n = 1$  or  $2$  as usual. Should some colored states be long-lived enough that hadrons would form around them, the coding strategy of 11g applies, with the initial two  $nn_r$  digits preserved in the combined code.
  - (i) Magnetic monopoles and dyons are assumed to have one unit of Dirac monopole charge and a variable integer number  $n_{q_1}n_{q_2}n_{q_3}$  units of electric charge. Codes  $411n_{q_1}n_{q_2}n_{q_3}0$  are then used when the magnetic and electrical charge sign agree and  $412n_{q_1}n_{q_2}n_{q_3}0$  when they disagree, with the overall sign of the particle set by the magnetic charge. For now no spin information is provided.
  - (j) The nature of Dark Matter (DM) is not known, and therefore a definitive classification is too early. Candidates within specific scenarios are classified therein, such as 1000022 for the lightest neutralino. Generic fundamental states can be given temporary codes in the range 51 - 60, with 51, 52 and 53 reserved for spin 0, 1/2 and 1 ones (this could also be an axion state). Generic mediators of s-channel DM pair creation or annihilation can be given codes 54 and 55 for spin 0 or 1 ones. Separate antiparticles, with negative codes, may or may not exist. More elaborate new scenarios should be constructed with  $n = 5$  and  $n_r = 9$ .
  - (k) Hidden Valley particles have  $n = 4$  and  $n_r = 9$ , and trailing numbers in agreement with their nearest-analog standard particles, as far as possible. Thus 4900021 is the gauge boson  $g_v$  of a confining gauge field,  $490000n_{q_v}$  and  $490001n_{\ell_v}$  fundamental constituents charged or not under this, 4900022 is the  $\gamma_v$  of a non-confining field, and  $4900n_{q_{v1}}n_{q_{v2}}n_J$  a Hidden Valley meson.
12. Occasionally program authors add their own states. To avoid confusion, these should be flagged by setting  $nn_r = 99$ .

13. Concerning the non-99 numbers, it may be noted that only quarks, excited quarks, squarks, and diquarks have  $n_{q_3} = 0$ ; only diquarks, baryons, and the odderon have  $n_{q_1} \neq 0$ ; and only mesons, the reggeon, and the pomeron have  $n_{q_1} = 0$  and  $n_{q_2} \neq 0$ . Concerning mesons (not antimesons), if  $n_{q_1}$  is odd then it labels a quark and an antiquark if even.
14. The 9-digit tetra-quark codes are  $\pm 1n_r n_L n_{q_1} n_{q_2} 0n_{q_3} n_{q_4} n_J$ . For the particle  $q_1 q_2$  is a diquark and  $\bar{q}_3 \bar{q}_4$  an antiquark, sorted such that  $n_{q_1} \geq n_{q_2}$ ,  $n_{q_3} \geq n_{q_4}$ ,  $n_{q_1} \geq n_{q_3}$ , and  $n_{q_2} \geq n_{q_4}$  if  $n_{q_1} = n_{q_3}$ . For the antiparticle, given with a negative sign,  $\bar{q}_1 \bar{q}_2$  is an antiquark and  $q_3 q_4$  a diquark, with the same sorting except that either  $n_{q_1} > n_{q_3}$  or  $n_{q_2} > n_{q_4}$  (so that flavour-diagonal states are particles). The  $n_r$ ,  $n_L$ , and  $n_J$  numbers have the same meaning as for ordinary hadrons.
15. The 9-digit penta-quark codes are  $\pm 1n_r n_L n_{q_1} n_{q_2} n_{q_3} n_{q_4} n_{q_5} n_J$ , sorted such that  $n_{q_1} \geq n_{q_2} \geq n_{q_3} \geq n_{q_4}$ . In the particle the first four are quarks and the fifth an antiquark while the opposite holds in the antiparticle, which is given with a negative sign. The  $n_r$ ,  $n_L$ , and  $n_J$  numbers have the same meaning as for ordinary hadrons.
16. Nuclear codes are given as 10-digit numbers  $\pm 10LZZZAAAI$ . For a (hyper)nucleus consisting of  $n_p$  protons,  $n_n$  neutrons and  $n_A$   $A$ 's,  $A = n_p + n_n + n_A$  gives the total baryon number,  $Z = n_p$  the total charge and  $L = n_A$  the total number of strange quarks.  $I$  gives the isomer level, with  $I = 0$  corresponding to the ground state and  $I > 0$  to excitations, see [2], where states denoted  $m, n, p, q$  translate to  $I = 1-4$ . As examples, the deuteron is 1000010020 and  $^{235}\text{U}$  is 1000922350. To avoid ambiguities, nuclear codes should not be applied to a single hadron, like  $p$ ,  $n$  or  $A^0$ , where quark-contents-based codes already exist.

QUARKS		GAUGE AND HIGGS BOSONS		SPECIAL PARTICLES		DIQUARKS		SUSY PARTICLES	
$d$	1					$(dd)_1$	1103	$\tilde{d}_L$	1000001
$u$	2			$G$ (graviton)	39	$(ud)_0$	2101	$\tilde{u}_L$	1000002
$s$	3	$g$	(9) 21	$R^0$	41	$(ud)_1$	2103	$\tilde{s}_L$	1000003
$c$	4	$\gamma$	22	$LQ^c$	42	$(uu)_1$	2203	$\tilde{c}_L$	1000004
$b$	5	$Z^0$	23	DM ( $S=0$ )	51	$(sd)_0$	3101	$\tilde{b}_1$	1000005 <sup>a</sup>
$t$	6	$W^+$	24	DM ( $S=\frac{1}{2}$ )	52	$(sd)_1$	3103	$\tilde{t}_1$	1000006 <sup>a</sup>
$b'$	7	$h^0/H_1^0$	25	DM ( $S=1$ )	53	$(su)_0$	3201	$\tilde{e}_L$	1000011
$t'$	8	$Z'/Z_2^0$	32	reggeon	110	$(su)_1$	3203	$\tilde{\nu}_{eL}$	1000012
		$Z''/Z_3^0$	33	pomeron	990	$(ss)_1$	3303	$\tilde{\mu}_L$	1000013
		$W'/W_2^+$	34	odderon	9990	$(cd)_0$	4101	$\tilde{\nu}_{\mu L}$	1000014
		$H^0/H_2^0$	35			$(cd)_1$	4103	$\tilde{\tau}_1$	1000015 <sup>a</sup>
		$A^0/H_3^0$	36	for MC internal		$(cu)_0$	4201	$\tilde{\nu}_{\tau L}$	1000016
		$H^+$	37	use 81–100,		$(cu)_1$	4203	$\tilde{d}_R$	2000001
$e^-$	11	$H^{++}$	38	901–930,		$(cs)_0$	4301	$\tilde{u}_R$	2000002
$\nu_e$	12	$a^0/H_4^0$	40	998–999,		$(cs)_1$	4303	$\tilde{s}_R$	2000003
$\mu^-$	13			1901–1930,		$(cc)_1$	4403	$\tilde{c}_R$	2000004
$\nu_\mu$	14			2901–2930, and		$(bd)_0$	5101	$\tilde{b}_2$	2000005 <sup>a</sup>
$\tau^-$	15			3901–3930		$(bd)_1$	5103	$\tilde{t}_2$	2000006 <sup>a</sup>
$\nu_\tau$	16					$(bu)_0$	5201	$\tilde{e}_R$	2000011
$\tau'^-$	17					$(bu)_1$	5203	$\tilde{\mu}_R$	2000013
$\nu_{\tau'}$	18					$(bs)_0$	5301	$\tilde{\tau}_2$	2000015 <sup>a</sup>
						$(bs)_1$	5303	$\tilde{g}$	1000021
						$(bc)_0$	5401		
						$(bc)_1$	5403		
						$(bb)_1$	5503		

$\tilde{\chi}_1^0$	1000022 <sup>b</sup>	$\rho_3(2250)^0$	9010117	$f_J(2220)$	9000229	<b>CHARMED MESONS</b>		<b><math>c\bar{c}</math> MESONS</b>	
$\tilde{\chi}_2^0$	1000023 <sup>b</sup>	$\rho_3(2250)^+$	9010217	$f_4(2300)$	9010229				
$\tilde{\chi}_1^+$	1000024 <sup>b</sup>	$a_4(2040)^0$	119	<b>STRANGE MESONS</b>		$D^+$	411	$\eta_c(1S)$	441
$\tilde{\chi}_3^0$	1000025 <sup>b</sup>	$a_4(2040)^+$	219			$D^0$	421	$\chi_{c0}(1P)$	10441
$\tilde{\chi}_4^0$	1000035 <sup>b</sup>	<b>LIGHT <math>I = 0</math> MESONS</b> ( $u\bar{u}, d\bar{d}, s\bar{s}$ admixtures)		$K_L^0$	130	$D_0^*(2400)^+$	10411	$\eta_c(2S)$	100441
$\tilde{\chi}_4^+$	1000037 <sup>b</sup>			$K_S^0$	310	$D_0^*(2400)^0$	10421	$J/\psi(1S)$	443
$\tilde{\chi}_2^+$	1000039	$K^0$	311	$D^*(2010)^+$	413	$h_c(1P)$	10443	$\chi_{c1}(1P)$	20443
$\tilde{G}$		$K^+$	321	$D^*(2007)^0$	423	$\psi(2S)$	100443	$\psi(3770)$	30443
<b>LIGHT <math>I = 1</math> MESONS</b>		$\eta$	221	$D_1(2420)^+$	10413	$\psi(4040)$	9000443	$\psi(4160)$	9010443
		$\eta'(958)$	331	$K_0^*(700)^0$	9000311	$D_1(2420)^0$	10423	$\psi(4415)$	9020443
$\pi^0$	111	$f_0(500)$	9000221	$K_0^*(700)^+$	9000321	$D_1(H)^+$	20413	$\chi_{c2}(1P)$	445
$\pi^+$	211	$f_0(980)$	9010221	$K_0^*(1430)^0$	10311	$D_1(2430)^0$	20423	$\chi_{c2}(3930)$	100445
$a_0(980)^0$	9000111	$\eta(1295)$	100221	$K_0^*(1430)^+$	10321	$D_2^*(2460)^+$	415	<b><math>b\bar{b}</math> MESONS</b>	
$a_0(980)^+$	9000211	$f_0(1370)$	10221	$K(1460)^0$	100311	$D_2^*(2460)^0$	425		
$\pi(1300)^0$	100111	$\eta(1405)$	9020221	$K(1460)^+$	100321	$D_s^+$	431	$\eta_b(1S)$	551
$\pi(1300)^+$	100211	$\eta(1475)$	100331	$K(1830)^0$	9010311	$D_{s0}^*(2317)^+$	10431	$\chi_{b0}(1P)$	10551
$a_0(1450)^0$	10111	$\eta(1760)$	9040221	$K(1830)^+$	9010321	$D_{s1}^+$	433	$\eta_b(2S)$	100551
$a_0(1450)^+$	10211	$f_0(1500)$	9030221	$K_0^*(1950)^0$	9020311	$D_{s1}(2536)^+$	10433	$\eta_b(3S)$	200551
$\pi(1800)^0$	9010111	$f_0(1710)$	10331	$K_0^*(1950)^+$	9020321	$D_{s1}(2460)^+$	20433	$\chi_{b0}(2P)$	110551
$\pi(1800)^+$	9010211	$\eta(1760)$	9040221	$K^*(892)^0$	313	$D_{s2}^*(2573)^+$	435	$\eta_b(3P)$	210551
$\rho(770)^0$	113	$f_0(2020)$	9050221	$K^*(892)^+$	323	<b>BOTTOM MESONS</b>		$\Upsilon(1S)$	553
$\rho(770)^+$	213	$f_0(2100)$	9060221	$K_1(1270)^0$	10313			$B^0$	511
$b_1(1235)^0$	10113	$f_0(2200)$	9070221	$K_1(1270)^+$	10323	$B^+$	521	$\chi_{b1}(1P)$	20553
$b_1(1235)^+$	10213	$\eta(2225)$	9080221	$K_1(1400)^0$	20313	$B_0^{*0}$	10511	$\Upsilon_1(1D)$	30553
$a_1(1260)^0$	20113	$\omega(782)$	223	$K_1(1400)^+$	20323	$B_0^{*+}$	10521	$\Upsilon(2S)$	100553
$a_1(1260)^+$	20213	$\phi(1020)$	333	$K^*(1410)^0$	100313	$B^{*0}$	513	$h_b(2P)$	110553
$\pi_1(1400)^0$	9000113	$h_1(1170)$	10223	$K^*(1410)^+$	100323	$B^{*+}$	523	$\chi_{b1}(2P)$	120553
$\pi_1(1400)^+$	9000213	$f_1(1285)$	20223	$K_1(1650)^0$	9000313	$B_1(L)^0$	10513	$\Upsilon_1(2D)$	130553
$\rho(1450)^0$	100113	$h_1(1380)$	10333	$K_1(1650)^+$	9000323	$B_1(L)^+$	10523	$\Upsilon(3S)$	200553
$\rho(1450)^+$	100213	$f_1(1420)$	20333	$K^*(1680)^0$	30313	$B_1(H)^0$	20513	$h_b(3P)$	210553
$\pi_1(1600)^0$	9010113	$\omega(1420)$	100223	$K^*(1680)^+$	30323	$B_1(H)^+$	20523	$\chi_{b1}(3P)$	220553
$\pi_1(1600)^+$	9010213	$f_1(1510)$	9000223	$K_2^*(1430)^0$	315	$B_2^{*0}$	515	$\Upsilon(4S)$	300553
$a_1(1640)^0$	9020113	$h_1(1595)$	9010223	$K_2^*(1430)^+$	325	$B_2^{*+}$	525	$\Upsilon(10860)$	9000553
$a_1(1640)^+$	9020213	$\omega(1650)$	30223	$K_2(1580)^0$	9000315	$B_s^0$	531	$\Upsilon(11020)$	9010553
$\rho(1700)^0$	30113	$\phi(1680)$	100333	$K_2(1580)^+$	9000325	$B_{s0}^0$	10531	$\chi_{b2}(1P)$	555
$\rho(1700)^+$	30213	$f_2(1270)$	225	$K_2(1770)^0$	10315	$B_s^{*0}$	533	$\eta_{b2}(1D)$	10555
$\rho(1900)^0$	9030113	$f_2(1430)$	9000225	$K_2(1770)^+$	10325	$B_{s1}(L)^0$	10533	$\Upsilon_2(1D)$	20555
$\rho(1900)^+$	9030213	$f_2'(1525)$	335	$K_2(1820)^0$	20315	$B_{s1}(H)^0$	20533	$\chi_{b2}(2P)$	100555
$\rho(2150)^0$	9040113	$f_2(1565)$	9010225	$K_2(1820)^+$	20325	$B_c^+$	541	$\eta_{b2}(2D)$	110555
$\rho(2150)^+$	9040213	$f_2(1640)$	9020225	$K_2^*(1980)^0$	9010315	$B_{c0}^{*+}$	10541	$\Upsilon_2(2D)$	120555
$a_2(1320)^0$	115	$\eta_2(1645)$	10225	$K_2^*(1980)^+$	9010325	$B_c^{*+}$	543	$\chi_{b2}(3P)$	200555
$a_2(1320)^+$	215	$f_2(1810)$	9030225	$K_2(2250)^0$	9020315	$B_{c1}(L)^+$	10543	$\Upsilon_3(1D)$	557
$\pi_2(1670)^0$	10115	$\eta_2(1870)$	10335	$K_2(2250)^+$	9020325	$B_{c1}(H)^+$	20543	$\Upsilon_3(2D)$	100557
$\pi_2(1670)^+$	10215	$f_2(1910)$	9040225	$K_3^*(1780)^0$	317	$B_{c2}^{*+}$	545		
$a_2(1700)^0$	9000115	$f_2(1950)$	9050225	$K_3^*(1780)^+$	327				
$a_2(1700)^+$	9000215	$f_2(2010)$	9060225	$K_3(2320)^0$	9010317				
$\pi_2(2100)^0$	9010115	$f_2(2150)$	9070225	$K_3(2320)^+$	9010327				
$\pi_2(2100)^+$	9010215	$f_2(2300)$	9080225	$K_4^*(2045)^0$	319				
$\rho_3(1690)^0$	117	$f_2(2340)$	9090225	$K_4^*(2045)^+$	329				
$\rho_3(1690)^+$	217	$\omega_3(1670)$	227	$K_4(2500)^0$	9000319				
$\rho_3(1990)^0$	9000117	$\phi_3(1850)$	337	$K_4(2500)^+$	9000329				
$\rho_3(1990)^+$	9000217	$f_4(2050)$	229						

<b>LIGHT BARYONS</b>		$\Xi^0$	3322	$\Xi_c^{\prime 0}$	4312	$\Sigma_b^+$	5222	$\Omega_{bc}^{\prime 0}$	5432
$p$	2212	$\Xi^-$	3312	$\Xi_c^{*+}$	4324	$\Sigma_b^{*-}$	5114	$\Omega_{bc}^{*0}$	5434
$n$	2112	$\Xi^{*0}$	3324 <sup>c</sup>	$\Xi_c^{*0}$	4314	$\Sigma_b^{*0}$	5214	$\Omega_{bcc}^+$	5442
$\Delta^{++}$	2224	$\Xi^{*-}$	3314 <sup>c</sup>	$\Omega_c^0$	4332	$\Sigma_b^{*+}$	5224	$\Omega_{bcc}^{*+}$	5444
$\Delta^+$	2214	$\Omega^-$	3334	$\Omega_c^{\prime 0}$	4334	$\Xi_b^-$	5132	$\Xi_{bb}^-$	5512
$\Delta^0$	2114	<b>CHARMED BARYONS</b>		$\Xi_c^+$	4412	$\Xi_b^0$	5232	$\Xi_{bb}^0$	5522
$\Delta^-$	1114	$\Lambda_c^+$	4122	$\Xi_{cc}^{++}$	4422	$\Xi_b^{\prime -}$	5312	$\Xi_{bb}^{*-}$	5514
<b>STRANGE BARYONS</b>		$\Sigma_c^{++}$	4222	$\Xi_{cc}^{*+}$	4414	$\Xi_b^{\prime 0}$	5322	$\Xi_{bb}^{*0}$	5524
$\Lambda$	3122	$\Sigma_c^+$	4212	$\Xi_{cc}^{*+}$	4424	$\Xi_b^{*-}$	5314	$\Omega_{bb}^-$	5532
$\Sigma^+$	3222	$\Sigma_c^0$	4112	$\Omega_{cc}^+$	4432	$\Xi_b^{*0}$	5324	$\Omega_{bb}^{*0}$	5534
$\Sigma^0$	3212	$\Sigma_c^{*+}$	4224	$\Omega_{cc}^{*+}$	4434	$\Omega_b^-$	5332	$\Omega_{bbc}^0$	5542
$\Sigma^-$	3112	$\Sigma_c^{*0}$	4214	$\Omega_{ccc}^{*+}$	4444	$\Omega_b^{*-}$	5334	$\Omega_{bbc}^{*0}$	5544
$\Sigma^{*+}$	3224 <sup>c</sup>	$\Sigma_c^{\prime 0}$	4112	<b>BOTTOM BARYONS</b>		$\Xi_{bc}^0$	5142	$\Omega_{bbb}^-$	5554
$\Sigma^{*0}$	3214 <sup>c</sup>	$\Sigma_c^{*+}$	4224	$\Lambda_b^0$	5122	$\Xi_{bc}^+$	5242	<b>PENTA-QUARKS</b>	
$\Sigma^{*-}$	3114 <sup>c</sup>	$\Sigma_c^{\prime 0}$	4214	$\Sigma_b^-$	5112	$\Xi_{bc}^{\prime 0}$	5412	$\Theta^+$	<b>100221132*</b>
		$\Xi_c^+$	4232	$\Sigma_b^0$	5212	$\Xi_{bc}^*$	5414	$\Phi^{--}$	<b>100331122*</b>
		$\Xi_c^0$	4132			$\Xi_{bc}^{*+}$	5424		
		$\Xi_c^{\prime +}$	4322			$\Omega_{bc}^0$	5342		

**Footnotes to the tables:**

- \*) Numbers or names in bold face are new or have changed since the 2020 *Review*.
- a) Particular in the third generation, the left and right sfermion states may mix, as shown. The lighter mixed state is given the smaller number.
- b) The physical  $\tilde{\chi}$  states are admixtures of the pure  $\tilde{\gamma}$ ,  $\tilde{Z}^0$ ,  $\tilde{W}^+$ ,  $\tilde{H}_1^0$ ,  $\tilde{H}_2^0$ , and  $\tilde{H}^+$  states.
- c)  $\Sigma^*$  and  $\Xi^*$  are alternate names for  $\Sigma(1385)$  and  $\Xi(1530)$ .

This text and full lists of particle numbers can be found online [3].

**References**

- [1] G. P. Yost *et al.* (Particle Data Group), *Phys. Lett.* **B204**, 1 (1988).
- [2] G. Audi *et al.*, *Nucl. Phys.* **A729**, 3 (2003).
- [3] <http://pdg.lbl.gov/current/mc-particle-id>.