## THE $\rho(1450)$ AND THE $\rho(1700)$

Updated May 2010 by S. Eidelman (Novosibirsk) and G. Venanzoni (Frascati).

In our 1988 edition, we replaced the  $\rho(1600)$  entry with two new ones, the  $\rho(1450)$  and the  $\rho(1700)$ , because there was emerging evidence that the 1600-MeV region actually contains two  $\rho$ -like resonances. Erkal [1] had pointed out this possibility with a theoretical analysis on the consistency of  $2\pi$  and  $4\pi$ electromagnetic form factors and the  $\pi\pi$  scattering length. Donnachie [2], with a full analysis of data on the  $2\pi$  and  $4\pi$ final states in  $e^+e^-$  annihilation and photoproduction reactions, had also argued that in order to obtain a consistent picture, two resonances were necessary. The existence of  $\rho(1450)$  was supported by the analysis of  $\eta\rho^0$  mass spectra obtained in photoproduction and  $e^+e^-$  annihilation [3], as well as that of  $e^+e^- \to \omega\pi$  [4].

The analysis of [2] was further extended by [5,6] to include new data on  $4\pi$ -systems produced in  $e^+e^-$  annihilation, and in  $\tau$ -decays ( $\tau$  decays to  $4\pi$ , and  $e^+e^-$  annihilation to  $4\pi$  can be related by the Conserved Vector Current assumption). These systems were successfully analyzed using interfering contributions from two  $\rho$ -like states, and from the tail of the  $\rho(770)$ decaying into two-body states. While specific conclusions on  $\rho(1450) \rightarrow 4\pi$  were obtained, little could be said about the  $\rho(1700)$ .

Independent evidence for two 1<sup>-</sup> states is provided by [7] in  $4\pi$  electroproduction at  $\langle Q^2 \rangle = 1$  (GeV/c)<sup>2</sup>, and by [8] in a high-statistics sample of the  $\eta\pi\pi$  system in  $\pi^-p$  charge exchange.

This scenario with two overlapping resonances is supported by other data. Bisello [9] measured the pion form factor in the interval 1.35–2.4 GeV, and observed a deep minimum around 1.6 GeV. The best fit was obtained with the hypothesis of  $\rho$ -like resonances at 1420 and 1770 MeV, with widths of about 250 MeV. Antonelli [10] found that the  $e^+e^- \rightarrow \eta \pi^+ \pi^-$  cross section is better fitted with two fully interfering Breit-Wigners, with parameters in fair agreement with those of [2] and [9]. These results can be considered as a confirmation of the  $\rho(1450)$ . Decisive evidence for the  $\pi\pi$  decay mode of both  $\rho(1450)$ and  $\rho(1700)$  comes from  $\overline{p}p$  annihilation at rest [11]. It has been shown that these resonances also possess a  $K\overline{K}$ decay mode [12–14]. High-statistics studies of the decays  $\tau \to \pi\pi\nu_{\tau}$  [15,16], and  $\tau \to 4\pi\nu_{\tau}$  [17] also require the  $\rho(1450)$ , but are not sensitive to the  $\rho(1700)$ , because it is too close to the  $\tau$  mass. A recent very-high-statistics study of the  $\tau \to \pi\pi\nu_{\tau}$ decay performed at Belle [18] reports the first observation of both  $\rho(1450)$  and  $\rho(1700)$  in  $\tau$  decays.

The structure of these  $\rho$  states is not yet completely clear. Barnes [19] and Close [20] claim that  $\rho(1450)$  has a mass consistent with radial 2*S*, but its decays show characteristics of hybrids, and suggest that this state may be a 2*S*-hybrid mixture. Donnachie [21] argues that hybrid states could have a  $4\pi$  decay mode dominated by the  $a_1\pi$ . Such behavior has been observed by [22] in  $e^+e^- \rightarrow 4\pi$  in the energy range 1.05–1.38 GeV, and by [17] in  $\tau \rightarrow 4\pi$  decays. Alexander [23] observes the  $\rho(1450) \rightarrow \omega\pi$  decay mode in *B*-meson decays, however, does not find  $\rho(1700) \rightarrow \omega\pi^0$ . A similar conclusion is made by [24], who studied the process  $e^+e^- \rightarrow \omega\pi^0$ . Various decay modes of the  $\rho(1450)$  and  $\rho(1700)$  are observed in  $\overline{p}n$  and  $\overline{p}p$ annihilation [25,26], but no definite conclusions can be drawn. More data should be collected to clarify the nature of the  $\rho$ states, particularly in the energy range above 1.6 GeV.

We now list under a separate entry the  $\rho(1570)$ , the  $\phi\pi$ state with  $J^{PC} = 1^{--}$  earlier observed by [27] (referred to as C(1480)) and recently confirmed by [28]. While [29] shows that it may be a threshold effect, [5] and [30] suggest two independent vector states with this decay mode. The C(1480)has not been seen in the  $\bar{p}p$  [31] and  $e^+e^-$  [32,33] experiments. However, the sensitivity of the two latter is an order of magnitude lower than that of [28]. Note that [28] can not exclude that their observation is due to an OZI-suppressed decay mode of the  $\rho(1700)$ .

Several observations on the  $\omega\pi$  system in the 1200-MeV region [34–40] may be interpreted in terms of either  $J^P = 1^ \rho(770) \rightarrow \omega\pi$  production [41], or  $J^P = 1^+ b_1(1235)$  production [39,40]. We argue that no special entry for a  $\rho(1250)$  is needed. The LASS amplitude analysis [42] showing evidence for  $\rho(1270)$  is preliminary and needs confirmation. For completeness, the relevant observations are listed under the  $\rho(1450)$ .

Recently [43] reported a very broad  $1^{--}$  resonance-like  $K^+K^-$  state in  $J/\psi \to K^+K^-\pi^0$  decays. Its pole position corresponds to mass of 1576 MeV and width of 818 MeV. [44–46] suggest its exotic structure (molecular or multiquark), while [47] and [48] explain it by the interference between the  $\rho(1450)$  and  $\rho(1700)$ . We quote [43] as X(1575) in the section "Further States."

Evidence for  $\rho$ -like mesons decaying into  $6\pi$  states was first noted by [49] in the analysis of  $6\pi$  mass spectra from  $e^+e^-$  annihilation [50,51] and diffractive photoproduction [52]. Clegg [49] argued that two states at about 2.1 and 1.8 GeV exist: while the former is a candidate for the  $\rho(2150)$ , the latter could be a manifestation of the  $\rho(1700)$  distorted by threshold effects. BaBar reported observations of the new decay modes of the  $\rho(2150)$  in the channels  $\eta'(958)\pi^+\pi^-$  and  $f_1(1285)\pi^+\pi^-$ [53]. The relativistic quark model [54] predicts the  $2^3D_1$  state with  $J^{PC} = 1^{--}$  at 2.15 GeV which can be identified with the  $\rho(2150)$ .

The E687 Collaboration at Fermilab reported an observation of a narrow-dip structure at 1.9 GeV in the  $3\pi^+3\pi^-$  diffractive photoproduction [55]. A similar effect of the dip in the cross section of  $e^+e^- \rightarrow 6\pi$  around 1.9 GeV has been earlier reported by DM2 [51], where  $6\pi$  included both  $3\pi^+3\pi^-$  and  $2\pi^+2\pi^-2\pi^0$ . Later the dip in the R value (the total cross section of  $e^+e^- \rightarrow$ hadrons divided by the cross section of  $e^+e^- \rightarrow \mu^+\mu^-$ ) was observed by [56], again around 1.9 GeV. This energy is close to the  $N\overline{N}$  threshold, which hints at the possible relation between the dip and  $N\overline{N}$ , e.q., the frequently discussed narrow  $N\overline{N}$  resonance or just a threshold effect. Such behaviour is also characteristic of exotic objects like vector  $q\overline{q}$  hybrids. Note that [57] failed to find this state in the reaction  $\overline{n}p \to 3\pi^+ 2\pi^- \pi^0$ . A reanalysis of the E687 data by [58] shows that a dip may arise due to interference of a narrow object with a broad  $\rho(1700)$ independently of the nature of the former. BaBar studied the processes  $e^+e^- \rightarrow 3\pi^+3\pi^-$  and  $e^+e^- \rightarrow 2\pi^+2\pi^-2\pi^0$  using the radiative return, and observed a structure around 1.9 GeV in both final states [59]. The data are not well described by a single Breit-Wigner state, and a good fit is achieved while taking into account the interference of such a structure with a Jacob-Slansky amplitude for continuum. The mass of this state obtained by BaBar is consistent with [56] and [55], but the width is substantially larger. Recently [28] observed a structure at 1.9 GeV in the radiative return to the  $\phi\pi$  final state, with a much smaller width of  $48 \pm 17$  MeV consistent with that of [56,58]. We list these observations under a separate particle  $\rho(1900)$ , which needs confirmation.

## References

- 1. C. Erkal, Z. Phys. C31, 615 (1986).
- 2. A. Donnachie and H. Mirzaie, Z. Phys. C33, 407 (1987).
- 3. A. Donnachie and A.B. Clegg, Z. Phys. C34, 257 (1987).
- 4. A. Donnachie and A.B. Clegg, Z. Phys. C51, 689 (1991).
- 5. A.B. Clegg and A. Donnachie, Z. Phys. C40, 313 (1988).
- 6. A.B. Clegg and A. Donnachie, Z. Phys. C62, 455 (1994).
- 7. T.J. Killian *et al.*, Phys. Rev. **D21**, 3005 (1980).
- 8. S. Fukui *et al.*, Phys. Lett. **B202**, 441 (1988).
- 9. D. Bisello et al., Phys. Lett. **B220**, 321 (1989).
- 10. A. Antonelli et al., Phys. Lett. B212, 133 (1988).
- 11. A. Abele *et al.*, Phys. Lett. **B391**, 191 (1997).
- 12. A. Abele et al., Phys. Rev. D57, 3860 (1998).
- 13. A. Bertin *et al.*, Phys. Lett. **B434**, 180 (1998).
- 14. A. Abele *et al.*, Phys. Lett. **B468**, 178 (1999).
- 15. R. Barate *et al.*, Z. Phys. **C76**, 15 (1997).
- 16. S. Anderson, Phys. Rev. **D61**, 112002 (2000).
- 17. K.W. Edwards *et al.*, Phys. Rev. **D61**, 072003 (2000).
- 18. M. Fujikawa et al., Phys. Rev. D78, 072006 (2008).
- 19. T. Barnes *et al.*, Phys. Rev. **D55**, 4157 (1997).
- 20. F.E. Close *et al.*, Phys. Rev. **D56**, 1584 (1997).
- A. Donnachie and Yu.S. Kalashnikova, Phys. Rev. D60, 114011 (1999).
- 22. R.R. Akhmetshin *et al.*, Phys. Lett. **B466**, 392 (1999).
- 23. J.P. Alexander *et al.*, Phys. Rev. **D64**, 092001 (2001).
- 24. R.R. Akhmetshin et al., Phys. Lett. B562, 173 (2003).
- 25. A. Abele *et al.*, Eur. Phys. J. **C21**, 261 (2001).

- 26. M. Bargiotti et al., Phys. Lett. B561, 233 (2003).
- 27. S.I. Bityukov et al., Phys. Lett. B188, 383 (1987).
- 28. B. Aubert *et al.*, Phys. Rev. **D77**, 092002 (2008).
- N.N. Achasov and G.N. Shestakov, Phys. Atom. Nucl. 59, 1262 (1996).
- 30. L.G. Landsberg, Sov. J. Nucl. Phys. 55, 1051 (1992).
- 31. A. Abele *et al.*, Phys. Lett. **B415**, 280 (1997).
- V.M. Aulchenko *et al.*, Sov. Phys. JETP Lett. 45, 145 (1987).
- 33. D. Bisello *et al.*, Z. Phys. **C52**, 227 (1991).
- 34. P. Frenkiel *et al.*, Nucl. Phys. **B47**, 61 (1972).
- 35. G. Cosme *et al.*, Phys. Lett. **B63**, 352 (1976).
- 36. D.P. Barber *et al.*, Z. Phys. C4, 169 (1980).
- 37. D. Aston, Phys. Lett. **B92**, 211 (1980).
- 38. M. Atkinson *et al.*, Nucl. Phys. **B243**, 1 (1984).
- 39. J.E. Brau *et al.*, Phys. Rev. **D37**, 2379 (1988).
- 40. C. Amsler *et al.*, Phys. Lett. **B311**, 362 (1993).
- 41. J. Layssac and F.M. Renard, Nuovo Cimento **6A**, 134 (1971).
- D. Aston *et al.*, Nucl. Phys. (Proc. Supp.) **B21**, 105 (1991).
- 43. M. Ablikim et al., Phys. Rev. Lett. 97, 142002 (2006).
- 44. G.-J. Ding and M.-L. Yan, Phys. Lett. **B643**, 33 (2006).
- 45. F.K. Guo et al., Nucl. Phys. A773, 78 (2006).
- 46. A. Zhang et al., Phys. Rev. D76, 036004 (2007).
- 47. B.A. Li, Phys. Rev. **D76**, 094016 (2007).
- 48. X. Liu *et al.*, Phys. Rev. **D75**, 074017 (2007).
- 49. A.B. Clegg and A. Donnachie, Z. Phys. C45, 677 (1990).
- 50. D. Bisello *et al.*, Phys. Lett. **107B**, 145 (1981).
- 51. A. Castro *et al.*, LAL-88-58(1988).
- 52. M. Atkinson *et al.*, Z. Phys. **C29**, 333 (1985).
- 53. B. Aubert *et al.*, Phys. Rev. **D76**, 092005 (2007).
- 54. S. Godfrey and N. Isgur, Phys. Rev. **D32**, 189 (1985).
- 55. P.L. Frabetti *et al.*, Phys. Lett. **B514**, 240 (2001).
- 56. A. Antonelli *et al.*, Phys. Lett. **B365**, 427 (1996).
- 57. M. Agnello *et al.*, Phys. Lett. **B527**, 39 (2002).
- 58. P.L. Frabetti *et al.*, Phys. Lett. **B578**, 290 (2004).
- 59. B. Aubert *et al.*, Phys. Rev. **D73**, 052003 (2006).