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## W MASS

The $W$-mass listed here corresponds to the mass parameter in a BreitWigner distribution with mass-dependent width. To obtain the world average, common systematic uncertainties between experiments are properly taken into account. The LEP-2 average $W$ mass based on published results is $80.376 \pm 0.033 \mathrm{GeV}$ [SCHAEL 13A]. The combined Tevatron data yields an average $W$ mass of $80.387 \pm 0.016 \mathrm{GeV}$ [AALTONEN 13N].

OUR FIT uses these average LEP and Tevatron mass values and combines them assuming no correlations.

| VALUE (GeV) | EVTS | DOCUMENT ID |  | TECN | COMMENT |
| :---: | :---: | :---: | :---: | :---: | :---: |
| 80.385 0.015 OUR FIT |  |  |  |  |  |
| $80.375 \pm 0.023$ | 2177k | 1 ABAZOV | 14N | D0 | $E_{\mathrm{cm}}^{p \bar{p}}=1.96 \mathrm{TeV}$ |
| $80.387 \pm 0.019$ | 1095k | 2 AALTONEN | 12E | CDF | $E_{\mathrm{cm}}^{p \bar{p}}=1.96 \mathrm{TeV}$ |
| $80.336 \pm 0.055 \pm 0.039$ | 10.3k | 3 ABDALLAH | 08A | DLPH | $E_{\mathrm{cm}}^{e \mathrm{ee}}=161-209 \mathrm{GeV}$ |
| $80.415 \pm 0.042 \pm 0.031$ | 11830 | 4 ABBIENDI | 06 | OPAL | $E_{\mathrm{cm}}^{\mathrm{ee}}=170-209 \mathrm{GeV}$ |
| $80.270 \pm 0.046 \pm 0.031$ | 9909 | ${ }^{5}$ ACHARD | 06 | L3 | $E_{\mathrm{cm}}^{e e}=161-209 \mathrm{GeV}$ |
| $80.440 \pm 0.043 \pm 0.027$ | 8692 | 6 SCHAEL | 06 | ALEP | $E_{\mathrm{cm}}^{\mathrm{ee}}=161-209 \mathrm{GeV}$ |
| $80.483 \pm 0.084$ | 49247 | 7 ABAZOV | 02D | D0 | $E_{\mathrm{cm}}^{p \bar{p}}=1.8 \mathrm{TeV}$ |
| $80.433 \pm 0.079$ | 53841 | 8 AFFOLDER | 01E | CDF | $E_{\mathrm{cm}}^{p \bar{p}}=1.8 \mathrm{TeV}$ |

-     - We do not use the following data for averages, fits, limits, etc.

| $80.367 \pm 0.026$ | 1677k | 9 ABAZOV | 12F | D0 | $E_{\mathrm{cm}}^{p \bar{p}}=1.96 \mathrm{TeV}$ |
| :---: | :---: | :---: | :---: | :---: | :---: |
| $80.401 \pm 0.043$ | 500k | 10 ABAZOV | 09AB | D0 | $E_{\mathrm{cm}}^{p \bar{p}}=1.96 \mathrm{TeV}$ |
| $80.413 \pm 0.034 \pm 0.034$ | 115k | 11 AALTONEN | 07F | CDF | $E_{\mathrm{cm}}^{p \bar{p}}=1.96 \mathrm{TeV}$ |
| $82.87 \pm 1.82 \begin{aligned} & +0.30 \\ & -0.16\end{aligned}$ | 1500 | 12 AKTAS | 06 | H1 | $\begin{aligned} e^{ \pm} p & \rightarrow \bar{\nu}_{e}\left(\nu_{e}\right) X \\ \sqrt{s} & \approx 300 \mathrm{GeV} \end{aligned}$ |
| $80.3 \pm 2.1 \pm 1.2 \pm 1.0$ | 645 | 13 CHEKANOV | 02C | ZEUS | $\underset{318 \mathrm{GeV}}{e^{-} p \rightarrow \sqrt{s}=}$ |
| $81.4{ }_{-2.6}^{+2.7} \pm 2.0_{-3.0}^{+3.3}$ | 1086 | 14 BREITWEG | 00D | ZEUS | $e^{+} \underset{300 \mathrm{GGV}_{e}}{p} \mathrm{X}, \sqrt{s} \approx$ |
| $80.84 \pm 0.22 \pm 0.83$ | 2065 | ${ }^{15}$ ALITTI | 92B | UA2 | See $W / Z$ ratio below |
| $80.79 \pm 0.31 \pm 0.84$ |  | 16 ALITTI | 90B | UA2 | $E_{\mathrm{cm}}^{p \bar{p}}=546,630 \mathrm{GeV}$ |
| $80.0 \pm 3.3 \pm 2.4$ | 22 | 17 ABE | 891 | CDF | $E_{\mathrm{cm}}^{p \bar{p}}=1.8 \mathrm{TeV}$ |
| $82.7 \pm 1.0 \pm 2.7$ | 149 | 18 ALBAJAR | 89 | UA1 | $E_{\mathrm{cm}}^{p \bar{p}}=546,630 \mathrm{GeV}$ |
| 81.8 +5.0 -5.3 | 46 | 19 ALBAJAR | 89 | UA1 | $E_{\mathrm{cm}}^{p \bar{p}}=546,630 \mathrm{GeV}$ |
| $89 \pm 3 \pm 6$ | 32 | 20 ALBAJAR | 89 | UA1 | $E_{\mathrm{cm}}^{p \bar{p}}=546,630 \mathrm{GeV}$ |
| 81. $\pm 5$. | 6 | ARNISON | 83 | UA1 | $E_{\mathrm{Cm}}^{e \mathrm{e}}$ = $=546 \mathrm{GeV}$ |
| 80. +10. | 4 | BANNER | 83B | UA2 | Repl. by ALITTI 90B |

${ }^{1}$ ABAZOV 14 N is a combination of $A B A Z O V 09 A B$ and $A B A Z O V 12 F$, also giving more details on the analysis.
${ }^{2}$ AALTONEN 12E select 470k $W \rightarrow e \nu$ decays and $625 \mathrm{k} W \rightarrow \mu \nu$ decays in $2.2 \mathrm{fb}^{-1}$ of Run-II data. The mass is determined using the transverse mass, transverse lepton momentum and transverse missing energy distributions, accounting for correlations. This result supersedes AALTONEN 07F. AALTONEN 14D gives more details on the procedures followed by the authors.
${ }^{3}$ ABDALLAH 08A use direct reconstruction of the kinematics of $W^{+} W^{-} \rightarrow q \bar{q} \ell \nu$ and $W^{+} W^{-} \rightarrow q \bar{q} q \bar{q}$ events for energies 172 GeV and above. The $W$ mass was also extracted from the dependence of the $W W$ cross section close to the production threshold and combined appropriately to obtain the final result. The systematic error includes $\pm 0.025 \mathrm{GeV}$ due to final state interactions and $\pm 0.009 \mathrm{GeV}$ due to LEP energy uncertainty.
${ }^{4}$ ABBIENDI 06 use direct reconstruction of the kinematics of $W^{+} W^{-} \rightarrow q \bar{q} \ell \nu_{\ell}$ and $W^{+} W^{-} \rightarrow q \bar{q} q \bar{q}$ events. The result quoted here is obtained combining this mass value with the results using $W^{+} W^{-} \rightarrow \ell \nu_{\ell} \ell^{\prime} \nu_{\ell^{\prime}}$ events in the energy range 183-207 GeV (ABBIENDI 03C) and the dependence of the $W W$ production cross-section on $m_{W}$ at threshold. The systematic error includes $\pm 0.009 \mathrm{GeV}$ due to the uncertainty on the LEP beam energy.
${ }^{5}$ ACHARD 06 use direct reconstruction of the kinematics of $W^{+} W^{-} \rightarrow q \bar{q} \ell \nu_{\ell}$ and $W^{+} W^{-} \rightarrow q \bar{q} q \bar{q}$ events in the C.M. energy range $189-209 \mathrm{GeV}$. The result quoted here is obtained combining this mass value with the results obtained from a direct $W$ mass reconstruction at 172 and 183 GeV and with those from the dependence of the $W W$ production cross-section on $m_{W}$ at 161 and 172 GeV (ACCIARRI 99).
${ }^{6}$ SCHAEL 06 use direct reconstruction of the kinematics of $W^{+} W^{-} \rightarrow q \bar{q} \ell \nu_{\ell}$ and $W^{+} W^{-} \rightarrow q \bar{q} q \bar{q}$ events in the C.M. energy range $183-209 \mathrm{GeV}$. The result quoted here is obtained combining this mass value with those obtained from the dependence of the $W$ pair production cross-section on $m_{W}$ at 161 and 172 GeV (BARATE 97 and BARATE 97S respectively). The systematic error includes $\pm 0.009 \mathrm{GeV}$ due to possible effects of final state interactions in the $q \bar{q} q \bar{q}$ channel and $\pm 0.009 \mathrm{GeV}$ due to the uncertainty on the LEP beam energy.
${ }^{7}$ ABAZOV 02D improve the measurement of the $W$-boson mass including $W \rightarrow$ e $\nu_{e}$ events in which the electron is close to a boundary of a central electromagnetic calorimeter module. Properly combining the results obtained by fitting $m_{T}(W), p_{T}(e)$, and $p_{T}(\nu)$, this sample provides a mass value of $80.574 \pm 0.405 \mathrm{GeV}$. The value reported here is a combination of this measurement with all previous $D \varnothing W$-boson mass measurements.
${ }^{8}$ AFFOLDER 01E fit the transverse mass spectrum of $30115 \mathrm{~W} \rightarrow e \nu_{e}$ events ( $M_{W}=$ $80.473 \pm 0.065 \pm 0.092 \mathrm{GeV})$ and of $14740 W \rightarrow \mu \nu_{\mu}$ events $\left(M_{W}=80.465 \pm 0.100 \pm\right.$ 0.103 GeV ) obtained in the run IB (1994-95). Combining the electron and muon results, accounting for correlated uncertainties, yields $M_{W}=80.470 \pm 0.089 \mathrm{GeV}$. They combine this value with their measurement of ABE 95P reported in run IA (1992-93) to obtain the quoted value.
${ }^{9}$ ABAZOV 12F select $1677 \mathrm{k} W \rightarrow e \nu$ decays in $4.3 \mathrm{fb}^{-1}$ of Run-II data. The mass is determined using the transverse mass and transverse lepton momentum distributions, accounting for correlations.
10 ABAZOV 09AB study the transverse mass, transverse electron momentum, and transverse missing energy in a sample of 0.5 million $W \rightarrow e \nu$ decays selected in Run-II data. The quoted result combines all three methods, accounting for correlations.
11 AALTONEN 07F obtain high purity $W \rightarrow e \nu_{e}$ and $W \rightarrow \mu \nu_{\mu}$ candidate samples totaling 63,964 and 51,128 events respectively. The $W$ mass value quoted above is derived by simultaneously fitting the transverse mass and the lepton, and neutrino $\mathrm{p}_{T}$ distributions.
12 AKTAS 06 fit the $\mathrm{Q}^{2}$ dependence $\left(300<\mathrm{Q}^{2}<30,000 \mathrm{GeV}^{2}\right)$ of the charged-current differential cross section with a propagator mass. The first error is experimental and the second corresponds to uncertainties due to input parameters and model assumptions.

13 CHEKANOV 02C fit the $Q^{2}$ dependence $\left(200<Q^{2}<60000 \mathrm{GeV}^{2}\right)$ of the charged-current differential cross sections with a propagator mass fit. The last error is due to the uncertainty on the probability density functions.
14 BREITWEG 00D fit the $Q^{2}$ dependence ( $200<\mathrm{Q}^{2}<22500 \mathrm{GeV}^{2}$ ) of the chargedcurrent differential cross sections with a propagator mass fit. The last error is due to the uncertainty on the probability density functions.
${ }^{15}$ ALITTI 92B result has two contributions to the systematic error $( \pm 0.83)$; one $( \pm 0.81)$ cancels in $m_{W} / m_{Z}$ and one ( $\pm 0.17$ ) is noncancelling. These were added in quadrature. We choose the ALITTI 92B value without using the LEP $m_{Z}$ value, because we perform our own combined fit.
16 There are two contributions to the systematic error $( \pm 0.84)$ : one $( \pm 0.81)$ which cancels in $m_{W} / m_{Z}$ and one ( $\pm 0.21$ ) which is non-cancelling. These were added in quadrature.
17 ABE 89l systematic error dominated by the uncertainty in the absolute energy scale.
18 ALBAJAR 89 result is from a total sample of $299 W \rightarrow e \nu$ events.
19 ALBAJAR 89 result is from a total sample of $67 \mathrm{~W} \rightarrow \mu \nu$ events.
20 ALBAJAR 89 result is from $W \rightarrow \tau \nu$ events.

| W/Z MASS RATIO |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: |
| Value | EVTS | DOCU | TECN | COMMENT |
| $0.88153 \pm 0.00017$ |  | 1 PDG |  |  |

-     - We do not use the following data for averages, fits, limits, etc.

| $0.8821 \pm 0.0011 \pm 0.0008$ | 28323 | 2 ABBOTT | 98 N D0 | $E_{\mathrm{cm}}^{p \bar{p}}=1.8 \mathrm{TeV}$ |  |
| :--- | :---: | :--- | :--- | :--- | :--- |
| $0.88114 \pm 0.00154 \pm 0.00252$ | 5982 | 3 ABBOTT | 98 P D0 | $E_{\mathrm{cm}}^{p \bar{p}}=1.8 \mathrm{TeV}$ |  |
| $0.8813 \pm 0.0036 \pm 0.0019$ | 156 | 4 ALITTI | 92 B | UA2 | $E_{\mathrm{cm}}^{p \bar{p}}=630 \mathrm{GeV}$ |

${ }^{1}$ PDG 16 is the PDG average using the world average $m_{W}$ and $m_{Z}$ values as quoted in this edition of Review of Particle Physics. The directly measured values of $m_{W} / m_{Z}$ are not used as their correlation with the Tevatron measured $m_{W}$ is unknown.
${ }^{2}$ ABBOTT 98N obtain this from a study of $28323 W \rightarrow e \nu_{e}$ and $3294 Z \rightarrow e^{+} e^{-}$ decays. Of this latter sample, 2179 events are used to calibrate the electron energy scale.
${ }^{3}$ ABBOTT 98P obtain this from a study of $5982 \mathrm{~W} \rightarrow e \nu_{e}$ events. The systematic error includes an uncertainty of $\pm 0.00175$ due to the electron energy scale.
${ }^{4}$ Scale error cancels in this ratio.
$m_{Z}=m_{W}$
$\frac{V A L U E(\mathrm{GeV})}{\mathbf{1 0 . 8 0 3} \pm \mathbf{0 . 0 1 5} \text { OUR AVERAGE }}$
$10.803 \pm 0.015$
$10.4 \pm 1.4 \pm 0.8$
1 PDG 16
ALBAJAR 89 UA1 $E_{\mathrm{Cm}}^{p \bar{p}}=546,630 \mathrm{GeV}$

DOCUMENT ID
TECN
COMMENT

-     - We do not use the following data for averages, fits, limits, etc.
$11.3 \pm 1.3 \pm 0.9$
ANSARI
87 UA2
$E_{\mathrm{Cm}}^{p \bar{p}}=546,630 \mathrm{GeV}$
${ }^{1}$ PDG 16 value was obtained using the world average values of $m_{Z}$ and $m_{W}$ as listed in this publication.

$$
m_{W^{+}}-m_{W^{-}}
$$

Test of CPT invariance.

| $\frac{\operatorname{VALUE}(\mathrm{GeV})}{\mathbf{- 0 . 1 9 \pm 0 . 5 8}} \frac{\text { EVTS }}{1722}$ | $\frac{\text { DOCUMENT ID }}{\text { ABE }} \frac{T E C N}{C D} \frac{\text { COMMENT }}{E_{\mathrm{cm}}^{p \bar{p}}=1.8 \mathrm{TeV}}$ |
| :--- | :--- | :--- | :--- |

## W WIDTH

The $W$ width listed here corresponds to the width parameter in a BreitWigner distribution with mass-dependent width. To obtain the world average, common systematic uncertainties between experiments are properly taken into account. The LEP-2 average $W$ width based on published results is $2.195 \pm 0.083 \mathrm{GeV}$ [SCHAEL 13A]. The combined Tevatron data yields an average $W$ width of $2.046 \pm 0.049 \mathrm{GeV}$ [FERMILAB-TM-2460-E].

OUR FIT uses these average LEP and Tevatron width values and combines them assuming no correlations.

| VALUE (GeV) | EVTS | DOCUMENT ID |  | TECN | COMMENT |
| :---: | :---: | :---: | :---: | :---: | :---: |
| $\mathbf{2 . 0 8 5} \pm 0.042$ OUR FIT |  |  |  |  |  |
| $2.028 \pm 0.072$ | 5272 | 1 ABAZOV | 09AK | D0 | $E_{\mathrm{cm}}^{p \bar{p}}=1.96 \mathrm{GeV}$ |
| $2.032 \pm 0.045 \pm 0.057$ | 6055 | 2 AALTONEN | 08B | CDF | $E_{\mathrm{cm}}^{p \bar{p}}=1.96 \mathrm{TeV}$ |
| $2.404 \pm 0.140 \pm 0.101$ | 10.3k | 3 ABDALLAH | 08A | DLPH | $E_{\mathrm{cm}}^{e e}=183-209 \mathrm{GeV}$ |
| $1.996 \pm 0.096 \pm 0.102$ | 10729 | 4 ABBIENDI | 06 | OPAL | $E_{\mathrm{cm}}^{e \mathrm{e}}=170-209 \mathrm{GeV}$ |
| $2.18 \pm 0.11 \pm 0.09$ | 9795 | 5 ACHARD | 06 | L3 | $E_{\mathrm{cm}}^{e \mathrm{e}}=172-209 \mathrm{GeV}$ |
| $2.14 \pm 0.09 \pm 0.06$ | 8717 | 6 SCHAEL | 06 | ALEP | $E_{\mathrm{cm}}^{e e}=183-209 \mathrm{GeV}$ |
| $2.23 \begin{array}{r}+0.15 \\ -0.14\end{array} \pm 0.10$ | 294 | 7 ABAZOV | 02E | D0 | $E_{\mathrm{cm}}^{p \bar{p}}=1.8 \mathrm{TeV}$ |
| $2.05 \pm 0.10 \pm 0.08$ | 662 | 8 AFFOLDER | 00M | CDF | $E_{\mathrm{cm}}^{p \bar{p}}=1.8 \mathrm{TeV}$ |

-     - We do not use the following data for averages, fits, limits, etc. - • -

| $2.152 \pm 0.066$ | 79176 | 9 <br> ABBOTT | 00B D0 | Extracted value |  |
| :--- | :--- | :--- | :--- | :--- | :--- |
| $2.064 \pm 0.060 \pm 0.059$ |  | 10 ABE | 95 W | CDF | Extracted value |
| 2.10+0.14 <br> -0.13 <br> 0.09 | 3559 | 11 ALITTI | 92 | UA2 | Extracted value |
| $2.18 \underset{-0.24}{+0.26} \pm 0.04$ |  | 12 ALBAJAR | 91 | UA1 | Extracted value |

1 ABAZOV 09AK obtain this result fitting the high-end tail (100-200 GeV) of the transverse mass spectrum in $W \rightarrow e \nu$ decays.
${ }^{2}$ AALTONEN 08B obtain this result fitting the high-end tail $(90-200 \mathrm{GeV})$ of the transverse mass spectrum in semileptonic $W \rightarrow e \nu_{e}$ and $W \rightarrow \mu \nu_{\mu}$ decays.
${ }^{3}$ ABDALLAH 08A use direct reconstruction of the kinematics of $W^{+} W^{-} \rightarrow q \bar{q} \ell \nu$ and $W^{+} W^{-} \rightarrow q \bar{q} q \bar{q}$ events. The systematic error includes $\pm 0.065 \mathrm{GeV}$ due to final state interactions.
${ }^{4}$ ABBIENDI 06 use direct reconstruction of the kinematics of $W^{+} W^{-} \rightarrow q \bar{q} \ell \nu_{\ell}$ and $W^{+} W^{-} \rightarrow q \bar{q} q \bar{q}$ events. The systematic error includes $\pm 0.003 \mathrm{GeV}$ due to the uncertainty on the LEP beam energy.
${ }^{5}$ ACHARD 06 use direct reconstruction of the kinematics of $W^{+} W^{-} \rightarrow q \bar{q} \ell \nu_{\ell}$ and $W^{+} W^{-} \rightarrow q \bar{q} q \bar{q}$ events in the C.M. energy range $189-209 \mathrm{GeV}$. The result quoted here is obtained combining this value of the width with the result obtained from a direct $W$ mass reconstruction at 172 and 183 GeV (ACCIARRI 99).
${ }^{6}$ SCHAEL 06 use direct reconstruction of the kinematics of $W^{+} W^{-} \rightarrow q \bar{q} \ell \nu_{\ell}$ and $W^{+} W^{-} \rightarrow q \bar{q} q \bar{q}$ events. The systematic error includes $\pm 0.05 \mathrm{GeV}$ due to possible effects of final state interactions in the $q \bar{q} q \bar{q}$ channel and $\pm 0.01 \mathrm{GeV}$ due to the uncertainty on the LEP beam energy.
7 ABAZOV 02E obtain this result fitting the high-end tail ( $90-200 \mathrm{GeV}$ ) of the transversemass spectrum in semileptonic $W \rightarrow e \nu_{e}$ decays.
${ }^{8}$ AFFOLDER 00M fit the high transverse mass $(100-200 \mathrm{GeV}) W \rightarrow$ e $\nu_{e}$ and $W \rightarrow$ $\mu \nu_{\mu}$ events to obtain $\Gamma(W)=2.04 \pm 0.11$ (stat) $\pm 0.09$ (syst) GeV . This is combined with the earlier CDF measurement (ABE 95C) to obtain the quoted result.
${ }^{9}$ ABBOTT 00B measure $R=10.43 \pm 0.27$ for the $W \rightarrow e \nu_{e}$ decay channel. They use the SM theoretical predictions for $\sigma(W) / \sigma(Z)$ and $\Gamma\left(W \rightarrow e \nu_{e}\right)$ and the world average for $\mathrm{B}(Z \rightarrow e e)$. The value quoted here is obtained combining this result ( $2.169 \pm 0.070$ GeV ) with that of ABBOTT 99H.
${ }^{10} \mathrm{ABE} 95 \mathrm{~W}$ measured $R=10.90 \pm 0.32 \pm 0.29$. They use $m_{W}=80.23 \pm 0.18 \mathrm{GeV}$, $\sigma(W) / \sigma(Z)=3.35 \pm 0.03, \Gamma(W \rightarrow e \nu)=225.9 \pm 0.9 \mathrm{MeV}, \Gamma\left(Z \rightarrow e^{+} e^{-}\right)=$ $83.98 \pm 0.18 \mathrm{MeV}$, and $\Gamma(Z)=2.4969 \pm 0.0038 \mathrm{GeV}$.
11 ALITTI 92 measured $R=10.4_{-0.6}^{+0.7} \pm 0.3$. The values of $\sigma(Z)$ and $\sigma(W)$ come from $O\left(\alpha_{s}^{2}\right)$ calculations using $m_{W}=80.14 \pm 0.27 \mathrm{GeV}$, and $m_{Z}=91.175 \pm 0.021 \mathrm{GeV}$ along with the corresponding value of $\sin ^{2} \theta_{W}=0.2274$. They use $\sigma(W) / \sigma(Z)=$ $3.26 \pm 0.07 \pm 0.05$ and $\Gamma(Z)=2.487 \pm 0.010 \mathrm{GeV}$.
12 ALBAJAR 91 measured $R=9.5_{-1.0}^{+1.1}$ (stat. + syst.). $\sigma(W) / \sigma(Z)$ is calculated in QCD at the parton level using $m_{W}=80.18 \pm 0.28 \mathrm{GeV}$ and $m_{Z}=91.172 \pm 0.031 \mathrm{GeV}$ along with $\sin ^{2} \theta_{W}=0.2322 \pm 0.0014$. They use $\sigma(W) / \sigma(Z)=3.23 \pm 0.05$ and $\Gamma(Z)$ $=2.498 \pm 0.020 \mathrm{GeV}$. This measurement is obtained combining both the electron and muon channels.

## $W^{+}$DECAY MODES

$W^{-}$modes are charge conjugates of the modes below.

|  | Mode | Fraction $\left(\Gamma_{i} / \Gamma\right)$ | Confidence level |  |
| :--- | :--- | :---: | :--- | :--- |
| $\Gamma_{1}$ | $\ell^{+} \nu$ | $[a]$ | $(10.86 \pm 0.09) \%$ |  |
| $\Gamma_{2}$ | $e^{+} \nu$ | $(10.71 \pm 0.16) \%$ |  |  |
| $\Gamma_{3}$ | $\mu^{+} \nu$ | $(10.63 \pm 0.15) \%$ |  |  |
| $\Gamma_{4}$ | $\tau^{+}{ }_{\nu}$ | $(11.38 \pm 0.21) \%$ |  |  |
| $\Gamma_{5}$ | hadrons | $(67.41 \pm 0.27) \%$ |  |  |
| $\Gamma_{6}$ | $\pi^{+} \gamma$ | $<7$ | $\times 10^{-6}$ | $95 \%$ |
| $\Gamma_{7}$ | $D_{s}^{+} \gamma$ | $<1.3$ | $\times 10^{-3}$ | $95 \%$ |
| $\Gamma_{8}$ | $c X$ | $(33.3 \pm 2.6) \%$ |  |  |
| $\Gamma_{9}$ | $c \bar{S}$ |  | $(31 \quad \pm 13$ | $) \%$ |

[a] $\ell$ indicates each type of lepton ( $e, \mu$, and $\tau$ ), not sum over them.
[b] This represents the width for the decay of the $W$ boson into a charged particle with momentum below detectability, $\mathrm{p}<200 \mathrm{MeV}$.

## W PARTIAL WIDTHS

## $\Gamma$ (invisible)

$\Gamma_{10}$
This represents the width for the decay of the $W$ boson into a charged particle with momentum below detectability, $\mathrm{p}<200 \mathrm{MeV}$.
$\frac{\text { VALUE }(\mathrm{MeV})}{\mathbf{3 0 + 5 2} \mathbf{+ 3 3}} \quad \frac{\text { DOCUMENT ID }}{1 \text { BARATE }} \quad$ 99। $\quad$ ALECN $\quad \frac{\text { COMMENT }}{E_{\mathrm{Cm}}^{e e}=161+172+183 \mathrm{GeV}}$

-     - We do not use the following data for averages, fits, limits, etc. - •


## 2 BARATE 99 L ALEP $\quad E_{\mathrm{Cm}}^{e e}=161+172+183 \mathrm{GeV}$

${ }^{1}$ BARATE 99I measure this quantity using the dependence of the total cross section $\sigma_{W} W$ upon a change in the total width. The fit is performed to the $W W$ measured cross sections at 161,172 , and 183 GeV . This partial width is $<139 \mathrm{MeV}$ at $95 \% \mathrm{CL}$.
2 BARATE 99L use $W$-pair production to search for effectively invisible $W$ decays, tagging with the decay of the other $W$ boson to Standard Model particles. The partial width for effectively invisible decay is $<27 \mathrm{MeV}$ at $95 \% \mathrm{CL}$.

## W BRANCHING RATIOS

Overall fits are performed to determine the branching ratios of the $W$ boson. Averages on $W \rightarrow e \nu, W \rightarrow \mu \nu$, and $W \rightarrow \tau \nu$, and their correlations are obtained by combining results from the four LEP experiments properly taking into account the common systematic uncertainties and their correlations [SCHAEL 13A]. A first fit determines the three individual leptonic braching ratios $\mathrm{B}(W \rightarrow e \nu), \mathrm{B}(W \rightarrow \mu \nu)$, and $\mathrm{B}(W \rightarrow$ $\tau \nu)$. This fit has a $\chi^{2}=6.3$ for 9 degrees of freedom. The correlation coefficients between the branching fractions are $0.14(e-\mu),-0.20(e-\tau)$, $-0.12(\mu-\tau)$. A second fit assumes lepton universality and determines the leptonic branching ratio $\operatorname{br} W \rightarrow \ell \nu$ and the hadronic branching ratio is derived as $\mathrm{B}(W \rightarrow$ hadrons $)=1-3 \mathrm{br} W \rightarrow \ell$. This fit has a $\chi^{2}=$ 15.4 for 11 degrees of freedom.
$\Gamma\left(\ell^{+} \nu\right) / \Gamma_{\text {total }}$
$\Gamma_{1} / \Gamma$
$\ell$ indicates average over $e, \mu$, and $\tau$ modes, not sum over modes.

| VALUE (units $10^{-2}$ ) | EVTS | DOCUMENT ID |  | TECN | COMMENT |
| :---: | :---: | :---: | :---: | :---: | :---: |
| $\mathbf{1 0 . 8 6} \pm 0.09$ OUR FIT |  |  |  |  |  |
| $10.86 \pm 0.12 \pm 0.08$ | 16438 | ABBIENDI | 07A | OPAL | $E_{\mathrm{cm}}^{e \mathrm{e}}$ = $=161-209 \mathrm{GeV}$ |
| $10.85 \pm 0.14 \pm 0.08$ | 13600 | ABDALLAH | 04G | DLPH | $E_{\mathrm{cm}}^{e e}=161-209 \mathrm{GeV}$ |
| $10.83 \pm 0.14 \pm 0.10$ | 11246 | ACHARD | 04J | L3 | $E_{\mathrm{cm}}^{e e}=161-209 \mathrm{GeV}$ |
| $10.96 \pm 0.12 \pm 0.05$ | 16116 | SCHAEL | 04A | ALEP | $E_{\mathrm{cm}}^{e e}=183-209 \mathrm{GeV}$ |

-     - We do not use the following data for averages, fits, limits, etc. - - -

| $11.02 \pm 0.52$ | 11858 | 1 | ABBOTT | 99 H | DO |
| :--- | :---: | :--- | :--- | :--- | :--- |
| $10.4 \pm 0.8$ | 3642 | 2 ABE | $E_{\mathrm{Cm}}^{p \bar{p}}=1.8 \mathrm{TeV}$ |  |  |
|  |  | 92 CDF | $E_{\mathrm{Cm}}^{p \bar{p}}=1.8 \mathrm{TeV}$ |  |  |

${ }^{1}$ ABBOTT 99н measure $R \equiv\left[\sigma_{W} \mathrm{~B}\left(W \rightarrow \ell \nu_{\ell}\right)\right] /\left[\sigma_{Z} \mathrm{~B}(Z \rightarrow \ell \ell)\right]=10.90 \pm 0.52$ combining electron and muon channels. They use $M_{W}=80.39 \pm 0.06 \mathrm{GeV}$ and the SM theoretical predictions for $\sigma(W) / \sigma(Z)$ and $\mathrm{B}(Z \rightarrow \ell \ell)$.
$21216 \pm 38_{-31}^{+27} W \rightarrow \mu \nu$ events from ABE 92 and $2426 W \rightarrow e \nu$ events of ABE 91C. ABE 92I give the inverse quantity as $9.6 \pm 0.7$ and we have inverted.

|  |  |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| VALUE (units $10^{-2}$ ) | EVTS | DOCUMENT ID |  | TECN | COMMENT |  |
| 10.71 $\pm 0.16$ OUR FIT |  |  |  |  |  |  |
| $10.71 \pm 0.25 \pm 0.11$ | 2374 | ABBIENDI | 07A | OPAL | $E_{\mathrm{cm}}^{e}{ }_{\mathrm{m}}^{e}=161-209$ | GeV |
| $10.55 \pm 0.31 \pm 0.14$ | 1804 | ABDALLAH | 04G | DLPH | $E_{\mathrm{cm}}^{\mathrm{ee}}=161-209$ | GeV |
| $10.78 \pm 0.29 \pm 0.13$ | 1576 | ACHARD | 04」 | L3 | $E_{c m}^{e e}=161-209$ | GeV |
| $10.78 \pm 0.27 \pm 0.10$ | 2142 | SCHAEL | 04A | ALEP | $E_{c m}^{e e}=183-209$ |  |

-     - We do not use the following data for averages, fits, limits, etc. - • -
$10.61 \pm 0.28$
${ }^{1}$ ABAZOV
04D TEVA $E_{\mathrm{cm}}^{p \bar{p}}=1.8 \mathrm{TeV}$

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${ }^{1}$ ABAZOV 04D take into account all correlations to properly combine the CDF (ABE 95W) and $D \varnothing$ ( $A B B O T T$ OOB) measurements of the ratio $R$ in the electron channel. The ratio R is defined as $\left[\sigma_{W} \cdot \mathrm{~B}\left(W \rightarrow e \nu_{e}\right)\right] /\left[\sigma_{Z} \cdot \mathrm{~B}(Z \rightarrow e e)\right]$. The combination gives $\mathrm{R}^{\text {Tevatron }}=10.59 \pm 0.23 . \sigma_{W} / \sigma_{Z}$ is calculated at next-to-next-to-leading order $(3.360 \pm 0.051)$. The branching fraction $\mathrm{B}(Z \rightarrow e e)$ is taken from this Review as $(3.363 \pm 0.004) \%$.

| $\Gamma\left(\boldsymbol{\mu}^{+} \boldsymbol{\nu}\right) / \Gamma_{\text {total }}$ |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: |
| VALUE (units $10^{-2}$ ) | EVTS | DOCUMENT ID |  | TECN | COMMENT |
| $10.63 \pm 0.15$ OUR FIT |  |  |  |  |  |
| $10.78 \pm 0.24 \pm 0.10$ | 2397 | ABBIENDI | 07A | OPAL | $E_{\mathrm{cm}}^{e \mathrm{em}}=161-209 \mathrm{GeV}$ |
| $10.65 \pm 0.26 \pm 0.08$ | 1998 | ABDALLAH | 04G | DLPH | $E_{\mathrm{Cm}}^{e \mathrm{e}}=161-209 \mathrm{GeV}$ |
| $10.03 \pm 0.29 \pm 0.12$ | 1423 | ACHARD | 04J | L3 | $E_{\mathrm{Cm}}^{e \mathrm{e}}=161-209 \mathrm{GeV}$ |
| $10.87 \pm 0.25 \pm 0.08$ | 2216 | SCHAEL | 04A | ALEP | $E_{\mathrm{Cm}}^{e \mathrm{e}}=183-209 \mathrm{GeV}$ |
| $\Gamma\left(\tau^{+} \nu\right) / \Gamma_{\text {total }}$ |  |  |  |  |  |
| VALUE (units $10^{-2}$ ) | EVTS | DOCUMENT ID |  | TECN | COMMENT |
| $11.38 \pm 0.21$ OUR FIT |  |  |  |  |  |
| $11.14 \pm 0.31 \pm 0.17$ | 2177 | ABBIENDI | 07A | OPAL | $E_{\mathrm{cm}}^{e \mathrm{e}}=161-209 \mathrm{GeV}$ |
| $11.46 \pm 0.39 \pm 0.19$ | 2034 | ABDALLAH | 04G | DLPH | $E_{\mathrm{cm}}^{e \mathrm{e}}=161-209 \mathrm{GeV}$ |
| $11.89 \pm 0.40 \pm 0.20$ | 1375 | ACHARD | 04J | L3 | $E_{\mathrm{Cm}}^{e \mathrm{e}}=161-209 \mathrm{GeV}$ |
| $11.25 \pm 0.32 \pm 0.20$ | 2070 | SCHAEL | 04A | ALEP | $E_{\mathrm{cm}}^{e e}=183-209 \mathrm{GeV}$ |

## $\Gamma($ hadrons $) / \Gamma_{\text {total }}$

OUR FIT value is obtained by a fit to the lepton branching ratio data assuming lepton universality.

| VALUE (units $10^{-2}$ ) | EVTS | DOCUMENT ID |  | TECN | COMMENT |
| :---: | :---: | :---: | :---: | :---: | :---: |
| $\mathbf{6 7 . 4 1} \pm 0.27$ OUR FIT |  |  |  |  |  |
| $67.41 \pm 0.37 \pm 0.23$ | 16438 | ABBIENDI | 07A | OPAL | $E_{\mathrm{cm}}^{e \mathrm{e}}$ ¢ $=161-209 \mathrm{GeV}$ |
| $67.45 \pm 0.41 \pm 0.24$ | 13600 | ABDALLAH | 04G | DLPH |  |
| $67.50 \pm 0.42 \pm 0.30$ | 11246 | ACHARD | 04J | L3 | $E_{\mathrm{cm}}^{e}{ }_{\text {el }}^{\text {en }}=161-209 \mathrm{GeV}$ |
| $67.13 \pm 0.37 \pm 0.15$ | 16116 | SCHAEL | 04A | ALEP | $E_{\mathrm{Cm}}^{e e}=183-209 \mathrm{GeV}$ |
| $\Gamma\left(\mu^{+} \nu\right) / \Gamma\left(e^{+} \nu\right)$ |  |  |  |  | $\Gamma_{3} / \Gamma_{2}$ |
| VALUE | EVTS | DOCUMENT ID |  | TECN | COMMENT |
| $\mathbf{0 . 9 9 1} \pm 0.018$ OUR AVERAGE |  |  |  |  |  |
| $0.993 \pm 0.019$ |  | SCHAEL | 13A | LEP | $E_{\mathrm{cm}}^{e e}=130-209 \mathrm{GeV}$ |
| $0.89 \pm 0.10$ | 13k | ${ }^{1}$ ABACHI | 95D | D0 | $E_{\mathrm{cm}}^{p \bar{p}}=1.8 \mathrm{TeV}$ |
| $1.02 \pm 0.08$ | 1216 | 2 ABE | 921 | CDF | $E_{\mathrm{cm}}^{p \bar{p}}=1.8 \mathrm{TeV}$ |
| $1.00 \pm 0.14 \pm 0.08$ | 67 | ALBAJAR | 89 | UA1 | $E_{\mathrm{cm}}^{p \bar{P}}=546,630 \mathrm{GeV}$ |

-     - We do not use the following data for averages, fits, limits, etc. • • -
$1.24{ }_{-0.4}^{+0.6}$
14 ARNISON
84D UA1
Repl. by ALBAJAR 89
${ }^{1}$ ABACHI 95D obtain this result from the measured $\sigma_{W} \mathrm{~B}(W \rightarrow \mu \nu)=2.09 \pm 0.23 \pm$ 0.11 nb and $\sigma_{W} \mathrm{~B}(W \rightarrow e \nu)=2.36 \pm 0.07 \pm 0.13 \mathrm{nb}$ in which the first error is the combined statistical and systematic uncertainty, the second reflects the uncertainty in the luminosity.
${ }^{2} \mathrm{ABE} 92$ I obtain $\sigma_{W} \mathrm{~B}(W \rightarrow \mu \nu)=2.21 \pm 0.07 \pm 0.21$ and combine with $\mathrm{ABE} 91 \mathrm{C} \sigma_{W}$ $\mathrm{B}((W \rightarrow e \nu))$ to give a ratio of the couplings from which we derive this measurement.

-     - We do not use the following data for averages, fits, limits, etc. - - -

| $0.995 \pm 0.112 \pm 0.083$ | 198 | ALITTI | 91C | UA2 | Repl. by ALITTI 92F |
| :--- | ---: | :--- | :--- | :--- | :--- |
| $1.02 \pm 0.20 \pm 0.10$ | 32 | ALBAJAR | 87 | UA1 | Repl. by ALBAJAR 89 |

${ }^{1}$ ABBOTT 00D measure $\sigma_{W} \times \mathrm{B}\left(W \rightarrow \tau \nu_{\tau}\right)=2.22 \pm 0.09 \pm 0.10 \pm 0.10 \mathrm{nb}$. Using the ABBOTT 00B result $\sigma_{W} \times \mathrm{B}\left(W \rightarrow e \nu_{e}\right)=2.31 \pm 0.01 \pm 0.05 \pm 0.10 \mathrm{nb}$, they quote the ratio of the couplings from which we derive this measurement.
${ }^{2}$ ABE 92E use two procedures for selecting $W \rightarrow \tau \nu_{\tau}$ events. The missing $\mathrm{E}_{T}$ trigger leads to $132 \pm 14 \pm 8$ events and the $\tau$ trigger to $47 \pm 9 \pm 4$ events. Proper statistical and systematic correlations are taken into account to arrive at $\sigma \mathrm{B}(W \rightarrow \tau \nu)=2.05 \pm 0.27$ nb. Combined with ABE 91 C result on $\sigma \mathrm{B}(W \rightarrow e \nu)$, ABE 92 E quote a ratio of the couplings from which we derive this measurement.
3 This measurement is derived by us from the ratio of the couplings of ALITTI 92F.
$\Gamma\left(\tau^{+} \nu\right) / \Gamma\left(\mu^{+} \nu\right)$
VALUE
$1.070 \pm 0.026$
$\Gamma\left(\pi^{+} \gamma\right) / \Gamma\left(e^{+} \nu\right)$

| VALUE | CL\% | DOCUMENT |
| :---: | :---: | :---: |
| $<6.4 \times 10^{-5}$ | 95 | AALTONE |
| $<7 \times 10^{-4}$ | 95 | ABE |
| $<4.9 \times 10^{-3}$ | 95 | 1 ALITTI |
| $<58 \times 10^{-3}$ | 95 | 2 ALBAJAR |
| $1^{1}$ ALITTI 92D 2 ALBAJAR | . $8 \times 10$ $<0.0$ | $\begin{aligned} & 3 \text { at } 90 \% \text { CL. } \\ & \text { at } 90 \% \mathrm{CL} \text {. } \end{aligned}$ |

$\frac{\text { VALUE }}{<\mathbf{1 . 2 \times 1 0 ^ { - 2 }}} \frac{C L \%}{95} \quad \frac{\text { DOCUMENT ID }}{\text { ABE }} \frac{\text { TECN }}{C} \frac{\text { COMMENT }}{E_{\mathrm{cm}}^{p \bar{p}}=1.8 \mathrm{TeV}}$
$\Gamma(c X) / \Gamma$ (hadrons)

| DOCUMENT ID |  | TECN | COMMENT | $\Gamma_{8} / \Gamma_{5}$ |
| :---: | :---: | :---: | :---: | :---: |
|  |  |  |  |  |
| 1 ABBIENDI | 00V | OPAL | $E_{\mathrm{cm}}^{e \mathrm{e}} \mathrm{m}=183$ | 9 GeV |
| 2 BARATE | 99M | ALEP | $E_{\mathrm{cm}}^{e \mathrm{e}} \mathrm{m}=172$ | 3 GeV |

$\frac{V A L U E}{0.49} \pm 0.04$ OUR AVE $\frac{E V T S}{R A G E}$

| $0.481 \pm 0.042 \pm 0.032$ | 3005 | 1 ABBIENDI | 00 V OPAL | $E_{\mathrm{Cm}}^{e e}=183+189 \mathrm{GeV}$ |
| :--- | ---: | :--- | :--- | :--- |
| $0.51 \pm 0.05 \pm 0.03$ | 746 | 2 BARATE | 99 m | ALEP |
| $E_{\mathrm{Cm}}^{e e}=172+183 \mathrm{GeV}$ |  |  |  |  |

${ }^{1}$ ABBIENDI $00 V$ tag $W \rightarrow c X$ decays using measured jet properties, lifetime information, and leptons produced in charm decays. From this result, and using the additional measurements of $\Gamma(W)$ and $\mathrm{B}(W \rightarrow$ hadrons $),\left|V_{c s}\right|$ is determined to be $0.969 \pm 0.045 \pm 0.036$.
2 BARATE 99M tag $c$ jets using a neural network algorithm. From this measurement $\left|V_{c s}\right|$ is determined to be $1.00 \pm 0.11 \pm 0.07$.
$\mathrm{R}_{c s}=\Gamma(c \bar{s}) / \Gamma$ (hadrons)
VALUE
$0.46+0.18 \pm 0.07$

DOCUMENT ID TECN
1 ABREU 98 N DLPH
$\Gamma_{9} / \Gamma_{5}$
${ }^{1}$ ABREU 98 N tag $c$ and $s$ jets by identifying a charged kaon as the highest momentum particle in a hadronic jet. They also use a lifetime tag to independently identify a $c$ jet, based on the impact parameter distribution of charged particles in a jet. From this measurement $\left|V_{c s}\right|$ is determined to be $0.94_{-0.26}^{+0.32} \pm 0.13$.

## AVERAGE PARTICLE MULTIPLICITIES IN HADRONIC $W$ DECAY

Summed over particle and antiparticle, when appropriate.
$\left\langle N_{\pi^{ \pm}}\right\rangle$
VALUE
$15.70 \pm 0.35$
$1 \frac{\text { DOCUMENT ID }}{1 \text { ABREU,P }} \quad$ OOF $\quad \frac{\text { TECN }}{\text { DLPH }} \frac{\text { COMMENT }}{E_{\mathrm{Cm}}^{e \mathrm{e}}=189 \mathrm{GeV}}$
${ }^{1}$ ABREU,P 00F measure $\left\langle N_{\pi^{ \pm}}\right\rangle=31.65 \pm 0.48 \pm 0.76$ and $15.51 \pm 0.38 \pm 0.40$ in the fully hadronic and semileptonic final states respectively. The value quoted is a weighted average without assuming any correlations.
$\left\langle\boldsymbol{N}_{\boldsymbol{K} \pm}\right\rangle$
VALUE DOCUMENT ID TECN COMMENT
$2.20 \pm 0.19$
1 ABREU,P 00F DLPH $E_{\mathrm{cm}}^{e e}=189 \mathrm{GeV}$
${ }^{1}$ ABREU,P 00F measure $\left\langle N_{K}{ }^{ \pm}\right\rangle=4.38 \pm 0.42 \pm 0.12$ and $2.23 \pm 0.32 \pm 0.17$ in the fully hadronic and semileptonic final states respectively. The value quoted is a weighted average without assuming any correlations.
$\left\langle N_{p}\right\rangle$

${ }^{1}$ ABREU,P 00F measure $\left\langle N_{p}\right\rangle=1.82 \pm 0.29 \pm 0.16$ and $0.94 \pm 0.23 \pm 0.06$ in the fully hadronic and semileptonic final states respectively. The value quoted is a weighted average without assuming any correlations.

## $\left\langle N_{\text {charged }}\right\rangle$

VALUE
$19.39 \pm 0.08$ OUR AVERAGE
$19.38 \pm 0.05 \pm 0.08$
$19.44 \pm 0.17$
$19.3 \pm 0.3 \pm 0.3$
$19.23 \pm 0.74$

| DOCUMENT ID |  | TECN | COMMENT |
| :---: | :---: | :---: | :---: |
| 1 ABBIENDI | 06A | OPAL | $E_{\mathrm{cm}}^{e e}=189-209 \mathrm{GeV}$ |
| 2 ABREU,P | 00F | DLPH | $E_{\mathrm{Cm}}^{e e}=183+189 \mathrm{GeV}$ |
| 3 ABBIENDI | 99N | OPAL | $E_{\text {cme }}^{e e}{ }_{\text {em }}^{e}=183 \mathrm{GeV}$ |
| 4 ABREU | 98C | DLPH | $E_{\mathrm{cm}}^{e e}=172 \mathrm{GeV}$ |

${ }^{1}$ ABBIENDI 06A measure $\left\langle N_{\text {charged }}\right\rangle=38.74 \pm 0.12 \pm 0.26$ when both $W$ bosons decay hadronically and $\left\langle N_{\text {charged }}\right\rangle=19.39 \pm 0.11 \pm 0.09$ when one $W$ boson decays semileptonically. The value quoted here is obtained under the assumption that there is no color reconnection between $W$ bosons; the value is a weighted average taking into account correlations in the systematic uncertainties.
${ }^{2}$ ABREU,P 00F measure $\left\langle N_{\text {charged }}\right\rangle=39.12 \pm 0.33 \pm 0.36$ and $38.11 \pm 0.57 \pm 0.44$ in the fully hadronic final states at 189 and 183 GeV respectively, and $\left\langle N_{\text {charged }}\right\rangle=$
$19.49 \pm 0.31 \pm 0.27$ and $19.78 \pm 0.49 \pm 0.43$ in the semileptonic final states. The value quoted is a weighted average without assuming any correlations.
${ }^{3}$ ABBIENDI 99 N use the final states $W^{+} W^{-} \rightarrow q \bar{q} \ell \bar{\nu}_{\ell}$ to derive this value.
${ }^{4}$ ABREU 98C combine results from both the fully hadronic as well semileptonic $W W$ final states after demonstrating that the $W$ decay charged multiplicity is independent of the topology within errors.

## TRIPLE GAUGE COUPLINGS (TGC'S)

## A REVIEW GOES HERE - Check our WWW List of Reviews

$g_{1}^{Z}$
OUR FIT below is taken from [SCHAEL 13A].

| VALUE | EVTS | DOCUMENT ID |  | TECN | COMMENT |
| :---: | :---: | :---: | :---: | :---: | :---: |
| $0.984 \underset{=0.020}{+0.018}$ OUR FIT |  |  |  |  |  |
| $\begin{array}{r} 0.975_{-0.030}^{+0.033} \end{array}$ | 7872 | 1 ABDALLAH | 10 | DLPH | $E_{\mathrm{Cm}}^{e e}=189-209 \mathrm{GeV}$ |
| $1.001 \pm 0.027 \pm 0.013$ | 9310 | 2 SCHAEL | 05A | ALEP | $E_{\mathrm{Cm}}^{e e}=183-209 \mathrm{GeV}$ |
| ${ }_{0.987}^{+0.034}{ }_{-0.033}$ | 9800 | 3 ABBIENDI | 04D | OPAL | $E_{\mathrm{cm}}^{e e}=183-209 \mathrm{GeV}$ |
| $0.966_{-0.032}^{+0.034} \pm 0.015$ | 8325 | ${ }^{4}$ ACHARD | 04D | L3 | $E_{\mathrm{cm}}^{e \mathrm{em}}=161-209 \mathrm{GeV}$ |

-     - We do not use the following data for averages, fits, limits, etc.


Each parameter is determined from a single-parameter fit in which the other parameters assume their Standard Model values.
${ }^{3}$ ABBIENDI 04D combine results from $W^{+} W^{-}$in all decay channels. Only CP-conserving couplings are considered and each parameter is determined from a single-parameter fit in which the other parameters assume their Standard Model values. The $95 \%$ confidence interval is $0.923<g_{1}^{Z}<1.054$.
${ }^{4}$ ACHARD 04D study $W W$-pair production, single- $W$ production and single-photon production with missing energy from 189 to 209 GeV . The result quoted here is obtained from the $W W$-pair production sample including data from 161 to 183 GeV , ACCIARRI 99Q. Each parameter is determined from a single-parameter fit in which the other parameters assume their Standard Model values.
${ }^{5}$ AAD 14 Y determine the electroweak $Z$-dijet cross section in $8 \mathrm{TeV} p p$ collisions. $Z \rightarrow$ $e e$ and $Z \rightarrow \mu \mu$ decays are selected with the di-lepton $p_{T}>20 \mathrm{GeV}$ and mass in the $81-101 \mathrm{GeV}$ range. Minimum two jets are required with $p_{T}>55$ and 45 GeV and no additional jets with $p_{T}>25 \mathrm{GeV}$ in the rapidity interval between them. The normalized $p_{T}$ balance between the $Z$ and the two jets is required to be $<0.15$. This leads to a selection of 900 events with dijet mass $>1 \mathrm{TeV}$. The number of signal and background events expected is 261 and 592 respectively. A Poisson likelihood method is used on an event by event basis to obtain the $95 \%$ CL limit $0.5<g_{1}^{Z}<1.26$ for a form factor value $\Lambda=\infty$.
${ }^{6}$ AAD 13AL study $W W$ production in $p p$ collisions and select $1325 W W$ candidates in decay modes with electrons or muons with an expected background of $369 \pm 61$ events. Assuming the LEP formulation and setting the form-factor $\Lambda=$ infinity, a fit to the transverse momentum distribution of the leading charged lepton, leads to a $95 \%$ C.L. range of $0.961<g_{1}^{Z}<1.052$. Supersedes AAD 12AC.
${ }^{7}$ CHATRCHYAN 13BF determine the $W^{+} W^{-}$production cross section using unlike sign di-lepton (e or $\mu$ ) events with high $\not p_{T}$. The leptons have $p_{T}>20 \mathrm{GeV} / \mathrm{c}$ and are isolated. 1134 candidate events are observed with an expected SM background of $247 \pm$ 34. The $p_{T}$ distribution of the leading lepton is fitted to obtain $95 \%$ C.L. limits of 0.905 $\leq g_{1}^{Z} \leq 1.095$.
${ }^{8}$ AAD 12 CD study $W Z$ production in $p p$ collisions and select $317 W Z$ candidates in three $\ell \nu$ decay modes with an expected background of $68.0 \pm 10.0$ events. The resulting $95 \%$ C.L. range is: $0.943<g_{1}^{Z}<1.093$. Supersedes AAD 12 V .
${ }^{9}$ AALTONEN 12AC study $W Z$ production in $p \bar{p}$ collisions and select $63 W Z$ candidates in three $\ell \nu$ decay modes with an expected background of $7.9 \pm 1.0$ events. Based on the cross section and shape of the $Z$ transverse momentum spectrum, the following $95 \%$ C.L. range is reported: $0.92<g_{1}^{Z}<1.20$ for a form factor of $\Lambda=2 \mathrm{TeV}$.

10 ABAZOV 12AG combine new results with already published results on $W \gamma, W W$ and $W Z$ production in order to determine the couplings with increased precision, superseding ABAZOV 08R, ABAZOV 11AC, ABAZOV 09AJ, ABAZOV 09AD. The 68\% C.L. result for a formfactor cutoff of $\Lambda=2 \mathrm{TeV}$ is $g_{1}^{Z}=1.022_{-0.030}^{+0.032}$.
11 ABAZOV 11 study the $p \bar{p} \rightarrow 3 \ell \nu$ process arising in $W Z$ production. They observe $34 W Z$ candidates with an estimated background of 6 events. An analysis of the $p_{T}$ spectrum of the $Z$ boson leads to a $95 \%$ C.L. limit of $0.944<g_{1}^{Z}<1.154$, for a form factor $\Lambda=2 \mathrm{TeV}$.
12 AALTONEN 10K study $p \bar{p} \rightarrow W^{+} W^{-}$with $W \rightarrow e / \mu \nu$. The $p_{T}$ of the leading (second) lepton is required to be $>20(10) \mathrm{GeV}$. The final number of events selected is 654 of which $320 \pm 47$ are estimated to be background. The $95 \%$ C.L. interval is 0.76 $<g_{1}^{Z}<1.34$ for $\Lambda=1.5 \mathrm{TeV}$ and $0.78<g_{1}^{Z}<1.30$ for $\Lambda=2 \mathrm{TeV}$.
13 ABAZOV 09AD study the $p \bar{p} \rightarrow \ell \nu 2$ jet process arising in $W W$ and $W Z$ production. They select $12,473(14,392)$ events in the electron (muon) channel with an expected di-boson signal of 436 (527) events. The results on the anomalous couplings are derived from an analysis of the $p_{T}$ spectrum of the $2-$ jet system and quoted at $68 \%$ C.L. and
for a form factor of 2 TeV . This measurement is not used for obtaining the mean as it is for a specific form factor. The $95 \%$ confidence interval is $0.88<g_{1}^{Z}<1.20$.
14 ABAZOV 09AJ study the $p \bar{p} \rightarrow 2 \ell 2 \nu$ process arising in $W W$ production. They select 100 events with an expected $W W$ signal of 65 events. An analysis of the $p_{T}$ spectrum of the two charged leptons leads to $95 \%$ C.L. limits of $0.86<g_{1}^{Z}<1.3$, for a form factor $\Lambda=2 \mathrm{TeV}$.
15 ABDALLAH 08C determine this triple gauge coupling from the measurement of the spin density matrix elements in $e^{+} e^{-} \rightarrow W^{+} W^{-} \rightarrow(q q)(\ell \nu)$, where $\ell=e$ or $\mu$. Values of all other couplings are fixed to their standard model values.
16 ABAZOV $07 Z$ set limits on anomalous TGCs using the measured cross section and $p_{T}(Z)$ distribution in $W Z$ production with both the $W$ and the $Z$ decaying leptonically into electrons and muons. Setting the other couplings to their standard model values, the $95 \%$ C.L. limit for a form factor scale $\Lambda=2 \mathrm{TeV}$ is $0.86<g_{1}^{Z}<1.35$.
17 ABAZOV 05S study $\bar{p} p \rightarrow W Z$ production with a subsequent trilepton decay to $\ell \nu \ell^{\prime} \bar{\ell}^{\prime}$ ( $\ell$ and $\ell^{\prime}=e$ or $\mu$ ). Three events (estimated background $0.71 \pm 0.08$ events) with $W Z$ decay characteristics are observed from which they derive limits on the anomalous $W W Z$ couplings. The $95 \%$ CL limit for a form factor scale $\Lambda=1.5 \mathrm{TeV}$ is $0.51<g_{1}^{Z}<$ 1.66, fixing $\lambda_{Z}$ and $\kappa_{Z}$ to their Standard Model values.

18 ABREU 01I combine results from $e^{+} e^{-}$interactions at 189 GeV leading to $W^{+} W^{-}$ and $W e \nu_{e}$ final states with results from ABREU 99 L at 183 GeV . The $95 \%$ confidence interval is $0.84<g_{1}^{Z}<1.13$.
${ }^{19}$ ABBOTT 99I perform a simultaneous fit to the $W \gamma, W W \rightarrow$ dilepton, $W W / W Z \rightarrow$ $e \nu j j, W W / W Z \rightarrow \mu \nu j j$, and $W Z \rightarrow$ trilepton data samples. For $\Lambda=2.0 \mathrm{TeV}$, the $95 \%$ CL limits are $0.63<g_{1}^{Z}<1.57$, fixing $\lambda_{Z}$ and $\kappa_{Z}$ to their Standard Model values, and assuming Standard Model values for the $W W \gamma$ couplings.
$\kappa_{\gamma}$
OUR FIT below is taken from [SCHAEL 13A].

| VALUE | EVTS | DOCUMENT ID |  | TECN | COMMENT |
| :---: | :---: | :---: | :---: | :---: | :---: |
| 0.982 $\pm \mathbf{\pm} 0.042$ OUR FIT |  |  |  |  |  |
| $\begin{gathered} 1.024_{-0.081}^{+0.077} \end{gathered}$ | 7872 | 1 ABDALLAH | 10 | DLPH | $E_{\mathrm{cm}}^{e e}=189-209 \mathrm{GeV}$ |
| $0.971 \pm 0.055 \pm 0.030$ | 10689 | ${ }^{2}$ SCHAEL | 05A | ALEP | $E_{\mathrm{cm}}^{e e}=183-209 \mathrm{GeV}$ |
| $0.88{ }_{-0.08}^{+0.09}$ | 9800 | 3 ABBIENDI | 04D | OPAL | $E_{\mathrm{cm}}^{e e}=183-209 \mathrm{GeV}$ |
| $1.013_{-0.064}^{+0.067} \pm 0.026$ | 10575 | 4 ACHARD | 04D | L3 | $E_{\mathrm{Cm}}^{e e}=161-209 \mathrm{GeV}$ |

-     - We do not use the following data for averages, fits, limits, etc. - - -


${ }^{1}$ ABDALLAH 10 use data on the final states $e^{+} e^{-} \rightarrow j j \ell \nu, j j j j, j j X, \ell X$, at center-of-mass energies between $189-209 \mathrm{GeV}$ at LEP2, where $j=$ jet, $\ell=$ lepton, and $X$ represents missing momentum. The fit is carried out keeping all other parameters fixed at their SM values.
${ }^{2}$ SCHAEL 05A study single-photon, single- $W$, and $W W$-pair production from 183 to 209 GeV . Each parameter is determined from a single-parameter fit in which the other parameters assume their Standard Model values.
3 ABBIENDI 04D combine results from $W^{+} W^{-}$in all decay channels. Only $C P$-conserving couplings are considered and each parameter is determined from a single-parameter fit in which the other parameters assume their Standard Model values. The 95\% confidence interval is $0.73<\kappa_{\gamma}<1.07$.
${ }^{4}$ ACHARD 04D study $W W$-pair production, single- $W$ production and single-photon production with missing energy from 189 to 209 GeV . The result quoted here is obtained including data from 161 to 183 GeV , ACCIARRI 99Q. Each parameter is determined from a single-parameter fit in which the other parameters assume their Standard Model values.
${ }^{5}$ CHATRCHYAN 14AB measure $W \gamma$ production cross section for $p_{T}^{\gamma}>15 \mathrm{GeV}$ and $\mathrm{R}(\ell \gamma)$ $>0.7$, which is the separation between the $\gamma$ and the final state charged lepton (e or $\mu$ ) in the azimuthal angle-pseudorapidity $(\phi-\eta)$ plane. After background subtraction the number of $e \nu \gamma$ and $\mu \nu \gamma$ events is determined to be $3200 \pm 325$ and $4970 \pm 543$ respectively, compatible with expectations from the SM. This leads to a $95 \%$ CL limit of $0.62<\kappa_{\gamma}<1.29$, assuming other parameters have SM values.
${ }^{6}$ AAD 13AN study $W \gamma$ production in $p p$ collisions. In events with no additional jet, 4449 (6578) W decays to electron (muon) are selected, with an expected background of $1662 \pm 262(2538 \pm 362)$ events. Analysing the photon $p_{T}$ spectrum above 100 GeV yields a $95 \%$ C.L. limit of $0.59<\kappa_{\gamma}<1.46$. Supersedes AAD 12BX.
${ }^{7}$ CHATRCHYAN $13 B F$ determine the $W^{+} W^{-}$production cross section using unlike sign di-lepton (e or $\mu$ ) events with high $\not p T_{T}$. The leptons have $p_{T}>20 \mathrm{GeV} / \mathrm{c}$ and are isolated. 1134 candidate events are observed with an expected SM background of $247 \pm$ 34. The $p_{T}$ distribution of the leading lepton is fitted to obtain $95 \%$ C.L. limits of 0.79 $\leq k_{\gamma} \leq 1.22$.
${ }^{8}$ ABAZOV 12AG combine new results with already published results on $W \gamma, W W$ and $W Z$ production in order to determine the couplings with increased precision, superseding ABAZOV 08R, ABAZOV 11AC, ABAZOV 09AJ, ABAZOV 09AD. The 68\% C.L. result for a formfactor cutoff of $\Lambda=2 \mathrm{TeV}$ is $\kappa_{\gamma}=1.048_{-0.105}^{+0.106}$.
${ }^{9}$ ABAZOV 11AC study $W \gamma$ production in $p \bar{p}$ collisions at 1.96 TeV , with the $W$ decay products containing an electron or a muon. They select 196 (363) events in the electron (muon) mode, with a SM expectation of 190 (372) events. A likelihood fit to the photon $E_{T}$ spectrum above 15 GeV yields at $95 \%$ C.L. the result: $0.6<\kappa_{\gamma}<1.4$ for a formfactor $\Lambda=2 \mathrm{TeV}$.
${ }^{10}$ CHATRCHYAN 11 M study $W \gamma$ production in $p p$ collisions at $\sqrt{s}=7 \mathrm{TeV}$ using $36 \mathrm{pb}^{-1}$ $p p$ data with the $W$ decaying to electron and muon. The total cross section is measured for photon transverse energy $E_{T}^{\gamma}>10 \mathrm{GeV}$ and spatial separation from charged leptons in the plane of pseudo rapidity and azimuthal angle $\Delta R(\ell, \gamma)>0.7$. The number of candidate (background) events is $452(228 \pm 21)$ for the electron channel and 520 $(277 \pm 25)$ for the muon channel. Setting other couplings to their standard model value, they derive a $95 \% \mathrm{CL}$ limit of $-0.11<\kappa_{\gamma}<2.04$.
11 AALTONEN 10K study $p \bar{p} \rightarrow W^{+} W^{-}$with $W \rightarrow e / \mu \nu$. The $p_{T}$ of the leading (second) lepton is required to be $>20(10) \mathrm{GeV}$. The final number of events selected is 654 of which $320 \pm 47$ are estimated to be background. The $95 \%$ C.L. interval is 0.37 $<\kappa_{\gamma}<1.72$ for $\Lambda=1.5 \mathrm{TeV}$ and $0.43<\kappa_{\gamma}<1.65$ for $\Lambda=2 \mathrm{TeV}$.
12 AARON 09B study single- $W$ production in ep collisions at 0.3 TeV C.M. energy. They select $53 \mathrm{~W} \rightarrow e / \mu$ events with a standard model expectation of $54.1 \pm 7.4$ events. Fitting the transverse momentum spectrum of the hadronic recoil system they obtain a $95 \%$ C.L. limit of $-3.7<\kappa_{\gamma}<-1.5$ or $0.3<\kappa_{\gamma}<1.5$, where the ambiguity is due to the quadratic dependence of the cross section to the coupling parameter.
13 ABAZOV 09AD study the $p \bar{p} \rightarrow \ell \nu 2$ jet process arising in $W W$ and $W Z$ production. They select $12,473(14,392)$ events in the electron (muon) channel with an expected di-boson signal of 436 (527) events. The results on the anomalous couplings are derived from an analysis of the $p_{T}$ spectrum of the 2 -jet system and quoted at $68 \%$ C.L. and for a form factor of 2 TeV . This measurement is not used for obtaining the mean as it is for a specific form factor. The $95 \%$ confidence interval is $0.56<\kappa_{\gamma}<1.55$.
14 ABAZOV 09AJ study the $p \bar{p} \rightarrow 2 \ell 2 \nu$ process arising in $W W$ production. They select 100 events with an expected $W W$ signal of 65 events. An analysis of the $p_{T}$ spectrum of the two charged leptons leads to $95 \%$ C.L. limits of $0.46<\kappa_{\gamma}<1.83$, for a form factor $\Lambda=2 \mathrm{TeV}$.
15 ABAZOV 08R use $0.7 \mathrm{fb}^{-1} p \bar{p}$ data at $\sqrt{s}=1.96 \mathrm{TeV}$ to select $263 W \gamma+X$ events, of which 187 constitute signal, with the $W$ decaying into an electron or a muon, which is required to be well separated from a photon with $E_{T}>9 \mathrm{GeV}$. A likelihood fit to the photon $E_{T}$ spectrum yields a $95 \%$ CL limit $0.49<\kappa_{\gamma}<1.51$ with other couplings fixed to their Standard Model values.
16 ABDALLAH 08C determine this triple gauge coupling from the measurement of the spin density matrix elements in $e^{+} e^{-} \rightarrow W^{+} W^{-} \rightarrow(q q)(\ell \nu)$, where $\ell=e$ or $\mu$. Values of all other couplings are fixed to their standard model values.
17 AALTONEN 07L set limits on anomalous TGCs using the $p_{T}(W)$ distribution in $W W$ and $W Z$ production with the $W$ decaying to an electron or muon and the $Z$ to 2 jets. Setting other couplings to their standard model value, the $95 \%$ C.L. limits are 0.54 $<\kappa_{\gamma}<1.39$ for a form factor scale $\Lambda=1.5 \mathrm{TeV}$.
18 ABAZOV 06 H study $\bar{p} p \rightarrow W W$ production with a subsequent decay $W W \rightarrow$ $e^{+} \nu_{e} e^{-} \bar{\nu}_{e}, W W \rightarrow e^{ \pm} \nu_{e} \mu^{\mp} \nu_{\mu}$ or $W W \rightarrow \mu^{+} \nu_{\mu} \mu^{-} \bar{\nu}_{\mu}$. The 95\% C.L. limit for a form factor scale $\Lambda=1 \mathrm{TeV}$ is $-0.05<\kappa_{\gamma}<2.29$, fixing $\lambda_{\gamma}=0$. With the assumption that the $W W \gamma$ and $W W Z$ couplings are equal the $95 \%$ C.L. one-dimensional limit ( $\Lambda$ $=2 \mathrm{TeV}$ ) is $0.68<\kappa<1.45$.
${ }^{19}$ ABAZOV 05J perform a likelihood fit to the photon $E_{T}$ spectrum of $W \gamma+\mathrm{X}$ events, where the $W$ decays to an electron or muon which is required to be well separated from the photon. For $\Lambda=2.0 \mathrm{TeV}$ the $95 \% \mathrm{CL}$ limits are $0.12<\kappa_{\gamma}<1.96$. In the fit $\lambda_{\gamma}$ is kept fixed to its Standard Model value.
${ }^{20}$ ABREU 01 combine results from $e^{+} e^{-}$interactions at 189 GeV leading to $W^{+} W^{-}$, $W e \nu_{e}$, and $\nu \bar{\nu} \gamma$ final states with results from ABREU 99L at 183 GeV . The $95 \%$ confidence interval is $0.87<\kappa_{\gamma}<1.68$.
21 BREITWEG 00 search for $W$ production in events with large hadronic $p_{T}$. For $p_{T}>20$ GeV , the upper limit on the cross section gives the $95 \% \mathrm{CL}$ limit $-3.7<\kappa_{\gamma}<2.5$ (for $\lambda_{\gamma}=0$ ).

[^0]$\lambda_{\gamma}$
OUR FIT below is taken from [SCHAEL 13A].

| VALUE | EVTS | DOCUMENT ID |  | TECN | COMMENT |
| :---: | :---: | :---: | :---: | :---: | :---: |
| -0.022 $\pm 0.019$ OUR FIT |  |  |  |  |  |
| $0.002 \pm 0.035$ | 7872 | ${ }^{1}$ ABDALLAH | 10 | DLPH | $E_{\mathrm{cm}}^{e e}=189-209 \mathrm{GeV}$ |
| $-0.012 \pm 0.027 \pm 0.011$ | 10689 | 2 SCHAEL | 05A | ALEP | $E_{\mathrm{cm}}^{e e}=183-209 \mathrm{GeV}$ |
| $-0.060+0.034$ | 9800 | 3 ABBIENDI | 04D | OPAL | $E_{\mathrm{Cm}}^{e e}=183-209 \mathrm{GeV}$ |
| $-0.021_{-0.034}^{+0.035} \pm 0.017$ | 10575 | 4 ACHARD | 04D | L3 | $E_{\mathrm{Cm}}^{e \mathrm{e}}=161-209 \mathrm{GeV}$ |

-     - We do not use the following data for averages, fits, limits, etc.

|  | 53 | ${ }^{5}$ CHATRCHYA | 14AB | CMS | $E_{\mathrm{cm}}^{p p}=7 \mathrm{TeV}$ |
| :---: | :---: | :---: | :---: | :---: | :---: |
|  |  | ${ }^{6}$ AAD | 13AN | ATLS | $E_{\mathrm{cm}}^{p p}=7 \mathrm{TeV}$ |
|  |  | 7 ABAZOV | 12AG | D0 | $E_{\mathrm{cm}}^{p \bar{p}}=1.96 \mathrm{TeV}$ |
|  |  | 8 ABAZOV | 11AC | D0 | $E_{\mathrm{cm}}^{p \bar{p}}=1.96 \mathrm{TeV}$ |
|  |  | 9 CHATRCHYA |  | CMS | $E_{\mathrm{cm}}^{p p}=7 \mathrm{TeV}$ |
|  |  | 10 AARON | 09B | H1 | $E_{\mathrm{cm}}^{e p}=0.3 \mathrm{TeV}$ |
| $0.00 \pm 0.06$ |  | 11 ABAZOV | 09AD | D0 | $E_{\mathrm{cm}}^{p \bar{p}}=1.96 \mathrm{TeV}$ |
|  |  | 12 ABAZOV | 09AJ | D0 | $E_{\mathrm{cm}}^{p \bar{p}}=1.96 \mathrm{TeV}$ |
|  |  | 13 ABAZOV | 08R | D0 | $E_{\mathrm{cm}}^{p \bar{p}}=1.96 \mathrm{TeV}$ |
| 0.16 ${ }_{-0.13}^{+0.12}$ | 1880 | 14 ABDALLAH | 08C | DLPH | Superseded by ABDAL- <br> LAH 10 |
|  | 1617 | 15 AALTONEN | 07L | CDF | $E_{\mathrm{cm}}^{p \bar{p}}=1.96 \mathrm{GeV}$ |
|  | 17 | 16 ABAZOV | 06H | D0 | $E_{\mathrm{cm}}^{p \bar{p}}=1.96 \mathrm{TeV}$ |
|  | 141 | 17 ABAZOV | 05J | D0 | $E_{\mathrm{cm}}^{p \bar{p}}=1.96 \mathrm{TeV}$ |
| $0.05 \pm 0.09 \pm 0.01$ | 2298 | 18 ABREU | 011 | DLPH | $E_{\mathrm{cm}}^{e \mathrm{e}}=183+189 \mathrm{GeV}$ |
|  |  | 19 BREITWEG | 00 | ZEUS | $\begin{gathered} e^{+} p \rightarrow e^{+} W^{ \pm} X \\ \sqrt{s} \approx 300 \mathrm{GeV} \end{gathered}$ |
| 0.00 $\begin{array}{r}+0.10 \\ -0.09\end{array}$ | 331 | 20 ABBOTT | 991 | D0 | $E_{\mathrm{cm}}^{p \bar{p}}=1.8 \mathrm{TeV}$ |

${ }^{1}$ ABDALLAH 10 use data on the final states $e^{+} e^{-} \rightarrow j j \ell \nu, j j j j, j j X, \ell X$, at center-of-mass energies between $189-209 \mathrm{GeV}$ at LEP2, where $j=$ jet, $\ell=$ lepton, and $X$ represents missing momentum. The fit is carried out keeping all other parameters fixed at their SM values.
2 SCHAEL 05A study single-photon, single- $W$, and $W W$-pair production from 183 to 209 GeV . Each parameter is determined from a single-parameter fit in which the other parameters assume their Standard Model values.
${ }^{3}$ ABBIENDI 04D combine results from $W^{+} W^{-}$in all decay channels. Only CP-conserving couplings are considered and each parameter is determined from a single-parameter fit in which the other parameters assume their Standard Model values. The $95 \%$ confidence interval is $-0.13<\lambda_{\gamma}<0.01$.
${ }^{4}$ ACHARD 04D study $W W$-pair production, single- $W$ production and single-photon production with missing energy from 189 to 209 GeV . The result quoted here is obtained including data from 161 to 183 GeV , ACCIARRI 99Q. Each parameter is determined
from a single-parameter fit in which the other parameters assume their Standard Model values.
${ }^{5}$ CHATRCHYAN 14AB measure $W \gamma$ production cross section for $p_{T}^{\gamma}>15 \mathrm{GeV}$ and $\mathrm{R}(\ell \gamma)$ $>0.7$, which is the separation between the $\gamma$ and the final state charged lepton (e or $\mu$ ) in the azimuthal angle-pseudorapidity $(\phi-\eta)$ plane. After background subtraction the number of e $\nu \gamma$ and $\mu \nu \gamma$ events is determined to be $3200 \pm 325$ and $4970 \pm 543$ respectively, compatible with expectations from the SM. This leads to a $95 \%$ CL limit of $-0.050<\lambda_{\gamma}<0.037$, assuming all other parameters have SM values.
${ }^{6}$ AAD 13AN study $W \gamma$ production in $p p$ collisions. In events with no additional jet, 4449 (6578) W decays to electron (muon) are selected, with an expected background of $1662 \pm 262(2538 \pm 362)$ events. Analysing the photon $p_{T}$ spectrum above 100 GeV yields a $95 \%$ C.L. limit of $-0.065<\lambda_{\gamma}<0.061$. Supersedes AAD 12BX.
${ }^{7}$ ABAZOV 12AG combine new results with already published results on $W \gamma, W W$ and $W Z$ production in order to determine the couplings with increased precision, superseding ABAZOV 08R, ABAZOV 11AC, ABAZOV 09AJ, ABAZOV 09AD. The 68\% C.L. result for a formfactor cutoff of $\Lambda=2 \mathrm{TeV}$ is $\lambda_{\gamma}=0.007_{-0.022}^{+0.021}$.
${ }^{8}$ ABAZOV 11AC study $W \gamma$ production in $p \bar{p}$ collisions at 1.96 TeV , with the $W$ decay products containing an electron or a muon. They select 196 (363) events in the electron (muon) mode, with a SM expectation of 190 (372) events. A likelihood fit to the photon $E_{T}$ spectrum above 15 GeV yields at $95 \%$ C.L. the result: $-0.08<\lambda_{\gamma}<0.07$ for a formfactor $\Lambda=2 \mathrm{TeV}$.
${ }^{9}$ CHATRCHYAN 11 M study $W \gamma$ production in $p p$ collisions at $\sqrt{s}=7 \mathrm{TeV}$ using $36 \mathrm{pb}^{-1}$ $p p$ data with the $W$ decaying to electron and muon. The total cross section is measured for photon transverse energy $E_{T}^{\gamma}>10 \mathrm{GeV}$ and spatial separation from charged leptons in the plane of pseudo rapidity and azimuthal angle $\Delta R(\ell, \gamma)>0.7$. The number of candidate (background) events is $452(228 \pm 21)$ for the electron channel and 520 $(277 \pm 25)$ for the muon channel. Setting other couplings to their standard model value, they derive a $95 \% \mathrm{CL}$ limit of $-0.18<\lambda_{\gamma}<0.17$.
10 AARON 09B study single- $W$ production in ep collisions at 0.3 TeV C.M. energy. They select $53 W \rightarrow e / \mu$ events with a standard model expectation of $54.1 \pm 7.4$ events. Fitting the transverse momentum spectrum of the hadronic recoil system they obtain a $95 \%$ C.L. limit of $-2.5<\lambda_{\gamma}<2.5$.
11 ABAZOV 09AD study the $p \bar{p} \rightarrow \ell \nu 2$ jet process arising in $W W$ and $W Z$ production. They select $12,473(14,392)$ events in the electron (muon) channel with an expected di-boson signal of 436 (527) events. The results on the anomalous couplings are derived from an analysis of the $p_{T}$ spectrum of the 2 -jet system and quoted at $68 \%$ C.L. and for a form factor of 2 TeV . This measurement is not used for obtaining the mean as it is for a specific form factor. The $95 \%$ confidence interval is $-0.10<\lambda_{\gamma}<0.11$.
12 ABAZOV 09AJ study the $p \bar{p} \rightarrow 2 \ell 2 \nu$ process arising in $W W$ production. They select 100 events with an expected $W W$ signal of 65 events. An analysis of the $p_{T}$ spectrum of the two charged leptons leads to $95 \%$ C.L. limits of $-0.14<\lambda_{\gamma}<0.18$, for a form factor $\Lambda=2 \mathrm{TeV}$.
13 ABAZOV 08R use $0.7 \mathrm{fb}^{-1} p \bar{p}$ data at $\sqrt{s}=1.96 \mathrm{TeV}$ to select $263 W \gamma+X$ events, of which 187 constitute signal, with the $W$ decaying into an electron or a muon, which is required to be well separated from a photon with $E_{T}>9 \mathrm{GeV}$. A likelihood fit to the photon $E_{T}$ spectrum yields a $95 \%$ CL limit $-0.12<\lambda_{\gamma}<0.13$ with other couplings fixed to their Standard Model values.
14 ABDALLAH 08C determine this triple gauge coupling from the measurement of the spin density matrix elements in $e^{+} e^{-} \rightarrow W^{+} W^{-} \rightarrow(q q)(\ell \nu)$, where $\ell=e$ or $\mu$. Values of all other couplings are fixed to their standard model values.
15 AALTONEN 07L set limits on anomalous TGCs using the $p_{T}(W)$ distribution in $W W$ and $W Z$ production with the $W$ decaying to an electron or muon and the $Z$ to 2 jets. Setting other couplings to their standard model value, the $95 \%$ C.L. limits are $-0.18<\lambda_{\gamma}<0.17$ for a form factor scale $\Lambda=1.5 \mathrm{TeV}$.
${ }^{16}$ ABAZOV 06 H study $\bar{p} p \rightarrow W W$ production with a subsequent decay $W W \rightarrow$ $e^{+} \nu_{e} e^{-} \bar{\nu}_{e}, W W \rightarrow e^{ \pm} \nu_{e} \mu^{\mp} \nu_{\mu}$ or $W W \rightarrow \mu^{+} \nu_{\mu} \mu^{-} \bar{\nu}_{\mu}$. The 95\% C.L. limit for a form factor scale $\Lambda=1 \mathrm{TeV}$ is $-0.97<\lambda_{\gamma}<1.04$, fixing $\kappa_{\gamma}=1$. With the assumption that the $W W \gamma$ and $W W Z$ couplings are equal the $95 \%$ C.L. one-dimensional limit ( $\Lambda$ $=2 \mathrm{TeV}$ ) is $-0.29<\lambda<0.30$.
17 ABAZOV 05J perform a likelihood fit to the photon $E_{T}$ spectrum of $W \gamma+\mathbf{X}$ events, where the $W$ decays to an electron or muon which is required to be well separated from the photon. For $\Lambda=2.0 \mathrm{TeV}$ the $95 \% \mathrm{CL}$ limits are $-0.20<\lambda_{\gamma}<0.20$. In the fit $\kappa_{\gamma}$ is kept fixed to its Standard Model value.
18 ABREU 01 combine results from $e^{+} e^{-}$interactions at 189 GeV leading to $W^{+} W^{-}$, $W e \nu_{e}$, and $\nu \bar{\nu} \gamma$ final states with results from ABREU 99L at 183 GeV . The $95 \%$ confidence interval is $-0.11<\lambda_{\gamma}<0.23$.
19 BREITWEG 00 search for $W$ production in events with large hadronic $p_{T}$. For $p_{T}>20$ GeV , the upper limit on the cross section gives the $95 \% \mathrm{CL}$ limit $-3.2<\lambda_{\gamma}<3.2$ for $\kappa_{\gamma}$ fixed to its Standard Model value.
${ }^{20}$ ABBOTT 99ı perform a simultaneous fit to the $W \gamma, W W \rightarrow$ dilepton, $W W / W Z \rightarrow$ $e \nu j j, W W / W Z \rightarrow \mu \nu j j$, and $W Z \rightarrow$ trilepton data samples. For $\Lambda=2.0 \mathrm{TeV}$, the $95 \% \mathrm{CL}$ limits are $-0.18<\lambda_{\gamma}<0.19$.

This coupling is $C P$-conserving ( $C$ - and $P$ - separately conserving).

| VALUE | EVTS | DOCUMENT |  | TECN | COMMENT |
| :---: | :---: | :---: | :---: | :---: | :---: |
| $0.924+0.059 \pm 0.024$ | 7171 | 1 ACHARD | 04D | L3 | $E_{\mathrm{Cm}}^{e \mathrm{ee}}=189-209 \mathrm{GeV}$ |

-     - We do not use the following data for averages, fits, limits, etc. - - -

|  | ${ }^{2}$ AAD | 13 AL ATLS | $E_{\mathrm{cm}}^{p p}=7 \mathrm{TeV}$ |
| :--- | :--- | :--- | :--- |
|  | ${ }^{3}$ AAD | 12 CD ATLS | $E_{\mathrm{cm}}^{p p}=7 \mathrm{TeV}$ |
|  | ${ }^{4}$ AALTONEN | 12 AC CDF | $E_{\mathrm{Cm}}^{p \bar{p}}=1.96 \mathrm{TeV}$ |
| 34 | ${ }^{5}$ ABAZOV | 11 DO | $E_{\mathrm{Cm}}^{p \bar{p}}=1.96 \mathrm{TeV}$ |
| 17 | ${ }^{6}$ ABAZOV | 06 H D0 | $E_{\mathrm{Cm}}^{p \bar{p}}=1.96 \mathrm{TeV}$ |
| 2.3 | ${ }^{7}$ ABAZOV | 05 S D0 | $E_{\mathrm{cm}}^{p \bar{p}}=1.96 \mathrm{TeV}$ |

1 ACHARD 04D study $W W$-pair production, single- $W$ production and single-photon production with missing energy from 189 to 209 GeV . The result quoted here is obtained using the $W W$-pair production sample. Each parameter is determined from a singleparameter fit in which the other parameters assume their Standard Model values.
${ }^{2}$ AAD 13 AL study $W W$ production in $p p$ collisions and select $1325 W W$ candidates in decay modes with electrons or muons with an expected background of $369 \pm 61$ events. Assuming the LEP formulation and setting the form-factor $\Lambda=$ infinity, a fit to the transverse momentum distribution of the leading charged lepton, leads to a 95\% C.L. range of $0.957<\kappa_{Z}<1.043$. Supersedes AAD 12AC.
${ }^{3}$ AAD 12CD study $W Z$ production in $p p$ collisions and select $317 W Z$ candidates in three $\ell \nu$ decay modes with an expected background of $68.0 \pm 10.0$ events. The resulting $95 \%$ C.L. range is: $0.63<\kappa_{Z}<1.57$. Supersedes AAD 12V.
${ }^{4}$ AALTONEN 12AC study $W Z$ production in $p \bar{p}$ collisions and select $63 W Z$ candidates in three $\ell \nu$ decay modes with an expected background of $7.9 \pm 1.0$ events. Based on the cross section and shape of the $Z$ transverse momentum spectrum, the following $95 \%$ C.L. range is reported: $0.61<\kappa_{Z}<1.90$ for a form factor of $\Lambda=2 \mathrm{TeV}$.
${ }^{5}$ ABAZOV 11 study the $p \bar{p} \rightarrow 3 \ell \nu$ process arising in $W Z$ production. They observe $34 W Z$ candidates with an estimated background of 6 events. An analysis of the $p_{T}$ spectrum of the $Z$ boson leads to a $95 \%$ C.L. limit of $0.600<\kappa_{Z}<1.675$, for a form factor $\Lambda=2 \mathrm{TeV}$.
${ }^{6}$ ABAZOV 06 H study $\bar{p} p \rightarrow W W$ production with a subsequent decay $W W \rightarrow$ $e^{+} \nu_{e} e^{-} \bar{\nu}_{e}, W W \rightarrow e^{ \pm} \nu_{e} \mu^{\mp} \nu_{\mu}$ or $W W \rightarrow \mu^{+} \nu_{\mu} \mu^{-} \bar{\nu}_{\mu}$. The 95\% C.L. limit for a form factor scale $\Lambda=2 \mathrm{TeV}$ is $0.55<\kappa_{Z}<1.55$, fixing $\lambda_{Z}=0$. With the assumption that the $W W \gamma$ and $W W Z$ couplings are equal the $95 \%$ C.L. one-dimensional limit ( $\Lambda$ $=2 \mathrm{TeV}$ ) is $0.68<\kappa<1.45$.
${ }^{7}$ ABAZOV 05S study $\bar{p} p \rightarrow W Z$ production with a subsequent trilepton decay to $\ell \nu \ell^{\prime} \bar{\ell}^{\prime}$ ( $\ell$ and $\ell^{\prime}=e$ or $\mu$ ). Three events (estimated background $0.71 \pm 0.08$ events) with $W Z$ decay characteristics are observed from which they derive limits on the anomalous $W W Z$ couplings. The $95 \%$ CL limit for a form factor scale $\Lambda=1 \mathrm{TeV}$ is $-1.0<\kappa_{Z}<3.4$, fixing $\lambda_{Z}$ and $g_{1}^{Z}$ to their Standard Model values.

## $\lambda_{z}$

This coupling is $C P$-conserving ( $C$ - and $P$ - separately conserving).

| VALUE | EVTS | DOCUMENT |  | TECN | COMMENT |
| :---: | :---: | :---: | :---: | :---: | :---: |
| $-0.088 \pm-0.060{ }_{-0.023}^{+0.023}$ | 7171 | 1 ACHARD | 04D | L3 | $E_{\mathrm{Cm}}^{e \mathrm{e}}$ ( ${ }^{\text {c }} 189-209 \mathrm{GeV}$ |

-     - We do not use the following data for averages, fits, limits, etc. - - -

|  | 2 AAD | 14Y | ATLS | $E_{\mathrm{cm}}^{p p}=8 \mathrm{TeV}$ |
| :---: | :---: | :---: | :---: | :---: |
|  | ${ }^{3} \mathrm{AAD}$ | 13AL | ATLS | $E_{\mathrm{cm}}^{p p}=7 \mathrm{TeV}$ |
|  | ${ }^{4}$ CHATRCHYAN | 13 BF | CMS | $E_{\mathrm{cm}}^{p p}=7 \mathrm{TeV}$ |
|  | ${ }^{5}$ AAD | 12CD | ATLS | $E_{\mathrm{cm}}^{p p}=7 \mathrm{TeV}$ |
|  | ${ }^{6}$ AALTONEN | 12AC | CDF | $E_{\mathrm{Cm}}^{p \bar{p}}=1.96 \mathrm{TeV}$ |
| 34 | 7 ABAZOV | 11 | D0 | $E_{\mathrm{Cm}}^{p \bar{p}}=1.96 \mathrm{TeV}$ |
| 334 | 8 AALTONEN | 10K | CDF | $E_{\mathrm{cm}}^{p \bar{p}}=1.96 \mathrm{TeV}$ |
| 13 | 9 ABAZOV | 072 | D0 | $E_{\mathrm{cm}}^{p \bar{p}}=1.96 \mathrm{TeV}$ |
| 17 | 10 ABAZOV | 06H | D0 | $E_{\mathrm{cm}}^{p \bar{p}}=1.96 \mathrm{TeV}$ |
| 2.3 | 11 ABAZOV | 05S | D0 | $E_{\mathrm{Cm}}^{p \bar{p}}=1.96 \mathrm{TeV}$ |

1 ACHARD 04D study $W W$-pair production, single- $W$ production and single-photon production with missing energy from 189 to 209 GeV . The result quoted here is obtained using the $W W$-pair production sample. Each parameter is determined from a singleparameter fit in which the other parameters assume their Standard Model values.
${ }^{2}$ AAD 14 Y determine the electroweak $Z$-dijet cross section in $8 \mathrm{TeV} p p$ collisions. $Z \rightarrow$ $e e$ and $Z \rightarrow \mu \mu$ decays are selected with the di-lepton $p_{T}>20 \mathrm{GeV}$ and mass in the $81-101 \mathrm{GeV}$ range. Minimum two jets are required with $p_{T}>55$ and 45 GeV and no additional jets with $p_{T}>25 \mathrm{GeV}$ in the rapidity interval between them. The normalized $p_{T}$ balance between the $Z$ and the two jets is required to be $<0.15$. This leads to a selection of 900 events with dijet mass $>1 \mathrm{TeV}$. The number of signal and background events expected is 261 and 592 respectively. A Poisson likelihood method is used on an event by event basis to obtain the $95 \%$ CL limit $-0.15<\lambda_{Z}<0.13$ for a form factor value $\Lambda=\infty$.
${ }^{3}$ AAD 13AL study $W W$ production in $p p$ collisions and select $1325 W W$ candidates in decay modes with electrons or muons with an expected background of $369 \pm 61$ events. Assuming the LEP formulation and setting the form-factor $\Lambda=$ infinity, a fit to the transverse momentum distribution of the leading charged lepton, leads to a $95 \%$ C.L. range of $-0.062<\lambda_{Z}<0.059$. Supersedes AAD 12AC.
${ }^{4}$ CHATRCHYAN 13BF determine the $W^{+} W^{-}$production cross section using unlike sign di-lepton ( $e$ or $\mu$ ) events with high $\not p_{T}$. The leptons have $p_{T}>20 \mathrm{GeV} / \mathrm{c}$ and are isolated. 1134 candidate events are observed with an expected SM background of $247 \pm$ 34. The $p_{T}$ distribution of the leading lepton is fitted to obtain $95 \%$ C.L. limits of $-0.048 \leq \lambda_{Z} \leq 0.048$.
${ }^{5}$ AAD 12CD study $W Z$ production in $p p$ collisions and select $317 W Z$ candidates in three $\ell \nu$ decay modes with an expected background of $68.0 \pm 10.0$ events. The resulting $95 \%$ C.L. range is: $-0.046<\lambda_{Z}<0.047$. Supersedes AAD 12V.
${ }^{6}$ AALTONEN 12AC study $W Z$ production in $p \bar{p}$ collisions and select $63 W Z$ candidates in three $\ell \nu$ decay modes with an expected background of $7.9 \pm 1.0$ events. Based on the cross section and shape of the $Z$ transverse momentum spectrum, the following 95\% C.L. range is reported: $-0.08<\lambda_{Z}<0.10$ for a form factor of $\Lambda=2 \mathrm{TeV}$.
${ }^{7}$ ABAZOV 11 study the $p \bar{p} \rightarrow 3 \ell \nu$ process arising in $W Z$ production. They observe $34 W Z$ candidates with an estimated background of 6 events. An analysis of the $p_{T}$ spectrum of the $Z$ boson leads to a $95 \%$ C.L. limit of $-0.077<\lambda_{Z}<0.093$, for a form factor $\Lambda=2 \mathrm{TeV}$.
${ }^{8}$ AALTONEN 10 K study $p \bar{p} \rightarrow W^{+} W^{-}$with $W \rightarrow e / \mu \nu$. The $p_{T}$ of the leading (second) lepton is required to be $>20(10) \mathrm{GeV}$. The final number of events selected is 654 of which $320 \pm 47$ are estimated to be background. The $95 \%$ C.L. interval is $-0.16<\lambda_{Z}<0.16$ for $\Lambda=1.5 \mathrm{TeV}$ and $-0.14<\lambda_{Z}<0.15$ for $\Lambda=2 \mathrm{TeV}$.
${ }^{9}$ ABAZOV $07 Z$ set limits on anomalous TGCs using the measured cross section and $p_{T}(Z)$ distribution in $W Z$ production with both the $W$ and the $Z$ decaying leptonically into electrons and muons. Setting the other couplings to their standard model values, the $95 \%$ C.L. limit for a form factor scale $\Lambda=2 \mathrm{TeV}$ is $-0.17<\lambda_{Z}<0.21$.
10 ABAZOV 06 H study $\bar{p} p \rightarrow W W$ production with a subsequent decay $W W \rightarrow$ $e^{+} \nu_{e} e^{-} \bar{\nu}_{e}, W W \rightarrow e^{ \pm} \nu_{e} \mu^{\mp} \nu_{\mu}$ or $W W \rightarrow \mu^{+} \nu_{\mu} \mu^{-} \bar{\nu}_{\mu}$. The 95\% C.L. limit for a form factor scale $\Lambda=2 \mathrm{TeV}$ is $-0.39<\lambda_{Z}<0.39$, fixing $\kappa_{Z}=1$. With the assumption that the $W W \gamma$ and $W W Z$ couplings are equal the $95 \%$ C.L. one-dimensional limit ( $\Lambda$ $=2 \mathrm{TeV})$ is $-0.29<\lambda<0.30$.
${ }^{11}$ ABAZOV 05S study $\bar{p} p \rightarrow W Z$ production with a subsequent trilepton decay to $\ell \nu \ell^{\prime} \bar{\ell}^{\prime}$ ( $\ell$ and $\ell^{\prime}=e$ or $\mu$ ). Three events (estimated background $0.71 \pm 0.08$ events) with $W Z$ decay characteristics are observed from which they derive limits on the anomalous $W W Z$ couplings. The $95 \%$ CL limit for a form factor scale $\Lambda=1.5 \mathrm{TeV}$ is $-0.48<\lambda_{Z}<$ 0.48 , fixing $g_{1}^{Z}$ and $\kappa_{Z}$ to their Standard Model values.

This coupling is $C P$-conserving but $C$ - and $P$-violating.
VALUE EVTS DOCUMENT ID TECN COMMENT
$\mathbf{- 0 . 0 7} \pm \mathbf{0 . 0 9}$ OUR AVERAGE Error includes scale factor of 1.1.

| $\begin{array}{r} -0.04+0.13 \\ -0.12 \end{array}$ | 9800 | 1 ABBIEND | 04D | OPAL | $E_{\mathrm{cm}}^{\mathrm{ee}}=183-209 \mathrm{GeV}$ |
| :---: | :---: | :---: | :---: | :---: | :---: |
| $0.00 \pm 0.13 \pm 0.05$ | 7171 | 2 ACHARD | 04D | L3 | $E_{\mathrm{cm}}^{e e}=189-209 \mathrm{GeV}$ |
| $-0.44_{-0.22}^{+0.23} \pm 0.12$ | 1154 | 3 ACCIARRI | 99Q | L3 | $E_{\mathrm{cm}}^{e e}=161+172+183 \mathrm{GeV}$ |

-     - We do not use the following data for averages, fits, limits, etc. - - -
$-0.31 \pm 0.23 \quad 4$ EBOLI 00 THEO LEP1, SLC + Tevatron
${ }^{1}$ ABBIENDI 04D combine results from $W^{+} W^{-}$in all decay channels. Only CP-conserving couplings are considered and each parameter is determined from a single-parameter fit in which the other parameters assume their Standard Model values. The 95\% confidence interval is $-0.28<g_{5}^{Z}<+0.21$.
2 ACHARD 04D study $W W$-pair production, single- $W$ production and single-photon production with missing energy from 189 to 209 GeV . The result quoted here is obtained using the $W W$-pair production sample. Each parameter is determined from a singleparameter fit in which the other parameters assume their Standard Model values.
3 ACCIARRI 99Q study $W$-pair, single- $W$, and single photon events.
${ }^{4}$ EBOLI 00 extract this indirect value of the coupling studying the non-universal one-loop contributions to the experimental value of the $Z \rightarrow b \bar{b}$ width ( $\Lambda=1 \mathrm{TeV}$ is assumed).

This coupling is $C P$-violating ( $C$-violating and $P$-conserving).
VALUE EVTS DOCUMENTID TECN COMMENT
$-0.30 \pm 0.17$ OUR AVERAGE
$-0.39_{-0.20}^{+0.19} \quad 1880$
$-0.02_{-0.33}^{+0.32} \quad 1065$
${ }^{1}$ ABDALLAH 08 C DLPH $E_{\mathrm{cm}}^{e e}=189-209 \mathrm{GeV}$
${ }^{2}$ ABBIENDI 01H OPAL $E_{\mathrm{cm}}^{e e}=189 \mathrm{GeV}$
${ }^{1}$ ABDALLAH 08 C determine this triple gauge coupling from the measurement of the spin density matrix elements in $e^{+} e^{-} \rightarrow W^{+} W^{-} \rightarrow(q q)(\ell \nu)$, where $\ell=e$ or $\mu$. Values of all other couplings are fixed to their standard model values.
${ }^{2}$ ABBIENDI 01H study $W$-pair events, with one leptonically and one hadronically decaying $W$. The coupling is extracted using information from the $W$ production angle together with decay angles from the leptonically decaying $W$.

## $\tilde{\kappa}_{\boldsymbol{Z}}$

This coupling is $C P$-violating ( $C$-conserving and $P$-violating).
VALUE EVTS DOCUMENTID TECN COMMENT

## $=0.12{ }_{-0.04}^{+0.06}$ OUR AVERAGE

$\begin{array}{lllll}-0.09_{-0.05}^{+0.08} & 1880 & { }^{1} \text { ABDALLAH } & 08 C & \text { DLPH } \\ E_{\mathrm{cm}}^{e e}=189-209 \mathrm{GeV} \\ -0.20_{-0.07}^{+0.10} & 1065 & { }^{2} \text { ABBIENDI } & 01 \mathrm{H} & \text { OPAL } \\ E_{\mathrm{cm}}^{e e}=189 \mathrm{GeV}\end{array}$

-     - We do not use the following data for averages, fits, limits, etc. - -

$$
{ }^{3} \text { BLINOV } \quad 11 \quad \text { LEP } \quad E_{\mathrm{cm}}^{e e}=183-207 \mathrm{GeV}
$$

${ }^{1}$ ABDALLAH 08C determine this triple gauge coupling from the measurement of the spin density matrix elements in $e^{+} e^{-} \rightarrow W^{+} W^{-} \rightarrow(q q)(\ell \nu)$, where $\ell=e$ or $\mu$. Values of all other couplings are fixed to their standard model values.
${ }^{2}$ ABBIENDI 01H study $W$-pair events, with one leptonically and one hadronically decaying $W$. The coupling is extracted using information from the $W$ production angle together with decay angles from the leptonically decaying $W$.
${ }^{3}$ BLINOV 11 use the LEP-average $e^{+} e^{-} \rightarrow W^{+} W^{-}$cross section data for $\sqrt{s}=$ $183-207 \mathrm{GeV}$ to determine an upper limit on the TGC $\widetilde{\kappa}_{Z}$. The average values of the cross sections as well as their correlation matrix, and standard model expectations of the cross sections are taken from the LEPEWWG note hep-ex/0612034. At 95\% confidence level $\left|\widetilde{\kappa}_{Z}\right|<0.13$.

## $\tilde{\lambda}_{z}$

This coupling is $C P$-violating ( $C$-conserving and $P$-violating).

| Value | EVTS | DOCUMENT ID |  | TECN | COMMENT |
| :---: | :---: | :---: | :---: | :---: | :---: |
| -0.09 $\pm 0.07$ OUR AVERAGE |  |  |  |  |  |
| $-0.08 \pm 0.07$ | 1880 | ${ }^{1}$ ABDALLAH | 08C | DLPH | $E_{\mathrm{cm}}^{e e}=189-209 \mathrm{GeV}$ |
| $-0.18{ }_{-0.16}^{+0.24}$ | 1065 | ${ }^{2}$ ABBIENDI | 01H | OPAL | $E_{\mathrm{cm}}^{e \mathrm{e}}=189 \mathrm{GeV}$ |

-     - We do not use the following data for averages, fits, limits, etc.

$$
3^{3} \text { BLINOV } \quad 11 \text { LEP } \quad E_{\mathrm{cm}}^{e e}=183-207 \mathrm{GeV}
$$

${ }^{1}$ ABDALLAH 08 C determine this triple gauge coupling from the measurement of the spin density matrix elements in $e^{+} e^{-} \rightarrow W^{+} W^{-} \rightarrow(q q)(\ell \nu)$, where $\ell=e$ or $\mu$. Values of all other couplings are fixed to their standard model values.
${ }^{2}$ ABBIENDI 01 H study $W$-pair events, with one leptonically and one hadronically decaying $W$. The coupling is extracted using information from the $W$ production angle together with decay angles from the leptonically decaying $W$.
${ }^{3}$ BLINOV 11 use the LEP-average $e^{+} e^{-} \rightarrow W^{+} W^{-}$cross section data for $\sqrt{s}=$ $183-207 \mathrm{GeV}$ to determine an upper limit on the TGC $\tilde{\lambda}_{Z}$. The average values of the cross sections as well as their correlation matrix, and standard model expectations of the cross sections are taken from the LEPEWWG note hep-ex/0612034. At 95\% confidence level $\left|\widetilde{\lambda}_{Z}\right|<0.31$.

## W ANOMALOUS MAGNETIC MOMENT

The full magnetic moment is given by $\mu_{W}=e(1+\kappa+\lambda) / 2 m_{W}$. In the Standard Model, at tree level, $\kappa=1$ and $\lambda=0$. Some papers have defined $\Delta \kappa=1-\kappa$ and assume that $\lambda=0$. Note that the electric quadrupole moment is given by $-e(\kappa-\lambda) / m_{W}^{2}$. A description of the parameterization of these moments and additional references can be found in HAGIWARA 87 and BAUR 88. The parameter $\Lambda$ appearing in the theoretical limits below is a regularization cutoff which roughly corresponds to the energy scale where the structure of the $W$ boson becomes manifest.

${ }^{1}$ ABREU 01I combine results from $e^{+} e^{-}$interactions at 189 GeV leading to $W^{+} W^{-}$, We $\nu_{e}$, and $\nu \bar{\nu} \gamma$ final states with results from ABREU 99L at 183 GeV to determine $\Delta g_{1}^{Z}, \Delta \kappa_{\gamma}$, and $\lambda_{\gamma} . \Delta \kappa_{\gamma}$ and $\lambda_{\gamma}$ are simultaneously floated in the fit to determine ${ }^{\mu} W$.
${ }^{2} \mathrm{ABE} 95 \mathrm{G}$ report $-1.3<\kappa<3.2$ for $\lambda=0$ and $-0.7<\lambda<0.7$ for $\kappa=1$ in $p \bar{p} \rightarrow e \nu_{e} \gamma \mathrm{X}$ and $\mu \nu_{\mu} \gamma \mathrm{X}$ at $\sqrt{s}=1.8 \mathrm{TeV}$.
${ }^{3}$ ALITTI 92C measure $\kappa=1_{-2.2}^{+2.6}$ and $\lambda=0_{-1.8}^{+1.7}$ in $p \bar{p} \rightarrow e \nu \gamma+X$ at $\sqrt{s}=630 \mathrm{GeV}$. At $95 \%$ CL they report $-3.5<\kappa<5.9$ and $-3.6<\lambda<3.5$.
${ }^{4}$ SAMUEL 92 use preliminary CDF and UA2 data and find $-2.4<\kappa<3.7$ at $96 \% \mathrm{CL}$ and $-3.1<\kappa<4.2$ at $95 \% \mathrm{CL}$ respectively. They use data for $W \gamma$ production and radiative $W$ decay.
${ }^{5}$ SAMUEL 91 use preliminary CDF data for $p \bar{p} \rightarrow W \gamma X$ to obtain $-11.3 \leq \Delta \kappa \leq$ 10.9. Note that their $\kappa=1-\Delta \kappa$.
${ }^{6}$ GRIFOLS 88 uses deviation from $\rho$ parameter to set limit $\Delta \kappa \lesssim 65\left(M_{W}^{2} / \Lambda^{2}\right)$.
7 GROTCH 87 finds the limit $-37<\Delta \kappa<73.5$ ( $90 \% \mathrm{CL}$ ) from the experimental limits on $e^{+} e^{-} \rightarrow \nu \bar{\nu} \gamma$ assuming three neutrino generations and $-19.5<\Delta \kappa<56$ for four generations. Note their $\Delta \kappa$ has the opposite sign as our definition.
8 VANDERBIJ 87 uses existing limits to the photon structure to obtain $|\Delta \kappa|<33$ $\left(m_{W} / \Lambda\right)$. In addition VANDERBIJ 87 discusses problems with using the $\rho$ parameter of the Standard Model to determine $\Delta \kappa$.
${ }^{9}$ GRAU 85 uses the muon anomaly to derive a coupled limit on the anomalous magnetic dipole and electric quadrupole $(\lambda)$ moments $1.05>\Delta \kappa \ln \left(\Lambda / m_{W}\right)+\lambda / 2>-2.77$. In the Standard Model $\lambda=0$.
${ }^{10}$ SUZUKI 85 uses partial-wave unitarity at high energies to obtain $|\Delta \kappa| \lesssim 190$ $\left(m_{W} / \Lambda\right)^{2}$. From the anomalous magnetic moment of the muon, SUZUKI 85 obtains $|\Delta \kappa| \lesssim 2.2 / \ln \left(\Lambda / m_{W}\right)$. Finally SUZUKI 85 uses deviations from the $\rho$ parameter and obtains a very qualitative, order-of-magnitude limit $|\Delta \kappa| \lesssim 150\left(m_{W} / \Lambda\right)^{4}$ if $|\Delta \kappa| \ll$ 1.
${ }^{11}$ HERZOG 84 consider the contribution of $W$-boson to muon magnetic moment including anomalous coupling of $W W \gamma$. Obtain a limit $-1<\Delta \kappa<3$ for $\Lambda \gtrsim 1 \mathrm{TeV}$.

## ANOMALOUS W/Z QUARTIC COUPLINGS

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## $a_{0} / \Lambda^{2}, a_{c} / \Lambda^{2}, a_{n} / \Lambda^{2}, \kappa_{0}^{W} / \Lambda^{2}, \kappa_{c}^{W} / \Lambda^{2}, f_{T, 0} / \Lambda^{4}, f_{M, i} / \Lambda^{4}, \alpha_{4}, \alpha_{5}$, $\mathrm{F}_{S, i} / \Lambda^{4}, \mathrm{~F}_{M, i} / \Lambda^{4}, \mathrm{~F}_{T, i} / \Lambda^{4}$

Anomalous $W$ quartic couplings are measured by the experiments at LEP, the Tevatron, and the LHC. Some of the recent results from the Tevatron and LHC experiments individually surpass the combined LEP-2 results in precision (see below). As discussed in the review on the "Anomalous $W / Z$ quartic couplings (QGCS)," the measurements are typically done using different operator expansions which then do not allow the results to be compared and averaged. At least one common framework should be agreed upon for the use in the future publications by the experiments.

## VALUE <br> DOCUMENT ID TECN

-     - We do not use the following data for averages, fits, limits, etc. • • -

|  | AAD | 15N | A |
| :---: | :---: | :---: | :---: |
|  | KHACHATR |  | CMS |
|  | AAD |  | ATLS |
|  | CHATRCHYAN | 14Q | CMS |
|  | ABAZOV | 13D | D0 |
|  | CHATRCHYA | 13 | CMS |
|  | ABBIENDI | 04B | OPAL |
|  | ABBIENDI | 04L | OP |
|  | HEISTER | 04A | AL |
|  | ABDALLAH | 031 | DL |
|  | ACHARD | 02F | L3 |

${ }^{1}$ AAD 15 N study $W \gamma \gamma$ events in $8 \mathrm{TeV} p p$ interactions, where the $W$ decays into an electron or a muon. The events are characterized by an isolated lepton, a missing transverse energy due to the decay neutrino, and two isolated photons, with the $p_{T}$ of the lepton and the photons being $>20 \mathrm{GeV}$. The number of candidate events observed in the electron channel for $N($ jet $) \geq 0$ and $N(j e t)=0$ is 47 and 15 , the corresponding numbers for the muon channel being 110 and 53. The backgrounds expected are $30.2 \pm 7.4$, $8.7 \pm 3.0,52.1 \pm 12.2$, and $24.4 \pm 8.3$ respectively. The $95 \%$ C.L. limits on the values of the parameters $f_{T, 0} / \Lambda^{4}, f_{M, 2} / \Lambda^{4}$ and $f_{M, 3} / \Lambda^{4}$ are $-0.9-0.9 \times 10^{2},-0.8-0.8 \times 10^{4}$, and $-1.5-1.4 \times 10^{4}$ respectively, without application of a form factor $\Lambda_{\mathrm{FF}}$.

2 KHACHATRYAN 15D study vector-boson-scattering tagged by two jets, requiring two same-sign charged leptons arising from $W^{ \pm} W^{ \pm}$production and decay. The two jets must have a transverse momentum larger than 30 GeV , while the leptons, electrons or muons, must have a transverse momentum $>20 \mathrm{GeV}$. The dijet mass is required to be $>$ 500 GeV , the dilepton mass $>50 \mathrm{GeV}$, with additional requirement of differing from the $Z$ mass by $>15 \mathrm{GeV}$. In the two categories $W^{+} W^{+}$and $W^{-} W^{-}, 10$ and 2 data events are observed in a data sample corresponding to an integrated luminosity of $19.4 \mathrm{fb}^{-1}$, with an expected background of $3.1 \pm 0.6$ and $2.6 \pm 0.5$ events. Analysing the distribution of the dilepton invariant mass, the following limits at 95\% C.L. are obtained, in units of $\mathrm{TeV}^{-4}:-38<\mathrm{F}_{S, 0} / \Lambda^{4}<40,-118<\mathrm{F}_{S, 1} / \Lambda^{4}<120,-33<\mathrm{F}_{M, 0} / \Lambda^{4}<32$, $-44<\mathrm{F}_{M, 1} / \Lambda^{4}<47,-65<\mathrm{F}_{M, 6} / \Lambda^{4}<63,-70<\mathrm{F}_{M, 7} / \Lambda^{4}<66,-4.2<$ $\mathrm{F}_{T, 0} / \Lambda^{4}<4.6,-1.9<\mathrm{F}_{T, 1} / \Lambda^{4}<2.2,-6.2<\mathrm{F}_{T, 2} / \Lambda^{4}<6.4$.
${ }^{3}$ AAD 14AM analyze electroweak production of $W W$ jet jet same-charge diboson plus two jets production, with the $W$ bosons decaying to electron or muon, to study the quartic $W W W W$ coupling. In a kinematic region enhancing the electroweak production over the strong production, 34 events are observed in the data while $29.8 \pm 2.4$ events are expected with a backgound of $15.9 \pm 1.9$ events. Assuming the other QGC coupling to have the SM value of zero, the observed event yield is used to determine $95 \%$ CL limits on the quartic gauge couplings: $-0.14<\alpha_{4}<0.16$ and $-0.23<\alpha_{5}<0.24$.
${ }^{4}$ CHATRCHYAN 14Q study $W V \gamma$ production in $8 \mathrm{TeV} p p$ collisions, in the single lepton final state, with $W \rightarrow \ell \nu, Z \rightarrow \operatorname{dijet}$ or $W \rightarrow \ell \nu, W \rightarrow$ dijet, the dijet mass resolution precluding differentiation between the $W$ and $Z . p_{T}$ and pseudo-rapidity cuts are put on the lepton, the photon and the two jets to minimize backgrounds. The dijet mass is required to be between $7-100 \mathrm{GeV}$ and $\left|\Delta \eta_{j j}\right|<1.4$. The selected number of muon (electron) events are 183 (139), with SM expectation being $194.2 \pm 11.5$ (147.9 $\pm 10.7$ ) including signal and background. The photon $E_{T}$ distribution is used to set limits on the anomalous quartic couplings. The following $95 \%$ CL limits are deduced (all in units of $\mathrm{TeV}^{-2}$ or $\left.\mathrm{TeV}^{-4}\right):-21<a_{0}^{W} / \Lambda^{2}<20,-34<a_{c}^{W} / \Lambda^{2}<32,-12<\kappa_{0}^{W} / \Lambda^{2}<$ 10 and $-18<\kappa_{c}^{W} / \Lambda^{2}<17$; and $-25<f_{T, 0} / \Lambda^{4}<24 \mathrm{TeV}^{-4}$.
${ }^{5}$ ABAZOV 13D searches for anomalous $W W \gamma \gamma$ quartic gauge couplings in the two-photon-mediated process $p p \rightarrow p p W W$, assuming the $W W \gamma$ triple gauge boson couplings to be at their Standard Model values. 946 events containing an $e^{+} e^{-}$pair with missing energy are selected in a total luminosity of $9.7 \mathrm{fb}^{-1}$, with an expectation of $983 \pm 108$ events from Standard-Model processes. The following 1-parameter limits at $95 \% \mathrm{CL}$ are otained: $\left|a_{0}^{W} / \Lambda^{2}\right|<4.3 \times 10^{-4} \mathrm{GeV}^{-2}\left(a_{c}^{W}=0\right),\left|a_{c}^{W} / \Lambda^{2}\right|<$ $1.5 \times 10^{-3} \mathrm{GeV}^{-2}\left(a_{0}^{W}=0\right)$.
${ }^{6}$ CHATRCHYAN 13AA searches for anomalous $W W \gamma \gamma$ quartic gauge couplings in the two-photon-mediated process $p p \rightarrow p p W W$, assuming the $W W \gamma$ triple gauge boson couplings to be at their Standard Model values. 2 events containing an $e^{ \pm} \mu^{\mp}$ pair with $p_{T}(e, \mu)>30 \mathrm{GeV}$ are selected in a total luminosity of $5.05 \mathrm{fb}^{-1}$, with an expected $p p W W$ signal of $2.2 \pm 0.4$ events and an expected background of $0.84 \pm 0.15$ events. The following 1-parameter limits at $95 \% \mathrm{CL}$ are otained from the $p_{T}(e, \mu)$ spectrum: $\left|a_{0}^{W} / \Lambda^{2}\right|<4.0 \times 10^{-6} \mathrm{GeV}^{-2}\left(a_{c}^{W}=0\right),\left|a_{c}^{W} / \Lambda^{2}\right|<1.5 \times 10^{-5} \mathrm{GeV}^{-2}\left(a_{0}^{W}\right.$ $=0$ ).
${ }^{7}$ ABBIENDI 04B select $187 e^{+} e^{-} \rightarrow W^{+} W^{-} \gamma$ events in the C.M. energy range $180-209 \mathrm{GeV}$, where $E_{\gamma}>2.5 \mathrm{GeV}$, the photon has a polar angle $\left|\cos \theta_{\gamma}\right|<0.975$ and is well isolated from the nearest jet and charged lepton, and the effective masses of both fermion-antifermion systems agree with the $W$ mass within $3 \Gamma_{W}$. The measured differential cross section as a function of the photon energy and photon polar angle is used to extract the $95 \% \mathrm{CL}$ limits: $-0.020 \mathrm{GeV}^{-2}<a_{0} / \Lambda^{2}<0.020 \mathrm{GeV}^{-2}$, $-0.053 \mathrm{GeV}^{-2}<a_{c} / \Lambda^{2}<0.037 \mathrm{GeV}^{-2}$ and $-0.16 \mathrm{GeV}^{-2}<a_{n} / \Lambda^{2}<0.15 \mathrm{GeV}^{-2}$.
${ }^{8}$ ABBIENDI 04L select $20 e^{+} e^{-} \rightarrow \nu \bar{\nu} \gamma \gamma$ acoplanar events in the energy range 180-209 GeV and $176 e^{+} e^{-} \rightarrow q \bar{q} \gamma \gamma$ events in the energy range $130-209 \mathrm{GeV}$. These samples are used to constrain possible anomalous $W^{+} W^{-} \gamma \gamma$ and $Z Z \gamma \gamma$ quartic couplings. Further combining with the $W^{+} W^{-} \gamma$ sample of ABBIENDI 04B the following oneparameter $95 \%$ CL limits are obtained: $-0.007<a_{0}^{Z} / \Lambda^{2}<0.023 \mathrm{GeV}^{-2},-0.029<$ $a_{c}^{Z} / \Lambda^{2}<0.029 \mathrm{GeV}^{-2},-0.020<a_{0}^{W} / \Lambda^{2}<0.020 \mathrm{GeV}^{-2},-0.052<a_{c}^{W} / \Lambda^{2}<$ $0.037 \mathrm{GeV}^{-2}$.
9 In the CM energy range 183 to 209 GeV HEISTER 04A select $30 e^{+} e^{-} \rightarrow \nu \bar{\nu} \gamma \gamma$ events with two acoplanar, high energy and high transverse momentum photons. The photonphoton acoplanarity is required to be $>5^{\circ}, E_{\gamma} / \sqrt{s}>0.025$ (the more energetic photon having energy $>0.2 \sqrt{s}$ ), $\mathrm{p}_{T_{\gamma}} / \mathrm{E}_{\text {beam }}>0.05$ and $\left|\cos \theta_{\gamma}\right|<0.94$. A likelihood fit to the photon energy and recoil missing mass yields the following one-parameter 95\% CL limits: $-0.012<a_{0}^{Z} / \Lambda^{2}<0.019 \mathrm{GeV}^{-2},-0.041<a_{c}^{Z} / \Lambda^{2}<0.044 \mathrm{GeV}^{-2}$, $-0.060<a_{0}^{W} / \Lambda^{2}<0.055 \mathrm{GeV}^{-2},-0.099<a_{c}^{W} / \Lambda^{2}<0.093 \mathrm{GeV}^{-2}$.
10 ABDALLAH 03 I select $122 e^{+} e^{-} \rightarrow W^{+} W^{-} \gamma$ events in the C.M. energy range 189-209 GeV, where $E_{\gamma}>5 \mathrm{GeV}$, the photon has a polar angle $\left|\cos \theta_{\gamma}\right|<0.95$ and is well isolated from the nearest charged fermion. A fit to the photon energy spectra yields $a_{c} / \Lambda^{2}=0.000_{-0.040}^{+0.019} \mathrm{GeV}^{-2}, a_{0} / \Lambda^{2}=-0.004_{-0.010}^{+0.018} \mathrm{GeV}^{-2}, \widetilde{a}_{0} / \Lambda^{2}=$ $-0.007_{-0.008}^{+0.019} \mathrm{GeV}^{-2}, a_{n} / \Lambda^{2}=-0.09_{-0.05}^{+0.16} \mathrm{GeV}^{-2}$, and $\widetilde{a}_{n} / \Lambda^{2}=+0.05_{-0.15}^{+0.07}$ $\mathrm{GeV}^{-2}$, keeping the other parameters fixed to their Standard Model values (0). The $95 \% \mathrm{CL}$ limits are: $-0.063 \mathrm{GeV}^{-2}<a_{C} / \Lambda^{2}<+0.032 \mathrm{GeV}^{-2},-0.020$ $\mathrm{GeV}^{-2}<a_{0} / \Lambda^{2}<+0.020 \mathrm{GeV}^{-2},-0.020 \mathrm{GeV}^{-2}<\widetilde{a}_{0} / \Lambda^{2}<+0.020 \mathrm{GeV}^{-2}$, $-0.18 \mathrm{GeV}^{-2}<a_{n} / \Lambda^{2}<+0.14 \mathrm{GeV}^{-2},-0.16 \mathrm{GeV}^{-2}<\widetilde{a}_{n} / \Lambda^{2}<+0.17 \mathrm{GeV}^{-2}$.
11 ACHARD 02F select $86 e^{+} e^{-} \rightarrow W^{+} W^{-} \gamma$ events at $192-207 \mathrm{GeV}$, where $E_{\gamma}>5$ GeV and the photon is well isolated. They also select 43 acoplanar $e^{+} e^{-} \rightarrow \nu \bar{\nu} \gamma \gamma$ events in this energy range, where the photon energies are $>5 \mathrm{GeV}$ and $>1 \mathrm{GeV}$ and the photon polar angles are between $14^{\circ}$ and $166^{\circ}$. All these 43 events are in the recoil mass region corresponding to the $Z(75-110 \mathrm{GeV})$. Using the shape and normalization of the photon spectra in the $W^{+} W^{-} \gamma$ events, and combining with the 42 event sample from 189 GeV data (ACCIARRI 00T), they obtain: $a_{0} / \Lambda^{2}=0.000 \pm 0.010 \mathrm{GeV}^{-2}, a_{c} / \Lambda^{2}=$ $-0.013 \pm 0.023 \mathrm{GeV}^{-2}$, and $a_{n} / \Lambda^{2}=-0.002 \pm 0.076 \mathrm{GeV}^{-2}$. Further combining the analyses of $W^{+} W^{-} \gamma$ events with the low recoil mass region of $\nu \bar{\nu} \gamma \gamma$ events (including samples collected at $183+189 \mathrm{GeV}$ ), they obtain the following one-parameter $95 \% \mathrm{CL}$ limits: $-0.015 \mathrm{GeV}^{-2}<a_{0} / \Lambda^{2}<0.015 \mathrm{GeV}^{-2},-0.048 \mathrm{GeV}^{-2}<a_{C} / \Lambda^{2}<0.026$ $\mathrm{GeV}^{-2}$, and $-0.14 \mathrm{GeV}^{-2}<a_{n} / \Lambda^{2}<0.13 \mathrm{GeV}^{-2}$.

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| KHACHATRY. | 15D | PRL 114051801 | V. Khachatryan et al. | (CMS Collab.) |
| AAD | 14AM | PRL 113141803 | G. Aad et al. | (ATLAS Collab.) |
| AAD | 14Y | JHEP 1404031 | G. Aad et al. | (ATLAS Collab.) |
| AALTONEN | 14D | PR D89 072003 | T. Aaltonen et al. | (CDF Collab.) |
| ABAZOV | 14N | PR D89 012005 | V.M. Abazov et al. | (D0 Collab.) |
| CHATRCHYAN | 14AB | PR D89 092005 | S. Chatrchyan et al. | (CMS Collab.) |
| CHATRCHYAN | 14Q | PR D90 032008 | S. Chatrchyan et al. | (CMS Collab.) |
| AAD | 13AL | PR D87 112001 | G. Aad et al. | (ATLAS Collab.) |
| Also |  | PR D88 079906 (errat.) | G. Aad et al. | (ATLAS Collab.) |
| AAD | 13AN | PR D87 112003 | G. Aad et al. | (ATLAS Collab.) |
| Also |  | PR D91 119901 (errat.) | G. Aad et al. | (ATLAS Collab.) |
| AALTONEN | 13N | PR D88 052018 | T. Aaltonen et al. | (CDF and D0 Collabs.) |


| ABAZOV | 13D | PR D88 012005 | V.M. Abazov et al. | (D0 Collab.) |
| :---: | :---: | :---: | :---: | :---: |
| CHATRCHYAN | 13AA | JHEP 1307116 | S. Chatrchyan et al. | (CMS Collab.) |
| CHATRCHYAN | 13BF | EPJ C73 2610 | S. Chatrchyan et al. | (CMS Collab.) |
| SCHAEL | 13A | PRPL 532119 | S. Schael et al. ( | (ALEPH Collab., DELPHI, L3+) |
| AAD | 12AC | PL B712 289 | G. Aad et al. | (ATLAS Collab.) |
| AAD | 12BX | PL B717 49 | G. Aad et al. | (ATLAS Collab.) |
| AAD | 12CD | EPJ C72 2173 | G. Aad et al. | (ATLAS Collab.) |
| AAD | 12V | PL B709 341 | G. Aad et al. | (ATLAS Collab.) |
| AALTONEN | 12AC | PR D86 031104 | T. Aaltonen et al. | (CDF Collab.) |
| AALTONEN | 12E | PRL 108151803 | T. Aaltonen et al. | (CDF Collab.) |
| AALTONEN | 12W | PR D85 032001 | T. Aaltonen et al. | (CDF Collab.) |
| ABAZOV | 12AG | PL B718 451 | V.M. Abazov et al. | (D0 Collab.) |
| ABAZOV | 12 F | PRL 108151804 | V.M. Abazov et al. | (D0 Collab.) |
| ABAZOV | 11 | PL B695 67 | V.M. Abazov et al. | (D0 Collab.) |
| ABAZOV | 11AC | PRL 107241803 | V.M. Abazov et al. | (D0 Collab.) |
| BLINOV | 11 | PL B699 287 | A.E. Blinov, A.S. Rudenko | o (NOVO) |
| CHATRCHYAN | 11M | PL B701 535 | S. Chatrchyan et al. | (CMS Collab.) |
| AALTONEN | 10K | PRL 104201801 | T. Aaltonen et al. | (CDF Collab.) |
| Also |  | PRL 105 019905(errat.) | T. Aaltonen et al. | (CDF Collab.) |
| ABDALLAH | 10 | EPJ C66 35 | J. Abdallah et al. | (DELPHI Collab.) |
| AARON | 09B | EPJ C64 251 | F.D. Aaron et al. | (H1 Collab.) |
| ABAZOV | 09AB | PRL 103141801 | V.M. Abazov et al. | (D0 Collab.) |
| ABAZOV | 09AD | PR D80 053012 | V.M. Abazov et al. | (D0 Collab.) |
| ABAZOV | 09AJ | PRL 103191801 | V.M. Abazov et al. | (D0 Collab.) |
| ABAZOV | 09AK | PRL 103231802 | V.M. Abazov et al. | (D0 Collab.) |
| AALTONEN | 08B | PRL 100071801 | T. Aaltonen et al. | (CDF Collab.) |
| ABAZOV | 08R | PRL 100241805 | V.M. Abazov et al. | (D0 Collab.) |
| ABDALLAH | 08A | EPJ C55 1 | J. Abdallah et al. | (DELPHI Collab.) |
| ABDALLAH | 08C | EPJ C54 345 | J. Abdallah et al. | (DELPHI Collab.) |
| AALTONEN | 07F | PRL 99151801 | T. Aaltonen et al. | (CDF Collab.) |
| Also |  | PR D77 112001 | T. Aaltonen et al. | (CDF Collab.) |
| AALTONEN | 07L | PR D76 111103 | T. Aaltonen et al. | (CDF Collab.) |
| ABAZOV | 07Z | PR D76 111104 | V.M. Abazov et al. | (D0 Collab.) |
| ABBIENDI | 07A | EPJ C52 767 | G. Abbiendi et al. | (OPAL Collab.) |
| ABAZOV | 06H | PR D74 057101 | V.M. Abazov et al. | (D0 Collab.) |
| Also |  | PR D74 059904(errat.) | V.M. Abazov et al. | (D0 Collab.) |
| ABBIENDI | 06 | EPJ C45 307 | G. Abbiendi et al. | (OPAL Collab.) |
| ABBIENDI | 06A | EPJ C45 291 | G. Abbiendi et al. | (OPAL Collab.) |
| ACHARD | 06 | EPJ C45 569 | P. Achard et al. | (L3 Collab.) |
| AKTAS | 06 | PL B632 35 | A. Aktas et al. | (H1 Collab.) |
| SCHAEL | 06 | EPJ C47 309 | S. Schael et al. | (ALEPH Collab.) |
| ABAZOV | 05J | PR D71 091108 | V.M. Abazov et al. | (D0 Collab.) |
| ABAZOV | 05S | PRL 95141802 | V.M. Abazov et al. | (D0 Collab.) |
| SCHAEL | 05A | PL B614 7 | S. Schael et al. | (ALEPH Collab.) |
| ABAZOV | 04D | PR D70 092008 | V.M. Abazov et al. | (CDF Collab., D0 Collab.) |
| ABBIENDI | 04B | PL B580 17 | G. Abbiendi et al. | (OPAL Collab.) |
| ABBIENDI | 04D | EPJ C33 463 | G. Abbiendi et al. | (OPAL Collab.) |
| ABBIENDI | 04L | PR D70 032005 | G. Abbiendi et al. | (OPAL Collab.) |
| ABDALLAH | 04G | EPJ C34 127 | J. Abdallah et al. | (DELPHI Collab.) |
| ACHARD | 04D | PL B586 151 | P. Achard et al. | (L3 Collab.) |
| ACHARD | 04J | PL B600 22 | P. Achard et al. | (L3 Collab.) |
| HEISTER | 04A | PL B602 31 | A. Heister et al. | (ALEPH Collab.) |
| SCHAEL | 04A | EPJ C38 147 | S. Schael et al. | (ALEPH Collab.) |
| ABBIENDI | 03C | EPJ C26 321 | G. Abbiendi et al. | (OPAL Collab.) |
| ABDALLAH | 031 | EPJ C31 139 | J. Abdallah et al. | (DELPHI Collab.) |
| ABAZOV | 02D | PR D66 012001 | V.M. Abazov et al. | (D0 Collab.) |
| ABAZOV | 02E | PR D66 032008 | V.M. Abazov et al. | (D0 Collab.) |
| ACHARD | 02F | PL B527 29 | P. Achard et al. | (L3 Collab.) |
| CHEKANOV | 02C | PL B539 197 | S. Chekanov et al. | (ZEUS Collab.) |
| ABBIENDI | 01H | EPJ C19 229 | G. Abbiendi et al. | (OPAL Collab.) |
| ABREU | 011 | PL B502 9 | P. Abreu et al. | (DELPHI Collab.) |
| AFFOLDER | 01E | PR D64 052001 | T. Affolder et al. | (CDF Collab.) |
| ABBIENDI | 00V | PL B490 71 | G. Abbiendi et al. | (OPAL Collab.) |
| ABBOTT | 00B | PR D61 072001 | B. Abbott et al. | (D0 Collab.) |
| ABBOTT | 00D | PRL 845710 | B. Abbott et al. | (D0 Collab.) |
| ABREU,P | 00F | EPJ C18 203 | P. Abreu et al. | (DELPHI Collab.) |
| Also |  | EPJ C25 493 (errat.) | P. Abreu et al. | (DELPHI Collab.) |
| ACCIARRI | 00T | PL B490 187 | M. Acciarri et al. | (L3 Collab.) |
| AFFOLDER | 00M | PRL 853347 | T. Affolder et al. | (CDF Collab.) |
| BREITWEG | 00 | PL B471 411 | J. Breitweg et al. | (ZEUS Collab.) |
| BREITWEG | 00D | EPJ C12 411 | J. Breitweg et al. | (ZEUS Collab.) |


| EBOLI | 00 | MPL A15 1 | O. Eboli, M. Gonzalez-Garcia, | aes |
| :---: | :---: | :---: | :---: | :---: |
| ABBIENDI | 99N | PL B453 153 | G. Abbiendi et al. | (OPAL Collab.) |
| ABBOTT | 99H | PR D60 052003 | B. Abbott et al. | (D0 Collab.) |
| ABBOTT | 991 | PR D60 072002 | B. Abbott et al. | (D0 Collab.) |
| ABREU | 99L | PL B459 382 | P. Abreu et al. | (DELPHI Collab.) |
| ACCIARRI | 99 | PL B454 386 | M. Acciarri et al. | (L3 Collab.) |
| ACCIARRI | 99Q | PL B467 171 | M. Acciarri et al. | (L3 Collab.) |
| BARATE | 991 | PL B453 107 | R. Barate et al. | (ALEPH Collab.) |
| BARATE | 99L | PL B462 389 | R. Barate et al. | (ALEPH Collab.) |
| BARATE | 99M | PL B465 349 | R. Barate et al. | (ALEPH Collab.) |
| ABBOTT | 98N | PR D58 092003 | B. Abbott et al. | (D0 Collab.) |
| ABBOTT | 98P | PR D58 012002 | B. Abbott et al. | (D0 Collab.) |
| ABE | 98 H | PR D58 031101 | F. Abe et al. | (CDF Collab.) |
| ABE | 98P | PR D58 091101 | F. Abe et al. | (CDF Collab.) |
| ABREU | 98C | PL B416 233 | P. Abreu et al. | (DELPHI Collab.) |
| ABREU | 98N | PL B439 209 | P. Abreu et al. | (DELPHI Collab.) |
| BARATE | 97 | PL B401 347 | R. Barate et al. | (ALEPH Collab.) |
| BARATE | 97S | PL B415 435 | R. Barate et al. | (ALEPH Collab.) |
| ABACHI | 95D | PRL 751456 | S. Abachi et al. | (D0 Collab.) |
| ABE | 95C | PRL 74341 | F. Abe et al. | (CDF Collab.) |
| ABE | 95G | PRL 741936 | F. Abe et al. | (CDF Collab.) |
| ABE | 95P | PRL 7511 | F. Abe et al. | (CDF Collab.) |
| Also |  | PR D52 4784 | F. Abe et al. | (CDF Collab.) |
| ABE | 95W | PR D52 2624 | F. Abe et al. | (CDF Collab.) |
| Also |  | PRL 73220 | F. Abe et al. | (CDF Collab.) |
| ABE | 92E | PRL 683398 | F. Abe et al. | (CDF Collab.) |
| ABE | 921 | PRL 6928 | F. Abe et al. | (CDF Collab.) |
| ALITTI | 92 | PL B276 365 | J. Alitti et al. | (UA2 Collab.) |
| ALITTI | 92B | PL B276 354 | J. Alitti et al. | (UA2 Collab.) |
| ALITTI | 92C | PL B277 194 | J. Alitti et al. | (UA2 Collab.) |
| ALITTI | 92D | PL B277 203 | J. Alitti et al. | (UA2 Collab.) |
| ALITTI | 92F | PL B280 137 | J. Alitti et al. | (UA2 Collab.) |
| SAMUEL | 92 | PL B280 124 | M.A. Samuel et al. | (OKSU, CARL) |
| ABE | 91 C | PR D44 29 | F. Abe et al. | (CDF Collab.) |
| ALBAJAR | 91 | PL B253 503 | C. Albajar et al. | (UA1 Collab.) |
| ALITTI | 91 C | ZPHY C52 209 | J. Alitti et al. | (UA2 Collab.) |
| SAMUEL | 91 | PRL 679 | M.A. Samuel et al. | (OKSU, CARL) |
| Also |  | PRL 672920 (erratum) | M.A. Samuel et al. |  |
| ABE | 90G | PRL 652243 | F. Abe et al. | (CDF Collab.) |
| Also |  | PR D43 2070 | F. Abe et al. | (CDF Collab.) |
| ALBAJAR | 90 | PL B241 283 | C. Albajar et al. | (UA1 Collab.) |
| ALITTI | 90B | PL B241 150 | J. Alitti et al. | (UA2 Collab.) |
| ABE | 891 | PRL 621005 | F. Abe et al. | (CDF Collab.) |
| ALBAJAR | 89 | ZPHY C44 15 | C. Albajar et al. | (UA1 Collab.) |
| BAUR | 88 | NP B308 127 | U. Baur, D. Zeppenfeld | (FSU, WISC) |
| GRIFOLS | 88 | IJMP A3 225 | J.A. Grifols, S. Peris, J. Sola | (BARC, DESY) |
| Also |  | PL B197 437 | J.A. Grifols, S. Peris, J. Sola | (BARC, DESY) |
| ALBAJAR | 87 | PL B185 233 | C. Albajar et al. | (UA1 Collab.) |
| ANSARI | 87 | PL B186 440 | R. Ansari et al. | (UA2 Collab.) |
| GROTCH | 87 | PR D36 2153 | H. Grotch, R.W. Robinett | (PSU) |
| HAGIWARA | 87 | NP B282 253 | K. Hagiwara et al. | (KEK, UCLA, FSU) |
| VANDERBIJ | 87 | PR D35 1088 | J.J. van der Bij | (FNAL) |
| GRAU | 85 | PL 154B 283 | A. Grau, J.A. Grifols | (BARC) |
| SUZUKI | 85 | PL 153B 289 | M. Suzuki | (LBL) |
| ARNISON | 84D | PL 134B 469 | G.T.J. Arnison et al. | (UA1 Collab.) |
| HERZOG | 84 | PL 148B 355 | F. Herzog | (WISC) |
| Also |  | PL 155B 468 (erratum) | F. Herzog | (WISC) |
| ARNISON | 83 | PL 122B 103 | G.T.J. Arnison et al. | (UA1 Collab.) |
| BANNER | 83B | PL 122B 476 | M. Banner et al. | (UA2 Collab.) |


[^0]:    ${ }^{22}$ ABBOTT 99I perform a simultaneous fit to the $W \gamma, W W \rightarrow$ dilepton, $W W / W Z \rightarrow$ $e \nu j j, W W / W Z \rightarrow \mu \nu j j$, and $W Z \rightarrow$ trilepton data samples. For $\Lambda=2.0 \mathrm{TeV}$, the $95 \%$ CL limits are $0.75<\kappa_{\gamma}<1.39$.

