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#### Z MASS

OUR FIT is obtained using the fit procedure and correlations as determined by the LEP Electroweak Working Group (see the note "The Z boson" and ref. LEP-SLC 06). The fit is performed using the Z mass and width, the Z hadronic pole cross section, the ratios of hadronic to leptonic partial widths, and the Z pole forward-backward lepton asymmetries. This set is believed to be most free of correlations.

The Z-boson mass listed here corresponds to the mass parameter in a Breit-Wigner distribution with mass dependent width. The value is 34 MeV greater than the real part of the position of the pole (in the energy-squared plane) in the Z-boson propagator. Also the LEP experiments have generally assumed a fixed value of the  $\gamma-Z$  interferences term based on the standard model. Keeping this term as free parameter leads to a somewhat larger error on the fitted Z mass. See ACCIARRI 00Q and ABBIENDI 04G for a detailed investigation of both these issues.

VALUE (	GeV)		EVTS	DOCUMENT ID		TECN	COMMENT
91.187	5±0.002	1 OUR FI	Γ				
91.1852	$2 \pm 0.003$	0	4.57M	<sup>1</sup> ABBIENDI	<b>01</b> A	OPAL	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
91.1863	$3 \pm 0.002$	8	4.08M	<sup>2</sup> ABREU	00F	DLPH	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
91.1898	$3 \pm 0.003$	1	3.96M	<sup>3</sup> ACCIARRI	<b>00</b> C	L3	$E_{\rm cm}^{\it ee}$ = 88–94 GeV
91.188	$5\pm0.003$	1	4.57M	<sup>4</sup> BARATE	<b>00</b> C	ALEP	$E_{\rm cm}^{\it ee}$ = 88–94 GeV
• • • '	We do no	ot use the	following o	data for averages, fi	ts, lin	nits, etc.	• • •
91.1872	$2 \pm 0.003$	3		<sup>5</sup> ABBIENDI	<b>04</b> G	OPAL	Eee LEP1 +
91.272	±0.032	±0.033		<sup>6</sup> ACHARD	<b>04</b> C	_	130–209 GeV <i>Eee</i> <sub>cm</sub> = 183–209 GeV
91.187	$5 \pm 0.003$	9	3.97M	<sup>7</sup> ACCIARRI	00Q	L3	$E_{\rm cm}^{\it ee} = {\sf LEP1} +$
91.151	±0.008			<sup>8</sup> MIYABAYASHI	95	TOPZ	CIII
91.74	$\pm  0.28$	$\pm 0.93$	156	<sup>9</sup> ALITTI	<b>92</b> B	UA2	$E_{\rm cm}^{p\overline{p}}$ = 630 GeV
90.9	$\pm 0.3$	$\pm 0.2$	188	<sup>10</sup> ABE	<b>89</b> C	CDF	$E_{cm}^{ar{p}} = 1.8 \; TeV$
91.14	$\pm  0.12$		480	<sup>11</sup> ABRAMS	<b>89</b> B	MRK2	E <sup>ee</sup> <sub>cm</sub> = 89–93 GeV
93.1	$\pm  1.0$	$\pm 3.0$	24	<sup>12</sup> ALBAJAR	89	UA1	$E_{\rm cm}^{p\overline{p}}$ = 546,630 GeV

<sup>&</sup>lt;sup>1</sup> ABBIENDI 01A error includes approximately 2.3 MeV due to statistics and 1.8 MeV due to LEP energy uncertainty.

<sup>&</sup>lt;sup>2</sup> The error includes 1.6 MeV due to LEP energy uncertainty.

<sup>&</sup>lt;sup>3</sup>The error includes 1.8 MeV due to LEP energy uncertainty.

<sup>&</sup>lt;sup>4</sup>BARATE 00C error includes approximately 2.4 MeV due to statistics, 0.2 MeV due to experimental systematics, and 1.7 MeV due to LEP energy uncertainty.

 $<sup>^5</sup>$  ABBIENDI 04G obtain this result using the S-matrix formalism for a combined fit to their cross section and asymmetry data at the Z peak and their data at 130–209 GeV. The authors have corrected the measurement for the 34 MeV shift with respect to the Breit-Wigner fits.

- <sup>6</sup> ACHARD 04C select  $e^+e^- \rightarrow Z\gamma$  events with hard initial–state radiation. Z decays to  $q \overline{q}$  and muon pairs are considered. The fit results obtained in the two samples are found consistent to each other and combined considering the uncertainty due to ISR modelling as fully correlated.
- <sup>7</sup> ACCIARRI 00Q interpret the s-dependence of the cross sections and lepton forward-backward asymmetries in the framework of the S-matrix formalism. They fit to their cross section and asymmetry data at high energies, using the results of S-matrix fits to Z-peak data (ACCIARRI 00C) as constraints. The 130–189 GeV data constrains the  $\gamma/Z$  interference term. The authors have corrected the measurement for the 34.1 MeV shift with respect to the Breit-Wigner fits. The error contains a contribution of  $\pm 2.3$  MeV due to the uncertainty on the  $\gamma Z$  interference.
- <sup>8</sup> MIYABAYASHI 95 combine their low energy total hadronic cross-section measurement with the ACTON 93D data and perform a fit using an S-matrix formalism. As expected, this result is below the mass values obtained with the standard Breit-Wigner parametrization.
- <sup>9</sup> Enters fit through W/Z mass ratio given in the W Particle Listings. The ALITTI 92B systematic error  $(\pm 0.93)$  has two contributions: one  $(\pm 0.92)$  cancels in  $m_W/m_Z$  and one  $(\pm 0.12)$  is noncancelling. These were added in quadrature.
- <sup>10</sup> First error of ABE 89 is combination of statistical and systematic contributions; second is mass scale uncertainty.
- $^{11}$  ABRAMS 89B uncertainty includes 35 MeV due to the absolute energy measurement.
- $^{12}$  ALBAJAR 89 result is from a total sample of 33  $Z \rightarrow e^+e^-$  events.

#### Z WIDTH

OUR FIT is obtained using the fit procedure and correlations as determined by the LEP Electroweak Working Group (see the note "The Z boson" and ref. LEP-SLC 06).

VALUE	(GeV)		<b>EVTS</b>	DOCUMENT ID		TECN	COMMENT	
2.4952±0.0023 OUR FIT								
2.494	$8 \pm 0.004$	1	4.57M	$^{ m 1}$ ABBIENDI	01A	OPAL	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV	
2.487	$6 \pm 0.004$	1	4.08M	<sup>2</sup> ABREU	00F	DLPH	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV	
2.502	$4 \pm 0.004$	2	3.96M	<sup>3</sup> ACCIARRI	<b>00</b> C	L3	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV	
2.495	$1 \pm 0.004$	3	4.57M	<sup>4</sup> BARATE	<b>00</b> C	ALEP	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV	
• • •	We do i	not use tl	he followir	ng data for averages	s, fits,	limits, e	etc. • • •	
2.494	$3 \pm 0.004$	1		<sup>5</sup> ABBIENDI	04G	OPAL	$E_{\rm cm}^{\rm ee} = {\sf LEP1} +$	
2.502	5±0.004	1	3.97M	<sup>6</sup> ACCIARRI	00Q	L3	130–209 GeV Eee LEP1 +	
2.50	$\pm 0.21$	$\pm 0.06$		<sup>7</sup> ABREU	<b>96</b> R	DLPH	130–189 GeV <i>E</i> <sup>ee</sup> <sub>cm</sub> = 91.2 GeV	
3.8	$\pm0.8$	$\pm 1.0$	188	ABE	89C	CDF	$E_{cm}^{p\overline{p}} = 1.8 \; TeV$	
2.42	$^{+0.45}_{-0.35}$		480	<sup>8</sup> ABRAMS	<b>89</b> B	MRK2	<i>E</i> <sup>ee</sup> <sub>cm</sub> = 89−93 GeV	
2.7	$^{+1.2}_{-1.0}$	$\pm 1.3$	24	<sup>9</sup> ALBAJAR	89	UA1	$E_{\rm cm}^{p\overline{p}}=$ 546,630 GeV	
2.7	$\pm2.0$	$\pm 1.0$	25	<sup>10</sup> ANSARI	87	UA2	$E_{cm}^{p\overline{p}} = 546,630 \; GeV$	

#### Z DECAY MODES

	Mode	I	Fraction	(Γ <sub>i</sub> /Γ)			cale factor/ dence level
Γ <sub>1</sub>	$e^+e^-$		( 3.363	3 ±0.004	) %		
-	$\mu^+\mu^-$		`	5 ±0.007	,		
Γ <sub>3</sub>	$\tau^+\tau^-$		( 3.370	±0.008	) %		
Γ <sub>4</sub>	$\ell^+\ell^-$	[ <i>a</i> ]	( 3.365	$58 \pm 0.002$	3) %		
Γ <sub>5</sub>	$\ell^+\ell^-\ell^+\ell^-$	[ <i>b</i> ]	( 4.2	$^{+0.9}_{-0.8}$	) × 10	<sub>0</sub> –6	
$\Gamma_6$	invisible		(20.00	$\pm0.06$	) %		
$\Gamma_7$	hadrons		(69.91	$\pm0.06$	) %		
Γ <sub>8</sub>	$(u\overline{u}+c\overline{c})/2$		(11.6	$\pm  0.6$	) %		
$\Gamma_9$	$(d\overline{d} + s\overline{s} + b\overline{b})/3$		(15.6	$\pm  0.4$	) %		
$\Gamma_{10}$	<u>c</u>		(12.03	$\pm 0.21$	) %		
	$b\overline{b}$		(15.12	$\pm  0.05$	) %		
$\Gamma_{12}$	<i>b</i> <del>b</del> b <del>b</del>		( 3.6	$\pm 1.3$	) × 10	)-4	
$\Gamma_{13}$	ggg		< 1.1		%		CL=95%
$\Gamma_{14}$	$\pi^{0}\gamma$		< 5.2		$\times$ 10	<sub>)</sub> –5	CL=95%
Γ <sub>15</sub>	$\eta\gamma$		< 5.1		$\times$ 10	ე–5	CL=95%
$\Gamma_{16}$	$\omega \gamma$		< 6.5				CL=95%
$\Gamma_{17}$	$\eta'(958)\gamma$		< 4.2		$\times$ 10	<sub>)</sub> –5	CL=95%
Γ <sub>18</sub>	$\gamma\gamma$		< 5.2		$\times$ 10	<sub>)</sub> –5	CL=95%
Γ <sub>19</sub>	$\gamma\gamma\gamma$		< 1.0		× 10	)-5	CL=95%
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 $<sup>^{</sup>m 1}$  ABBIENDI 01A error includes approximately 3.6 MeV due to statistics, 1 MeV due to event selection systematics, and 1.3 MeV due to LEP energy uncertainty.

<sup>&</sup>lt;sup>2</sup>The error includes 1.2 MeV due to LEP energy uncertainty.

<sup>&</sup>lt;sup>3</sup>The error includes 1.3 MeV due to LEP energy uncertainty.

<sup>&</sup>lt;sup>4</sup>BARATE 00C error includes approximately 3.8 MeV due to statistics, 0.9 MeV due to experimental systematics, and 1.3 MeV due to LEP energy uncertainty.

<sup>&</sup>lt;sup>5</sup> ABBIENDI 04G obtain this result using the S-matrix formalism for a combined fit to their cross section and asymmetry data at the Z peak and their data at 130-209 GeV. The authors have corrected the measurement for the 1 MeV shift with respect to the Breit-Wigner fits.

<sup>&</sup>lt;sup>6</sup> ACCIARRI 00Q interpret the s-dependence of the cross sections and lepton forwardbackward asymmetries in the framework of the S-matrix formalism. They fit to their cross section and asymmetry data at high energies, using the results of S-matrix fits to Z-peak data (ACCIARRI 00C) as constraints. The 130–189 GeV data constrains the  $\gamma/Z$ interference term. The authors have corrected the measurement for the 0.9 MeV shift with respect to the Breit-Wigner fits.

<sup>&</sup>lt;sup>7</sup> ABREU 96R obtain this value from a study of the interference between initial and final state radiation in the process  $e^+e^- \rightarrow Z \rightarrow \mu^+\mu^-$ .

<sup>&</sup>lt;sup>8</sup> ABRAMS 89B uncertainty includes 50 MeV due to the miniSAM background subtraction error.

 $<sup>^9</sup>$ ALBAJAR 89 result is from a total sample of 33  $Z \rightarrow e^+e^-$  events.

 $<sup>^{10}</sup>$  Quoted values of ANSARI 87 are from direct fit. Ratio of Z and W production gives either  $\Gamma(Z)<(1.09\pm0.07)\times\Gamma(W)$ , CL = 90% or  $\Gamma(Z)=(0.82^{+0.19}_{-0.14}\pm0.06)\times\Gamma(W)$ . Assuming Standard-Model value  $\Gamma(W)=2.65$  GeV then gives  $\Gamma(Z)<2.89\pm0.19$  or  $=2.17^{+0.50}_{-0.37}\pm0.16.$ 

```
\pi^{\pm}W^{\mp}
                                                                                                    \times 10^{-5}
                                                                 [c] < 7
                                                                                                                    CL=95%
\Gamma_{21} \quad \rho^{\pm} W^{\mp}
                                                                                                    \times 10^{-5}
                                                                 [c] < 8.3
                                                                                                                    CL=95%
                                                                                     ^{+\,0.23}_{-\,0.25}
\Gamma_{22} = J/\psi(1S)X
                                                                                                  ) \times 10^{-3}
                                                                        ( 3.51
                                                                                                                        S = 1.1
\Gamma_{23}
        \psi(2S)X
                                                                                                 ) \times 10^{-3}
                                                                                     \pm 0.29
                                                                        ( 1.60
        \chi_{c1}(1P)X
                                                                                                 ) \times 10^{-3}
\Gamma_{24}
                                                                        ( 2.9
                                                                                     \pm 0.7
\Gamma_{25}
         \chi_{c2}(1P)X
                                                                                                    \times 10^{-3}
                                                                      < 3.2
                                                                                                                    CL=90%
          \Upsilon(1S) \times + \Upsilon(2S) \times
                                                                        ( 1.0
                                                                                                  ) \times 10^{-4}
                                                                                     \pm 0.5
               +\Upsilon(3S) X
               \Upsilon(1S)X
                                                                                                    \times 10^{-5}
                                                                                                                    CL=95%
                                                                      < 4.4
                                                                                                    \times 10^{-4}
              \Upsilon(2S)X
\Gamma_{28}
                                                                                                                    CL=95%
                                                                      < 1.39
               \Upsilon(3S)X
\Gamma_{29}
                                                                                                    \times 10^{-5}
                                                                                                                    CL=95%
                                                                      < 9.4
          (D^0/\overline{D}^0) X
\Gamma_{30}
                                                                                                  ) %
                                                                         (20.7)
                                                                                     \pm 2.0
          D^{\pm}X
\Gamma_{31}
                                                                                                 ) %
                                                                        (12.2)
                                                                                     \pm 1.7
\Gamma_{32}
          D^*(2010)^{\pm}X
                                                                                                 ) %
                                                                 [c] (11.4
                                                                                     \pm 1.3
                                                                                                  ) \times 10^{-3}
\Gamma_{33}
          D_{s1}(2536)^{\pm} X
                                                                         ( 3.6
                                                                                     \pm 0.8
                                                                                                 ) \times 10^{-3}
          D_{s,I}(2573)^{\pm} X
\Gamma_{34}
                                                                         (5.8
                                                                                     \pm 2.2
          D^{*'}(2629)^{\pm}X
\Gamma_{35}
                                                                     searched for
\Gamma_{36}
          BX
\Gamma_{37}
          B^*X
          B^{+}X
\Gamma_{38}
                                                                                     \pm 0.13 )%
                                                                 [d] (6.08
          B_{\varepsilon}^{0}X
\Gamma_{39}
                                                                                     \pm 0.13 )%
                                                                 [d] (1.59
\Gamma_{40}
          B^+_{\cdot}X
                                                                     searched for
          \Lambda^{+}X
\Gamma_{41}
                                                                                     \pm 0.33 ) %
                                                                         ( 1.54
\Gamma_{42}
                                                                          seen
\Gamma_{43}
          \Xi_b X
                                                                          seen
\Gamma_{44}
          b-baryon X
                                                                 [d] ( 1.38
                                                                                     \pm 0.22 )%
                                                                                                    \times 10^{-3}
          anomalous \gamma + hadrons
                                                                 [e] < 3.2
                                                                                                                    CL=95%
\Gamma_{46}
          e^+e^-\gamma
                                                                                                    \times 10^{-4}
                                                                                                                    CL=95%
                                                                 [e] < 5.2
          \mu^+\mu^-\gamma
\Gamma_{47}
                                                                                                    \times 10^{-4}
                                                                                                                    CL=95%
                                                                 [e] < 5.6
          \tau^+\tau^-\gamma
\Gamma_{48}
                                                                                                    \times 10^{-4}
                                                                                                                    CL=95%
                                                                 [e] < 7.3
\Gamma_{49}
          \ell^+\ell^-\gamma\gamma
                                                                                                    \times 10^{-6}
                                                                 [f] < 6.8
                                                                                                                    CL=95%
\Gamma_{50}
                                                                                                    \times 10^{-6}
          q \overline{q} \gamma \gamma
                                                                 [f] < 5.5
                                                                                                                    CL=95%
                                                                                                    \times 10^{-6}
\Gamma_{51}
          \nu \overline{\nu} \gamma \gamma
                                                                                                                    CL=95%
                                                                 [f] < 3.1
          e^{\pm}\mu^{\mp}
\Gamma_{52}
                                                                                                    \times 10^{-6}
                                                      LF
                                                                 [c] < 1.7
                                                                                                                    CL=95%
          e^{\pm} \tau^{\mp}
\Gamma_{53}
                                                                                                    \times 10^{-6}
                                                      LF
                                                                                                                    CL=95%
                                                                 [c] <
                                                                           9.8
\Gamma_{54}
          \mu^{\pm} \tau^{\mp}
                                                                                                    \times 10^{-5}
                                                      LF
                                                                                                                    CL=95%
                                                                 [c] <
                                                                          1.2
                                                                                                    \times 10^{-6}
\Gamma_{55}
          ре
                                                      L,B
                                                                      <
                                                                           1.8
                                                                                                                    CL=95%
                                                                                                    \times 10^{-6}
\Gamma_{56}
                                                                                                                    CL=95%
          p\mu
                                                      L,B
                                                                      < 1.8
```

- [a]  $\ell$  indicates each type of lepton  $(e, \mu, \text{ and } \tau)$ , not sum over them.
- [b] Here  $\ell$  indicates e or  $\mu$ .
- [c] The value is for the sum of the charge states or particle/antiparticle states indicated.

- [d] This value is updated using the product of (i) the  $Z \rightarrow b \overline{b}$  fraction from this listing and (ii) the b-hadron fraction in an unbiased sample of weakly decaying b-hadrons produced in Z-decays provided by the Heavy Flavor Averaging Group (HFAG, http://www.slac.stanford.edu/xorg/hfag/osc/PDG\_2009/#FRACZ).
- [e] See the Particle Listings below for the  $\gamma$  energy range used in this measurement.
- [f] For  $m_{\gamma\gamma}=$  (60  $\pm$  5) GeV.

#### Z PARTIAL WIDTHS

 $\Gamma(e^+e^-)$ For the LEP experiments, this parameter is not directly used in the overall fit but is

derived using the fit results; see the note "The Z boson" and ref. LEP-SLC 06.

 VALUE (MeV)
 EVTS
 DOCUMENT ID
 TECN
 COMMENT

 83.91±0.12 OUR FIT
 83.66±0.20
 137.0K
 ABBIENDI
 01A
 OPAL
 Eee = 88-94 GeV

 83.54±0.27
 117.8k
 ABBELL
 00E DLPH
 Fee = 88-94 GeV

 $83.66 \pm 0.20$ 00F DLPH  $E_{cm}^{ee} = 88-94 \text{ GeV}$  $83.54 \pm 0.27$ 117.8k **ABREU**  $E_{cm}^{ee} = 88-94 \text{ GeV}$  $84.16 \pm 0.22$ 124.4k **ACCIARRI** 00C L3 00C ALEP *E*<sub>cm</sub><sup>ee</sup> = 88–94 GeV  $83.88 \pm 0.19$ **BARATE** <sup>1</sup> ABE  $82.89 \pm 1.20 \pm 0.89$  $E_{\rm cm}^{\rm ee}=91.31~{\rm GeV}$ 95J SLD

 $\Gamma(\mu^+\mu^-)$ This parameter is not directly used in the overall fit but is derived using the fit results; see the note "The Z boson" and ref. LEP-SLC 06.

VALUE (MeV)	<i>EVTS</i>	DOCUMENT ID		TECN	COMMENT
83.99±0.18 OUR FIT					
$84.03 \pm 0.30$	182.8K	ABBIENDI	<b>01</b> A	OPAL	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
$84.48 \pm 0.40$	157.6k	ABREU	00F	DLPH	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
$83.95 \pm 0.44$	113.4k	ACCIARRI	<b>00</b> C	L3	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
$84.02 \!\pm\! 0.28$		BARATE	<b>00</b> C	ALEP	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV

 $\Gamma( au^+ au^-)$ 

This parameter is not directly used in the overall fit but is derived using the fit results; see the note "The Z boson" and ref. LEP-SLC 06.

<i>VALUE</i> (MeV)	EVTS	DOCUMENT ID		TECN	COMMENT
84.08±0.22 OUR FIT					
$83.94 \pm 0.41$	151.5K	ABBIENDI	01A	OPAL	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
$83.71 \pm 0.58$	104.0k	ABREU	00F	DLPH	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
$84.23 \pm 0.58$	103.0k	ACCIARRI	<b>00</b> C	L3	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
$84.38 \!\pm\! 0.31$		BARATE	00C	ALEP	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV

 $<sup>^1</sup>$  ABE 95J obtain this measurement from Bhabha events in a restricted fiducial region to improve systematics. They use the values 91.187 and 2.489 GeV for the Z mass and total decay width to extract this partial width.

 $\Gamma(\ell^+\ell^-)$ 

In our fit  $\Gamma(\ell^+\ell^-)$  is defined as the partial Z width for the decay into a pair of massless charged leptons. This parameter is not directly used in the 5-parameter fit assuming lepton universality but is derived using the fit results. See the note "The Z boson" and ref. LEP-SLC 06.

VALUE (MeV)	EVTS	DOCUMENT ID		TECN	COMMENT
83.984±0.086 OUR F	IT				
$83.82 \pm 0.15$	471.3K	ABBIENDI	01A	OPAL	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
83.85 $\pm 0.17$	379.4k	ABREU	00F	DLPH	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
84.14 $\pm 0.17$	340.8k	ACCIARRI	<b>00</b> C	L3	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
$84.02 \pm 0.15$	500k	BARATE	<b>00</b> C	ALEP	$E_{\rm cm}^{\it ee} = 88 - 94 \; {\rm GeV}$

# Γ(invisible)

We use only direct measurements of the invisible partial width using the single photon channel to obtain the average value quoted below. OUR FIT value is obtained as a difference between the total and the observed partial widths assuming lepton universality.

 $\Gamma_6$ 

Γ<sub>7</sub>

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VALUE (MeV)	EVTS	DOCUMENT ID		TECN	COMMENT
499.0± 1.5 OUR FIT					
503 $\pm$ 16 OUR AVER	RAGE Erro	or includes scale f	actor	of 1.2.	
$498\pm12\pm12$	1791	ACCIARRI	98G	L3	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
$539 \pm 26 \pm 17$	410	AKERS	<b>95</b> C	OPAL	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
$450$ $\pm 34$ $\pm 34$	258	BUSKULIC	93L	ALEP	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
$540$ $\pm 80$ $\pm 40$	52	ADEVA	92	L3	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
• • • We do not use th	e following	data for averages	s, fits,	limits, e	etc. • • •
498.1± 2.6		<sup>1</sup> ABBIENDI	01A	OPAL	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
498.1± 3.2		<sup>1</sup> ABREU	00F	DLPH	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
$499.1 \pm 2.9$		<sup>1</sup> ACCIARRI	<b>00</b> C	L3	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
$499.1 \pm 2.5$		<sup>1</sup> BARATE	<b>00</b> C	ALEP	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV

<sup>&</sup>lt;sup>1</sup> This is an indirect determination of  $\Gamma$ (invisible) from a fit to the visible Z decay modes.

### Γ(hadrons)

This parameter is not directly used in the 5-parameter fit assuming lepton universality, but is derived using the fit results. See the note "The Z boson" and ref. LEP-SLC 06.

VALUE (MeV)	EVTS	DOCUMENT ID		TECN	COMMENT
1744.4±2.0 OUR FIT					
$1745.4 \pm 3.5$	4.10M	ABBIENDI	<b>01</b> A	OPAL	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
$1738.1\!\pm\!4.0$	3.70M	ABREU	00F	DLPH	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
$1751.1 \pm 3.8$	3.54M	ACCIARRI	<b>00</b> C	L3	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
$1744.0 \pm 3.4$	4.07M	BARATE	00C	ALEP	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV

#### **Z BRANCHING RATIOS**

OUR FIT is obtained using the fit procedure and correlations as determined by the LEP Electroweak Working Group (see the note "The Z boson" and ref. LEP-SLC 06).

$\Gamma(\text{hadrons})/\Gamma(e^+e^-)$					$\Gamma_7/\Gamma_1$
VALUE	<b>EVTS</b>	DOCUMENT ID		TECN	COMMENT
20.804± 0.050 OUR FIT					
$20.902 \pm \ 0.084$	137.0K	$^{ m 1}$ ABBIENDI	<b>01</b> A	OPAL	$E_{\mathrm{cm}}^{\mathrm{ee}} = 88-94 \; \mathrm{GeV}$
$20.88 \pm 0.12$	117.8k	ABREU	00F	DLPH	$E_{\rm cm}^{\it ee}=$ 88–94 GeV
$20.816 \pm \ 0.089$	124.4k	ACCIARRI	<b>00</b> C	L3	$E_{\rm cm}^{\it ee}=$ 88–94 GeV
$20.677 \pm \ 0.075$		<sup>2</sup> BARATE	<b>00</b> C	ALEP	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
• • • We do not use the fo	ollowing da	ita for averages, fit	ts, lim	its, etc.	• • •
$27.0 \begin{array}{c} +11.7 \\ -8.8 \end{array}$	12	<sup>3</sup> ABRAMS	<b>89</b> D	MRK2	E <sup>ee</sup> <sub>cm</sub> = 89–93 GeV

<sup>&</sup>lt;sup>1</sup> ABBIENDI 01A error includes approximately 0.067 due to statistics, 0.040 due to event selection systematics, 0.027 due to the theoretical uncertainty in *t*-channel prediction, and 0.014 due to LEP energy uncertainty.

# $\Gamma(\text{hadrons})/\Gamma(\mu^+\mu^-)$

 $\Gamma_7/\Gamma_2$ 

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OUR FIT is obtained using the fit procedure and correlations as determined by the LEP Electroweak Working Group (see the note "The Z boson" and ref. LEP-SLC 06).

VALUE	<u>EVTS</u>	DOCUMENT ID		TECN	COMMENT
$20.785\pm0.033$ OUR FIT					
$20.811 \!\pm\! 0.058$	182.8K	<sup>1</sup> ABBIENDI	01A	OPAL	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
$20.65 \pm 0.08$	157.6k	ABREU	00F	DLPH	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
$20.861\!\pm\!0.097$	113.4k	ACCIARRI	<b>00</b> C	L3	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
$20.799 \pm 0.056$		<sup>2</sup> BARATE	<b>00</b> C	ALEP	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
• • • We do not use the fe	ollowing da	ata for averages, fi	ts, lim	its, etc.	• • •
$18.9  {}^{+7.1}_{-5.3}$	13	<sup>3</sup> ABRAMS	<b>89</b> D	MRK2	E <sup>ee</sup> <sub>cm</sub> = 89–93 GeV

 $<sup>^{1}</sup>$ ABBIENDI 01A error includes approximately 0.050 due to statistics and 0.027 due to event selection systematics.

<sup>&</sup>lt;sup>2</sup> BARATE 00C error includes approximately 0.062 due to statistics, 0.033 due to experimental systematics, and 0.026 due to the theoretical uncertainty in *t*-channel prediction.

<sup>&</sup>lt;sup>3</sup> ABRAMS 89D have included both statistical and systematic uncertainties in their quoted errors.

<sup>&</sup>lt;sup>2</sup>BARATE 00C error includes approximately 0.053 due to statistics and 0.021 due to experimental systematics.

<sup>&</sup>lt;sup>3</sup> ABRAMS 89D have included both statistical and systematic uncertainties in their quoted errors.

 $\Gamma(\text{hadrons})/\Gamma(\tau^+\tau^-)$ 

 $\Gamma_7/\Gamma_3$ 

OUR FIT is obtained using the fit procedure and correlations as determined by the LEP Electroweak Working Group (see the note "The Z boson" and ref. LEP-SLC 06).

VALUE	<b>EVTS</b>	DOCUMENT ID		TECN	COMMENT
20.764 ± 0.045 OUR FIT					
$20.832 \!\pm\! 0.091$	151.5K	$^{ m 1}$ ABBIENDI	01A	OPAL	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
$20.84 \pm 0.13$	104.0k	ABREU	00F	DLPH	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
$20.792\!\pm\!0.133$	103.0k	ACCIARRI	<b>00</b> C	L3	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
$20.707 \pm 0.062$		<sup>2</sup> BARATE	<b>00</b> C	ALEP	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
• • • We do not use the fe	ollowing da	ata for averages, fi	ts, lim	its, etc.	• • •
15.2 +4.8	21	3 аррамс	000	MDKO	ree on 02 CaV

<sup>15.2</sup>  $^{+4.0}_{-3.9}$  21 3 ABRAMS 89D MRK2  $E_{cm}^{ee} = 89-93 \text{ GeV}$ 

# $\Gamma(\text{hadrons})/\Gamma(\ell^+\ell^-)$

 $\Gamma_7/\Gamma_4$ 

 $\ell$  indicates each type of lepton  $(e, \mu, \text{ and } \tau)$ , not sum over them.

Our fit result is obtained requiring lepton universality.

<u>VALUE</u>	<u>EVTS</u>	<u>DOCUMENT ID</u>		TECN	COMMENT
20.767±0.025 OUR	FIT				
$20.823 \pm 0.044$	471.3K	$^{ m 1}$ ABBIENDI	<b>01</b> A	OPAL	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
$20.730 \pm 0.060$	379.4k	ABREU	00F	DLPH	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
$20.810 \pm 0.060$	340.8k	ACCIARRI	<b>00</b> C	L3	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
$20.725 \!\pm\! 0.039$	500k	<sup>2</sup> BARATE	<b>00</b> C	ALEP	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
• • • We do not us	e the follov	ving data for avera	ges, f	its, limit	s, etc. • • •
18.9 + 3.6	46	ABRAMS	<b>89</b> B	MRK2	E <sub>cm</sub> <sup>ee</sup> = 89–93 GeV

<sup>&</sup>lt;sup>1</sup> ABBIENDI 01A error includes approximately 0.034 due to statistics and 0.027 due to event selection systematics.

# $\Gamma(\text{hadrons})/\Gamma_{\text{total}}$

 $\Gamma_7/\Gamma$ 

This parameter is not directly used in the overall fit but is derived using the fit results; see the note "The Z boson" and ref. LEP-SLC 06.

\*\*DOCUMENT ID\*\*

69.911±0.056 OUR FIT

 $\Gamma(e^+e^-)/\Gamma_{\text{total}}$ 

 $\Gamma_1/\Gamma$ 

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This parameter is not directly used in the overall fit but is derived using the fit results; see the note "The Z boson" and ref. LEP-SLC 06.

VALUE (%) DOCUMENT ID

 $(3363.2\pm4.2)\times10^{-3}$  OUR FIT

<sup>&</sup>lt;sup>1</sup> ABBIENDI 01A error includes approximately 0.055 due to statistics and 0.071 due to event selection systematics.

<sup>&</sup>lt;sup>2</sup>BARATE 00C error includes approximately 0.054 due to statistics and 0.033 due to experimental systematics.

<sup>&</sup>lt;sup>3</sup> ABRAMS 89D have included both statistical and systematic uncertainties in their quoted errors.

 $<sup>^2</sup>$  BARATE 00C error includes approximately 0.033 due to statistics, 0.020 due to experimental systematics, and 0.005 due to the theoretical uncertainty in t-channel prediction.

 $\Gamma(\mu^+\mu^-)/\Gamma_{\text{total}}$ 

This parameter is not directly used in the overall fit but is derived using the fit results; see the note "The Z boson" and ref. LEP-SLC 06. DOCUMENT ID

 $(3366.2\pm6.6)\times10^{-3}$  OUR FIT

 $\Gamma(\mu^+\mu^-)/\Gamma(e^+e^-)$ 

 $\Gamma_2/\Gamma_1$ This parameter is not directly used in the overall fit but is derived using the fit results; see the note "The Z boson" and ref. LEP-SLC 06.  $\underline{DOCUMENT\ ID}$ 

1.0009±0.0028 OUR FIT

 $\Gamma(\tau^+\tau^-)/\Gamma_{\text{total}}$  $\Gamma_3/\Gamma$ 

This parameter is not directly used in the overall fit but is derived using the fit results; see the note "The Z boson" and ref. LEP-SLC 06.

DOCUMENT ID

 $(3369.6\pm8.3)\times10^{-3}$  OUR FIT

 $\Gamma \big(\tau^+ \, \tau^-\big)/\Gamma \big(e^+ \, e^-\big)$  $\Gamma_3/\Gamma_1$ 

This parameter is not directly used in the overall fit but is derived using the fit results; see the note "The Z boson" and ref. LEP-SLC 06.

DOCUMENT ID

1.0019 ± 0.0032 OUR FIT

 $\Gamma(\ell^+\ell^-)/\Gamma_{\text{total}}$  $\Gamma_4/\Gamma$ 

 $\ell$  indicates each type of lepton  $(e, \mu, \text{ and } \tau)$ , not sum over them.

Our fit result assumes lepton universality.

This parameter is not directly used in the overall fit but is derived using the fit results; see the note "The Z boson" and ref. LEP-SLC 06.

DOCUMENT ID

 $(3365.8\pm2.3)\times10^{-3}$  OUR FIT

 $\Gamma(\ell^+\ell^-\ell^+\ell^-)/\Gamma_{ ext{total}}$ Here  $\ell$  indicates either e or  $\mu$ .  $\Gamma_5/\Gamma$ 

VALUE (units  $10^{-6}$ ) DOCUMENT ID TECN COMMENT  $4.2^{+0.9}_{-0.8}\pm0.2$ CHATRCHYAN 12BN CMS  $E_{cm}^{pp} = 7 \text{ TeV}$ 28

 $\Gamma(\text{invisible})/\Gamma_{\text{total}}$  $\Gamma_6/\Gamma$ 

See the data, the note, and the fit result for the partial width,  $\Gamma_6$ , above.

DOCUMENT ID

20.000 ± 0.055 OUR FIT

### $\Gamma((u\overline{u}+c\overline{c})/2)/\Gamma(\text{hadrons})$

 $\Gamma_8/\Gamma_7$ 

This quantity is the branching ratio of  $Z \to$  "up-type" quarks to  $Z \to$  hadrons. Except ACKERSTAFF 97T the values of  $Z \to$  "up-type" and  $Z \to$  "down-type" branchings are extracted from measurements of  $\Gamma(\text{hadrons})$ , and  $\Gamma(Z \to \gamma + \text{jets})$  where  $\gamma$  is a high-energy (>5 or 7 GeV) isolated photon. As the experiments use different procedures and slightly different values of  $M_Z$ ,  $\Gamma(\text{hadrons})$  and  $\alpha_S$  in their extraction procedures, our average has to be taken with caution.

VALUE	DOCUMENT ID		TECN	COMMENT
$0.166 \pm 0.009$ OUR AVERAGE				
$0.172^{igoplus 0.011}_{igoplus 0.010}$	<sup>1</sup> ABBIENDI	04E	OPAL	$E_{cm}^{\mathit{ee}} = 91.2 \; GeV$
$0.160\pm0.019\pm0.019$	<sup>2</sup> ACKERSTAFF	97T	OPAL	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
$0.137 ^{igoplus 0.038}_{-0.054}$	<sup>3</sup> ABREU	95x	DLPH	Eee = 88–94 GeV
$0.137 \pm 0.033$	<sup>4</sup> ADRIANI	93	L3	$E_{ m cm}^{\it ee} = 91.2 \; { m GeV}$

 $<sup>^1</sup>$  ABBIENDI 04E select photons with energy > 7 GeV and use  $\Gamma({\rm hadrons})=1744.4\pm2.0$  MeV and  $\alpha_{\rm S}=0.1172\pm0.002$  to obtain  $\Gamma_u=300^{+19}_{-18}$  MeV.

### $\Gamma((d\overline{d}+s\overline{s}+b\overline{b})/3)/\Gamma(hadrons)$

 $\Gamma_0/\Gamma_7$ 

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This quantity is the branching ratio of  $Z \to$  "down-type" quarks to  $Z \to$  hadrons. Except ACKERSTAFF 97T the values of  $Z \to$  "up-type" and  $Z \to$  "down-type" branchings are extracted from measurements of  $\Gamma(\text{hadrons})$ , and  $\Gamma(Z \to \gamma + \text{jets})$  where  $\gamma$  is a high-energy (>5 or 7 GeV) isolated photon. As the experiments use different procedures and slightly different values of  $M_Z$ ,  $\Gamma(\text{hadrons})$  and  $\alpha_S$  in their extraction procedures, our average has to be taken with caution.

VALUE	DOCUMENT ID		TECN	COMMENT
$0.223\pm0.006$ OUR AVERAGE				
$0.218 \pm 0.007$	<sup>1</sup> ABBIENDI	04E	OPAL	$E_{cm}^{\mathit{ee}} = 91.2 \; GeV$
$0.230 \pm 0.010 \pm 0.010$	<sup>2</sup> ACKERSTAFF	97T	OPAL	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
$0.243^{+0.036}_{-0.026}$	<sup>3</sup> ABREU	95X	DLPH	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
$0.243 \pm 0.022$	<sup>4</sup> ADRIANI	93	L3	$E_{cm}^{\mathit{ee}} = 91.2 \; GeV$

 $<sup>^1</sup>$  ABBIENDI 04E select photons with energy > 7 GeV and use  $\Gamma({\rm hadrons})=1744.4\pm2.0$  MeV and  $\alpha_{\rm S}=0.1172\pm0.002$  to obtain  $\Gamma_{\rm d}=381\pm12$  MeV.

<sup>&</sup>lt;sup>2</sup> ACKERSTAFF 97T measure  $\Gamma_{u\overline{u}}/(\Gamma_{d\overline{d}}+\Gamma_{u\overline{u}}+\Gamma_{s\overline{s}})=0.258\pm0.031\pm0.032$ . To obtain this branching ratio authors use  $R_c+R_b=0.380\pm0.010$ . This measurement is fully negatively correlated with the measurement of  $\Gamma_{d\overline{d},s\overline{s}}/(\Gamma_{d\overline{d}}+\Gamma_{u\overline{u}}+\Gamma_{s\overline{s}})$  given in the next data block.

<sup>&</sup>lt;sup>3</sup> ABREU 95x use  $M_Z = 91.187 \pm 0.009$  GeV,  $\Gamma(\text{hadrons}) = 1725 \pm 12$  MeV and  $\alpha_s = 0.123 \pm 0.005$ . To obtain this branching ratio we divide their value of  $C_{2/3} = 0.91 + 0.25 + 0.005$  by their value of  $(3C_{1/3} + 2C_{2/3}) = 6.66 \pm 0.05$ .

 $<sup>^4</sup>$  ADRIANI 93 use  $M_Z=91.181\pm0.022$  GeV,  $\Gamma({\rm hadrons})=1742\pm19$  MeV and  $\alpha_s=0.125\pm0.009$ . To obtain this branching ratio we divide their value of  $C_{2/3}=0.92\pm0.22$  by their value of  $(3C_{1/3}+2C_{2/3})=6.720\pm0.076$ .

<sup>&</sup>lt;sup>2</sup> ACKERSTAFF 97T measure  $\Gamma_{d\overline{d},s\overline{s}}/(\Gamma_{d\overline{d}}+\Gamma_{u\overline{u}}+\Gamma_{s\overline{s}})=0.371\pm0.016\pm0.016$ . To obtain this branching ratio authors use  $R_c+R_b=0.380\pm0.010$ . This measurement is fully negatively correlated with the measurement of  $\Gamma_{u\overline{u}}/(\Gamma_{d\overline{d}}+\Gamma_{u\overline{u}}+\Gamma_{s\overline{s}})$  presented in the previous data block.

- <sup>3</sup> ABREU 95X use  $M_Z=91.187\pm0.009$  GeV, Γ(hadrons) = 1725 ± 12 MeV and  $\alpha_S=0.123\pm0.005$ . To obtain this branching ratio we divide their value of  $C_{1/3}=1.62^{+0.24}_{-0.17}$  by their value of  $(3C_{1/3}+2C_{2/3})=6.66\pm0.05$ .
- <sup>4</sup> ADRIANI 93 use  $M_Z = 91.181 \pm 0.022$  GeV, Γ(hadrons) = 1742 ± 19 MeV and  $\alpha_s = 0.125 \pm 0.009$ . To obtain this branching ratio we divide their value of  $C_{1/3} = 1.63 \pm 0.15$  by their value of  $(3C_{1/3} + 2C_{2/3}) = 6.720 \pm 0.076$ .

 $R_c = \Gamma(c\overline{c})/\Gamma(\text{hadrons})$  OUR FIT is obtained by a simultaneous fit to several c- and b-quark measurements

OUR FIT is obtained by a simultaneous fit to several c- and b-quark measurements as explained in the note "The Z boson" and ref. LEP-SLC 06.

The Standard Model predicts  $R_c=0.1723$  for  $m_t=174.3$  GeV and  $M_H=150$  GeV.

VALUE	DOCUMENT ID		TECN	COMMENT
0.1721±0.0030 OUR FIT				
$0.1744 \pm 0.0031 \pm 0.0021$	<sup>1</sup> ABE	05F	SLD	<i>E</i> <sup>ee</sup> <sub>cm</sub> =91.28 GeV
$0.1665 \pm 0.0051 \pm 0.0081$	<sup>2</sup> ABREU			<i>E</i> <sup>ee</sup> <sub>cm</sub> = 88−94 GeV
$0.1698 \pm 0.0069$	<sup>3</sup> BARATE	<b>00</b> B	ALEP	<i>E</i> <sup>ee</sup> <sub>cm</sub> = 88−94 GeV
$0.180 \pm 0.011 \pm 0.013$	<sup>4</sup> ACKERSTAFF	98E	OPAL	<i>E</i> <sup>ee</sup> <sub>cm</sub> = 88−94 GeV
$0.167\ \pm0.011\ \pm0.012$	<sup>5</sup> ALEXANDER	<b>96</b> R	OPAL	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
• • • We do not use the fo	llowing data for a	verag	es, fits, I	imits, etc. • • •
$0.1623 \pm 0.0085 \pm 0.0209$	<sup>6</sup> ABREU	<b>95</b> D	DLPH	$E_{cm}^{ee} = 88-94 \text{ GeV}$

- $^1$  ABE 05F use hadronic Z decays collected during 1996–98 to obtain an enriched sample of  $c\overline{c}$  events using a double tag method. The single c–tag is obtained with a neural network trained to perform flavor discrimination using as input several signatures (corrected secondary vertex mass, vertex decay length, multiplicity and total momentum of the hemisphere). A multitag approach is used, defining 4 regions of the output value of the neural network and  $R_c$  is extracted from a simultaneous fit to the count rates of the 4 different tags. The quoted systematic error includes an uncertainty of  $\pm 0.0006$  due to the uncertainty on  $R_b$ .
- $^2$  ABREU 00 obtain this result properly combining the measurement from the  $D^{*+}$  production rate ( $R_c = 0.1610 \pm 0.0104 \pm 0.0077 \pm 0.0043$  (BR)) with that from the overall charm counting ( $R_c = 0.1692 \pm 0.0047 \pm 0.0063 \pm 0.0074$  (BR)) in  $c\,\overline{c}$  events. The systematic error includes an uncertainty of  $\pm 0.0054$  due to the uncertainty on the charmed hadron branching fractions.
- <sup>3</sup> BARATE 00B use exclusive decay modes to independently determine the quantities  $R_c \times \mathrm{f}(c \to \mathrm{X}), \ \mathrm{X} = D^0, \ D^+, \ D_s^+, \ \mathrm{and} \ \Lambda_c.$  Estimating  $R_c \times \mathrm{f}(c \to \Xi_c / \Omega_c) = 0.0034$ , they simply sum over all the charm decays to obtain  $R_c = 0.1738 \pm 0.0047 \pm 0.0088 \pm 0.0075 (\mathrm{BR})$ . This is combined with all previous ALEPH measurements (BARATE 98T and BUSKULIC 94G,  $R_c = 0.1681 \pm 0.0054 \pm 0.0062$ ) to obtain the quoted value.
- <sup>4</sup> ACKERSTAFF 98E use an inclusive/exclusive double tag. In one jet  $D^{*\pm}$  mesons are exclusively reconstructed in several decay channels and in the opposite jet a slow pion (opposite charge inclusive  $D^{*\pm}$ ) tag is used. The b content of this sample is measured by the simultaneous detection of a lepton in one jet and an inclusively reconstructed  $D^{*\pm}$  meson in the opposite jet. The systematic error includes an uncertainty of  $\pm 0.006$  due to the external branching ratios.
- <sup>5</sup> ALEXANDER 96R obtain this value via direct charm counting, summing the partial contributions from  $D^0$ ,  $D^+$ ,  $D_s^+$ , and  $\Lambda_c^+$ , and assuming that strange-charmed baryons account for the 15% of the  $\Lambda_c^+$  production. An uncertainty of  $\pm 0.005$  due to the uncertainties in the charm hadron branching ratios is included in the overall systematics.
- <sup>6</sup> ABREU 95D perform a maximum likelihood fit to the combined p and  $p_T$  distributions of single and dilepton samples. The second error includes an uncertainty of  $\pm 0.0124$  due to models and branching ratios.

### $R_b = \Gamma(b\overline{b})/\Gamma(\text{hadrons})$

 $\Gamma_{11}/\Gamma_7$ 

OUR FIT is obtained by a simultaneous fit to several c- and b-quark measurements as explained in the note "The Z boson" and ref. LEP-SLC 06.

The Standard Model predicts  $R_b$ =0.21581 for  $m_t$ =174.3 GeV and  $M_H$ =150 GeV.

VALUE	DOCUMENT ID		TECN	COMMENT
0.21629±0.00066 OUR FIT	·			
$0.21594 \pm 0.00094 \pm 0.00075$	$^{ m 1}$ ABE	05F	SLD	E <sup>ee</sup> <sub>cm</sub> =91.28 GeV
$0.2174\ \pm0.0015\ \pm0.0028$	<sup>2</sup> ACCIARRI	00	L3	<i>E</i> <sup>ee</sup> <sub>cm</sub> = 89−93 GeV
$0.2178 \pm 0.0011 \pm 0.0013$	<sup>3</sup> ABBIENDI	<b>99</b> B	OPAL	<i>E</i> <sup>ee</sup> <sub>cm</sub> = 88−94 GeV
$0.21634 \pm 0.00067 \pm 0.00060$	<sup>4</sup> ABREU	<b>99</b> B	DLPH	<i>E</i> <sup>ee</sup> <sub>cm</sub> = 88−94 GeV
$0.2159\ \pm0.0009\ \pm0.0011$	<sup>5</sup> BARATE	97F	ALEP	<i>E</i> <sup>ee</sup> <sub>cm</sub> = 88−94 GeV
• • • We do not use the following	ng data for averag	es, fit	s, limits,	etc. • • •
$0.2145\ \pm0.0089\ \pm0.0067$	<sup>6</sup> ABREU	<b>95</b> D	DLPH	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
$0.219 \pm 0.006 \pm 0.005$	<sup>7</sup> BUSKULIC	<b>94</b> G	ALEP	<i>E</i> <sup>ee</sup> <sub>cm</sub> = 88−94 GeV
$0.251 \pm 0.049 \pm 0.030$	<sup>8</sup> JACOBSEN	91	MRK2	$E_{\rm cm}^{\rm ee} = 91~{\rm GeV}$

- $^1$  ABE 05F use hadronic Z decays collected during 1996–98 to obtain an enriched sample of  $b\overline{b}$  events using a double tag method. The single b–tag is obtained with a neural network trained to perform flavor discrimination using as input several signatures (corrected secondary vertex mass, vertex decay length, multiplicity and total momentum of the hemisphere; the key tag is obtained requiring the secondary vertex corrected mass to be above the  $D{\rm -meson}$  mass). ABE 05F obtain  $R_b=0.21604\pm0.00098\pm0.00074$  where the systematic error includes an uncertainty of  $\pm0.00012$  due to the uncertainty on  $R_c$ . The value reported here is obtained properly combining with ABE 98D. The quoted systematic error includes an uncertainty of  $\pm0.00012$  due to the uncertainty on  $R_c$ .
- $^2$  ACCIARRI 00 obtain this result using a double-tagging technique, with a high  $p_T$  lepton tag and an impact parameter tag in opposite hemispheres.
- <sup>3</sup> ABBIENDI 99B tag  $Z \rightarrow b\overline{b}$  decays using leptons and/or separated decay vertices. The b-tagging efficiency is measured directly from the data using a double-tagging technique.
- <sup>4</sup> ABREU 99B obtain this result combining in a multivariate analysis several tagging methods (impact parameter and secondary vertex reconstruction, complemented by event shape variables). For  $R_c$  different from its Standard Model value of 0.172,  $R_b$  varies as  $-0.024 \times (R_c 0.172)$ .
- $^5$  BARATE 97F combine the lifetime-mass hemisphere tag (BARATE 97E) with event shape information and lepton tag to identify  $Z \to b \, \overline{b}$  candidates. They further use c- and  $u \, d \, s\text{-}$  selection tags to identify the background. For  $R_{C}$  different from its Standard Model value of 0.172,  $R_{D}$  varies as  $-0.019 \times (R_{C}-0.172)$ .
- <sup>6</sup> ABREU 95D perform a maximum likelihood fit to the combined p and  $p_T$  distributions of single and dilepton samples. The second error includes an uncertainty of  $\pm 0.0023$  due to models and branching ratios.
- $^{7}$  BUSKULIC 94G perform a simultaneous fit to the  $\it p$  and  $\it p_T$  spectra of both single and dilepton events.
- <sup>8</sup> JACOBSEN 91 tagged  $b\overline{b}$  events by requiring coincidence of  $\geq$  3 tracks with significant impact parameters using vertex detector. Systematic error includes lifetime and decay uncertainties ( $\pm 0.014$ ).

# $\Gamma(b\overline{b}b\overline{b})/\Gamma(hadrons)$

 $\Gamma_{12}/\Gamma_7$ 

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VALUE (units $10^{-4}$ )	DOCUMENT ID		TECN	COMMENT
5.2±1.9 OUR AVERAGE				
$3.6 \pm 1.7 \pm 2.7$	<sup>1</sup> ABBIENDI	<b>01</b> G	OPAL	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
$6.0 \pm 1.9 \pm 1.4$	<sup>2</sup> ABREU	<b>99</b> U	DLPH	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV

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 $\Gamma(ggg)/\Gamma(hadrons)$  $\Gamma_{13}/\Gamma_{7}$ 96s DLPH *E*<sup>ee</sup><sub>cm</sub>= 88–94 GeV

<sup>&</sup>lt;sup>1</sup> This branching ratio is slightly dependent on the jet-finder algorithm. The value we quote is obtained using the JADE algorithm, while using the DURHAM algorithm ABREU 96S obtain an upper limit of  $1.5 \times 10^{-2}$ .

$\Gamma(\pi^0\gamma)/\Gamma_{ m total}$					Γ <sub>14</sub> /Γ
VALUE	<u>CL%</u>	DOCUMENT ID		TECN	COMMENT
$< 5.2 \times 10^{-5}$	95	<sup>1</sup> ACCIARRI	<b>95</b> G	L3	<i>E</i> <sup>ee</sup> <sub>cm</sub> = 88−94 GeV
$< 5.5 \times 10^{-5}$	95	ABREU	<b>94</b> B	DLPH	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
$< 2.1 \times 10^{-4}$	95	DECAMP	92	ALEP	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
$< 1.4 \times 10^{-4}$	95	AKRAWY	91F	OPAL	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV

<sup>&</sup>lt;sup>1</sup> This limit is for both decay modes  $Z \to \pi^0 \gamma/\gamma \gamma$  which are indistinguishable in ACCIA-RRI 95G.

$\Gamma(\eta\gamma)/\Gamma_{total}$						Γ <sub>15</sub> /Γ
VALUE	CL%	DOCUMENT ID		TECN	COMMENT	
$< 7.6 \times 10^{-5}$	95	ACCIARRI	95G	L3	$E_{\rm cm}^{ee} = 88-94$	GeV
$< 8.0 \times 10^{-5}$	95	ABREU	<b>94</b> B	DLPH	$E_{cm}^{ee} = 88-94$	GeV
$< 5.1 \times 10^{-5}$	95	DECAMP	92	ALEP	$E_{cm}^{ee} = 88-94$	GeV
$< 2.0 \times 10^{-4}$	95	AKRAWY	91F	OPAL	$E_{\rm cm}^{\it ee} = 88-94$	GeV
$\Gamma(\omega\gamma)/\Gamma_{ m total}$						Γ <sub>16</sub> /Γ
VALUE	<u>CL%</u>	DOCUMENT ID		TECN	COMMENT	
$<6.5 \times 10^{-4}$	95	ABREU	<b>94</b> B	DLPH	$E_{\rm cm}^{ee} = 88-94$	GeV
$\Gamma(\eta'(958)\gamma)/\Gamma_{total}$						Γ <sub>17</sub> /Γ
VALUE	CL%	DOCUMENT ID		TECN	COMMENT	
$<4.2 \times 10^{-5}$	95	DECAMP	92	ALEP	$E_{\rm cm}^{\rm ee} = 88-94$	GeV
$\Gamma(\gamma\gamma)/\Gamma_{ m total}$						Γ <sub>18</sub> /Γ
This decay would		-	eorem			
VALUE	<u>CL%</u>	DOCUMENT ID		<u>TECN</u>	COMMENT	
$< 5.2 \times 10^{-5}$	95	<sup>1</sup> ACCIARRI	<b>95</b> G	L3	$E_{\rm cm}^{ee} = 88-94$	GeV
$< 5.5 \times 10^{-5}$	95	ABRFU	<b>94</b> B	DI PH	$F_{em}^{ee} = 88-94$	GeV

<sup>94</sup>B DLPH  $E_{cm}^{ee} = 88-94 \text{ GeV}$  $< 5.5 \times 10^{\circ}$ 91F OPAL  $E_{cm}^{ee}$ = 88–94 GeV  $<1.4 \times 10^{-4}$ AKRAWY

 $<sup>^{</sup>m 1}$  ABBIENDI 01G use a sample of four-jet events from hadronic Z decays. To enhance the  $b\overline{b}b\overline{b}$  signal, at least three of the four jets are required to have a significantly detached

 $<sup>^2</sup>$ ABREU 990 force hadronic Z decays into 3 jets to use all the available phase space and require a btag for every jet. This decay mode includes primary and secondary 4b production, e.g, from gluon splitting to  $b\overline{b}$ .

 $<sup>^1</sup>$  This limit is for both decay modes  $Z 
ightarrow ~\pi^0 \gamma/\gamma\gamma$  which are indistinguishable in ACCIA-RRI 95G.

$\Gamma(\gamma\gamma\gamma)/\Gamma_{total}$					Γ <sub>19</sub> /Γ
VALUE	CL%	DOCUMENT ID		TECN	COMMENT
$< 1.0 \times 10^{-5}$	95	<sup>1</sup> ACCIARRI	<b>95</b> C	L3	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
$< 1.7 \times 10^{-5}$	95	$^{ m 1}$ ABREU	<b>94</b> B	DLPH	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
$< 6.6 \times 10^{-5}$	95	AKRAWY	91F	OPAL	$E_{\rm cm}^{\it ee} = 88 - 94 \; {\rm GeV}$

 $<sup>^{1}</sup>$ Limit derived in the context of composite Z model.

 $\Gamma(\pi^{\pm}W^{\mp})/\Gamma_{\text{total}}$ 

 $\Gamma_{20}/\Gamma$ 

The value is for the sum of the charge states indicated.

VALUE	CL%	DOCUMENT ID		TECN	COMMENT
$< 7 \times 10^{-5}$	95	DECAMP	92	ALEP	Eee = 88–94 GeV

# $\Gamma(\rho^{\pm}W^{\mp})/\Gamma_{\text{total}}$

 $\Gamma_{21}/\Gamma$ 

The value is for the sum of the charge states indicated.

VALUE	CL%	DOCUMENT ID		TECN	COMMENT
<8.3 × 10 <sup>-5</sup>	95	DECAMP	92	ALEP	Eee = 88–94 GeV

### $\Gamma(J/\psi(1S)X)/\Gamma_{\text{total}}$

 $\Gamma_{22}/\Gamma$ 

VALUE (units 10<sup>-3</sup>) EVTS DOCUMENT ID TECN COMMENT

# $3.51^{+0.23}_{-0.25}$ **OUR AVERAGE** Error includes scale factor of 1.1.

$3.21 \pm 0.21 {+0.19 \atop -0.28}$	553	<sup>1</sup> ACCIARRI	99F	L3	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
$3.9 \pm 0.2 \pm 0.3$	511	<sup>2</sup> ALEXANDER	<b>96</b> B	OPAL	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
$3.73 \pm 0.39 \pm 0.36$	153	<sup>3</sup> ABREU	<b>94</b> P	DLPH	$E_{\rm cm}^{ee} = 88-94 {\rm GeV}$

 $<sup>^1</sup>$  ACCIARRI 99F combine  $\mu^+\,\mu^-$  and  $e^+\,e^-\,J/\psi(1S)$  decay channels. The branching ratio for prompt  $J/\psi(1S)$  production is measured to be  $(2.1\pm0.6\pm0.4^{+0.4}_{-0.2}(\text{theor.}))\times10^{-4}$  .

# $\Gamma(\psi(2S)X)/\Gamma_{\text{total}}$

 $\Gamma_{23}/\Gamma$ 

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VALUE (units $10^{-3}$ )	EVTS	DOCUMENT ID		TECN	COMMENT
1.60±0.29 OUR AVER	AGE				
$1.6 \pm 0.5 \pm 0.3$	39	<sup>1</sup> ACCIARRI	97J	L3	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
$1.6 \pm 0.3 \pm 0.2$	46.9	<sup>2</sup> ALEXANDER	<b>96</b> B	OPAL	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
$1.60 \pm 0.73 \pm 0.33$	5.4	<sup>3</sup> ABREU	94P	DLPH	$E_{cm}^{ee} = 88-94 \text{ GeV}$

<sup>&</sup>lt;sup>1</sup> ACCIARRI 97J measure this branching ratio via the decay channel  $\psi(2S) \rightarrow \ell^+\ell^-$  ( $\ell = \mu, e$ ).

<sup>&</sup>lt;sup>2</sup> ALEXANDER 96B identify  $J/\psi(1S)$  from the decays into lepton pairs. (4.8  $\pm$  2.4)% of this branching ratio is due to prompt  $J/\psi(1S)$  production (ALEXANDER 96N).

<sup>&</sup>lt;sup>3</sup> Combining  $\mu^+\mu^-$  and  $e^+e^-$  channels and taking into account the common systematic errors.  $(7.7^{+6.3}_{-5.4})\%$  of this branching ratio is due to prompt  $J/\psi(1S)$  production.

<sup>&</sup>lt;sup>2</sup> ALEXANDER 96B measure this branching ratio via the decay channel  $\psi(2S) \to J/\psi \, \pi^+ \, \pi^-$ , with  $J/\psi \to \ell^+ \ell^-$ .

<sup>&</sup>lt;sup>3</sup> ABREU 94P measure this branching ratio via decay channel  $\psi(2S) \to J/\psi \pi^+ \pi^-$ , with  $J/\psi \to \mu^+ \mu^-$ .

$\Gamma(\chi_{c1}(1P)X)/\Gamma_{tc}$	otal					$\Gamma_{24}/\Gamma$
$VALUE$ (units $10^{-3}$ )	EVTS	DOCUMENT ID		TECN	COMMENT	
2.9±0.7 OUR AVER	AGE					
07106105	22	1 ACCIADDI	07.	1.0	ree 00 04	C 1/

 $^{1}$  ACCIARRI 97J L3  $E_{cm}^{ee}$  = 88–94 GeV  $2.7 \pm 0.6 \pm 0.5$  $5.0\pm2.1^{+1.5}_{-0.9}$  $^2$  ABREU 94P DLPH  $E_{cm}^{ee} = 88-94$  GeV

### $\Gamma(\chi_{c2}(1P)X)/\Gamma_{total}$

 $\Gamma_{25}/\Gamma$ 

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<u>VALUE</u>	CL%	DOCUMENT ID		TECN	<u>COMMENT</u>	
$< 3.2 \times 10^{-3}$	90	<sup>1</sup> ACCIARRI	97J	L3	Eee = 88–94 GeV	,

 $<sup>^1</sup>$  ACCIARRI 97J derive this limit via the decay channel  $\chi_{c2} 
ightarrow ~J/\psi + ~\gamma$ , with  $J/\psi 
ightarrow$  $\ell^+\ell^-$  ( $\ell=\mu$ , e). The  $M(\ell^+\ell^-\gamma)$ – $M(\ell^+\ell^-)$  mass difference spectrum is fitted with two gaussian shapes for  $\chi_{c1}$  and  $\chi_{c2}$ .

### $\Gamma(\Upsilon(1S)X)/\Gamma_{\text{total}}$

 $\Gamma_{27}/\Gamma$ 

<u>VALUE</u>	CL%	DOCUMENT ID		TECN	COMMENT
$<4.4 \times 10^{-5}$	95	<sup>1</sup> ACCIARRI	99F	L3	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV

<sup>&</sup>lt;sup>1</sup> ACCIARRI 99F search for  $\Upsilon(1S)$  through its decay into  $\ell^+\ell^-$  ( $\ell=e$  or  $\mu$ ).

# $\Gamma(\Upsilon(2S)X)/\Gamma_{total}$

 $\Gamma_{28}/\Gamma$ 

<u>VALUE</u>	CL%	DOCUMENT ID		TECN	COMMENT
$<13.9 \times 10^{-5}$	95	<sup>1</sup> ACCIARRI	<b>97</b> R	L3	Eee = 88–94 GeV

<sup>&</sup>lt;sup>1</sup> ACCIARRI 97R search for  $\Upsilon(2S)$  through its decay into  $\ell^+\ell^-$  ( $\ell=e$  or  $\mu$ ).

# $\Gamma(\Upsilon(3S)X)/\Gamma_{\text{total}}$

 $\Gamma_{29}/\Gamma$ 

VALUE	CL%	DOCUMENT ID		TECN	COMMENT
$< 9.4 \times 10^{-5}$	95	<sup>1</sup> ACCIARRI	<b>97</b> R	L3	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV

<sup>&</sup>lt;sup>1</sup> ACCIARRI 97R search for  $\Upsilon(3S)$  through its decay into  $\ell^+\ell^-$  ( $\ell=e$  or  $\mu$ ).

# $\Gamma((D^0/\overline{D}^0)X)/\Gamma(\text{hadrons})$

 $\Gamma_{30}/\Gamma_{7}$ 

<u>VALUE</u>	<u>EVTS</u>	<u>DOCUMENT ID</u>		TECN	<u>COMMENT</u>
$0.296\pm0.019\pm0.021$	369	<sup>1</sup> ABREU	931	DLPH	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV

<sup>&</sup>lt;sup>1</sup>The  $(D^0/\overline{D}^0)$  states in ABREU 931 are detected by the  $K\pi$  decay mode. This is a corrected result (see the erratum of ABREU 931).

 $<sup>^1</sup>$ ACCIARRI 97J measure this branching ratio via the decay channel  $\chi_{c1} 
ightarrow ~J/\psi + ~\gamma$ , with  $J/\psi \to \ell^+\ell^-$  ( $\ell=\mu$ , e). The  $M(\ell^+\ell^-\gamma)-M(\ell^+\ell^-)$  mass difference spectrum is fitted with two gaussian shapes for  $\chi_{c1}$  and  $\chi_{c2}$ .

 $<sup>^2</sup>$  This branching ratio is measured via the decay channel  $\chi_{c1} o J/\psi + \gamma$ , with  $J/\psi o$  $\mu^{+}\mu^{-}$ .

 $<sup>^1</sup>$  ALEXANDER 96F identify the  $\Upsilon$  (which refers to any of the three lowest bound states) through its decay into  $e^+e^-$  and  $\mu^+\mu^-$ . The systematic error includes an uncertainty of  $\pm 0.2$  due to the production mechanism.

$\Gamma(D^{\pm}X)/\Gamma(\text{hadrons})$	)				$\Gamma_{31}/\Gamma_7$
VALUE	<b>EVTS</b>	DOCUMENT ID		TECN	COMMENT
$0.174 \pm 0.016 \pm 0.018$	539	<sup>1</sup> ABREU	931	DLPH	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV

<sup>&</sup>lt;sup>1</sup> The  $D^{\pm}$  states in ABREU 93I are detected by the  $K\pi\pi$  decay mode. This is a corrected result (see the erratum of ABREU 93I).

### $\Gamma(D^*(2010)^{\pm}X)/\Gamma(hadrons)$

 $\Gamma_{32}/\Gamma_{7}$ 

The value is for the sum of the charge states indicated.

<u>VALUE</u>	<b>EVTS</b>	DOCUMENT ID		TECN	COMMENT
0.163 ± 0.019 OUR AVE	RAGE	Error includes scal	e facto	r of 1.3.	
$0.155 \pm 0.010 \pm 0.013$	358	$^{ m 1}$ ABREU	931	DLPH	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
$0.21 \pm 0.04$	362	<sup>2</sup> DECAMP	<b>91</b> J	ALEP	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV

 $<sup>^1</sup>D^*(2010)^\pm$  in ABREU 93I are reconstructed from  $D^0\pi^\pm$ , with  $D^0\to K^-\pi^+$ . The new CLEO II measurement of B( $D^{*\pm}\to D^0\pi^\pm$ ) = (68.1  $\pm$  1.6) % is used. This is a corrected result (see the erratum of ABREU 93I).

# $\Gamma(D_{s1}(2536)^{\pm}X)/\Gamma(hadrons)$

 $\Gamma_{33}/\Gamma_{7}$ 

 $D_{\rm s1}(2536)^{\pm}$  is an expected orbitally-excited state of the  $D_{\rm s}$  meson.

VALUE (%)	<b>EVTS</b>	DOCUMENT ID		TECN	COMMENT
$0.52 \pm 0.09 \pm 0.06$	92	$^{ m 1}$ HEISTER	<b>02</b> B	ALEP	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV

 $<sup>^1</sup>$  HEISTER 02B reconstruct this meson in the decay modes  $D_{s1}(2536)^\pm\to D^{*\pm}\,K^0$  and  $D_{s1}(2536)^\pm\to D^{*0}\,K^\pm.$  The quoted branching ratio assumes that the decay width of the  $D_{s1}(2536)$  is saturated by the two measured decay modes.

# $\Gamma(D_{s,J}(2573)^{\pm}X)/\Gamma(hadrons)$

 $\Gamma_{34}/\Gamma_{7}$ 

 $D_{sJ}$ (2573) $^{\pm}$  is an expected orbitally-excited state of the  $D_{
m S}$  meson.

VALUE (%)	EVTS	DOCUMENT ID		TECN	COMMENT
$0.83 \pm 0.29 ^{+0.07}_{-0.13}$	64	<sup>1</sup> HEISTER	<b>02</b> B	ALEP	Eee = 88–94 GeV

 $<sup>^1</sup>$  HEISTER 02B reconstruct this meson in the decay mode  $D_{s2}^*(2573)^\pm \to D^0 \, K^\pm$ . The quoted branching ratio assumes that the detected decay mode represents 45% of the full decay width.

# $\Gamma(D^{*\prime}(2629)^{\pm}X)/\Gamma(hadrons)$

 $\Gamma_{35}/\Gamma_{7}$ 

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 $D^{*\prime}(2629)^{\pm}$  is a predicted radial excitation of the  $D^{*}(2010)^{\pm}$  meson.

VALUEDOCUMENT IDTECNCOMMENTsearched for $^1$  ABBIENDI01NOPAL $E_{\rm cm}^{ee}=88-94$  GeV

<sup>&</sup>lt;sup>2</sup> DECAMP 91J report B( $D^*(2010)^+ \to D^0\pi^+$ ) B( $D^0 \to K^-\pi^+$ )  $\Gamma(D^*(2010)^\pm X)$  /  $\Gamma(\text{hadrons}) = (5.11 \pm 0.34) \times 10^{-3}$ . They obtained the above number assuming B( $D^0 \to K^-\pi^+$ ) = (3.62  $\pm$  0.34  $\pm$  0.44)% and B( $D^*(2010)^+ \to D^0\pi^+$ ) = (55  $\pm$  4)%. We have rescaled their original result of 0.26  $\pm$  0.05 taking into account the new CLEO II branching ratio B( $D^*(2010)^+ \to D^0\pi^+$ ) = (68.1  $\pm$  1.6)%.

<sup>&</sup>lt;sup>1</sup> ABBIENDI 01N searched for the decay mode  $D^{*\prime}(2629)^{\pm} \rightarrow D^{*\pm}\pi^{+}\pi^{-}$  with  $D^{*+} \rightarrow D^{0}\pi^{+}$ , and  $D^{0} \rightarrow K^{-}\pi^{+}$ . They quote a 95% CL limit for  $Z \rightarrow D^{*\prime}(2629)^{\pm} \times B(D^{*\prime}(2629)^{+} \rightarrow D^{*+}\pi^{+}\pi^{-}) < 3.1 \times 10^{-3}$ .

### $\Gamma(B^*X)/|\Gamma(BX)+\Gamma(B^*X)|$

 $\Gamma_{37}/(\Gamma_{36}+\Gamma_{37})$ 

As the experiments assume different values of the b-baryon contribution, our average should be taken with caution.

VALUE	<b>EVTS</b>	DOCUMENT ID		TECN	COMMENT
$0.75 \pm 0.04$ OUR AVE	RAGE				
$0.760 \pm 0.036 \pm 0.083$		<sup>1</sup> ACKERSTAFF	97M	OPAL	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
$0.771 \pm 0.026 \pm 0.070$		<sup>2</sup> BUSKULIC	<b>96</b> D	ALEP	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
$0.72\ \pm0.03\ \pm0.06$		<sup>3</sup> ABREU	<b>95</b> R	DLPH	<i>E</i> <sup>ee</sup> <sub>cm</sub> = 88−94 GeV
$0.76\ \pm0.08\ \pm0.06$	1378	<sup>4</sup> ACCIARRI	<b>95</b> B	L3	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV

 $<sup>^1</sup>$ ACKERSTAFF 97M use an inclusive B reconstruction method and assume a (13.2  $\pm$ 4.1)% b-baryon contribution. The value refers to a b-flavored meson mixture of  $B_{\mu}$ ,  $B_{d}$ ,

### $\Gamma(B^+X)/\Gamma(hadrons)$

 $\Gamma_{38}/\Gamma_{7}$ 

"OUR EVALUATION" is obtained using our current values for f( $\overline{b} 
ightarrow B^+$ ) and R $_b$  $=\Gamma(b\overline{b})/\Gamma(\text{hadrons})$ . We calculate  $\Gamma(B^+ X)/\Gamma(\text{hadrons}) = R_b \times f(\overline{b} \to B^+)$ . The decay fraction  $f(\overline{b} \to B^+)$  was provided by the Heavy Flavor Averaging Group (HFAG, http://www.slac.stanford.edu/xorg/hfag/osc/PDG\_2009/#FRACZ).

#### DOCUMENT ID TECN COMMENT $0.0869\pm0.0019$ OUR EVALUATION

 $0.0887 \pm 0.0030$ 

 $^{1}$  ABDALLAH 03K DLPH  $E_{cm}^{ee} = 88-94$  GeV

# $\Gamma(B_s^0X)/\Gamma(hadrons)$

 $\Gamma_{39}/\Gamma_{7}$ 

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"OUR EVALUATION" is obtained using our current values for f( $\overline{b} 
ightarrow \ B_s^0$ ) and R $_b$  $= \Gamma(b\,\overline{b})/\Gamma(\text{hadrons}). \text{ We calculate } \Gamma(B_{s}^{0})/\Gamma(\text{hadrons}) = R_{b} \times f(\overline{b} \to B_{s}^{0}). \text{ The }$ decay fraction f( $\overline{b} \to B^0_s$ ) was provided by the Heavy Flavor Averaging Group (HFAG,  $http://www.slac.stanford.edu/xorg/hfag/osc/PDG\_2009/\#FRACZ)$ 

VALUE	DOCUMENT ID		TECN	COMMENT
0.0227±0.0019 OUR EVALUATIO	N			
seen	<sup>1</sup> ABREU	92M	DLPH	<i>E</i> <sup>ee</sup> <sub>cm</sub> = 88−94 GeV
seen	<sup>2</sup> ACTON	92N	OPAL	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
seen	<sup>3</sup> BUSKULIC	92E	ALEP	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV

 $<sup>^1</sup>$  ABREU 92M reported value is  $\Gamma(B_s^0{\rm X})*{\rm B}(B_s^0\to\ D_s\,\mu\nu_\mu{\rm X})*{\rm B}(D_s\to\ \phi\pi)/\Gamma({\rm hadrons})$  $= (18 \pm 8) \times 10^{-5}$ .

 $<sup>^2</sup>$ BUSKULIC 96D use an inclusive reconstruction of B hadrons and assume a (12.2  $\pm$ 4.3)% b-baryon contribution. The value refers to a b-flavored mixture of  $B_{\mu\nu}$ ,  $B_{cd}$ , and

 $<sup>^3</sup>$  ABREU 95R use an inclusive *B*-reconstruction method and assume a  $(10\pm4)\%$  *b*-baryon contribution. The value refers to a b-flavored meson mixture of  $B_u$ ,  $B_d$ , and  $B_s$ .

 $<sup>^4</sup>$  ACCIARRI 95B assume a 9.4% *b*-baryon contribution. The value refers to a *b*-flavored mixture of  $B_{\mu}$ ,  $B_{d}$ , and  $B_{s}$ .

 $<sup>^{1}</sup>$ ABDALLAH 03K measure the production fraction of  $B^{+}$  mesons in hadronic Z decays  $f(B^+) = (40.99 \pm 0.82 \pm 1.11)\%$ . The value quoted here is obtained multiplying this production fraction by our value of  $R_b = \Gamma(\overline{b}b)/\Gamma(\text{hadrons})$ .

<sup>&</sup>lt;sup>2</sup> ACTON 92N find evidence for  $B_s^0$  production using  $D_s$ - $\ell$  correlations, with  $D_s^+ o \phi \pi^+$ and  $K^*(892)K^+$ . Assuming  $R_b$  from the Standard Model and averaging over the e and

 $\mu$  channels, authors measure the product branching fraction to be  $f(\overline{b} \to B_s^0) \times B(B_s^0 \to D_s^- \ell^+ \nu_\ell X) \times B(D_s^- \to \phi \pi^-) = (3.9 \pm 1.1 \pm 0.8) \times 10^{-4}$ .

 $^3$  BUSKULIC 92E find evidence for  $B_s^0$  production using  $D_s$ - $\ell$  correlations, with  $D_s^+ \to \phi \pi^+$  and  $K^*(892)\,K^+$ . Using B( $D_s^+ \to \phi \pi^+$ ) = (2.7  $\pm$  0.7)% and summing up the e and  $\mu$  channels, the weighted average product branching fraction is measured to be B( $\overline{b} \to B_s^0$ )×B( $B_s^0 \to D_s^- \ell^+ \nu_\ell X$ ) = 0.040  $\pm$  0.011  $^{+0.010}_{-0.012}$ .

# $\Gamma(B_c^+X)/\Gamma(hadrons)$

 $\Gamma_{40}/\Gamma_{7}$ 

VALUE	DOCUMENT ID	TECN	COMMENT
searched for	<sup>1</sup> ACKERSTAFF 980	OPAL	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
searched for	<sup>2</sup> ABREU 97E	DLPH	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
searched for	<sup>3</sup> BARATE 97H	ALEP	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV

<sup>1</sup> ACKERSTAFF 980 searched for the decay modes  $B_c \to J/\psi \pi^+$ ,  $J/\psi a_1^+$ , and  $J/\psi \ell^+ \nu_\ell$ , with  $J/\psi \to \ell^+ \ell^-$ ,  $\ell = e, \mu$ . The number of candidates (background) for the three decay modes is 2 (0.63 ± 0.2), 0 (1.10 ± 0.22), and 1 (0.82 ± 0.19) respectively. Interpreting the 2  $B_c \to J/\psi \pi^+$  candidates as signal, they report  $\Gamma(B_c^+ X) \times B(B_c \to J/\psi \pi^+)/\Gamma(\text{hadrons}) = (3.8^{+5.0}_{-2.4} \pm 0.5) \times 10^{-5}$ . Interpreted as background, the 90% CL bounds are  $\Gamma(B_c^+ X) * B(B_c \to J/\psi \pi^+)/\Gamma(\text{hadrons}) < 1.06 \times 10^{-4}$ ,  $\Gamma(B_c^+ X) * B(B_c \to J/\psi a_1^+)/\Gamma(\text{hadrons}) < 5.29 \times 10^{-4}$ ,  $\Gamma(B_c^+ X) * B(B_c \to J/\psi \ell^+ \nu_\ell)/\Gamma(\text{hadrons}) < 6.96 \times 10^{-5}$ .

<sup>2</sup> ABREU 97E searched for the decay modes  $B_C \to J/\psi \pi^+$ ,  $J/\psi \ell^+ \nu_\ell$ , and  $J/\psi (3\pi)^+$ , with  $J/\psi \to \ell^+ \ell^-$ ,  $\ell = e, \mu$ . The number of candidates (background) for the three decay modes is 1 (1.7), 0 (0.3), and 1 (2.3) respectively. They report the following 90% CL limits:  $\Gamma(B_C^+ X)*B(B_C \to J/\psi \pi^+)/\Gamma(\text{hadrons}) < (1.05-0.84) \times 10^{-4}$ ,  $\Gamma(B_C^+ X)*B(B_C \to J/\psi \ell^+)/\Gamma(\text{hadrons}) < (5.8-5.0) \times 10^{-5}$ ,  $\Gamma(B_C^+ X)*B(B_C \to J/\psi (3\pi)^+)/\Gamma(\text{hadrons}) < 1.75 \times 10^{-4}$ , where the ranges are due to the predicted  $B_C$  lifetime (0.4–1.4) ps.

<sup>3</sup>BARATE 97H searched for the decay modes  $B_c \to J/\psi \pi^+$  and  $J/\psi \ell^+ \nu_\ell$  with  $J/\psi \to \ell^+ \ell^-$ ,  $\ell = e,\mu$ . The number of candidates (background) for the two decay modes is 0 (0.44) and 2 (0.81) respectively. They report the following 90% CL limits:  $\Gamma(B_c^+ X)*B(B_c \to J/\psi \pi^+)/\Gamma(\text{hadrons}) < 3.6 \times 10^{-5}$  and  $\Gamma(B_c^+ X)*B(B_c \to J/\psi \ell^+ \nu_\ell)/\Gamma(\text{hadrons}) < 5.2 \times 10^{-5}$ .

### $\Gamma(\Lambda_c^+ X)/\Gamma(hadrons)$

 $\Gamma_{41}/\Gamma_{7}$ 

Created: 7/12/2013 14:51

VALUE	DOCUMENT ID		TECN	COMMENT
0.022±0.005 OUR AVERAGE				
$0.024 \pm 0.005 \pm 0.006$	$^{ m 1}$ ALEXANDER	<b>96</b> R	OPAL	$E_{\mathrm{cm}}^{\mathit{ee}} = 88-94 \; \mathrm{GeV}$
$0.021 \pm 0.003 \pm 0.005$	<sup>2</sup> BUSKULIC	96Y	ALEP	$E_{ m cm}^{\it ee}=$ 88–94 GeV

 $^1$  ALEXANDER 96R measure R $_b \times {\rm f}(b \to \Lambda_c^+ X) \times {\rm B}(\Lambda_c^+ \to p \, K^- \, \pi^+) = (0.122 \pm 0.023 \pm 0.010)\%$  in hadronic Z decays; the value quoted here is obtained using our best value B( $\Lambda_c^+ \to p \, K^- \, \pi^+$ ) = (5.0  $\pm$  1.3)%. The first error is the total experiment's error and the second error is the systematic error due to the branching fraction uncertainty.

<sup>2</sup>BUSKULIC 96Y obtain the production fraction of  $\Lambda_c^+$  baryons in hadronic Z decays  $f(b \to \Lambda_c^+ X) = 0.110 \pm 0.014 \pm 0.006$  using  $B(\Lambda_c^+ \to p \, K^- \, \pi^+) = (4.4 \pm 0.6)\%$ ; we have rescaled using our best value  $B(\Lambda_c^+ \to p \, K^- \, \pi^+) = (5.0 \pm 1.3)\%$  obtaining  $f(b \to p \, K^- \, \pi^+) = (5.0 \pm 1.3)\%$ 

 $\Lambda_c^+ X) = 0.097 \pm 0.013 \pm 0.025$  where the first error is their total experiment's error and the second error is the systematic error due to the branching fraction uncertainty. The value quoted here is obtained multiplying this production fraction by our value of  $R_b = \Gamma(b\,\overline{b})/\Gamma(\text{hadrons})$ .

# $\Gamma(\Xi_c^0 X)/\Gamma(hadrons)$

 $\Gamma_{42}/\Gamma_{7}$ 

VALUE DOCUMENT ID TECN COMMENT

• • We do not use the following data for averages, fits, limits, etc.

seen  $^{1}$  ABDALLAH 05C DLPH  $E_{cm}^{ee} = 88-94$  GeV

<sup>1</sup> ABDALLAH 05C searched for the charmed strange baryon  $\Xi_c^0$  in the decay channel  $\Xi_c^0 \to \Xi^- \pi^+ (\Xi^- \to \Lambda \pi^-)$ . The production rate is measured to be  $f_{\Xi_c^0} \times B(\Xi_c^0 \to \Xi^- \pi^+) = (4.7 \pm 1.4 \pm 1.1) \times 10^{-4}$  per hadronic Z decay.

# $\Gamma(\Xi_b X)/\Gamma(hadrons)$

 $\Gamma_{43}/\Gamma_{7}$ 

Here  $\Xi_b$  is used as a notation for the strange *b*-baryon states  $\Xi_b^-$  and  $\Xi_b^0$ .

• • We do not use the following data for averages, fits, limits, etc. • •

seen  $^1$  ABDALLAH 05C DLPH  $E_{\rm cm}^{ee}=88-94~{\rm GeV}$  seen  $^2$  BUSKULIC 96T ALEP  $E_{\rm cm}^{ee}=88-94~{\rm GeV}$  seen  $^3$  ABREU 95V DLPH  $E_{\rm cm}^{ee}=88-94~{\rm GeV}$ 

- <sup>1</sup> ABDALLAH 05C searched for the beauty strange baryon  $\Xi_b$  in the inclusive semileptonic decay channel  $\Xi_b \to \Xi^- \ell^- \overline{\nu}_\ell X$ . Evidence for the  $\Xi_b$  production is seen from the observation of  $\Xi^\mp$  production accompanied by a lepton of the same sign. From the excess of "right-sign" pairs  $\Xi^\mp \ell^\mp$  compared to "wrong-sign" pairs  $\Xi^\mp \ell^\pm$  the production rate is measured to be B( $b \to \Xi_b$ )  $\times$  B( $\Xi_b \to \Xi^- \ell^- X$ ) = (3.0  $\pm$  1.0  $\pm$  0.3)  $\times$  10<sup>-4</sup> per lepton species, averaged over electrons and muons.
- <sup>2</sup> BUSKULIC 96T investigate  $\Xi$ -lepton correlations and find a significant excess of "right-sign" pairs  $\Xi^{\mp}\ell^{\mp}$  compared to "wrong-sign" pairs  $\Xi^{\mp}\ell^{\pm}$ . This excess is interpreted as evidence for  $\Xi_b$  semileptonic decay. The measured product branching ratio is B( $b \to \Xi_b$ )  $\times$  B( $\Xi_b \to X_c X \ell^- \overline{\nu}_\ell$ )  $\times$  B( $X_c \to \Xi^- X'$ ) = (5.4  $\pm$  1.1  $\pm$  0.8)  $\times$  10<sup>-4</sup> per lepton species, averaged over electrons and muons, with  $X_c$  a charmed baryon.
- <sup>3</sup> ABREU 95V observe an excess of "right-sign" pairs  $\Xi^{\mp}\ell^{\mp}$  compared to "wrong-sign" pairs  $\Xi^{\mp}\ell^{\pm}$  in jets: this excess is interpreted as evidence for the beauty strange baryon  $\Xi_b$  production, with  $\Xi_b \to \Xi^-\ell^-\overline{\nu}_\ell X$ . They find that the probability for this signal to come from non b-baryon decays is less than  $5\times 10^{-4}$  and that  $\Lambda_b$  decays can account for less than 10% of these events. The  $\Xi_b$  production rate is then measured to be B( $b\to\Xi_b$ )  $\times$  B( $\Xi_b\to\Xi^-\ell^- X$ ) =  $(5.9\pm 2.1\pm 1.0)\times 10^{-4}$  per lepton species, averaged over electrons and muons.

### $\Gamma(b\text{-baryon X})/\Gamma(\text{hadrons})$

 $\Gamma_{44}/\Gamma_{7}$ 

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"OUR EVALUATION" is obtained using our current values for  $f(b \rightarrow b\text{-baryon})$  and  $R_b = \Gamma(b\overline{b})/\Gamma(\text{hadrons})$ . We calculate  $\Gamma(b\text{-baryon X})/\Gamma(\text{hadrons}) = R_b \times f(b \rightarrow b\text{-baryon})$ . The decay fraction  $f(b \rightarrow b\text{-baryon})$  was provided by the Heavy Flavor Averaging Group (HFAG, http://www.slac.stanford.edu/xorg/hfag/osc/PDG\_2009).

VALUE <u>DOCUMENT ID TECN</u> COMMENT

0.0197±0.0032 OUR EVALUATION

 $^{
m 1}$  BARATE 98V use the overall number of identified protons in  $\it b$ -hadron decays to measure  $f(b \rightarrow b\text{-baryon}) = 0.102 \pm 0.007 \pm 0.027$ . They assume  $BR(b\text{-baryon} \rightarrow pX) = 0.007 \pm 0.007$ (58  $\pm$  6)% and BR( $B_s^0 \to pX$ ) = (8.0  $\pm$  4.0)%. The value quoted here is obtained multiplying this production fraction by our value of R<sub>b</sub> =  $\Gamma(b\overline{b})/\Gamma(\text{hadrons})$ .

#### $\Gamma(\text{anomalous } \gamma + \text{hadrons})/\Gamma_{\text{total}}$

Limits on additional sources of prompt photons beyond expectations for final-state bremsstrahlung.

<u>VALUE</u>	CL%	DOCUMENT ID		TECN	COMMENT
$< 3.2 \times 10^{-3}$	95	<sup>1</sup> AKRAWY	90J	OPAL	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV

 $<sup>^1</sup>$ AKRAWY 90J report  $\Gamma(\gamma {
m X}) < 8.2$  MeV at 95%CL. They assume a three-body  $\gamma q \overline{q}$ distribution and use  $E(\gamma) > 10$  GeV.

# $\Gamma(e^+e^-\gamma)/\Gamma_{\text{total}}$

 $\Gamma_{46}/\Gamma$ 

VALUE	CL%	DOCUMENT ID		TECN	COMMENT
$< 5.2 \times 10^{-4}$	95	<sup>1</sup> ACTON	<b>91</b> B	OPAL	$E_{\rm cm}^{ee} = 91.2$

<sup>&</sup>lt;sup>1</sup> ACTON 91B looked for isolated photons with E>2% of beam energy (> 0.9 GeV).

$$\Gamma(\mu^+\mu^-\gamma)/\Gamma_{\mathsf{total}}$$

 $\Gamma_{47}/\Gamma$ 

VALUE	CL%	DOCUMENT ID		TECN	COMMENT
<5.6 × 10 <sup>-4</sup>	95	<sup>1</sup> ACTON	<b>91</b> B	OPAL	<i>E</i> <sup>ee</sup> <sub>cm</sub> = 91.2 GeV

<sup>&</sup>lt;sup>1</sup> ACTON 91B looked for isolated photons with E>2% of beam energy (> 0.9 GeV).

$$\Gamma(\tau^+\tau^-\gamma)/\Gamma_{\rm total}$$

 $\Gamma_{48}/\Gamma$ 

VALUE	CL%	DOCUMENT ID		TECN	COMMENT
$< 7.3 \times 10^{-4}$	95	$^{ m 1}$ ACTON	<b>91</b> B	OPAL	E <sup>ee</sup> <sub>cm</sub> = 91.2 GeV

<sup>&</sup>lt;sup>1</sup> ACTON 91B looked for isolated photons with E>2% of beam energy (> 0.9 GeV).

# $\Gamma(\ell^+\ell^-\gamma\gamma)/\Gamma_{\text{total}}$

 $\Gamma_{49}/\Gamma$ 

The value is the sum over  $\ell = e, \mu, \tau$ .

VALUE	<u>CL%</u>	DOCUMENT ID		TECN	COMMENT
$<6.8 \times 10^{-6}$	95	<sup>1</sup> ACTON	93E	OPAL	Eee = 88-94 GeV
1 _					

<sup>&</sup>lt;sup>1</sup> For  $m_{\gamma\gamma}=60\pm 5$  GeV.

 $\Gamma_{50}/\Gamma$ 

$$^{1}$$
 For  $m_{\gamma\gamma}=$  60  $\pm$  5 GeV.

$$^{1}$$
 For  $m_{\gamma\gamma}=$  60  $\pm$  5 GeV.

 $\Gamma(e^{\pm}\mu^{\mp})/\Gamma_{\text{total}}$ 

Test of lepton family number conservation. The value is for the sum of the charge states indicated.

VALUE	<u>CL%</u>	DOCUMENT ID		TECN	COMMENT
$< 2.5 \times 10^{-6}$	95	ABREU	97C	DLPH	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
$< 1.7 \times 10^{-6}$	95	AKERS	95W	OPAL	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
$< 0.6 \times 10^{-5}$	95	ADRIANI	931	L3	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
$< 2.6 \times 10^{-5}$	95	DECAMP	92	ALEP	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV

 $\Gamma(e^{\pm}\mu^{\mp})/\Gamma(e^{+}e^{-})$   $\Gamma_{52}/\Gamma_{1}$ 

Test of lepton family number conservation. The value is for the sum of the charge states indicated.

VALUE	CL%	DOCUMENT ID		DOCUMENT ID TECN COMMENT	
<0.07	90	ALBAJAR	89	UA1	$E_{\rm cm}^{p\overline{p}}$ = 546,630 GeV

 $\Gamma(e^{\pm} au^{\mp})/\Gamma_{\mathsf{total}}$   $\Gamma_{\mathsf{53}}/\Gamma$ 

Test of lepton family number conservation. The value is for the sum of the charge states indicated.

VALUE	CL%	DOCUMENT ID		TECN	COMMENT
$< 2.2 \times 10^{-5}$	95	ABREU	<b>97</b> C	DLPH	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
$< 9.8 \times 10^{-6}$	95	AKERS	95W	OPAL	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
$< 1.3 \times 10^{-5}$	95	ADRIANI	931	L3	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
$< 1.2 \times 10^{-4}$	95	DECAMP	92	ALEP	$E_{\rm cm}^{\it ee}=$ 88–94 GeV

 $\Gamma(\mu^{\pm}\tau^{\mp})/\Gamma_{ ext{total}}$   $\Gamma_{54}/\Gamma$ 

Test of lepton family number conservation. The value is for the sum of the charge states indicated.

VALUE	CL%	DOCUMENT ID		TECN	COMMENT
$<1.2 \times 10^{-5}$	95	ABREU	<b>97</b> C	DLPH	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
$< 1.7 \times 10^{-5}$	95	AKERS	95W	OPAL	$E_{\rm cm}^{\it ee}=$ 88–94 GeV
$< 1.9 \times 10^{-5}$	95	ADRIANI	931	L3	$E_{\rm cm}^{\it ee}=$ 88–94 GeV
$< 1.0 \times 10^{-4}$	95	DECAMP	92	ALEP	$E_{\rm cm}^{ee} = 88-94 \; {\rm GeV}$

 $\Gamma(pe)/\Gamma_{\text{total}}$   $\Gamma_{55}/\Gamma$ 

Test of baryon number and lepton number conservations. Charge conjugate states are implied.

<u>VALUE</u>	CL%	DOCUMENT ID		TECN	COMMENT
<1.8 × 10 <sup>-6</sup>	95	<sup>1</sup> ABBIENDI	991	OPAL	Eee = 88–94 GeV

<sup>&</sup>lt;sup>1</sup> ABBIENDI 991 give the 95%CL limit on the partial width  $\Gamma(Z^0 \to pe)$ < 4.6 KeV and we have transformed it into a branching ratio.

 $\Gamma(p\mu)/\Gamma_{
m total}$   $\Gamma_{
m 56}/\Gamma$ 

Test of baryon number and lepton number conservations. Charge conjugate states are implied.

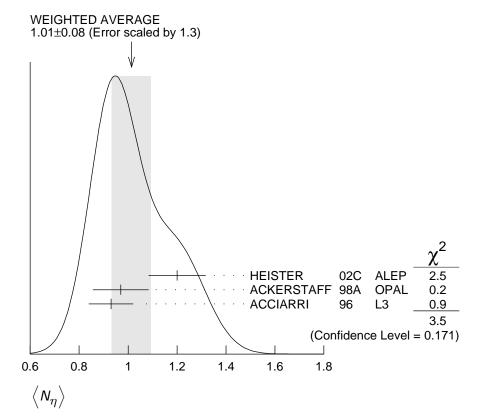
VALUE	CL%	CL%DOCUMENT ID		TECN	COMMENT
$<1.8 \times 10^{-6}$	95	<sup>1</sup> ABBIENDI	991	OPAL	$E_{cm}^{ee} = 88-94 \text{ GeV}$

 $<sup>^1</sup>$  ABBIENDI 991 give the 95%CL limit on the partial width  $\Gamma(Z^0\to p\mu)<$  4.4 KeV and we have transformed it into a branching ratio.

### AVERAGE PARTICLE MULTIPLICITIES IN HADRONIC Z DECAY

Summed over particle and antiparticle, when appropriate.

$\langle N_{\gamma} \rangle$				
VALUE	DOCUMENT ID		TECN	COMMENT
$20.97 \pm 0.02 \pm 1.15$	ACKERSTAFF	98A	OPAL	E <sup>ee</sup> <sub>cm</sub> = 91.2 GeV
///				
$\langle N_{\pi^{\pm}} \rangle$				
<u>VALUE</u> 17.03 ±0.16 OUR AVERAGE	DOCUMENT ID		<u>TECN</u>	COMMENT
17.007±0.209	ABE	046	SI D	$E_{\rm cm}^{\it ee}=91.2~{\rm GeV}$
				****
$17.26 \pm 0.10 \pm 0.88$	ABREU	98L		$E_{\rm cm}^{\rm ee} = 91.2 \text{ GeV}$
$17.04 \pm 0.31$	BARATE	98∨		$E_{cm}^{ee} = 91.2 \; GeV$
$17.05 \pm 0.43$	AKERS	<b>94</b> P	OPAL	$E_{ m cm}^{ m ee}=91.2~{ m GeV}$
$\langle N_{\pi^0}  angle$				
VALUE	DOCUMENT ID		TECN	COMMENT
9.76±0.26 OUR AVERAGE				
$9.55 \pm 0.06 \pm 0.75$	ACKERSTAFF	98A	OPAL	$E_{cm}^{ee} = 91.2 \; GeV$
$9.63 \pm 0.13 \pm 0.63$	BARATE	97J	ALEP	$E_{\rm cm}^{\it ee}=91.2~{\rm GeV}$
$9.90\pm0.02\pm0.33$	ACCIARRI	96	L3	$E_{\rm cm}^{\it ee}=91.2~{\rm GeV}$
$9.2 \pm 0.2 \pm 1.0$	ADAM	96	DLPH	$E_{ m cm}^{ m ee}=$ 91.2 GeV
$\langle N_{\eta}  angle$				
VALUE	DOCUMENT ID			COMMENT
1.01±0.08 OUR AVERAGE Error	includes scale fa	ctor o	of 1.3. So	ee the ideogram below.
$1.20 \pm 0.04 \pm 0.11$	HEISTER	<b>02</b> C	ALEP	$E_{ m cm}^{\it ee}=$ 91.2 GeV
$0.97 \pm 0.03 \pm 0.11$	ACKERSTAFF	98A	OPAL	$E_{\mathrm{cm}}^{\mathrm{ee}} = 91.2 \; \mathrm{GeV}$
$0.93 \pm 0.01 \pm 0.09$	ACCIARRI	96	L3	$E_{ m cm}^{ m ee}=91.2~{ m GeV}$



# $\langle N_{ ho^{\pm}} angle$

2.57±0.15 OUR AVERAGE							
$2.59\!\pm\!0.03\!\pm\!0.16$	$^{ m 1}$ BEDDALL	09	ALEPH archive,	E <sup>ee</sup> <sub>cm</sub> = 91.2 GeV			
$2.40\pm0.06\pm0.43$	ACKERSTAFF	98A OP	AL $E_{\rm cm}^{\rm ee}$ = 91.2 GeV	/			

TECN COMMENT

DOCUMENT ID

# $\langle N_{\rho^0} \rangle$

VALUE	DOCUMENT ID		TECN	COMMENT
1.24±0.10 OUR AVERAGE	Error includes scale fac	ctor o	f 1.1.	
$1.19 \pm 0.10$	ABREU	99J	DLPH	$E_{ m cm}^{\it ee}=$ 91.2 GeV
$1.45\!\pm\!0.06\!\pm\!0.20$	BUSKULIC	96н	ALEP	$E_{cm}^{\mathit{ee}} = 91.2 \; GeV$
$\langle \textit{N}_{\omega}  angle$				
VALUE	<u>DOCUMENT ID</u>		TECN	COMMENT
1.02±0.06 OUR AVERAGE				
$1.00 \pm 0.03 \pm 0.06$	HEISTER	<b>02</b> C	ALEP	$E_{ m cm}^{\it ee}=$ 91.2 GeV
$1.04 \pm 0.04 \pm 0.14$	ACKERSTAFF	98A	OPAL	$E_{ m cm}^{\it ee}=$ 91.2 GeV
$1.17 \pm 0.09 \pm 0.15$	ACCIARRI	<b>97</b> D	L3	$E_{\rm cm}^{\rm ee}=91.2~{\rm GeV}$

 $<sup>^1</sup>$  BEDDALL 09 analyse 3.2 million hadronic Z decays as archived by ALEPH collaboration and report a value of  $2.59\pm0.03\pm0.15\pm0.04$ . The first error is statistical, the second systematic, and the third arises from extrapolation to full phase space. We combine the systematic errors in quadrature.

<	N <sub>n</sub> ,	>
•	7/	,

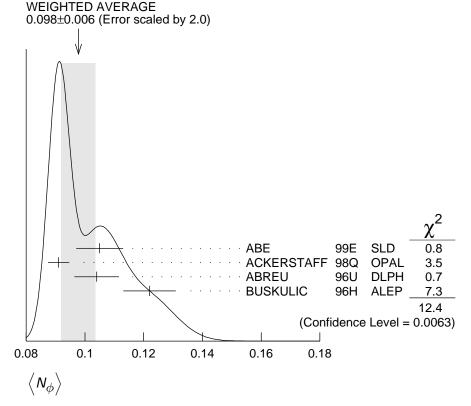
VALUE	DOCUMENT ID	TECN	COMMENT		
$0.17 \pm 0.05$ OUR AVERAGE	Error includes scale fac	tor of 2.4.			
$0.14\ \pm0.01\ \pm0.02$	ACKERSTAFF 98	BA OPAL	$E_{ m cm}^{\it ee}=91.2~{ m GeV}$		
$0.25 \pm 0.04$	<sup>1</sup> ACCIARRI 97	'D <b>L3</b>	$E_{ m cm}^{\it ee}=91.2~{ m GeV}$		
<ul> <li>• • We do not use the following data for averages, fits, limits, etc.</li> <li>• •</li> </ul>					
$0.068\!\pm\!0.018\!\pm\!0.016$	<sup>2</sup> BUSKULIC 92	D ALEP	$E_{cm}^{ee} = 91.2 \; GeV$		
$^1$ ACCIARRI 97D obtain this value averaging over the two decay channels $\eta'  ightarrow \ \pi^+\pi^-\eta$					

and  $\eta' \to \rho^0 \gamma$ . <sup>2</sup> BUSKULIC 92D obtain this value for x>0.1.

# $\langle N_{f_2(080)} \rangle$

/10(aon)/			
VALUE	DOCUMENT ID	TECN	COMMENT
$0.147 \pm 0.011$ OUR AVERAGE			
$0.164 \pm 0.021$	ABREU 99J	DLPH	$E_{cm}^{ee} = 91.2 \; GeV$
$0.141 \pm 0.007 \pm 0.011$	ACKERSTAFF 98Q	OPAL	$E_{ m cm}^{\it ee} = 91.2 \; { m GeV}$
$\langle N_{a_0(980)^{\pm}} \rangle$	DOCUMENT ID	TECN	COMMENT
0.27±0.04±0.10	ACKERSTAFF 98A		
0.21 ± 0.04 ± 0.10	ACKERSTAIT 90A	OLAL	2cm = 91.2 GeV
$\langle \textit{N}_{m{\phi}}  angle$			
VALUE	DOCUMENT ID	TECN	COMMENT
0.008±0.006.OUR AVERAGE	Error includes scale facto	r of 2 0	See the ideogram below

VALUL	DOCUMENT ID		TLCIV	COMMENT
$0.098 \pm 0.006$ OUR AVERAGE	Error includes scale factor		of 2.0.	See the ideogram below.
$0.105 \pm 0.008$	ABE	99E	SLD	$E_{ m cm}^{\it ee}=91.2~{ m GeV}$
$0.091 \pm 0.002 \pm 0.003$	ACKERSTAFF	98Q	OPAL	$E_{ m cm}^{ m ee}=91.2~{ m GeV}$
$0.104 \pm 0.003 \pm 0.007$	ABREU	<b>96</b> U	DLPH	$E_{ m cm}^{\it ee}=91.2~{ m GeV}$
$0.122\pm0.004\pm0.008$	BUSKULIC	96н	ALEP	$E_{cm}^{ee} = 91.2 \; GeV$



# $\langle N_{f_2(1270)} \rangle$

VALUE	DOCUMENT ID	TECN	COMMENT
$0.169\pm0.025$ OUR AVERAGE	Error includes scale fac	tor of 1.4.	
$0.214 \pm 0.038$	ABREU 99	J DLPH	$E_{\rm cm}^{\rm ee}=91.2~{\rm GeV}$
$0.155 \!\pm\! 0.011 \!\pm\! 0.018$	ACKERSTAFF 98	Q OPAL	E <sup>ee</sup> <sub>cm</sub> = 91.2 GeV
/*/			

# $\langle N_{f_1(1285)} \rangle$

VALUEDOCUMENT IDTECNCOMMENT $0.165 \pm 0.051$ 1 ABDALLAH03HDLPH $E_{cm}^{ee} = 91.2 \text{ GeV}$ 

# $\langle N_{f_1(1420)} \rangle$

VALUEDOCUMENT IDTECNCOMMENT0.056  $\pm$  0.0121 ABDALLAH03HDLPH $E_{cm}^{ee} = 91.2 \text{ GeV}$ 

# $\langle N_{f_2'(1525)} \rangle$

VALUEDOCUMENT IDTECNCOMMENT $0.012 \pm 0.006$ ABREU99JDLPH $E_{cm}^{ee} = 91.2 \text{ GeV}$ 

 $<sup>^{1}</sup>$  ABDALLAH 03H assume a  $K\overline{K}\pi$  branching ratio of (9.0  $\pm$  0.4)%.

 $<sup>^1</sup>$  ABDALLAH 03H assume a  $K\overline{K}\pi$  branching ratio of 100%.

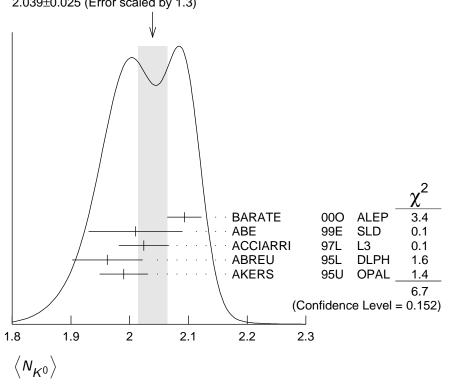
<	$N_{\kappa^{\pm}}$	. >
•		•

VALUE	DOCUMENT ID		TECN	COMMENT
2.24 $\pm$ 0.04 OUR AVERAGE				
$2.203 \pm 0.071$	ABE	<b>0</b> 4C	SLD	$E_{ m cm}^{\it ee}=91.2~{ m GeV}$
$2.21 \pm 0.05 \pm 0.05$	ABREU	98L	DLPH	$E_{ m cm}^{\it ee}=91.2~{ m GeV}$
$2.26 \pm 0.12$	BARATE	98V	ALEP	$E_{ m cm}^{\it ee}=91.2~{ m GeV}$
$2.42 \pm 0.13$	AKERS	<b>94</b> P	OPAL	$E_{cm}^{\mathit{ee}} = 91.2 \; GeV$

# $\langle \mathit{N_{K^0}} \rangle$

VALUE		DOCUMENT ID	DOCUMENT ID		COMMENT	
	2.039 ± 0.025 OUR AVERAGE	Error includes scale	factor	of 1.3.	See the ideogram below.	
	$2.093 \pm 0.004 \pm 0.029$	BARATE	000	ALEP	$E_{ m cm}^{\it ee}=91.2~{ m GeV}$	
	$2.01 \pm 0.08$	ABE	99E	SLD	$E_{cm}^{ee} = 91.2 \; GeV$	
	$2.024 \pm 0.006 \pm 0.042$	ACCIARRI	97L	L3	$E_{cm}^{ee} = 91.2 \; GeV$	
	$1.962 \pm 0.022 \pm 0.056$	ABREU	95L	DLPH	$E_{cm}^{ee} = 91.2 \; GeV$	
	$1.99 \pm 0.01 \pm 0.04$	AKERS	95∪	OPAL	$E_{\rm cm}^{\rm ee}=91.2~{\rm GeV}$	

# WEIGHTED AVERAGE 2.039±0.025 (Error scaled by 1.3)



# $\langle N_{K^*(892)^{\pm}} \rangle$

<u>VALUE</u>	DOCUMENT ID		TECN	COMMENT
$0.72 \pm 0.05$ OUR AVERAGE				
$0.712 \pm 0.031 \pm 0.059$	ABREU	95L	DLPH	$E_{\mathrm{cm}}^{\mathrm{ee}} = 91.2 \; \mathrm{GeV}$
$0.72\ \pm0.02\ \pm0.08$	ACTON	93	OPAL	$E_{\rm cm}^{\it ee}=91.2~{\rm GeV}$

# $\langle N_{K^*(892)^0} \rangle$

<u>VALUE</u>	DOCUMENT ID		TECN	COMMENT
$0.739 \pm 0.022$ OUR AVERAGE				
$0.707 \pm 0.041$	ABE	99E	SLD	$E_{ m cm}^{ m ee}=$ 91.2 GeV
$0.74 \pm 0.02 \pm 0.02$	ACKERSTAFF	<b>97</b> S	OPAL	$E_{ m cm}^{\it ee}=$ 91.2 GeV
$0.77 \pm 0.02 \pm 0.07$	ABREU	<b>96</b> U	DLPH	$E_{ m cm}^{\it ee}=$ 91.2 GeV
$0.83 \pm 0.01 \pm 0.09$	BUSKULIC	96н	ALEP	$E_{ m cm}^{\it ee}=$ 91.2 GeV
$0.97\ \pm0.18\ \pm0.31$	ABREU	93	DLPH	$E_{ m cm}^{\it ee}=$ 91.2 GeV

# $\left< N_{K_2^*(1430)} \right>$

<u>VALUE</u>	DOCUMENT ID		TECN	<u>COMMENT</u>
$0.073 \pm 0.023$	ABREU	99J	DLPH	$E_{\rm cm}^{\rm ee} = 91.2 \; {\rm GeV}$

ullet ullet We do not use the following data for averages, fits, limits, etc. ullet ullet

 $0.19\ \pm0.04\ \pm0.06$ 

<sup>1</sup> AKERS

95x OPAL  $E_{cm}^{ee}$  = 91.2 GeV

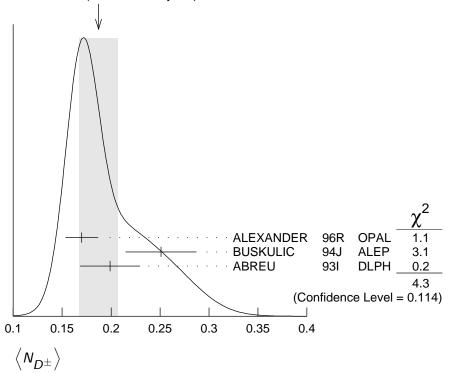
Created: 7/12/2013 14:51

# $\left< N_{D^{\pm}} \right>$

VALUE	DOCUMENT ID		TECN	COMMENT
$0.187 \pm 0.020$ OUR AVERAGE	Error includes scale	factor	of 1.5.	See the ideogram below.
$0.170\pm0.009\pm0.014$	ALEXANDER	<b>96</b> R	OPAL	$E_{ m cm}^{\it ee}=91.2~{ m GeV}$
$0.251 \pm 0.026 \pm 0.025$	BUSKULIC	94J	ALEP	$E_{\rm cm}^{\it ee}=91.2~{\rm GeV}$
$0.199 \pm 0.019 \pm 0.024$	$^{ m 1}$ ABREU	931	DLPH	$E_{cm}^{\mathit{ee}} = 91.2 \; GeV$

<sup>&</sup>lt;sup>1</sup> See ABREU 95 (erratum).

WEIGHTED AVERAGE 0.187±0.020 (Error scaled by 1.5)



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 $<sup>^{1}</sup>$  AKERS 95X obtain this value for x< 0.3.

$\langle N_{D^0} \rangle$					
VALUE	DOCUMENT ID		<u>TECN</u>	COMMENT	
0.462 ± 0.026 OUR AVERAGE	ALEXANDER	060	ODAL	E66 _ 01 2 CoV	
$0.465 \pm 0.017 \pm 0.027$ $0.518 \pm 0.052 \pm 0.035$	BUSKULIC			$E_{\text{cm}}^{\text{ee}} = 91.2 \text{ GeV}$ $E_{\text{cm}}^{\text{ee}} = 91.2 \text{ GeV}$	
$0.403 \pm 0.032 \pm 0.044$	<sup>1</sup> ABREU	94)		$E_{\rm cm}^{ee} = 91.2 \text{ GeV}$	
	ABILLO	931	DLITI	Lcm— 91.2 GeV	
<sup>1</sup> See ABREU 95 (erratum).					
$\langle N_{D_s^{\pm}} \rangle$					
VALUE	DOCUMENT ID		TECN	COMMENT	
0.131±0.010±0.018				E <sub>cm</sub> = 91.2 GeV	
$\langle N_{D^*(2010)^{\pm}} \rangle$					
<u>VALUE</u> 0.183 ±0.008 OUR AVERAGE	DOCUMENT ID		<u>TECN</u>	COMMENT	
$0.1854 \pm 0.0041 \pm 0.0091$	<sup>1</sup> ACKERSTAFF	98E	OPAL	E <sup>ee</sup> <sub>cm</sub> = 91.2 GeV	
$0.187 \pm 0.015 \pm 0.013$				$E_{\rm cm}^{ee} = 91.2 \text{ GeV}$	
$0.171\ \pm0.012\ \pm0.016$				$E_{\rm cm}^{ee} = 91.2 \text{ GeV}$	
<sup>1</sup> ACKERSTAFF 98E systematic error includes an uncertainty of $\pm 0.0069$ due to the branching ratios B( $D^{*+} \rightarrow D^0 \pi^+$ ) = $0.683 \pm 0.014$ and B( $D^0 \rightarrow K^- \pi^+$ ) = $0.0383 \pm 0.0012$ . <sup>2</sup> See ABREU 95 (erratum).					
$VALUE$ (units $10^{-3}$ )	DOCUMENT ID		TECN	COMMENT	
• • • We do not use the following					
$2.9^{+0.7}_{-0.6}\pm0.2$				E <sup>ee</sup> <sub>cm</sub> = 91.2 GeV	
<sup>1</sup> ACKERSTAFF 97W obtain this width is saturated by the <i>D* h</i>	s value for $x > 0.6$ final states.	and w	vith the	assumption that its decay	
$\langle N_{B^*} \rangle$					
VALUE				COMMENT	
$0.28 \pm 0.01 \pm 0.03$	<sup>1</sup> ABREU	<b>95</b> R	DLPH	$E_{cm}^{\mathit{ee}} = 91.2 \; GeV$	
$^{ m 1}$ ABREU 95R quote this value f	or a flavor-average	ed exc	ited stat	e.	
$\langle N_{J/\psi(1S)}  angle$					
VALUE	1 ALEXANDED	0.6-	<u>TECN</u>	<u>COMMENT</u>	
$0.0056 \pm 0.0003 \pm 0.0004$				E <sup>ee</sup> <sub>cm</sub> = 91.2 GeV	
$^{1}$ ALEXANDER 96B identify $J/v$	$\psi(1S)$ from the de	cays i	nto lepto	on pairs.	
/A/ \					
$\langle N_{\psi(2S)} \rangle$	DOCUMENT ID		TECN	COMMENT	
$(N_{\psi(2S)})$ <u>VALUE</u> 0.0023 $\pm$ 0.0004 $\pm$ 0.0003				COMMENT  Eee = 91.2 GeV	

/	M	١
1	'¶p	/

VALUE	DOCUMENT ID		TECN	COMMENT
1.046±0.026 OUR AVERAGE				
$1.054 \pm 0.035$	ABE	04C	SLD	$E_{ m cm}^{ m ee}=$ 91.2 GeV
$1.08 \pm 0.04 \pm 0.03$	ABREU	98L	DLPH	$E_{ m cm}^{ m ee}=$ 91.2 GeV
$1.00 \pm 0.07$	BARATE	98V	ALEP	$E_{ m cm}^{ m ee}=$ 91.2 GeV
$0.92\ \pm0.11$	AKERS	94P	OPAL	$E_{ m cm}^{ m ee}=$ 91.2 GeV

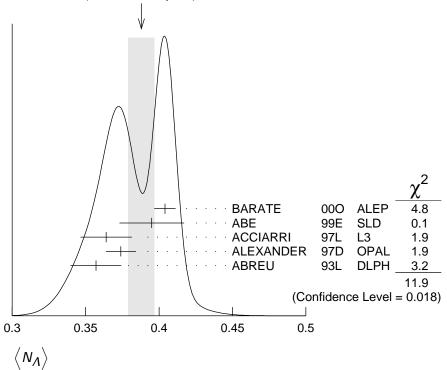
# $\langle N_{\Delta(1232)^{++}} \rangle$

VALUE	<u>DOCUMENT ID</u>	<u> IECN</u>	COMMENT
$0.087 \pm 0.033$ OUR AVERAGE	Error includes scale	factor of 2.4.	
$0.079 \pm 0.009 \pm 0.011$	ABREU	95W DLPH	$E_{ m cm}^{\it ee}=91.2~{ m GeV}$
$0.22 \pm 0.04 \pm 0.04$	ALEXANDER	95D OPAL	$E_{\rm cm}^{\it ee}=91.2~{\rm GeV}$

# $\langle N_A \rangle$

VALUE	DOCUMENT ID		TECN	COMMENT
$0.388 \pm 0.009$ OUR AVERAGE	Error includes scale	factor	of 1.7.	See the ideogram below.
$0.404\pm0.002\pm0.007$	BARATE	000	ALEP	$E_{ m cm}^{ m ee}=91.2~{ m GeV}$
$0.395 \pm 0.022$	ABE	99E	SLD	$E_{ m cm}^{\it ee}=91.2~{ m GeV}$
$0.364 \pm 0.004 \pm 0.017$	ACCIARRI	97L	L3	$E_{cm}^{ee} = 91.2 \; GeV$
$0.374 \pm 0.002 \pm 0.010$	ALEXANDER	<b>97</b> D	OPAL	$E_{ m cm}^{ m ee}=91.2~{ m GeV}$
$0.357 \pm 0.003 \pm 0.017$	ABREU	93L	DLPH	$E_{\rm cm}^{\rm ee} = 91.2 \; {\rm GeV}$

# WEIGHTED AVERAGE 0.388±0.009 (Error scaled by 1.7)



$\langle N_{\Lambda(1520)} \rangle$				
VALUE	DOCUMENT ID		TECN	COMMENT
$0.0224\pm0.0027$ OUR AVERAGE				
$0.029 \pm 0.005 \pm 0.005$	ABREU	00P	DLPH	$E_{ m cm}^{\it ee}=91.2~{ m GeV}$
$0.0213 \pm 0.0021 \pm 0.0019$	ALEXANDER	<b>97</b> D	OPAL	E <sup>ee</sup> <sub>cm</sub> = 91.2 GeV
$\langle N_{\Sigma^+}  angle$				
VALUE	DOCUMENT ID		TECN	COMMENT
0.107±0.010 OUR AVERAGE				
$0.114 \pm 0.011 \pm 0.009$	ACCIARRI	001	L3	$E_{\mathrm{cm}}^{\mathrm{ee}} = 91.2 \; \mathrm{GeV}$
$0.099 \pm 0.008 \pm 0.013$	ALEXANDER	97E	OPAL	$E_{ m cm}^{\it ee}=$ 91.2 GeV
$\langle N_{oldsymbol{\Sigma}^-}  angle$				
VALUE	DOCUMENT ID		TECN	COMMENT
0.082±0.007 OUR AVERAGE				
$0.081 \pm 0.002 \pm 0.010$	ABREU	<b>00</b> P	DLPH	$E_{cm}^{\mathit{ee}} = 91.2 \; GeV$
$0.083 \pm 0.006 \pm 0.009$	ALEXANDER	97E	OPAL	$E_{cm}^{ee} = 91.2 \; GeV$
$\langle N_{\Sigma^+ + \Sigma^-}  angle$				
VALUE	DOCUMENT ID		TECN	COMMENT
$0.181 \pm 0.018$ OUR AVERAGE	1			
$0.182 \pm 0.010 \pm 0.016$	<sup>1</sup> ALEXANDER			$E_{ m cm}^{\it ee}=$ 91.2 GeV
$0.170\pm0.014\pm0.061$	ABREU	950	DLPH	$E_{\rm cm}^{\it ee}=91.2~{\rm GeV}$

 $<sup>^1\</sup>text{We}$  have combined the values of  $\langle \textit{N}_{\sum^+}\rangle$  and  $\langle \textit{N}_{\sum^-}\rangle$  from ALEXANDER 97E adding the statistical and systematic errors of the two final states separately in quadrature. If isospin symmetry is assumed this value becomes 0.174  $\pm$  0.010  $\pm$  0.015.

# $\langle N_{\Sigma^0} \rangle$

\ Z \ /	DOCUMENT ID		TECN	COMMENT
VALUE	DOCUMENT ID		<u>TECN</u>	COMMENT
$0.076\pm0.010$ OUR AVERAGE				
$0.095 \pm 0.015 \pm 0.013$	ACCIARRI	001	L3	$E_{\rm cm}^{\it ee}=91.2~{\rm GeV}$
$0.071 \pm 0.012 \pm 0.013$	ALEXANDER	97E	OPAL	$E_{ m cm}^{\it ee}=91.2~{ m GeV}$
$0.070 \pm 0.010 \pm 0.010$	ADAM	<b>96</b> B	DLPH	$E_{cm}^{\mathit{ee}} = 91.2 \; GeV$
$\langle \mathit{N}_{(\Sigma^{+}+\Sigma^{-}+\Sigma^{0})/3}  angle$				
VALUE	DOCUMENT ID		TECN	COMMENT
$0.084 \pm 0.005 \pm 0.008$	ALEXANDER	97E	OPAL	E <sup>ee</sup> <sub>cm</sub> = 91.2 GeV
$\langle N_{\Sigma(1385)^+}  angle$				
VALUE	DOCUMENT ID		TECN	COMMENT
$0.0239 \pm 0.0009 \pm 0.0012$	ALEXANDER	<b>97</b> D	OPAL	$E_{\mathrm{cm}}^{\mathit{ee}} = 91.2 \; \mathrm{GeV}$
$\langle N_{oldsymbol{\Sigma}(1385)^-}  angle$				
VALUE	DOCUMENT ID		TECN	COMMENT
$0.0240 \pm 0.0010 \pm 0.0014$	ALEXANDER	<b>97</b> D	OPAL	$E_{cm}^{ee} = 91.2 \; GeV$

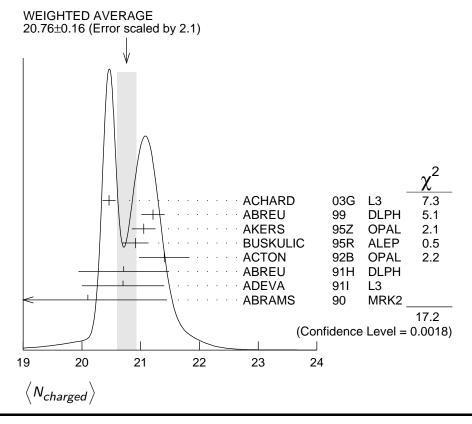
$\langle N_{\Sigma(1385)^++\Sigma(1385)^-}  angle$				
VALUE	DOCUMENT ID			
	Error includes sca			
$0.0479 \pm 0.0013 \pm 0.0026$				$E_{\rm cm}^{ee} = 91.2 \text{ GeV}$
$0.0382 \pm 0.0028 \pm 0.0045$	ABREU	950	DLPH	E <sup>ee</sup> <sub>cm</sub> = 91.2 GeV
⟨ <b>N</b> <sub>=</sub> -⟩  VALUE	DOCUMENT ID		TECN	COMMENT
0.0258±0.0009 OUR AVERAGE	DOCOMENT ID		TECIV	COMMILINI
$0.0247 \pm 0.0009 \pm 0.0025$	ABDALLAH	06E	DLPH	$E_{\rm cm}^{\rm ee}=91.2~{\rm GeV}$
$0.0259 \pm 0.0004 \pm 0.0009$				$E_{\rm cm}^{\rm ee} = 91.2 \; {\rm GeV}$
				CIII
$\langle N_{\equiv (1530)^0} \rangle$				
VALUE	DOCUMENT ID		TECN	COMMENT
0.0059±0.0011 OUR AVERAGE	Frror includes sca			
$0.0045 \pm 0.0005 \pm 0.0006$	ABDALLAH	<b>05</b> C	DLPH	$E_{cm}^{ee} = 91.2 \; GeV$
$0.0068\!\pm\!0.0005\!\pm\!0.0004$	ALEXANDER	<b>97</b> D	OPAL	$E_{cm}^{ee} = 91.2 \; GeV$
1				
$\langle N_{\mathcal{Q}^-}  angle$				
VALUE	DOCUMENT ID		TECN	COMMENT
0.00164±0.00028 OUR AVERAGE	AL EVANDED	07-	0041	T-P- 01 0 C V
$0.0018 \pm 0.0003 \pm 0.0002$				$E_{\rm cm}^{\rm ee} = 91.2 \; {\rm GeV}$
$0.0014 \pm 0.0002 \pm 0.0004$	ADAM	<b>96</b> B	DLPH	E <sup>ee</sup> <sub>cm</sub> = 91.2 GeV
$\langle N_{A_c^+} \rangle$				
VALUE	DOCUMENT ID		TECN	COMMENT
$0.078 \pm 0.012 \pm 0.012$	ALEXANDER	<b>96</b> R	OPAL	$E_{ m cm}^{ m ee}=$ 91.2 GeV
$\langle N_{\overline{D}}  angle$				
$VALUE$ (units $10^{-6}$ )	DOCUMENT ID		TECN	COMMENT
• • • We do not use the following of				
				$E_{ m cm}^{ m ee}=91.2~{ m GeV}$
1	JOHALL	JUA	/ \LLI	-cm - 31.2 GeV

 $<sup>^{1}\,\</sup>mathrm{SCHAEL}$  06A obtain this anti-deuteron production rate per hadronic Z decay in the anti-deuteron momentum range from 0.62 to 1.03 GeV/c.

# $\langle N_{charged} \rangle$

VALUE	DOCUMENT ID	TECN	COMMENT
20.76±0.16 OUR AVERAGE	Error includes scale factor	ot 2.1.	See the ideogram below.
$20.46 \pm 0.01 \pm 0.11$	ACHARD 03G	L3	$E_{cm}^{ee} = 91.2 \; GeV$
$21.21\!\pm\!0.01\!\pm\!0.20$	ABREU 99	DLPH	$E_{cm}^{ee} = 91.2 \; GeV$
$21.05 \!\pm\! 0.20$	AKERS 95z	OPAL	$E_{cm}^{ee} = 91.2 \; GeV$
$20.91\!\pm\!0.03\!\pm\!0.22$	BUSKULIC 95R	ALEP	$E_{cm}^{ee} = 91.2 \; GeV$
$21.40 \pm 0.43$	ACTON 92B	OPAL	$E_{cm}^{ee} = 91.2 \; GeV$
$20.71 \pm 0.04 \pm 0.77$	ABREU 91H	DLPH	$E_{cm}^{ee} = 91.2 \; GeV$
$20.7 \pm 0.7$	ADEVA 91ı	L3	$E_{ m cm}^{ m ee}=91.2~{ m GeV}$
$20.1 \pm 1.0 \pm 0.9$	ABRAMS 90	MRK2	$E_{cm}^{ee} = 91.1 \; GeV$

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#### Z HADRONIC POLE CROSS SECTION

OUR FIT is obtained using the fit procedure and correlations as determined by the LEP Electroweak Working Group (see the note "The Z boson" and ref. LEP-SLC 06). This quantity is defined as

$$\sigma_h^0 = \frac{12\pi}{M_Z^2} \, \frac{\Gamma(e^+ e^-) \, \Gamma(\text{hadrons})}{\Gamma_Z^2}$$

It is one of the parameters used in the Z lineshape fit.

VALUE (nb)	EVTS	DOCUMENT ID		TECN	COMMENT
41.541±0.037 OUR F	IT.				
$41.501 \pm 0.055$	4.10M	$^{ m 1}$ ABBIENDI	<b>01</b> A	OPAL	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
$41.578 \pm 0.069$	3.70M	ABREU	00F	DLPH	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
$41.535 \pm 0.055$	3.54M	ACCIARRI	<b>00</b> C	L3	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
$41.559 \pm 0.058$	4.07M	<sup>2</sup> BARATE	<b>00</b> C	ALEP	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
ullet $ullet$ We do not use	the followin	g data for averages	s, fits,	limits, e	etc. • • •
42 ±4	450	ABRAMS	<b>89</b> B	MRK2	$E_{\rm cm}^{ee} = 89.2 - 93.0  {\rm GeV}$

<sup>&</sup>lt;sup>1</sup>ABBIENDI 01A error includes approximately 0.031 due to statistics, 0.033 due to event selection systematics, 0.029 due to uncertainty in luminosity measurement, and 0.011 due to LEP energy uncertainty.

<sup>&</sup>lt;sup>2</sup> BARATE 00C error includes approximately 0.030 due to statistics, 0.026 due to experimental systematics, and 0.025 due to uncertainty in luminosity measurement.

#### Z VECTOR COUPLINGS

These quantities are the effective vector couplings of the Z to charged leptons. Their magnitude is derived from a measurement of the Z lineshape and the forward-backward lepton asymmetries as a function of energy around the Z mass. The relative sign among the vector to axial-vector couplings is obtained from a measurement of the Z asymmetry parameters,  $A_e$ ,  $A_\mu$ , and  $A_\tau$ . By convention the sign of  $g_A^e$  is fixed to be negative (and opposite to that of  $g^{\nu e}$  obtained using  $\nu_e$  scattering measurements).

For the light quarks, the sign of the couplings is assigned consistently with this assumption. The fit values quoted below correspond to global nine- or five-parameter fits to lineshape, lepton forward-backward asymmetry, and  $A_e$ ,  $A_\mu$ , and  $A_\tau$  measurements. See the note "The Z boson" and ref. LEP-SLC 06 for details. Where  $p\overline{p}$  and ep data is quoted, OUR FIT value corresponds to a weighted average of this with the LEP/SLD fit result.

g	e V
•	v

VALUE	<b>EVTS</b>	DOCUMENT ID		TECN	COMMENT
-0.03817±0.00047 OUR FI	T				
$-0.058$ $\pm 0.016$ $\pm 0.007$	5026	$^{ m 1}$ ACOSTA	05м	CDF	$E_{cm}^{p\overline{p}} = 1.96 \; TeV$
$-0.0346 \pm 0.0023$	137.0K	<sup>2</sup> ABBIENDI	010	OPAL	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
$-0.0412 \pm 0.0027$	124.4k	<sup>3</sup> ACCIARRI	<b>00</b> C	L3	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
$-0.0400 \pm 0.0037$		BARATE	<b>00</b> C	ALEP	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
$-0.0414\ \pm0.0020$		<sup>4</sup> ABE	<b>95</b> J	SLD	$E_{cm}^{\mathit{ee}} = 91.31 \; GeV$

<sup>&</sup>lt;sup>1</sup> ACOSTA 05M determine the forward–backward asymmetry of  $e^+e^-$  pairs produced via  $q \, \overline{q} \to Z/\gamma^* \to e^+e^-$  in 15 M( $e^+e^-$ ) effective mass bins ranging from 40 GeV to 600 GeV. These results are used to obtain the vector and axial–vector couplings of the Z to  $e^+e^-$ , assuming the quark couplings are as predicted by the standard model. Higher order radiative corrections have not been taken into account.

order radiative corrections have not been taken into account.  $^2$  ABBIENDI 010 use their measurement of the  $\tau$  polarization in addition to the lineshape and forward-backward lepton asymmetries.

# $g_V^\mu$

<u>VALUE</u>	<b>EVTS</b>	DOCUMENT ID		TECN	COMMENT
$-0.0367\pm0.0023$ OUF	R FIT				
$-0.0388 {}^{+ 0.0060}_{- 0.0064}$	182.8K	<sup>1</sup> ABBIENDI	010	OPAL	<i>E</i> <sup>ee</sup> <sub>cm</sub> = 88−94 GeV
$-0.0386 \pm 0.0073$	113.4k	<sup>2</sup> ACCIARRI	<b>00</b> C	L3	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
$-0.0362\!\pm\!0.0061$		BARATE	<b>00</b> C	ALEP	<i>E</i> <sup>ee</sup> <sub>cm</sub> = 88−94 GeV
$\bullet$ $\bullet$ We do not use t	he following	g data for averages	s, fits,	limits, e	etc. • • •
$-0.0413\!\pm\!0.0060$	66143	<sup>3</sup> ABBIENDI	<b>01</b> K	OPAL	<i>E</i> <sup>ee</sup> <sub>cm</sub> = 89−93 GeV

 $<sup>^1</sup>$ ABBIENDI 010 use their measurement of the au polarization in addition to the lineshape and forward-backward lepton asymmetries.

 $<sup>^3</sup>$  ACCIARRI 00C use their measurement of the au polarization in addition to forward-backward lepton asymmetries.

 $<sup>^4</sup>$  ABE 95J obtain this result combining polarized Bhabha results with the  $A_{LR}$  measurement of ABE 94C. The Bhabha results alone give  $-0.0507\pm0.0096\pm0.0020$ .

<sup>&</sup>lt;sup>2</sup> ACCIARRI 00C use their measurement of the  $\tau$  polarization in addition to forward-backward lepton asymmetries.

<sup>&</sup>lt;sup>3</sup> ABBIENDI 01K obtain this from an angular analysis of the muon pair asymmetry which takes into account effects of initial state radiation on an event by event basis and of initial-final state interference.

$g_V^{\tau}$
5 V

VALUE	<b>EVTS</b>	DOCUMENT ID		TECN	COMMENT
$-0.0366 \pm 0.0010$ OUR	FIT				
$-0.0365\!\pm\!0.0023$	151.5K	$^{ m 1}$ ABBIENDI	010	OPAL	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
$-0.0384\!\pm\!0.0026$	103.0k	<sup>2</sup> ACCIARRI	<b>00</b> C	L3	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
$-0.0361\!\pm\!0.0068$		BARATE	<b>00</b> C	ALEP	$E_{cm}^{ee} = 88-94 \text{ GeV}$

<sup>&</sup>lt;sup>1</sup> ABBIENDI 010 use their measurement of the  $\tau$  polarization in addition to the lineshape and forward-backward lepton asymmetries.

# $g_V^\ell$

VALUE	EVTS	DOCUMENT ID		TECN	COMMENT
-0.03783±0.00041 OL					
$-0.0358 \pm 0.0014$	471.3K	<sup>1</sup> ABBIENDI	010	OPAL	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
$-0.0397\ \pm0.0020$	379.4k	<sup>2</sup> ABREU	00F	DLPH	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
$-0.0397\ \pm0.0017$	340.8k	<sup>3</sup> ACCIARRI	<b>00</b> C	L3	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
$-0.0383 \pm 0.0018$	500k	BARATE	00C	ALEP	$E_{\rm cm}^{\rm ee} = 88-94 \; {\rm GeV}$

 $<sup>^1</sup>$  ABBIENDI 010 use their measurement of the au polarization in addition to the lineshape and forward-backward lepton asymmetries.

# $g_V^u$

VALUE	<u>EVTS</u>	DOCUMENT ID		TECN	COMMENT
$0.25 \begin{array}{c} +0.07 \\ -0.06 \end{array}$ OUR AV	ERAGE				
$0.201 \pm 0.112$	156k	$^{ m 1}$ ABAZOV	<b>11</b> D	D0	$E_{cm}^{oldsymbol{\overline{p}}}=1.97\;TeV$
0.27 ±0.13	1500	<sup>2</sup> AKTAS	06	H1	$e^{\pm} p  ightarrow \; \overline{ u}_{m{e}}( u_{m{e}}) X, \ \sqrt{s} pprox 300 \; {\sf GeV}$
$0.24 \begin{array}{l} +0.28 \\ -0.11 \end{array}$		<sup>3</sup> LEP-SLC	06		$E_{cm}^{ee} = 88-94 \; GeV$
$0.399^{+0.152}_{-0.188} \pm 0.066$	5026	<sup>4</sup> ACOSTA	05м	CDF	$E_{cm}^{ar{p}} = 1.96 \; TeV$

 $<sup>^1</sup>$  ABAZOV 11D study  $p\overline{p}\to Z/\gamma^*\,e^+\,e^-$  events using 5 fb $^{-1}$  data at  $\sqrt{s}=1.96$  TeV. The candidate events are selected by requiring two isolated electromagnetic showers with  $E_T>25$  GeV, at least one electron in the central region and the di-electron mass in the range 50–1000 GeV. From the forward-backward asymmetry, determined as a function of the di-electron mass, they derive the axial and vector couplings of the u- and d- quarks and the value of  $\sin^2\!\theta_{eff}^\ell=0.2309\pm0.0008(\text{stat})\pm0.0006(\text{syst}).$ 

<sup>3</sup> LEP-SLC 06 is a combination of the results from LEP and SLC experiments using light quark tagging. s- and d-quark couplings are assumed to be identical.

 $<sup>^2</sup>$  ACCIARRI 00C use their measurement of the au polarization in addition to forward-backward lepton asymmetries.

<sup>&</sup>lt;sup>2</sup> Using forward-backward lepton asymmetries.

 $<sup>^3</sup>$ ACCIARRI 00C use their measurement of the au polarization in addition to forward-backward lepton asymmetries.

 $<sup>^2</sup>$  AKTAS 06 fit the neutral current (1.5  $\leq$  Q $^2$   $\leq$  30,000 GeV $^2$ ) and charged current (1.5  $\leq$  Q $^2$   $\leq$  15,000 GeV $^2$ ) differential cross sections. In the determination of the *u*-quark couplings the electron and *d*-quark couplings are fixed to their standard model values.

<sup>&</sup>lt;sup>4</sup> ACOSTA 05M determine the forward-backward asymmetry of  $e^+e^-$  pairs produced via  $q \overline{q} \rightarrow Z/\gamma^* \rightarrow e^+e^-$  in 15 M( $e^+e^-$ ) effective mass bins ranging from 40 GeV to 600 GeV. These results are used to obtain the vector and axial-vector couplings of the Z to the light quarks, assuming the electron couplings are as predicted by the Standard Model. Higher order radiative corrections have not been taken into account.

gv					
VALUE	<b>EVTS</b>	DOCUMENT ID		TECN	COMMENT
$-0.33 \begin{array}{c} +0.05 \\ -0.06 \end{array}$ OUR AN	/ERAGE				
$-0.351\!\pm\!0.251$	156k	<sup>1</sup> ABAZOV	<b>11</b> D	D0	$E_{cm}^{ar{p}} = 1.97 \; TeV$
$-0.33 \pm 0.33$	1500	<sup>2</sup> AKTAS	06	H1	$e^{\pm} p \rightarrow \overline{\nu}_e(\nu_e) X$ ,
					$\sqrt{s}pprox$ 300 GeV
$-0.33 \begin{array}{l} +0.05 \\ -0.07 \end{array}$		<sup>3</sup> LEP-SLC	06		$E_{cm}^{\mathit{ee}} = 88 – 94 \; GeV$
$-0.226^{+0.635}_{-0.290}\pm0.090$	5026	<sup>4</sup> ACOSTA	05м	CDF	$E_{cm}^{ar{p}} = 1.96 \; TeV$

- $^1$  ABAZOV 11D study  $p\overline{p}\to Z/\gamma^*\,e^+\,e^-$  events using 5 fb $^{-1}$  data at  $\sqrt{s}=1.96$  TeV. The candidate events are selected by requiring two isolated electromagnetic showers with  $E_T>25$  GeV, at least one electron in the central region and the di-electron mass in the range 50–1000 GeV. From the forward-backward asymmetry, determined as a function of the di-electron mass, they derive the axial and vector couplings of the u- and d- quarks and the value of  $\sin^2\!\theta_{eff}^\ell=0.2309\pm0.0008(\text{stat})\pm0.0006(\text{syst}).$
- $^2$  AKTAS 06 fit the neutral current (1.5  $\leq$  Q $^2$   $\leq$  30,000 GeV $^2$ ) and charged current (1.5  $\leq$  Q $^2$   $\leq$  15,000 GeV $^2$ ) differential cross sections. In the determination of the d-quark couplings the electron and u-quark couplings are fixed to their standard model values.
- <sup>3</sup> LEP-SLC 06 is a combination of the results from LEP and SLC experiments using light quark tagging. s- and d-quark couplings are assumed to be identical.
- <sup>4</sup> ACOSTA 05M determine the forward-backward asymmetry of  $e^+e^-$  pairs produced via  $q \overline{q} \rightarrow Z/\gamma^* \rightarrow e^+e^-$  in 15 M( $e^+e^-$ ) effective mass bins ranging from 40 GeV to 600 GeV. These results are used to obtain the vector and axial-vector couplings of the Z to the light quarks, assuming the electron couplings are as predicted by the Standard Model. Higher order radiative corrections have not been taken into account.

#### Z AXIAL-VECTOR COUPLINGS

These quantities are the effective axial-vector couplings of the Z to charged leptons. Their magnitude is derived from a measurement of the Z line-shape and the forward-backward lepton asymmetries as a function of energy around the Z mass. The relative sign among the vector to axial-vector couplings is obtained from a measurement of the Z asymmetry parameters,  $A_e$ ,  $A_\mu$ , and  $A_\tau$ . By convention the sign of  $g_A^e$  is fixed to be negative (and opposite to that of  $g^{\nu_e}$  obtained using  $\nu_e$  scattering measurements). For the light quarks, the sign of the couplings is assigned consistently with this assumption. The fit values quoted below correspond to global nine- or five-parameter fits to lineshape, lepton forward-backward asymmetry, and  $A_e$ ,  $A_\mu$ , and  $A_\tau$  measurements. See the note "The Z boson" and ref. LEP-SLC 06 for details. Where  $p\overline{p}$  and ep data is quoted, OUR FIT value corresponds to a weighted average of this with the LEP/SLD fit result.

$g_A^e$				
VALUE	<b>EVTS</b>	DOCUMENT ID	TECN	COMMENT
$-0.50111\pm0.00035$ OUR FI	Т			
$-0.528$ $\pm 0.123$ $\pm 0.059$	5026	$^{ m 1}$ ACOSTA	05м CDF	$E_{cm}^{ar{p}} = 1.96 \; TeV$
$-0.50062\!\pm\!0.00062$	137.0K	<sup>2</sup> ABBIENDI	010 OPAL	$E_{\rm cm}^{ee} = 88-94 \; {\rm GeV}$
$-0.5015 \pm 0.0007$	124.4k	<sup>3</sup> ACCIARRI	00C L3	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
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$-0.50166\pm0.00057$	BARATE	00c ALEP	$E_{\rm cm}^{ee}$ = 88–94 GeV
$-0.4977 \pm 0.0045$	<sup>4</sup> ABE	95」 SLD	$E_{cm}^{ee} = 91.31 \text{ GeV}$

 $<sup>^{1}</sup>$  ACOSTA 05M determine the forward–backward asymmetry of  $e^{+}e^{-}$  pairs produced via  $q \overline{q} \rightarrow Z/\gamma^* \rightarrow e^+ e^-$  in 15 M( $e^+ e^-$ ) effective mass bins ranging from 40 GeV to 600 GeV. These results are used to obtain the vector and axial-vector couplings of the Z to  $e^+e^-$ , assuming the quark couplings are as predicted by the standard model. Higher

order radiative corrections have not been taken into account. <sup>2</sup>ABBIENDI 010 use their measurement of the  $\tau$  polarization in addition to the lineshape

and forward-backward lepton asymmetries.

 $^3$  ACCIARRI 00C use their measurement of the au polarization in addition to forwardbackward lepton asymmetries.

 $^4$  ABE 95J obtain this result combining polarized Bhabha results with the  $A_{LR}$  measurement of ABE 94C. The Bhabha results alone give  $-0.4968 \pm 0.0039 \pm 0.0027$ .

# $g^{\mu}_{A}$

EVTS	DOCUMENT ID		TECN	COMMENT			
-0.50120±0.00054 OUR FIT							
099 182.8K	$^{ m 1}$ ABBIENDI	010	OPAL	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV			
14 113.4k	<sup>2</sup> ACCIARRI	<b>00</b> C	L3	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV			
093	BARATE	<b>00</b> C	ALEP	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV			
<ul> <li>• • We do not use the following data for averages, fits, limits, etc.</li> </ul>							
66143	<sup>3</sup> ABBIENDI	01K	OPAL	E <sup>ee</sup> <sub>cm</sub> = 89–93 GeV			
	054 OUR FIT 099 182.8K 14 113.4k 093 : use the following	054 OUR FIT 099 182.8K <sup>1</sup> ABBIENDI 14 113.4k <sup>2</sup> ACCIARRI 093 BARATE : use the following data for averag	054 OUR FIT 099 182.8K	054 OUR FIT 099 182.8K <sup>1</sup> ABBIENDI 010 OPAL 14 113.4k <sup>2</sup> ACCIARRI 00C L3 093 BARATE 00C ALEP 1 use the following data for averages, fits, limits, e			

 $<sup>^{1}</sup>$  ABBIENDI 010 use their measurement of the au polarization in addition to the lineshape and forward-backward lepton asymmetries.

 $^2$  ACCIARRI 00C use their measurement of the au polarization in addition to forwardbackward lepton asymmetries.

 $^3$  ABBIENDI  $^{01}$ K obtain this from an angular analysis of the muon pair asymmetry which takes into account effects of initial state radiation on an event by event basis and of initial-final state interference.

### $g_A^{ au}$

**							
VALUE	<u>EVTS</u>	DOCUMENT ID		TECN	COMMENT		
-0.50204±0.00064 OUR FIT							
$-0.50165 \pm 0.00124$	151.5K	<sup>1</sup> ABBIENDI	010	OPAL	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV		
$-0.5023\ \pm0.0017$	103.0k	<sup>2</sup> ACCIARRI	<b>00</b> C	L3	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV		
$-0.50216 \pm 0.00100$		BARATE	<b>00</b> C	ALEP	$E_{\rm cm}^{\it ee} = 88 - 94 \; {\rm GeV}$		

 $<sup>^{1}</sup>$  ABBIENDI 010 use their measurement of the au polarization in addition to the lineshape

# $g_A^\ell$

<u>VALUE</u>	<u>EVTS</u>	DOCUMENT ID		TECN	COMMENT		
-0.50123±0.00026 OUR FIT							
$-0.50089\!\pm\!0.00045$	471.3K	<sup>1</sup> ABBIENDI	010	OPAL	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV		
$-0.5007 \pm 0.0005$	379.4k	ABREU	00F	DLPH	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV		
$-0.50153\!\pm\!0.00053$	340.8k	<sup>2</sup> ACCIARRI	<b>00</b> C	L3	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV		
$-0.50150 \pm 0.00046$	500k	BARATE	<b>00</b> C	ALEP	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV		

 $<sup>^1</sup>$ ABBIENDI 010 use their measurement of the au polarization in addition to the lineshape and forward-backward lepton asymmetries.

and forward-backward lepton asymmetries.  $^2$  ACCIARRI 00C use their measurement of the  $\tau$  polarization in addition to forwardbackward lepton asymmetries.

 $<sup>^2</sup>$ ACCIARRI 00C use their measurement of the au polarization in addition to forwardbackward lepton asymmetries.

<b>8 Å</b> VALUE	<u>EVTS</u>	DOCUMENT ID		<u>TECN</u>	<u>COMMENT</u>
$0.50 \begin{array}{l} +0.04 \\ -0.06 \end{array}$ OUR AVE	RAGE				
$0.501 \pm 0.110$	156k	$^{ m 1}$ ABAZOV	<b>11</b> D	D0	$E_{cm}^{ar{p}}=1.97\;TeV$
0.57 ±0.08	1500	<sup>2</sup> AKTAS	06	H1	$e^{\pm} p  ightarrow  \overline{ u}_{m{e}}( u_{m{e}}) X, \ \sqrt{s} pprox 300 \; {\sf GeV}$
$0.47 \begin{array}{l} +0.05 \\ -0.33 \end{array}$		<sup>3</sup> LEP-SLC	06		$E_{cm}^{\mathit{ee}} = 88 – 94 \; GeV$
$0.441^{+0.207}_{-0.173}{\pm0.067}$	5026	<sup>4</sup> ACOSTA	05м	CDF	$E_{cm}^{p\overline{p}} = 1.96 \; TeV$

- $^1$  ABAZOV 11D study  $p\overline{p}\to Z/\gamma^*\,e^+\,e^-$  events using 5 fb $^{-1}$  data at  $\sqrt{s}=1.96$  TeV. The candidate events are selected by requiring two isolated electromagnetic showers with  $E_T>25$  GeV, at least one electron in the central region and the di-electron mass in the range 50–1000 GeV. From the forward-backward asymmetry, determined as a function of the di-electron mass, they derive the axial and vector couplings of the u- and d- quarks and the value of  $\sin^2\!\theta_{eff}^\ell=0.2309\pm0.0008(\text{stat})\pm0.0006(\text{syst}).$
- $^2$  AKTAS 06 fit the neutral current (1.5  $\leq$  Q $^2$   $\leq$  30,000 GeV $^2$ ) and charged current (1.5  $\leq$  Q $^2$   $\leq$  15,000 GeV $^2$ ) differential cross sections. In the determination of the u-quark couplings the electron and d-quark couplings are fixed to their standard model values.
- <sup>3</sup>LEP-SLC 06 is a combination of the results from LEP and SLC experiments using light quark tagging. s- and d-quark couplings are assumed to be identical.
- <sup>4</sup> ACOSTA 05M determine the forward-backward asymmetry of  $e^+e^-$  pairs produced via  $q \overline{q} \rightarrow Z/\gamma^* \rightarrow e^+e^-$  in 15 M( $e^+e^-$ ) effective mass bins ranging from 40 GeV to 600 GeV. These results are used to obtain the vector and axial-vector couplings of the Z to the light quarks, assuming the electron couplings are as predicted by the Standard Model. Higher order radiative corrections have not been taken into account.

"d	
B A	

<u>VALUE</u>	<u>EVTS</u>	DOCUMENT ID		TECN	COMMENT		
-0.523 <sup>+0.050</sup> <sub>-0.029</sub> OUR A	<b>VERAGE</b>						
$-0.497 \pm 0.165$	156k	<sup>1</sup> ABAZOV	<b>11</b> D	D0	$E_{cm}^{ar{p}}=1.97\;TeV$		
$-0.80 \pm 0.24$	1500	<sup>2</sup> AKTAS	06	H1	${ m e}^{\pm} p  ightarrow \; \overline{ u}_e( u_e) X, \ \sqrt{s} pprox 300 \; { m GeV}$		
$-0.52 \begin{array}{l} +0.05 \\ -0.03 \end{array}$		<sup>3</sup> LEP-SLC	06		$E_{cm}^{ee} = 88 – 94 \; GeV$		
$-0.016^{+0.346}_{-0.536}\pm0.091$	5026	<sup>4</sup> ACOSTA	05м	CDF	$E_{cm}^{oldsymbol{p}\overline{oldsymbol{p}}}$ = 1.96 TeV		

- $^1$  ABAZOV 11D study  $p\overline{p}\to Z/\gamma^*\,e^+\,e^-$  events using  $5~{\rm fb}^{-1}$  data at  $\sqrt{s}=1.96$  TeV. The candidate events are selected by requiring two isolated electromagnetic showers with  $E_T>25~{\rm GeV}$ , at least one electron in the central region and the di-electron mass in the range 50–1000 GeV. From the forward-backward asymmetry, determined as a function of the di-electron mass, they derive the axial and vector couplings of the u- and d- quarks and the value of  $\sin^2\!\theta_{eff}^\ell=0.2309\pm0.0008({\rm stat})\pm0.0006({\rm syst})$ .
- $^2$  AKTAS 06 fit the neutral current (1.5  $\leq$  Q $^2$   $\leq$  30,000 GeV $^2$ ) and charged current (1.5  $\leq$  Q $^2$   $\leq$  15,000 GeV $^2$ ) differential cross sections. In the determination of the d-quark couplings the electron and u-quark couplings are fixed to their standard model values.
- <sup>3</sup> LEP-SLC 06 is a combination of the results from LEP and SLC experiments using light quark tagging. s- and d-quark couplings are assumed to be identical.

<sup>4</sup> ACOSTA 05M determine the forward-backward asymmetry of  $e^+e^-$  pairs produced via  $q \overline{q} \rightarrow Z/\gamma^* \rightarrow e^+e^-$  in 15 M( $e^+e^-$ ) effective mass bins ranging from 40 GeV to 600 GeV. These results are used to obtain the vector and axial-vector couplings of the Z to the light quarks, assuming the electron couplings are as predicted by the Standard Model. Higher order radiative corrections have not been taken into account.

### Z COUPLINGS TO NEUTRAL LEPTONS

Averaging over neutrino species, the invisible Z decay width determines the effective neutrino coupling  $g^{\nu\ell}$ . For  $g^{\nu}_{e}$  and  $g^{\nu\mu}$ ,  $\nu_{e}e$  and  $\nu_{\mu}e$  scattering results are combined with  $g^{e}_{A}$  and  $g^{e}_{V}$  measurements at the Z mass to obtain  $g^{\nu}_{e}$  and  $g^{\nu}_{\mu}$  following NOVIKOV 93C.

$oldsymbol{g}^{oldsymbol{ u}_{oldsymbol{\ell}}}$			
VALUE	<u>DOCUMENT ID</u>	)	COMMENT
$0.50076 \pm 0.00076$	$^{ m 1}$ LEP-SLC	06	$E_{cm}^{ee} = 88-94 \; GeV$
$^{ m 1}$ From invisible $Z$ -decay	width.		
$g^{ u_{\mathbf{e}}}$			
VALUE	DOCUMENT ID	TECN	COMMENT
0.528±0.085			From $\nu_{\mu} e$ and $\nu_{e} e$ scattering
$^1$ VILAIN 94 derive this $1.05 {+0.15 \atop -0.18}.$	value from their valu	ie of <i>g</i>	$g^{ u}\mu$ and their ratio $g^{ u}e/g^{ u}\mu$ =
${\sf g}^{\nu_\mu}$			
<u>VALUE</u>	<u>DOCUMENT ID</u>		TECN COMMENT
$0.502 \pm 0.017$	$^{1}$ VILAIN	94	CHM2 From $ u_{\mu}  e$ scattering
1			$e\nu_{\mu}$

 $<sup>^1</sup>$  VILAIN 94 derive this value from their measurement of the couplings  $g_A^{e\,\nu_\mu}=-0.503\pm0.017$  and  $g_V^{e\,\nu_\mu}=-0.035\pm0.017$  obtained from  $\nu_\mu\,e$  scattering. We have re-evaluated this value using the current PDG values for  $g_A^e$  and  $g_V^e$ .

#### Z ASYMMETRY PARAMETERS

For each fermion-antifermion pair coupling to the  ${\it Z}$  these quantities are defined as

$$A_f = \frac{2g_V^f g_A^f}{(g_V^f)^2 + (g_A^f)^2}$$

where  $g_V^f$  and  $g_A^f$  are the effective vector and axial-vector couplings. For their relation to the various lepton asymmetries see the note "The Z boson" and ref. LEP-SLC 06.

 $A_e$ 

Using polarized beams, this quantity can also be measured as  $(\sigma_L - \sigma_R)/(\sigma_L + \sigma_R)$ , where  $\sigma_L$  and  $\sigma_R$  are the  $e^+e^-$  production cross sections for Z bosons produced with left-handed and right-handed electrons respectively.

VALUE	<b>EVTS</b>	DOCUMENT ID		TECN	COMMENT
0.1515±0.0019 OUR AVERA	AGE				
$0.1454 \pm 0.0108 \pm 0.0036$	144810	$^{ m 1}$ ABBIENDI	010	OPAL	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
$0.1516 \pm 0.0021$	559000	<sup>2</sup> ABE	<b>01</b> B	SLD	$E_{\rm cm}^{\rm ee} = 91.24~{\rm GeV}$
$0.1504 \pm 0.0068 \pm 0.0008$		<sup>3</sup> HEISTER	01	ALEP	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
$0.1382 \pm 0.0116 \pm 0.0005$	105000	<sup>4</sup> ABREU	00E	DLPH	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
$0.1678 \pm 0.0127 \pm 0.0030$	137092	<sup>5</sup> ACCIARRI	98H	L3	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
$0.162\ \pm0.041\ \pm0.014$	89838	<sup>6</sup> ABE	97	SLD	$E_{cm}^{ee} = 91.27 \; GeV$
$0.202\ \pm0.038\ \pm0.008$		<sup>7</sup> ABE	95J	SLD	$E_{\rm cm}^{ee} = 91.31 \; {\rm GeV}$

 $<sup>^1</sup>$  ABBIENDI 010 fit for  $A_e$  and  $A_\tau$  from measurements of the  $\tau$  polarization at varying  $\tau$  production angles. The correlation between  $A_e$  and  $A_\tau$  is less than 0.03.

 $<sup>^7</sup>$ ABE 95J obtain this result from polarized Bhabha scattering.



This quantity is directly extracted from a measurement of the left-right forward-backward asymmetry in  $\mu^+\mu^-$  production at SLC using a polarized electron beam. This double asymmetry eliminates the dependence on the *Z-e-e* coupling parameter  $A_e$ .

<u>VALUE</u>	<u>EVTS</u>	<u>DOCUMENT ID</u>		TECN	COMMENT
$0.142 \pm 0.015$	16844	$^{ m 1}$ ABE	<b>01</b> B	SLD	$E_{cm}^{ee} = 91.24 \text{ GeV}$

<sup>&</sup>lt;sup>1</sup> ABE 01B obtain this direct measurement using the left-right production and left-right forward-backward polar angle asymmetries in  $\mu^+\mu^-$  decays of the Z boson obtained with a polarized electron beam.



The LEP Collaborations derive this quantity from the measurement of the  $\tau$  polarization in  $Z\to \tau^+\tau^-$ . The SLD Collaboration directly extracts this quantity from its measured left-right forward-backward asymmetry in  $Z\to \tau^+\tau^-$  produced using a polarized  $e^-$  beam. This double asymmetry eliminates the dependence on the Z-e-e coupling parameter  $A_e$ .

<u>VALUE</u>	<u>EVTS</u>	DOCUMENT ID TECN		TECN	COMMENT			
0.143 ±0.004 OUR AVERAGE								
$0.1456 \pm 0.0076 \pm 0.0057$	144810	$^{ m 1}$ abbiendi	010	OPAL	<i>E</i> <sup>ee</sup> <sub>cm</sub> = 88–94 GeV			
$0.136 \pm 0.015$	16083	<sup>2</sup> ABE	<b>01</b> B	SLD	$E_{\rm cm}^{\it ee}=91.24~{\rm GeV}$			
$0.1451 \pm 0.0052 \pm 0.0029$		<sup>3</sup> HEISTER	01	ALEP	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV			
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 $<sup>^2</sup>$  ABE 01B use the left-right production and left-right forward-backward decay asymmetries in leptonic Z decays to obtain a value of 0.1544  $\pm$  0.0060. This is combined with left-right production asymmetry measurement using hadronic Z decays (ABE 00B) to obtain the quoted value.

<sup>&</sup>lt;sup>3</sup> HEISTER 01 obtain this result fitting the  $\tau$  polarization as a function of the polar production angle of the  $\tau$ .

ABREU 00E obtain this result fitting the  $\tau$  polarization as a function of the polar  $\tau$  production angle. This measurement is a combination of different analyses (exclusive  $\tau$  decay modes, inclusive hadronic 1-prong reconstruction, and a neural network analysis).

 $<sup>^{5}</sup>$  Derived from the measurement of forward-backward  $\tau$  polarization asymmetry.

 $<sup>^6</sup>$  ABE 97 obtain this result from a measurement of the observed left-right charge asymmetry,  $A_Q^{\rm obs}=0.225\pm0.056\pm0.019,$  in hadronic Z decays. If they combine this value of  $A_Q^{\rm obs}$  with their earlier measurement of  $A_{LR}^{\rm obs}$  they determine  $A_e$  to be  $-0.1574\pm0.0197\pm0.0067$  independent of the beam polarization.

 $0.1359 \pm 0.0079 \pm 0.0055$  105000 <sup>4</sup> ABREU 00E DLPH  $E_{\text{cm}}^{\text{ee}} = 88-94 \text{ GeV}$   $0.1476 \pm 0.0088 \pm 0.0062$  137092 ACCIARRI 98H L3  $E_{\text{cm}}^{\text{ee}} = 88-94 \text{ GeV}$ 

<sup>&</sup>lt;sup>4</sup>ABREU 00E obtain this result fitting the  $\tau$  polarization as a function of the polar  $\tau$  production angle. This measurement is a combination of different analyses (exclusive  $\tau$  decay modes, inclusive hadronic 1-prong reconstruction, and a neural network analysis).



The SLD Collaboration directly extracts this quantity by a simultaneous fit to four measured s-quark polar angle distributions corresponding to two states of  $e^-$  polarization (positive and negative) and to the  $K^+$   $K^-$  and  $K^\pm$   $K_S^0$  strange particle tagging modes in the hadronic final states.

		. 504.000.			
VALUE	EVTS	DOCUMENT ID		TECN	COMMENT
$0.895 \pm 0.066 \pm 0.062$	2870	<sup>1</sup> ABE	<b>00</b> D	SLD	$E_{\rm cm}^{\rm ee} = 91.2 \text{ GeV}$

<sup>&</sup>lt;sup>1</sup> ABE 00D tag  $Z \to s\overline{s}$  events by an absence of B or D hadrons and the presence in each hemisphere of a high momentum  $K^{\pm}$  or  $K_{S}^{0}$ .



This quantity is directly extracted from a measurement of the left-right forward-backward asymmetry in  $c\overline{c}$  production at SLC using polarized electron beam. This double asymmetry eliminates the dependence on the Z-e-e coupling parameter  $A_e$ . OUR FIT is obtained by a simultaneous fit to several c- and b-quark measurements as explained in the note "The Z boson" and ref. LEP-SLC 06.

VALUE	DOCUMENT ID		TECN	COMMENT
$0.670 \pm 0.027$ OUR FIT				
$0.6712 \pm 0.0224 \pm 0.0157$	<sup>1</sup> ABE	05	SLD	$E_{\rm cm}^{\it ee}=91.24~{\rm GeV}$
• • • We do not use the following	data for averages	, fits,	limits, e	etc. • • •
$0.583 \pm 0.055 \pm 0.055$	<sup>2</sup> ABE	02G	SLD	E <sup>ee</sup> <sub>cm</sub> = 91.24 GeV
$0.688 \pm 0.041$	<sup>3</sup> ABE	010	SLD	$E_{\rm cm}^{\it ee}=91.25~{\rm GeV}$

 $<sup>^1</sup>$  ABE 05 use hadronic Z decays collected during 1996–98 to obtain an enriched sample of  $c\,\overline{c}$  events tagging on the invariant mass of reconstructed secondary decay vertices. The charge of the underlying c–quark is obtained with an algorithm that takes into account the net charge of the vertex as well as the charge of tracks emanating from the vertex and identified as kaons. This yields (9970 events)  $A_c=0.6747\pm0.0290\pm0.0233$ . Taking into account all correlations with earlier results reported in ABE 02G and ABE 01C, they obtain the quoted overall SLD result.

 $<sup>^1</sup>$  ABBIENDI 010 fit for  $A_e$  and  $A_\tau$  from measurements of the  $\tau$  polarization at varying  $\tau$  production angles. The correlation between  $A_e$  and  $A_\tau$  is less than 0.03.

<sup>&</sup>lt;sup>2</sup> ABE 01B obtain this direct measurement using the left-right production and left-right forward-backward polar angle asymmetries in  $\tau^+\tau^-$  decays of the Z boson obtained with a polarized electron beam.

<sup>&</sup>lt;sup>3</sup> HEISTER 01 obtain this result fitting the  $\tau$  polarization as a function of the polar production angle of the  $\tau$ .

 $<sup>^2</sup>$  ABE 02G tag b and c quarks through their semileptonic decays into electrons and muons. A maximum likelihood fit is performed to extract simultaneously  $A_{\mbox{\it b}}$  and  $A_{\mbox{\it c}}$  .

<sup>&</sup>lt;sup>3</sup> ABE 01C tag  $Z \to c \overline{c}$  events using two techniques: exclusive reconstruction of  $D^{*+}$ ,  $D^+$  and  $D^0$  mesons and the soft pion tag for  $D^{*+} \to D^0 \pi^+$ . The large background from D mesons produced in  $b \overline{b}$  events is separated efficiently from the signal using precision vertex information. When combining the  $A_c$  values from these two samples, care is taken to avoid double counting of events common to the two samples, and common systematic errors are properly taken into account.



This quantity is directly extracted from a measurement of the left-right forwardbackward asymmetry in  $b\overline{b}$  production at SLC using polarized electron beam. This double asymmetry eliminates the dependence on the Z-e-e coupling parameter  $A_e$ . OUR FIT is obtained by a simultaneous fit to several c- and b-quark measurements as explained in the note "The  $\it Z$  boson" and ref. LEP-SLC 06.

VALUE	<u>EVTS</u>	DOCUMENT ID		TECN	COMMENT			
0.923 ±0.020 OUR FIT								
$0.9170 \pm 0.0147 \pm 0.0145$		$^{ m 1}$ ABE	05	SLD	$E_{ m cm}^{ee} = 91.24 \; { m GeV}$			
• • • We do not use the following data for averages, fits, limits, etc. • •								
$0.907\ \pm0.020\ \pm0.024$	48028	<sup>2</sup> ABE	03F	SLD	$E_{\rm cm}^{\rm ee}=91.24~{\rm GeV}$			
$0.919 \pm 0.030 \pm 0.024$		<sup>3</sup> ABE	02G	SLD	$E_{cm}^{ee} = 91.24 \; GeV$			

<sup>4</sup> ABE

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- 99L SLD  $E_{\text{cm}}^{\text{ee}} = 91.27 \text{ GeV}$  $0.855 \pm 0.088 \pm 0.102$  $^{
  m 1}$  ABE 05 use hadronic Z decays collected during 1996–98 to obtain an enriched sample of  $b\overline{b}$  events tagging on the invariant mass of reconstructed secondary decay vertices. The charge of the underlying b-quark is obtained with an algorithm that takes into account the net charge of the vertex as well as the charge of tracks emanating from the vertex and identified as kaons. This yields (25917 events)  $A_b=0.9173\pm0.0184\pm0.0173$ . Taking into account all correlations with earlier results reported in ABE 03F, ABE 02G
  - and ABE 99L, they obtain the quoted overall SLD result.  $^2$  ABE 03F obtain an enriched sample of  $b\overline{b}$  events tagging on the invariant mass of a 3-dimensional topologically reconstructed secondary decay. The charge of the underlying b quark is obtained using a self-calibrating track-charge method. For the 1996–1998 data sample they measure  $A_b=0.906\pm0.022\pm0.023$ . The value quoted here is obtained combining the above with the result of ABE 98I (1993–1995 data sample).
  - $^3$  ABE 02G tag b and c quarks through their semileptonic decays into electrons and muons. A maximum likelihood fit is performed to extract simultaneously  $A_b$  and  $A_c$ .
  - $^4$  ABE 99L obtain an enriched sample of  $b\overline{b}$  events tagging with an inclusive vertex mass cut. For distinguishing b and  $\overline{b}$  quarks they use the charge of identified  $K^{\pm}$ .

### TRANSVERSE SPIN CORRELATIONS IN $Z \rightarrow \tau^+ \tau^-$

The correlations between the transverse spin components of  $\tau^+\tau^-$  produced in Z decays may be expressed in terms of the vector and axial-vector couplings:

$$C_{TT} = \frac{|g_A^{\tau}|^2 - |g_V^{\tau}|^2}{|g_A^{\tau}|^2 + |g_V^{\tau}|^2}$$

$$C_{TN} = -2 \frac{|g_A^{\tau}| |g_V^{\tau}|}{|g_A^{\tau}|^2 + |g_V^{\tau}|^2} \sin(\Phi_{g_V^{\tau}} - \Phi_{g_A^{\tau}})$$

 $C_{TT}$  refers to the transverse-transverse (within the collision plane) spin correlation and  $C_{TN}$  refers to the transverse-normal (to the collision plane) spin correlation.

The longitudinal  $\tau$  polarization  $P_{\tau}$   $(=-A_{\tau})$  is given by:

$$P_{\tau} = -2 \frac{|g_A^{\tau}||g_V^{\tau}|}{|g_A^{\tau}|^2 + |g_V^{\tau}|^2} \cos(\Phi_{g_V^{\tau}} - \Phi_{g_A^{\tau}})$$

Here  $\Phi$  is the phase and the phase difference  $\Phi_{{\mathcal g}_V^{\mathcal T}} - \Phi_{{\mathcal g}_A^{\mathcal T}}$  can be obtained using both the measurements of  $\mathcal C_{TN}$  and  $\mathcal P_{\mathcal T}.$ 

$c_{TT}$					
VALUE	<b>EVTS</b>	DOCUMENT ID		TECN	COMMENT
1.01±0.12 OUR AVERA	AGE				
$0.87\!\pm\!0.20{}^{+0.10}_{-0.12}$	9.1k	ABREU	<b>97</b> G	DLPH	Eee 91.2 GeV
$1.06\!\pm\!0.13\!\pm\!0.05$	120k	BARATE	<b>97</b> D	ALEP	E <sup>ee</sup> <sub>cm</sub> = 91.2 GeV
$C_{TN}$					
VALUE	<u>EVTS</u>	DOCUMENT ID		TECN	COMMENT
$0.08 \pm 0.13 \pm 0.04$	120k	<sup>1</sup> BARATE	<b>97</b> D	ALEP	$E_{\rm cm}^{\it ee}=91.2~{\rm GeV}$

 $<sup>^1</sup>$  BARATE 97D combine their value of  $C_{TN}$  with the world average  $P_{\tau}=-0.140\pm0.007$  to obtain  $\tan(\Phi_{g_{A}^{\tau}}-\Phi_{g_{A}^{\tau}})=-0.57\pm0.97.$ 

### FORWARD-BACKWARD $e^+e^- \rightarrow f \overline{f}$ CHARGE ASYMMETRIES

These asymmetries are experimentally determined by tagging the respective lepton or quark flavor in  $e^+\,e^-$  interactions. Details of heavy flavor (c- or b-quark) tagging at LEP are described in the note on "The Z boson" and ref. LEP-SLC 06. The Standard Model predictions for LEP data have been (re)computed using the ZFITTER package (version 6.36) with input parameters  $M_Z{=}91.187~{\rm GeV},~M_{\rm top}{=}174.3~{\rm GeV},~M_{\rm Higgs}{=}150~{\rm GeV},~\alpha_s{=}0.119,~\alpha^{(5)}~(M_Z){=}~1/128.877$  and the Fermi constant  $G_F{=}~1.16637\times 10^{-5}~{\rm GeV}^{-2}$  (see the note on "The Z boson" for references). For non-LEP data the Standard Model predictions are as given by the authors of the respective publications.

### $A_{FR}^{(0,e)}$ CHARGE ASYMMETRY IN $e^+e^- \rightarrow e^+e^-$

OUR FIT is obtained using the fit procedure and correlations as determined by the LEP Electroweak Working Group (see the note "The Z boson" and ref. LEP-SLC 06). For the Z peak, we report the pole asymmetry defined by  $(3/4)A_e^2$  as determined by the nine-parameter fit to cross-section and lepton forward-backward asymmetry data.

ASYMMETRY (%)	STD. MODEL	$\sqrt{s}$ (GeV)	DOCUMENT ID		TECN
1.45±0.25 OUR FIT					
$0.89 \pm 0.44$	1.57	91.2	<sup>1</sup> ABBIENDI	01A	OPAL
$1.71 \pm 0.49$	1.57	91.2	ABREU	00F	DLPH
$1.06 \pm 0.58$	1.57	91.2	ACCIARRI	<b>00</b> C	L3
$1.88 \pm 0.34$	1.57	91.2	<sup>2</sup> BARATE	<b>00</b> C	ALEP

## ---- $A^{(0,\mu)}_{FB}$ CHARGE ASYMMETRY IN $e^+e^ightarrow~\mu^+\mu^-$ -

OUR FIT is obtained using the fit procedure and correlations as determined by the LEP Electroweak Working Group (see the note "The Z boson" and ref. LEP-SLC 06). For the Z peak, we report the pole asymmetry defined by  $(3/4)A_{\rm e}A_{\mu}$  as determined by the nine-parameter fit to cross-section and lepton forward-backward asymmetry data.

ASYMMETRY (%)	STD. MODEL	$\sqrt{s}$ (GeV)	DOCUMENT ID		TECN
1.69± 0.13 OUR FIT					
$1.59 \pm 0.23$	1.57	91.2	<sup>1</sup> ABBIENDI	01A	OPAL
$1.65 \pm 0.25$	1.57	91.2	ABREU	00F	DLPH
$1.88 \pm 0.33$	1.57	91.2	ACCIARRI	<b>00</b> C	L3
$1.71 \pm 0.24$	1.57	91.2	<sup>2</sup> BARATE	00C	ALEP
• • • We do not use the follow	ving data for	averages, fi	ts, limits, etc. • •	•	
9 ±30	-1.3	20	<sup>3</sup> ABREU	95м	DLPH
$7 \pm 26$	-8.3	40	<sup>3</sup> ABREU	95M	DLPH
$-11$ $\pm 33$	-24.1	57	<sup>3</sup> ABREU	95M	DLPH
$-62$ $\pm 17$	-44.6	69	<sup>3</sup> ABREU	95M	DLPH
$-56$ $\pm 10$	-63.5	79	<sup>3</sup> ABREU	95M	DLPH
$-13$ $\pm$ $5$	-34.4	87.5	<sup>3</sup> ABREU	95M	DLPH
$-29.0 \ \ {}^{+}_{-}\ \ {}^{5.0}_{-}\ \ \pm 0.5$	-32.1	56.9	<sup>4</sup> ABE	901	VNS
$-$ 9.9 $\pm$ 1.5 $\pm$ 0.5	-9.2	35	HEGNER	90	JADE
$0.05 \pm 0.22$	0.026	91.14	<sup>5</sup> ABRAMS	89D	MRK2
$-43.4 \pm 17.0$	-24.9	52.0	<sup>6</sup> BACALA	89	AMY
$-11.0 \pm 16.5$	-29.4	55.0	<sup>6</sup> BACALA	89	AMY
$-30.0 \pm 12.4$	-31.2	56.0	<sup>6</sup> BACALA	89	AMY
$-46.2 \pm 14.9$	-33.0	57.0	<sup>6</sup> BACALA	89	AMY
$-29$ $\pm 13$	-25.9	53.3	ADACHI	<b>88</b> C	TOPZ
$+$ 5.3 $\pm$ 5.0 $\pm$ 0.5	-1.2	14.0	ADEVA	88	MRKJ
$-10.4~\pm~1.3~\pm0.5$	-8.6	34.8	ADEVA	88	MRKJ
$-12.3~\pm~5.3~\pm0.5$	-10.7	38.3	ADEVA	88	MRKJ
$-15.6~\pm~3.0~\pm0.5$	-14.9	43.8	ADEVA	88	MRKJ
$-$ 1.0 $\pm$ 6.0	-1.2	13.9	BRAUNSCH	88D	TASS
$-$ 9.1 $\pm$ 2.3 $\pm$ 0.5	-8.6	34.5	BRAUNSCH	88D	TASS
$-10.6 \ \ \begin{array}{c} + \ \ 2.2 \\ - \ \ 2.3 \end{array} \ \pm 0.5$	-8.9	35.0	BRAUNSCH	88D	TASS
$-17.6 \ \ \begin{array}{c} + \ \ 4.4 \\ - \ \ 4.3 \end{array} \ \pm 0.5$	-15.2	43.6	BRAUNSCH	88D	TASS
$-$ 4.8 $\pm$ 6.5 $\pm$ 1.0	-11.5	39	BEHREND	87C	CELL
$-18.8~\pm~4.5~\pm1.0$	-15.5	44	BEHREND	87C	CELL
$+\ 2.7\ \pm\ 4.9$	-1.2	13.9	BARTEL	86C	JADE
$-11.1 ~\pm~ 1.8 ~\pm 1.0$	-8.6	34.4	BARTEL	86C	JADE
$-17.3 \pm 4.8 \pm 1.0$	-13.7	41.5	BARTEL	<b>86</b> C	JADE
$-22.8 \pm 5.1 \pm 1.0$	-16.6	44.8	BARTEL	<b>86</b> C	JADE

 $<sup>^{1}</sup>$  ABBIENDI 01A error includes approximately 0.38 due to statistics, 0.16 due to event selection systematics, and 0.18 due to the theoretical uncertainty in t-channel prediction.

 $<sup>^2</sup>$  BARATE 00C error includes approximately 0.31 due to statistics, 0.06 due to experimental systematics, and 0.13 due to the theoretical uncertainty in t-channel prediction.

$-$ 6.3 $\pm$ 0.8 $\pm$ 0.2	-6.3	29	ASH	85	MAC
$-$ 4.9 $\pm$ 1.5 $\pm$ 0.5	-5.9	29	DERRICK	85	HRS
$-$ 7.1 $\pm$ 1.7	-5.7	29	LEVI	83	MRK2
-16.1 + 3.2	-9.2	34.2	BRANDELIK	82C	TASS

 $<sup>^{</sup>m 1}$  ABBIENDI 01A error is almost entirely on account of statistics.

## - $A^{(0, au)}_{FB}$ CHARGE ASYMMETRY IN $e^+\,e^ightarrow~ au^+\, au^-$ -

OUR FIT is obtained using the fit procedure and correlations as determined by the LEP Electroweak Working Group (see the note "The Z boson" and ref. LEP-SLC 06). For the Z peak, we report the pole asymmetry defined by  $(3/4)A_{\rho}A_{\tau}$  as determined by the nine-parameter fit to cross-section and lepton forward-backward asymmetry data.

	STD.	$\sqrt{s}$			
ASYMMETRY (%)	MODEL	(GeV)	DOCUMENT ID		TECN
1.88± 0.17 OUR FIT					
$1.45 \pm 0.30$	1.57	91.2	<sup>1</sup> ABBIENDI	01A	OPAL
$2.41 \pm 0.37$	1.57	91.2	ABREU	00F	DLPH
$2.60 \pm 0.47$	1.57	91.2	ACCIARRI	00C	L3
$1.70 \pm 0.28$	1.57	91.2	<sup>2</sup> BARATE	<b>00</b> C	ALEP
• • • We do not use the follo	wing data fo	r averages, f	fits, limits, etc. • •	• •	
$-32.8   {+}  {}^{\displaystyle 6.4}_{\displaystyle -6.2}   \pm 1.5$	-32.1	56.9	<sup>3</sup> ABE	901	VNS
$-\ 8.1\ \pm\ 2.0\ \pm0.6$	-9.2	35	HEGNER	90	JADE
$-18.4\ \pm 19.2$	-24.9	52.0	<sup>4</sup> BACALA	89	AMY
$-17.7 \pm 26.1$	-29.4	55.0	<sup>4</sup> BACALA	89	AMY
$-45.9 \pm 16.6$	-31.2	56.0	<sup>4</sup> BACALA	89	AMY
$-49.5\ \pm 18.0$	-33.0	57.0	<sup>4</sup> BACALA	89	AMY
$-20$ $\pm 14$	-25.9	53.3	ADACHI	88C	TOPZ
$-10.6~\pm~3.1~\pm1.5$	-8.5	34.7	ADEVA	88	MRKJ
$-$ 8.5 $\pm$ 6.6 $\pm$ 1.5	-15.4	43.8	ADEVA	88	MRKJ
$-$ 6.0 $\pm$ 2.5 $\pm$ 1.0	8.8	34.6	BARTEL	85F	JADE
$-11.8 \pm 4.6 \pm 1.0$	14.8	43.0	BARTEL	85F	JADE
$-$ 5.5 $\pm$ 1.2 $\pm$ 0.5	-0.063	29.0	FERNANDEZ	85	MAC
$-$ 4.2 $\pm$ 2.0	0.057	29	LEVI	83	MRK2
$-10.3 \pm 5.2$	-9.2	34.2	BEHREND	82	CELL
$-$ 0.4 $\pm$ 6.6	-9.1	34.2	BRANDELIK	82C	TASS

<sup>&</sup>lt;sup>1</sup> ABBIENDI 01A error includes approximately 0.26 due to statistics and 0.14 due to event selection systematics.

<sup>&</sup>lt;sup>2</sup>BARATE 00C error is almost entirely on account of statistics.

<sup>&</sup>lt;sup>3</sup> ABREU 95M perform this measurement using radiative muon-pair events associated with high-energy isolated photons.

<sup>&</sup>lt;sup>4</sup> ABE 901 measurements in the range 50  $< \sqrt{s} <$  60.8 GeV.

<sup>&</sup>lt;sup>5</sup> ABRAMS 89D asymmetry includes both 9  $\mu^+\mu^-$  and 15  $\tau^+\tau^-$  events.

<sup>&</sup>lt;sup>6</sup>BACALA 89 systematic error is about 5%.

<sup>&</sup>lt;sup>2</sup>BARATE 00C error includes approximately 0.26 due to statistics and 0.11 due to experimental systematics.

<sup>&</sup>lt;sup>3</sup> ABE 901 measurements in the range 50  $< \sqrt{s} <$  60.8 GeV.

<sup>&</sup>lt;sup>4</sup>BACALA 89 systematic error is about 5%.

### 

For the Z peak, we report the pole asymmetry defined by  $(3/4)A_\ell^2$  as determined by the five-parameter fit to cross-section and lepton forward-backward asymmetry data assuming lepton universality. For details see the note "The Z boson" and ref. LEP-SLC 06.

ASYMMETRY (%)	STD. MODEL	$\frac{\sqrt{s}}{(\text{GeV})}$	DOCUMENT ID		TECN
1.71±0.10 OUR FIT					
$1.45 \pm 0.17$	1.57	91.2	$^{ m 1}$ ABBIENDI	<b>01</b> A	OPAL
$1.87 \pm 0.19$	1.57	91.2	ABREU	00F	DLPH
$1.92 \pm 0.24$	1.57	91.2	ACCIARRI	00C	L3
$1.73 \pm 0.16$	1.57	91.2	<sup>2</sup> BARATE	<b>00</b> C	ALEP

 $<sup>^{1}</sup>$  ABBIENDI 01A error includes approximately 0.15 due to statistics, 0.06 due to event selection systematics, and 0.03 due to the theoretical uncertainty in t-channel prediction.

## ——— $A_{FB}^{(0,u)}$ CHARGE ASYMMETRY IN $e^+e^- \rightarrow u\overline{u}$

ASYMMETRY (%)	MODEL	(GeV)	DOCUMENT ID	TECN
$4.0\pm6.7\pm2.8$	7.0	91.2	1 ACKERSTAFE 97T	OPAL

<sup>&</sup>lt;sup>1</sup> ACKERSTAFF 97⊤ measure the forward-backward asymmetry of various fast hadrons made of light quarks. Then using SU(2) isospin symmetry and flavor independence for down and strange quarks authors solve for the different quark types.

# 

The *s*-quark asymmetry is derived from measurements of the forward-backward asymmetry of fast hadrons containing an *s* quark.

ASYMMETRY (%)	STD. MODEL	$\sqrt{s}$ (GeV)	DOCUMENT ID	TECN
9.8 ±1.1 OUR AVERAGE				
$10.08\!\pm\!1.13\!\pm\!0.40$	10.1	91.2		B DLPH
$6.8 \pm 3.5 \pm 1.1$	10.1	91.2	<sup>2</sup> ACKERSTAFF 97	T OPAL

<sup>&</sup>lt;sup>1</sup> ABREU 00B tag the presence of an *s* quark requiring a high-momentum-identified charged kaon. The *s*-quark pole asymmetry is extracted from the charged-kaon asymmetry taking the expected *d*- and *u*-quark asymmetries from the Standard Model and using the measured values for the *c*- and *b*-quark asymmetries.

## 

OUR FIT, which is obtained by a simultaneous fit to several c- and b-quark measurements as explained in the note "The Z boson" and ref. LEP-SLC 06, refers to the Z pole asymmetry. The experimental values,

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<sup>&</sup>lt;sup>2</sup> BARATE 00C error includes approximately 0.15 due to statistics, 0.04 due to experimental systematics, and 0.02 due to the theoretical uncertainty in *t*-channel prediction.

<sup>&</sup>lt;sup>2</sup> ACKERSTAFF 97T measure the forward-backward asymmetry of various fast hadrons made of light quarks. Then using SU(2) isospin symmetry and flavor independence for down and strange quarks authors solve for the different quark types. The value reported here corresponds then to the forward-backward asymmetry for "down-type" quarks.

on the other hand, correspond to the measurements carried out at the respective energies.

ASYMMETRY (%)	STD. MODEL	$\sqrt{s}$ (GeV)	DOCUMENT ID		TECN
7.07± 0.35 OUR FIT					
$6.31 \pm 0.93 \pm 0.65$	6.35	91.26	<sup>1</sup> ABDALLAH	04F	DLPH
$5.68 \pm 0.54 \pm 0.39$	6.3	91.25	<sup>2</sup> ABBIENDI	<b>03</b> P	OPAL
$6.45 \pm 0.57 \pm 0.37$	6.10	91.21	<sup>3</sup> HEISTER	02H	ALEP
$6.59 \pm 0.94 \pm 0.35$	6.2	91.235	<sup>4</sup> ABREU	99Y	DLPH
$6.3 \pm 0.9 \pm 0.3$	6.1	91.22	<sup>5</sup> BARATE	980	ALEP
$6.3 \pm 1.2 \pm 0.6$	6.1	91.22	<sup>6</sup> ALEXANDER	<b>97</b> C	OPAL
$8.3 \pm 3.8 \pm 2.7$	6.2	91.24	<sup>7</sup> ADRIANI	<b>92</b> D	L3
• • • We do not use the follow	ving data for	averages, fit	ts, limits, etc. • •	•	
$3.1 \pm 3.5 \pm 0.5$	-3.5	89.43	<sup>1</sup> ABDALLAH	04F	DLPH
$11.0 \pm 2.8 \pm 0.7$	12.3	92.99	<sup>1</sup> ABDALLAH	04F	DLPH
$-$ 6.8 $\pm$ 2.5 $\pm$ 0.9	-3.0	89.51	<sup>2</sup> ABBIENDI	<b>03</b> P	OPAL
$14.6 \pm 2.0 \pm 0.8$	12.2	92.95	<sup>2</sup> ABBIENDI	<b>03</b> P	OPAL
$-12.4 \pm 15.9 \pm 2.0$	-9.6	88.38	<sup>3</sup> HEISTER	02H	ALEP
$-$ 2.3 $\pm$ 2.6 $\pm$ 0.2	-3.8	89.38	<sup>3</sup> HEISTER	02H	ALEP
$-$ 0.3 $\pm$ 8.3 $\pm$ 0.6	0.9	90.21	<sup>3</sup> HEISTER	02H	ALEP
$10.6 \pm 7.7 \pm 0.7$	9.6	92.05	<sup>3</sup> HEISTER	02H	ALEP
$11.9 \pm 2.1 \pm 0.6$	12.2	92.94	<sup>3</sup> HEISTER	02H	ALEP
$12.1 \pm 11.0 \pm 1.0$	14.2	93.90	<sup>3</sup> HEISTER	02H	ALEP
$-4.96\pm3.68\pm0.53$	-3.5	89.434	<sup>4</sup> ABREU	99Y	DLPH
$11.80 \pm \ 3.18 \pm 0.62$	12.3	92.990	<sup>4</sup> ABREU	99Y	DLPH
$-$ 1.0 $\pm$ 4.3 $\pm$ 1.0	-3.9	89.37	<sup>5</sup> BARATE	980	ALEP
$11.0 \pm 3.3 \pm 0.8$	12.3	92.96	<sup>5</sup> BARATE	980	ALEP
$3.9 \pm 5.1 \pm 0.9$	-3.4	89.45	<sup>6</sup> ALEXANDER	97C	OPAL
$15.8 \pm 4.1 \pm 1.1$	12.4	93.00	<sup>6</sup> ALEXANDER	97C	OPAL
$-12.9~\pm~7.8~\pm5.5$	-13.6	35	BEHREND	<b>90</b> D	CELL
$7.7\ \pm 13.4\ \pm 5.0$	-22.1	43	BEHREND	<b>90</b> D	CELL
$-12.8 \pm 4.4 \pm 4.1$	-13.6	35	ELSEN	90	JADE
$-10.9 \pm 12.9 \pm 4.6$	-23.2	44	ELSEN	90	JADE
$-14.9 \pm 6.7$	-13.3	35	OULD-SAADA	89	JADE

<sup>&</sup>lt;sup>1</sup> ABDALLAH 04F tag b- and c-quarks using semileptonic decays combined with charge flow information from the hemisphere opposite to the lepton. Enriched samples of  $c\overline{c}$  and  $b\overline{b}$  events are obtained using lifetime information.

## 

<sup>&</sup>lt;sup>2</sup> ABBIENDI 03P tag heavy flavors using events with one or two identified leptons. This allows the simultaneous fitting of the b and c quark forward-backward asymmetries as well as the average  $B^0 - \overline{B}{}^0$  mixing.

 $<sup>^3</sup>$  HEISTER 02H measure simultaneously b and c quark forward-backward asymmetries using their semileptonic decays to tag the quark charge. The flavor separation is obtained with a discriminating multivariate analysis.

<sup>&</sup>lt;sup>4</sup> ABREU 99Y tag  $Z \rightarrow b\overline{b}$  and  $Z \rightarrow c\overline{c}$  events by an exclusive reconstruction of several D meson decay modes ( $D^{*+}$ ,  $D^0$ , and  $D^+$  with their charge-conjugate states).

 $<sup>^5</sup>$  BARATE 980 tag  $Z\to c\overline{c}$  events requiring the presence of high-momentum reconstructed  $D^{*+},\,D^+,$  or  $D^0$  mesons.

<sup>&</sup>lt;sup>6</sup> ALEXANDER 97C identify the b and c events using a  $D/D^*$  tag.

<sup>&</sup>lt;sup>7</sup> ADRIANI 92D use both electron and muon semileptonic decays.

OUR FIT, which is obtained by a simultaneous fit to several c- and b-quark measurements as explained in the note "The Z boson" and ref. LEP-SLC 06, refers to the  $\bf Z$  pole asymmetry. The experimental values, on the other hand, correspond to the measurements carried out at the respective energies.

	STD.	$\sqrt{s}$ (GeV)		
ASYMMETRY (%)  9.92 ± 0.16 OUR FIT	MODEL	(GeV)	DOCUMENT ID	TECN
9.58 $\pm$ 0.32 $\pm$ 0.14	9.68	91.231	<sup>1</sup> ABDALLAH 05	DLPH
$10.04\pm \ 0.56\pm \ 0.25$	9.69	91.251	<sup>2</sup> ABDALLAH 04F	
$9.72\pm 0.42\pm 0.15$	9.67	91.25	<sup>3</sup> ABBIENDI 03F	
$9.72\pm 0.42\pm 0.13$ $9.77\pm 0.36\pm 0.18$	9.69	91.25	<sup>4</sup> ABBIENDI 021	OPAL
$9.52\pm 0.41\pm 0.17$	9.59	91.20	5 HEISTER 02	
$10.00\pm \ 0.27\pm \ 0.11$	9.63	91.21	6 HEISTER 010	
$7.62\pm 1.94\pm 0.85$	9.64	91.235	7 ABREU 99Y	
$9.60\pm 0.66\pm 0.33$	9.69	91.255	8 ACCIARRI 990	
$9.31\pm \ 1.01\pm \ 0.55$	9.65	91.24	9 ACCIARRI 98L	
$9.4 \pm 2.7 \pm 2.2$	9.61	91.24	10 ALEXANDER 970	
• • • We do not use the follow				OFAL
				DI DII
$6.37 \pm 1.43 \pm 0.17$	5.8	89.449	<sup>1</sup> ABDALLAH 05	DLPH
$10.41\pm \ 1.15\pm \ 0.24$	12.1	92.990	<sup>1</sup> ABDALLAH 05	DLPH
$6.7 \pm 2.2 \pm 0.2$	5.7	89.43	<sup>2</sup> ABDALLAH 04F	
$11.2 \pm 1.8 \pm 0.2$	12.1	92.99	<sup>2</sup> ABDALLAH 04F	
$4.7 \pm 1.8 \pm 0.1$	5.9	89.51	<sup>3</sup> ABBIENDI 03F	
$10.3 \pm 1.5 \pm 0.2$	12.0	92.95	<sup>3</sup> ABBIENDI 03F	
$5.82\pm\ 1.53\pm\ 0.12$	5.9	89.50	<sup>4</sup> ABBIENDI 021	OPAL
$12.21 \pm 1.23 \pm 0.25$	12.0	92.91	<sup>4</sup> ABBIENDI 02ı	OPAL
$-13.1 \pm 13.5 \pm 1.0$	3.2	88.38	<sup>5</sup> HEISTER 02h	
$5.5 \pm 1.9 \pm 0.1$	5.6	89.38	<sup>5</sup> HEISTER 02	
$-$ 0.4 $\pm$ 6.7 $\pm$ 0.8	7.5	90.21	<sup>5</sup> HEISTER 02	
$11.1 ~\pm~ 6.4 ~\pm~ 0.5$	11.0	92.05	<sup>5</sup> HEISTER 02	
$10.4 \pm 1.5 \pm 0.3$	12.0	92.94	<sup>5</sup> HEISTER 02	
$13.8 \pm 9.3 \pm 1.1$	12.9	93.90	<sup>5</sup> HEISTER 021	
$4.36\pm \ 1.19\pm \ 0.11$	5.8	89.472	<sup>6</sup> HEISTER 01	
$11.72 \pm 0.97 \pm 0.11$	12.0	92.950	<sup>6</sup> HEISTER 01	
$5.67 \pm 7.56 \pm 1.17$	5.7	89.434	<sup>7</sup> ABREU 99Y	
$8.82\pm \ 6.33\pm \ 1.22$	12.1	92.990	<sup>7</sup> ABREU 99Y	
$6.11\pm\ 2.93\pm\ 0.43$	5.9	89.50	8 ACCIARRI 990	
$13.71\pm\ 2.40\pm\ 0.44$	12.2	93.10	8 ACCIARRI 990	
$4.95\pm \ 5.23\pm \ 0.40$	5.8	89.45	<sup>9</sup> ACCIARRI 981	
$11.37 \pm \ 3.99 \pm \ 0.65$	12.1	92.99	<sup>9</sup> ACCIARRI 980	
$-$ 8.6 $\pm 10.8$ $\pm$ 2.9	5.8	89.45	<sup>10</sup> ALEXANDER 970	
$-$ 2.1 $\pm$ 9.0 $\pm$ 2.6	12.1	93.00	<sup>10</sup> ALEXANDER 970	OPAL
$-71$ $\pm 34$ $+ 7$ $- 8$	-58	58.3	SHIMONAKA 91	TOPZ
$-22.2 \pm 7.7 \pm 3.5$	-26.0	35	BEHREND 900	CELL
$-49.1 \pm 16.0 \pm 5.0$	-39.7	43	BEHREND 900	
$-28$ $\pm 11$	-23	35	BRAUNSCH 90	TASS
$-16.6~\pm~7.7~\pm~4.8$	-24.3	35	ELSEN 90	JADE
$-33.6 \pm 22.2 \pm 5.2$	-39.9	44	ELSEN 90	JADE
$3.4 \pm 7.0 \pm 3.5$	-16.0	29.0	BAND 89	MAC
$-72$ $\pm 28$ $\pm 13$	-56	55.2	SAGAWA 89	AMY

- <sup>1</sup> ABDALLAH 05 obtain an enriched samples of  $b\bar{b}$  events using lifetime information. The quark (or antiquark) charge is determined with a neural network using the secondary vertex charge, the jet charge and particle identification.
- <sup>2</sup> ABDALLAH 04F tag b- and c-quarks using semileptonic decays combined with charge flow information from the hemisphere opposite to the lepton. Enriched samples of  $c\overline{c}$  and  $b\overline{b}$  events are obtained using lifetime information.
- <sup>3</sup> ABBIENDI 03P tag heavy flavors using events with one or two identified leptons. This allows the simultaneous fitting of the b and c quark forward-backward asymmetries as well as the average  $B^0 \overline{B}{}^0$  mixing.
- <sup>4</sup> ABBIENDI 02I tag  $Z^0 \rightarrow b\overline{b}$  decays using a combination of secondary vertex and lepton tags. The sign of the *b*-quark charge is determined using an inclusive tag based on jet, vertex, and kaon charges.
- $^{5}$  HEISTER 02H measure simultaneously b and c quark forward-backward asymmetries using their semileptonic decays to tag the quark charge. The flavor separation is obtained with a discriminating multivariate analysis.
- <sup>6</sup> HEISTER 01D tag  $Z \rightarrow b\overline{b}$  events using the impact parameters of charged tracks complemented with information from displaced vertices, event shape variables, and lepton identification. The *b*-quark direction and charge is determined using the hemisphere charge method along with information from fast kaon tagging and charge estimators of primary and secondary vertices. The change in the quoted value due to variation of  $A_{FB}^c$  and  $R_b$  is given as +0.103 ( $A_{FB}^c 0.0651$ ) -0.440 ( $R_b 0.21585$ ).
- <sup>7</sup> ABREU 99Y tag  $Z \rightarrow b\overline{b}$  and  $Z \rightarrow c\overline{c}$  events by an exclusive reconstruction of several D meson decay modes ( $D^{*+}$ ,  $D^0$ , and  $D^+$  with their charge-conjugate states).
- <sup>8</sup> ACCIARRI 99D tag  $Z \to b \, \overline{b}$  events using high p and p<sub>T</sub> leptons. The analysis determines simultaneously a mixing parameter  $\chi_b = 0.1192 \pm 0.0068 \pm 0.0051$  which is used to correct the observed asymmetry.
- <sup>9</sup> ACCIARRI 980 tag  $Z \rightarrow b\overline{b}$  events using lifetime and measure the jet charge using the hemisphere charge.
- <sup>10</sup> ALEXANDER 97C identify the b and c events using a  $D/D^*$  tag.

### CHARGE ASYMMETRY IN $e^+e^- \rightarrow q \overline{q}$

Summed over five lighter flavors.

Experimental and Standard Model values are somewhat event-selection dependent. Standard Model expectations contain some assumptions on  $B^0 - \overline{B}{}^0$  mixing and on other electroweak parameters.

ASYMMETRY (%)	STD. MODEL	$\frac{\sqrt{s}}{(\text{GeV})}$	DOCUMENT ID		TECN
• • • We do not use the following	owing data for	averages, fi	its, limits, etc. • •	•	
$-\ 0.76\pm0.12\pm0.15$		91.2	<sup>1</sup> ABREU	921	DLPH
$4.0 \pm 0.4 \pm 0.63$	4.0	91.3	<sup>2</sup> ACTON	92L	OPAL
$9.1\ \pm 1.4\ \pm 1.6$	9.0	57.9	ADACHI	91	TOPZ
$-0.84\pm0.15\pm0.04$		91	DECAMP	<b>91</b> B	ALEP
$8.3 \pm 2.9 \pm 1.9$	8.7	56.6	STUART	90	AMY
$11.4 \pm 2.2 \pm 2.1$	8.7	57.6	ABE	89L	VNS
$6.0 \pm 1.3$	5.0	34.8	GREENSHAW	89	JADE
$8.2 \pm 2.9$	8.5	43.6	GREENSHAW	89	JADE

<sup>1</sup> ABREU 921 has 0.14 systematic error due to uncertainty of quark fragmentation.

### CHARGE ASYMMETRY IN $p\overline{p} \rightarrow Z \rightarrow e^+e^-$

ASYMMETRY (%)	STD. MODEL	$\frac{\sqrt{s}}{(\text{GeV})}$	DOCUMENT ID		TECN
• • • We do not use the follow	wing data for	averages, fits	s, limits, etc. •	• •	
$5.2 \!\pm\! 5.9 \!\pm\! 0.4$		91	ABE	91E	CDF

### ANOMALOUS $ZZ\gamma$ , $Z\gamma\gamma$ , AND ZZV COUPLINGS

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Combining the LEP results properly taking into account the correlations the following 95% CL limits are derived (CERN-PH-EP/2005-051 or hep-ex/0511027):

$$\begin{array}{ll} -0.13 < h_1^Z < +0.13, & -0.078 < h_2^Z < +0.071, \\ -0.20 < h_3^Z < +0.07, & -0.05 < h_4^Z < +0.12, \\ -0.056 < h_1^\gamma < +0.055, & -0.045 < h_2^\gamma < +0.025, \\ -0.049 < h_3^\gamma < -0.008, & -0.002 < h_4^\gamma < +0.034. \end{array}$$

VALUE <u>DOCUMENT ID</u> <u>TECN</u> <u>COMMENT</u>

• • • We do not use the following data for averages, fits, limits, etc. • • •

1 AAD
 12BX ATLS
 
$$E_{\rm cm}^{pp} = 7 \, {\rm TeV}$$

 2 ABAZOV
 12S
 D0
  $E_{\rm cm}^{p\overline{p}} = 1.96 \, {\rm TeV}$ 

 3 AALTONEN
 11S
 CDF
  $E_{\rm cm}^{p\overline{p}} = 1.96 \, {\rm TeV}$ 

 4 CHATRCHYAN 11M
 CMS
  $E_{\rm cm}^{p\overline{p}} = 7 \, {\rm TeV}$ 

 5 ABAZOV
 09L
 D0
  $E_{\rm cm}^{p\overline{p}} = 1.96 \, {\rm TeV}$ 

 6 ABAZOV
 07M
 D0
  $E_{\rm cm}^{p\overline{p}} = 1.96 \, {\rm TeV}$ 

 7 ABDALLAH
 07C
 DLPH
  $E_{\rm cm}^{ee} = 183-208 \, {\rm GeV}$ 

 8 ACHARD
 04H
 L3
  $E_{\rm cm}^{ee} = 183-208 \, {\rm GeV}$ 

 9 ABBIENDI,G
 00C
 OPAL
  $E_{\rm cm}^{ee} = 189 \, {\rm GeV}$ 

 10 ABBOTT
 98M
 D0
  $E_{\rm cm}^{p\overline{p}} = 1.8 \, {\rm TeV}$ 

 11 ABREU
 98K
 DLPH
  $E_{\rm cm}^{ee} = 161, 172 \, {\rm GeV}$ 

 $<sup>^2</sup>$  ACTON 92L use the weight function method on 259k selected  $Z \to \text{hadrons}$  events. The systematic error includes a contribution of 0.2 due to  $B^0$ - $\overline{B}{}^0$  mixing effect, 0.4 due to Monte Carlo (MC) fragmentation uncertainties and 0.3 due to MC statistics. ACTON 92L derive a value of  $\sin^2\!\theta_W^{\text{eff}}$  to be 0.2321  $\pm$  0.0017  $\pm$  0.0028.

 $<sup>^1</sup>$  AAD 12BX study  $Z\gamma$  production in pp collisions and select 142  $Z\gamma$  candidates where the Z decays to electron or muon pairs, and the photon has a transverse energy larger than 60 GeV. The expected background is 12.9  $\pm$  4.9 events. The resulting 95% C.L. limits ranges are:  $-0.028 < h_{\gamma}^3 < 0.027, \, -0.022 < h_{Z}^3 < 0.026, \, -0.00021 < h_{\gamma}^4 < 0.00021, \, -0.00022 < h_{Z}^4 < 0.00021.$ 

- $^2$  ABAZOV 12S study  $Z\gamma$  production in  $p\overline{p}$  collisions at  $\sqrt{s}=1.96$  TeV using 6.2 fb $^{-1}$  of data where the Z decays to electron (muon) pairs and the photon has at least 10 GeV of transverse momentum. In data, 304 (308) di-electron (di-muon) events are observed with an expected background of 255  $\pm$  16 (285  $\pm$  24) events. Based on the photon  $p_T$  spectrum, and including also earlier data and the  $Z\to\nu\overline{\nu}$  decay mode (from ABAZOV 09L), the following 95% C.L. limits are reported:  $\left|h_{03}^Z\right| < 0.026, \left|h_{04}^Z\right| < 0.0013, \left|h_{03}^{\gamma}\right| < 0.027, \left|h_{04}^{\gamma}\right| < 0.0014$  for a form factor scale of  $\Lambda=1.5$  TeV.
- <sup>3</sup> AALTONEN 11s study  $Z\gamma$  events in  $p\overline{p}$  interactions at  $\sqrt{s}=1.96$  TeV with integrated luminosity 5.1 fb $^{-1}$  for  $Z\to e^+e^-/\mu^+\mu^-$  and 4.9 fb $^{-1}$  for  $Z\to \nu\overline{\nu}$ . For the charged lepton case, the two leptons must be of the same flavor with the transverse momentum/energy of one >20 GeV and the other >10 GeV. The isolated photon must have  $E_T>50$  GeV. They observe 91 events with  $87.2\pm7.8$  events expected from standard model processes. For the  $\nu\overline{\nu}$  case they require solitary photons with  $E_T>25$  GeV and missing  $E_T>25$  GeV and observe 85 events with standard model expectation of  $85.9\pm5.6$  events. Taking the form factor  $\Lambda=1.5$  TeV they derive 95% C.L. limits as  $|h_3^{\gamma}, Z|<0.022$  and  $|h_4^{\gamma}, Z|<0.0009$ .
- $^4$  CHATRCHYAN 11M study  $Z\gamma$  production in pp collisions at  $\sqrt{s}=7$  TeV using  $36~{\rm pb}^{-1}$  pp data, where the Z decays to  $e^+\,e^-$  or  $\mu^+\mu^-$ . The total cross sections are measured for photon transverse energy  $E_T^\gamma>10$  GeV and spatial separation from charged leptons in the plane of pseudo rapidity and azimuthal angle  $\Delta R(\ell,\gamma)>0.7$  with the dilepton invariant mass requirement of  $M_{\ell\ell}>50$  GeV. The number of  $e^+e^-\gamma$  and  $\mu^+\mu^-\gamma$  candidates is 81 and 90 with estimated backgrounds of  $20.5\pm2.5$  and  $27.3\pm3.2$  events respectively. The 95% CL limits for  $ZZ\gamma$  couplings are -0.05<  $h_3^Z<0.06$  and -0.0005<  $h_4^Z<0.0005$ , and for  $Z\gamma\gamma$  couplings are -0.07<  $h_3^\gamma<0.07$  and -0.0005<  $h_4^\gamma<0.0006$ .
- $^5$  ABAZOV 09L study  $Z\gamma,~Z\to\nu\overline{\nu}$  production in  $p\overline{p}$  collisions at 1.96 TeV C.M. energy. They select 51 events with a photon of transverse energy  $E_T$  larger than 90 GeV, with an expected background of 17 events. Based on the photon  $E_T$  spectrum and including also Z decays to charged leptons (from ABAZOV 07M), the following 95% CL limits are reported:  $\left|h_{30}^{\gamma}\right|<0.033,~\left|h_{40}^{\gamma}\right|<0.0017,~\left|h_{30}^{Z}\right|<0.033,~\left|h_{40}^{Z}\right|<0.0017.$
- <sup>6</sup> ABAZOV 07M use 968  $p\overline{p} \to e^+e^-/\mu^+\mu^-\gamma X$  candidates, at 1.96 TeV center of mass energy, to tag  $p\overline{p} \to Z\gamma$  events by requiring  $E_T(\gamma)>$  7 GeV, lepton-gamma separation  $\Delta R_{\ell\gamma}>0.7$ , and di-lepton invariant mass > 30 GeV. The cross section is in agreement with the SM prediction. Using these  $Z\gamma$  events they obtain 95% C.L. limits on each  $h_i^V$ , keeping all others fixed at their SM values. They report:  $-0.083 < h_{30}^Z < 0.082$ ,  $-0.0053 < h_{40}^Z < 0.0054$ ,  $-0.085 < h_{30}^\gamma < 0.084$ ,  $-0.0053 < h_{40}^\gamma < 0.0054$ , for the form factor scale Λ= 1.2 TeV.
- 7 Using data collected at  $\sqrt{s}=183$ –208, ABDALLAH 07C select 1,877  $e^+e^- \to Z\gamma$  events with  $Z \to q\overline{q}$  or  $\nu\overline{\nu}$ , 171  $e^+e^- \to ZZ$  events with  $Z \to q\overline{q}$  or lepton pair (except an explicit  $\tau$  pair), and 74  $e^+e^- \to Z\gamma^*$  events with a  $q\overline{q}\mu^+\mu^-$  or  $q\overline{q}e^+e^-$  signature, to derive 95% CL limits on  $h_i^V$ . Each limit is derived with other parameters set to zero. They report:  $-0.23 < h_1^Z < 0.23$ ,  $-0.30 < h_3^Z < 0.16$ ,  $-0.14 < h_1^\gamma < 0.14$ ,  $-0.049 < h_3^\gamma < 0.044$ .
- $^{8}$  ACHARD 04H select 3515  $e^{+}\,e^{-}\to Z\gamma$  events with  $Z\to q\,\overline{q}$  or  $\nu\overline{\nu}$  at  $\sqrt{s}=189$ –209 GeV to derive 95% CL limits on  $h_{i}^{V}$  . For deriving each limit the other parameters are fixed at zero. They report:  $-0.153 < h_{1}^{Z} < 0.141, \, -0.087 < h_{2}^{Z} < 0.079, \, -0.220 < h_{3}^{Z} < 0.112, \, -0.068 < h_{4}^{Z} < 0.148, \, -0.057 < h_{1}^{\gamma} < 0.057, \, -0.050 < h_{2}^{\gamma} < 0.023, \, -0.059 < h_{3}^{\gamma} < 0.004, \, -0.004 < h_{4}^{\gamma} < 0.042.$

- <sup>9</sup>ABBIENDI,G 00C study  $e^+e^- \rightarrow Z\gamma$  events (with  $Z \rightarrow q\overline{q}$  and  $Z \rightarrow \nu\overline{\nu}$ ) at 189 GeV to obtain the central values (and 95% CL limits) of these couplings:  $h_1^Z=0.000\pm0.100~(-0.190,0.190),~h_2^Z=0.000\pm0.068~(-0.128,0.128),~h_3^Z=-0.074^{+0.102}_{-0.103}~(-0.269,0.119),~h_4^Z=0.046\pm0.068~(-0.084,0.175),~h_1^{\gamma}=0.000\pm0.061~(-0.115,0.115),~h_2^{\gamma}=0.000\pm0.041~(-0.077,0.077),~h_3^{\gamma}=-0.080^{+0.039}_{-0.041}~(-0.164,-0.006),~h_4^{\gamma}=0.064^{+0.033}_{-0.030}~(+0.007,+0.134).~$  The results are derived assuming that only one coupling at a time is different from zero.
- <sup>10</sup> ABBOTT 98M study  $p \, \overline{p} \to Z \, \gamma + X$ , with  $Z \to e^+ \, e^-$ ,  $\mu^+ \, \mu^-$ ,  $\overline{\nu} \, \nu$  at 1.8 TeV, to obtain 95% CL limits at  $\Lambda = 750$  GeV:  $|h_{30}^Z| < 0.36$ ,  $|h_{40}^Z| < 0.05$  (keeping  $h_i^{\gamma} = 0$ ), and  $|h_{30}^{\gamma}| < 0.37$ ,  $|h_{40}^{\gamma}| < 0.05$  (keeping  $h_i^{\gamma} = 0$ ). Limits on the *CP*-violating couplings are  $|h_{10}^Z| < 0.36$ ,  $|h_{20}^Z| < 0.05$  (keeping  $h_i^{\gamma} = 0$ ), and  $|h_{10}^{\gamma}| < 0.37$ ,  $|h_{20}^{\gamma}| < 0.05$  (keeping  $h_i^{\gamma} = 0$ ).
- <sup>11</sup> ABREU 98K determine a 95% CL upper limit on  $\sigma(e^+e^- \to \gamma + \text{invisible particles}) < 2.5 pb using 161 and 172 GeV data. This is used to set 95% CL limits on <math>|h_{30}^{\gamma}| < 0.8$  and  $|h_{30}^{Z}| < 1.3$ , derived at a scale Λ=1 TeV and with n=3 in the form factor representation.



Combining the LEP results properly taking into account the correlations the following 95% CL limits are derived (CERN-PH-EP/2005-051 or hep-ex/0511027):

$$-0.30 < f_4^Z < +0.30,$$
  $-0.34 < f_5^Z < +0.38,$   $-0.17 < f_4^{\gamma} < +0.19,$   $-0.32 < f_5^{\gamma} < +0.36.$ 

**VALUE** 

DOCUMENT ID TECN COMMENT

• • • We do not use the following data for averages, fits, limits, etc. • • •

$$^{1}$$
 AAD  $^{1}$  13Z ATLS  $E_{\rm cm}^{pp}=7~{\rm TeV}$   $^{2}$  CHATRCHYAN 13B CMS  $E_{\rm cm}^{pp}=7~{\rm TeV}$   $^{3}$  SCHAEL 09 ALEP  $E_{\rm cm}^{ee}=192$ –209 GeV  $^{4}$  ABAZOV 08K D0  $E_{\rm cm}^{p\overline{p}}=1.96~{\rm TeV}$   $^{5}$  ABDALLAH 07C DLPH  $E_{\rm cm}^{ee}=183$ –208 GeV  $^{6}$  ABBIENDI 04C OPAL  $^{7}$  ACHARD 03D L3

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<sup>1</sup>AAD 13Z study ZZ production in pp collisions at  $\sqrt{s}=7$  TeV. In the  $ZZ\to \ell^+\ell^-\ell'^+\ell'^-$  final state they observe a total of 66 events with an expected background of  $0.9\pm1.3$ . In the  $ZZ\to \ell^+\ell^-\nu\nu$  final state they observe a total of 87 events with an expected background of  $46.9\pm5.2$ . The limits on anomalous TGCs are determined using the observed and expected numbers of these ZZ events binned in  $p_T^Z$ . The 95% C.L. are as follows: for form factor scale  $\Lambda=\infty$ , -0.015<  $f_4^\gamma<0.015$ , -0.013<  $f_4^Z<0.013$ , -0.016<  $f_5^\gamma<0.015$ , -0.013,  $f_5^Z<0.013$ ; for form factor scale  $\Lambda=3$  TeV, -0.022<  $f_4^\gamma<0.023$ , -0.019.

- $^2$ CHATRCHYAN 13B study ZZ production in pp collisions and select 54 ZZ candidates in the Z decay channel with electrons or muons with an expected background of  $1.4 \pm 0.5$ events. The resulting 95% C.L. ranges are:  $-0.013~<~f_{f 4}^{\gamma}~<$  0.015,  $-0.011~<~f_{f 4}^{Z}$  $0.012,\, -0.014 \; < \; f_{\, 5}^{\, \gamma} \; < 0.014, \, -0.012 \; < \; f_{\, 5}^{\, Z} \; < 0.012.$
- $^3$  Using data collected in the center of mass energy range 192–209 GeV, SCHAEL 09 select 318  $e^+e^- \rightarrow ZZ$  events with 319.4 expected from the standard model. Using this data they derive the following 95% CL limits:  $-0.321 < f_A^{\gamma} < 0.318, -0.534 < f_A^{Z} < 0.318$ 0.534,  $-0.724 < f_{\overline{5}}^{\gamma} < 0.733$ ,  $-1.194 < f_{\overline{5}}^{Z} < 1.190$ .
- $^4$  ABAZOV 08K search for ZZ and  $Z\gamma^*$  events with 1 fb $^{-1}$   $p\overline{p}$  data at  $\sqrt{s}=1.96$  TeV in (ee)(ee),  $(\mu\mu)(\mu\mu)$ ,  $(ee)(\mu\mu)$  final states requiring the lepton pair masses to be > 30GeV. They observe 1 event, which is consistent with an expected signal of 1.71  $\pm$  0.15 events and a background of 0.13  $\pm$  0.03 events. From this they derive the following limits, for a form factor (A) value of 1.2 TeV:  $-0.28 < f_{40}^Z < 0.28, \, -0.31 < f_{50}^Z < 0.28$  $0.29, -0.26 < f_{40}^{\gamma} < 0.26, -0.30 < f_{50}^{\gamma} < 0.28.$
- $^{5}$  Using data collected at  $\sqrt{s}=$  183–208 GeV, ABDALLAH 07C select 171  $e^{+}e^{-} 
  ightarrow ZZ$ events with  $Z \to q \overline{q}$  or lepton pair (except an explicit  $\tau$  pair), and 74  $e^+e^- \to Z\gamma^*$ events with a  $q \overline{q} \mu^+ \mu^-$  or  $q \overline{q} e^+ e^-$  signature, to derive 95% CL limits on  $f_i^V$ . Each limit is derived with other parameters set to zero. They report:  $-0.40 < f_{\Delta}^{Z'} < 0.42$ ,  $-0.38 < f_{5}^{Z} < 0.62, -0.23 < f_{4}^{\gamma} < 0.25, -0.52 < f_{5}^{\gamma} < 0.48.$
- $^6$  ABBIENDI 04C study ZZ production in  $e^+e^-$  collisions in the C.M. energy range 190–209 GeV. They select 340 events with an expected background of 180 events. Including the ABBIENDI 00N data at 183 and 189 GeV (118 events with an expected background of 65 events) they report the following 95% CL limits:  $-0.45 < f_4^Z < 0.58$ , -0.94 <  $f_{5}^{Z} < 0.25, \, -0.32 <$   $f_{4}^{\gamma} < 0.33, \, ext{and} \, -0.71 <$   $f_{5}^{\gamma} < 0.59.$
- $^7$ ACHARD 03D study Z-boson pair production in  $e^+e^-$  collisions in the C.M. energy range 200-209 GeV. They select 549 events with an expected background of 432 events. Including the ACCIARRI 99G and ACCIARRI 99O data (183 and 189 GeV respectively, 286 events with an expected background of 241 events) and the 192-202 GeV ACCIARRI 011 results (656 events, expected background of 512 events), they report the following 95% CL limits:  $-0.48 \le f_4^Z \le 0.46$ ,  $-0.36 \le f_5^Z \le 1.03$ ,  $-0.28 \le f_4^{\gamma} \le 0.28$ , and  $-0.40 \le 0.28$  $f_{\rm E}^{\gamma} \leq 0.47.$

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 $a_0/\Lambda^2$ ,  $a_c/\Lambda^2$ 

Combining published and unpublished preliminary LEP results the following 95% CL intervals for the QGCs associated with the  $ZZ\gamma\gamma$  vertex are derived (CERN-PH-EP/2005-051 or hep-ex/0511027):

$$-0.008 < a_0^Z/\Lambda^2 < +0.021$$
$$-0.029 < a_0^Z/\Lambda^2 < +0.039$$

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- ullet ullet We do not use the following data for averages, fits, limits, etc. ullet
  - <sup>1</sup> ABBIENDI 04L OPAL <sup>2</sup> HEISTER 04A ALEP
  - <sup>3</sup> ACHARD 02G L3

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- ABBIENDI 04L select 20  $e^+e^- \rightarrow \nu \overline{\nu} \gamma \gamma$  acoplanar events in the energy range 180–209 GeV and 176  $e^+e^- \rightarrow q \overline{q} \gamma \gamma$  events in the energy range 130–209 GeV. These samples are used to constrain possible anomalous  $W^+W^-\gamma \gamma$  and  $ZZ\gamma\gamma$  quartic couplings. Further combining with the  $W^+W^-\gamma$  sample of ABBIENDI 04B the following one-parameter 95% CL limits are obtained:  $-0.007 < a_0^Z/\Lambda^2 < 0.023 \ {\rm GeV}^{-2}, -0.029 < a_0^Z/\Lambda^2 < 0.029 \ {\rm GeV}^{-2}, -0.020 < a_0^W/\Lambda^2 < 0.020 \ {\rm GeV}^{-2}, -0.052 < a_0^W/\Lambda^2 < 0.037 \ {\rm GeV}^{-2}$
- In the CM energy range 183 to 209 GeV HEISTER 04A select  $30~e^+e^- \rightarrow \nu \overline{\nu} \gamma \gamma$  events with two acoplanar, high energy and high transverse momentum photons. The photon-photon acoplanarity is required to be > 5°,  $E_{\gamma}/\sqrt{s}>$  0.025 (the more energetic photon having energy > 0.2  $\sqrt{s}$ ), p $_{T_{\gamma}}/E_{\rm beam}>$  0.05 and  $|\cos\theta_{\gamma}|<$  0.94. A likelihood fit to the photon energy and recoil missing mass yields the following one–parameter 95% CL limits:  $-0.012 < a_0^Z/\Lambda^2 < 0.019~{\rm GeV}^{-2}$ ,  $-0.041 < a_c^Z/\Lambda^2 < 0.044~{\rm GeV}^{-2}$ ,  $-0.060 < a_0^W/\Lambda^2 < 0.055~{\rm GeV}^{-2}$ ,  $-0.099 < a_c^W/\Lambda^2 < 0.093~{\rm GeV}^{-2}$ .
- $^3$  ACHARD 02G study  $e^+\,e^-\to Z\gamma\gamma\to q\overline{q}\gamma\gamma$  events using data at center-of-mass energies from 200 to 209 GeV. The photons are required to be isolated, each with energy >5 GeV and  $|\cos\theta|<0.97$ , and the di-jet invariant mass to be compatible with that of the Z boson (74–111 GeV). Cuts on Z velocity ( $\beta<0.73$ ) and on the energy of the most energetic photon reduce the backgrounds due to non-resonant production of the  $q\overline{q}\gamma\gamma$  state and due to ISR respectively, yielding a total of 40 candidate events of which 8.6 are expected to be due to background. The energy spectra of the least energetic photon are fitted for all ten center-of-mass energy values from 130 GeV to 209 GeV (as obtained adding to the present analysis 130–202 GeV data of ACCIARRI 01E, for a total of 137 events with an expected background of 34.1 events) to obtain the fitted values  $a_0/\Lambda^2=0.00^{+0.02}_{-0.01}$  GeV $^{-2}$  and  $a_c/\Lambda^2=0.03^{+0.01}_{-0.02}$  GeV $^{-2}$ , where the other parameter is kept fixed to its Standard Model value (0). A simultaneous fit to both parameters yields the 95% CL limits -0.02 GeV $^{-2}< a_0/\Lambda^2<0.03$  GeV $^{-2}$  and -0.07 GeV $^{-2}< a_c/\Lambda^2<0.05$  GeV $^{-2}$ .

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BUSKULIC   93L   PL B313 520   D. Buskulic et al.   Okun, M.I. Vysotsky   (TTEP)					
NOVIKOV   93C   PL 8298 453   V.A. Novikov, L.B. Okun, M.I. Vysotsky   (TTEP)   ABREU   92M   PL 8289 199   P. Abreu et al.   (DELPHI Collab.)   ACTON   92B   ZPHY C53 539   D.P. Acton et al.   (OPAL Collab.)   ACTON   92B   ZPHY C53 539   P.P. Acton et al.   (OPAL Collab.)   ACTON   92N   PL 8294 436   P.D. Acton et al.   (OPAL Collab.)   ACTON   92N   PL 8295 357   P.D. Acton et al.   (OPAL Collab.)   ADEVA   92   PL 8275 299   B. Adeva et al.   (L3 Collab.)   ADEVA   92   PL 8276 354   D. Acton et al.   (UA2 Collab.)   ALITTI   92B   PL 8276 354   D. Buskulic et al.   (L4 Collab.)   BUSKULIC   92D   PL 8292 210   D. Buskulic et al.   (ALEPH Collab.)   BUSKULIC   92E   PL 8294 145   D. Buskulic et al.   (ALEPH Collab.)   DECAMP   92   PRPL 216 253   D. Decamp et al.   (ALEPH Collab.)   ABREU   91H   ZPHY C50 185   P. Abreu et al.   (ALEPH Collab.)   ACTON   91B   PL 8273 338   D.P. Acton et al.   (OPAL Collab.)   ACTON   91B   PL 8273 338   D.P. Acton et al.   (OPAL Collab.)   ACTON   91B   PL 8259 513   I. Addachi et al.   (DELPHI Collab.)   ADEVAN   911   PL 8259 513   M.Z. Akrawy et al.   (OPAL Collab.)   DECAMP   91B   PL 8277 531   M.Z. Akrawy et al.   (OPAL Collab.)   DECAMP   91B   PL 8277 537   D. Decamp et al.   (ALEPH Collab.)   ACOSSEN   91   PRL 67 3347   R.G. Jacobsen et al.   (ALEPH Collab.)   ACOSSEN   91   PRL 67 3347   R.G. Jacobsen et al.   (ALEPH Collab.)   ACOSSEN   91   PRL 67 3347   R.G. Jacobsen et al.   (Mark II Collab.)   ACOSSEN   91   PRL 67 3347   R.G. Jacobsen et al.   (Mark II Collab.)   ACOSSEN   91   PRL 67 3347   R.G. Jacobsen et al.   (Mark II Collab.)   ACOSSEN   91   PRL 67 3347   R.G. Jacobsen et al.   (Mark II Collab.)   ACOSSEN   91   PRL 67 3347   R.G. Jacobsen et al.   (Mark II Collab.)   ACOSSEN   91   PRL 67 3347   R.G. Jacobsen et al.   (Mark II Collab.)   ACOSSEN   91   PRL 67 3347   R.G. Jacobsen et al.   (Mark II Collab.)   ACOSSEN   91   PRL 67 3347   R.G. Jacobsen et al.   (Mark II Collab.)   ACOSSEN   91   PRL 67 3347   R.G. Jacobsen et al.   (Mark II Collab					
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AKRAWY         90J         PL B246 285         M.Z. Akrawy et al.         (OPAL Collab.)           BEHREND         90D         ZPHY C47 3333         H.J. Behrend et al.         (CELLO Collab.)           BRAUNSCH         90         ZPHY C48 433         W. Braunschweig et al.         (TASSO Collab.)           ELSEN         90         ZPHY C46 547         S. Hegner et al.         (JADE Collab.)           HEGNER         90         ZPHY C46 547         S. Hegner et al.         (JADE Collab.)           STUART         90         PRL 64 983         D. Stuart et al.         (AMY Collab.)           ABE         89         PRL 63 720         F. Abe et al.         (CDF Collab.)           ABE         89C         PRL 63 720         F. Abe et al.         (VENUS Collab.)           ABRAMS         89B         PRL 63 2173         G.S. Abrams et al.         (Mark II Collab.)           ABRAMS         89B         PRL 63 2780         G.S. Abrams et al.         (Mark II Collab.)           ALBAJAR         89         PL B218 112         A. Bacala et al.         (Mark II Collab.)           BAND         89         PL B218 369         H.R. Band et al.         (MAC Collab.)           GREENSHAW         89         PPH C42 1         T. Greenshaw et al.	ABE	901	ZPHY C48 13	K. Abe <i>et al.</i>	(VENUS Collab.)
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BEHREND         90D         ZPHY C47 333         H.J. Behrend et al.         (ČELLO Collab.)           BRAUNSCH         90         ZPHY C48 433         W. Braunschweig et al.         (TASSO Collab.)           ELSEN         90         ZPHY C46 547         S. Hegner et al.         (JADE Collab.)           HEGNER         90         ZPHY C46 547         S. Hegner et al.         (JADE Collab.)           STUART         90         PRL 64 983         D. Stuart et al.         (AMY Collab.)           ABE         89         PRL 63 720         F. Abe et al.         (CDF Collab.)           ABE         89C         PRL 63 720         F. Abe et al.         (VENUS Collab.)           ABRAMS         89B         PRL 63 2173         G.S. Abrams et al.         (Mark II Collab.)           ABRAMS         89D         PRL 63 2780         G.S. Abrams et al.         (Mark II Collab.)           ALBAJAR         89         PRL 63 2780         G.S. Abrams et al.         (Mark II Collab.)           BACALA         89         PL B218 112         A. Bacala et al.         (Mark II Collab.)           BACALA         89         PL B218 369         H.R. Band et al.         (MAC Collab.)           GREENSHAW         89         ZPHY C42 1         T. Greenshaw et al.	AKRAWY	90.J	PL B246 285		`
BRAUNSCH         90         ZPHY C48 433         W. Braunschweig et al.         (TASSO Collab.)           ELSEN         90         ZPHY C46 349         E. Elsen et al.         (JADE Collab.)           HEGNER         90         ZPHY C46 547         S. Hegner et al.         (JADE Collab.)           STUART         90         PRL 64 983         D. Stuart et al.         (AMY Collab.)           ABE         89         PRL 62 613         F. Abe et al.         (CDF Collab.)           ABE         89C         PRL 63 720         F. Abe et al.         (CDF Collab.)           ABE         89L         PL B232 425         K. Abe et al.         (VENUS Collab.)           ABRAMS         89B         PRL 63 2173         G.S. Abrams et al.         (Mark II Collab.)           ABRAMS         89D         PRL 63 2780         G.S. Abrams et al.         (Mark II Collab.)           ALBAJAR         89         ZPHY C44 15         C. Albajar et al.         (UAI Collab.)           BAND         89         PL B218 369         H.R. Band et al.         (MAC Collab.)           GREENSHAW         89         ZPHY C42 1         T. Greenshaw et al.         (JADE Collab.)           SAGAWA         89         PRL 63 2341         H. Sagawa et al.         (MAC Collab.)				,	
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HEGNER         90         ZPHY C46 547         S. Hegner et al.         (JADE Collab.)           STUART         90         PRL 64 983         D. Stuart et al.         (AMY Collab.)           ABE         89         PRL 62 613         F. Abe et al.         (CDF Collab.)           ABE         89C         PRL 63 720         F. Abe et al.         (VENUS Collab.)           ABE         89L         PL B232 425         K. Abe et al.         (VENUS Collab.)           ABRAMS         89B         PRL 63 2173         G.S. Abrams et al.         (Mark II Collab.)           ABRAMS         89D         PRL 63 2780         G.S. Abrams et al.         (Mark II Collab.)           ABRAMS         89D         PRL 63 2780         G.S. Abrams et al.         (Mark II Collab.)           ALBAJAR         89         ZPHY C44 15         C. Albajar et al.         (UA1 Collab.)           BACALA         89         PL B218 112         A. Bacala et al.         (AMY Collab.)           BAND         89         PL B218 369         H.R. Band et al.         (MAC Collab.)           GREENSHAW         90         ZPHY C44 167         F. Ould-Saada et al.         (JADE Collab.)           SAGAWA         89         PRL 63 2341         H. Sagawa et al.         (Mary Collab.)<					(IADE Collab.)
STUART         90         PRL 64 983         D. Stuart et al.         (AMY Collab.)           ABE         89         PRL 62 613         F. Abe et al.         (CDF Collab.)           ABE         89C         PRL 63 720         F. Abe et al.         (VENUS Collab.)           ABE         89L         PL B232 425         K. Abe et al.         (VENUS Collab.)           ABRAMS         89B         PRL 63 2173         G.S. Abrams et al.         (Mark II Collab.)           ABRAMS         89D         PRL 63 2780         G.S. Abrams et al.         (Mark II Collab.)           ALBAJAR         89         ZPHY C44 15         C. Albajar et al.         (UA1 Collab.)           BACALA         89         PL B218 369         H.R. Band et al.         (MAC Collab.)           GREENSHAW         89         PLB 218 369         H.R. Band et al.         (JADE Collab.)           OULD-SAADA         89         ZPHY C42 1         T. Greenshaw et al.         (JADE Collab.)           OULD-SAADA         89         PRL 63 2341         H. Sagawa et al.         (AMY Collab.)           ADACHI         88C         PL B208 319         I. Adachi et al.         (TOPAZ Collab.)           ADEVA         88         PR D38 2665         B. Adeva et al.         (Mark.)					
ABE         89         PRL 62 613         F. Abe et al.         (CDF Collab.)           ABE         89C         PRL 63 720         F. Abe et al.         (CDF Collab.)           ABE         89L         PL B232 425         K. Abe et al.         (VENUS Collab.)           ABRAMS         89B         PRL 63 2173         G.S. Abrams et al.         (Mark II Collab.)           ABRAMS         89D         PRL 63 2780         G.S. Abrams et al.         (Mark II Collab.)           ABRAMS         89D         PRL 63 2780         G.S. Abrams et al.         (Mark II Collab.)           ALBAJAR         89         PL B218 112         A. Bacala et al.         (UAI Collab.)           BACALA         89         PL B218 369         H.R. Band et al.         (MAC Collab.)           GREENSHAW         89         PLPHY C42 1         T. Greenshaw et al.         (JADE Collab.)           OULD-SAADA         89         ZPHY C44 567         F. Ould-Saada et al.         (JADE Collab.)           SAGAWA         89         PRL 63 2341         H. Sagawa et al.         (AMY Collab.)           ADACHI         88C         PL B208 319         I. Adachi et al.         (TOPAZ Collab.)           ADEVA         88         PR D38 2665         B. Adeva et al.         (Mark.)<				<u> </u>	
ABE         89C         PRL 63 720         F. Abe et al.         (CDF Collab.)           ABE         89L         PL B232 425         K. Abe et al.         (VENUS Collab.)           ABRAMS         89B         PRL 63 2173         G.S. Abrams et al.         (Mark II Collab.)           ABRAMS         89D         PRL 63 2780         G.S. Abrams et al.         (Mark II Collab.)           ALBAJAR         89         PRL 63 2780         G.S. Abrams et al.         (UA1 Collab.)           ALBAJAR         89         PRL 63 2780         G.S. Abrams et al.         (UA1 Collab.)           BACALA         89         PL B218 112         A. Bacala et al.         (UAY Collab.)           BAND         89         PL B218 369         H.R. Band et al.         (MAC Collab.)           GREENSHAW         89         ZPHY C42 1         T. Greenshaw et al.         (JADE Collab.)           OULD-SAADA         89         ZPHY C44 567         F. Ould-Saada et al.         (JADE Collab.)           SAGAWA         89         PRL 63 2341         H. Sagawa et al.         (AMY Collab.)           ADACHI         88C         PL B208 319         I. Adachi et al.         (TOPAZ Collab.)           ADEVA         88         PR D38 2665         B. Adeva et al.         (Mark-					
ABE         89L         PL B232 425         K. Abe et al.         (VÈNUS Collab.)           ABRAMS         89B         PRL 63 2173         G.S. Abrams et al.         (Mark II Collab.)           ABRAMS         89D         PRL 63 2780         G.S. Abrams et al.         (Mark II Collab.)           ALBAJAR         89         PPH C44 15         C. Albajar et al.         (UA1 Collab.)           BACALA         89         PL B218 112         A. Bacala et al.         (AMY Collab.)           BAND         89         PL B218 369         H.R. Band et al.         (MAC Collab.)           GREENSHAW         89         ZPHY C42 1         T. Greenshaw et al.         (JADE Collab.)           OULD-SAADA         89         ZPHY C44 567         F. Ould-Saada et al.         (JADE Collab.)           SAGAWA         89         PRL 63 2341         H. Sagawa et al.         (AMY Collab.)           ADACHI         88C         PL B208 319         I. Adachi et al.         (TOPAZ Collab.)           ADEVA         88         PR D38 2665         B. Adeva et al.         (Mark-J Collab.)           BRAUNSCH         88D         ZPHY C40 163         W. Braunschweig et al.         (TASSO Collab.)           BEHREND         87         P B 1919 209         H.J. Behrend et al.					
ABRAMS         89B         PRL 63 2173         G.S. Abrams et al.         (Mark II Collab.)           ABRAMS         89D         PRL 63 2780         G.S. Abrams et al.         (Mark II Collab.)           ALBAJAR         89         ZPHY C44 15         C. Albajar et al.         (UA1 Collab.)           BACALA         89         PL B218 112         A. Bacala et al.         (AMY Collab.)           BAND         89         PL B218 369         H.R. Band et al.         (MAC Collab.)           GREENSHAW         89         PL PY C42 1         T. Greenshaw et al.         (JADE Collab.)           OULD-SAADA         89         ZPHY C44 567         F. Ould-Saada et al.         (JADE Collab.)           SAGAWA         89         PRL 63 2341         H. Sagawa et al.         (AMY Collab.)           ADACHI         88C         PL B208 319         I. Adachi et al.         (TOPAZ Collab.)           ADEVA         88         PR D38 2665         B. Adeva et al.         (Mark-J Collab.)           BRAUNSCH         88D         ZPHY C40 163         W. Braunschweig et al.         (TASSO Collab.)           BEHREND         87C         PL B186 440         R. Ansari et al.         (UA2 Collab.)           BARTEL         86C         ZPHY C30 371         W. Bartel et					
ABRAMS       89D       PRL 63 2780       G.S. Abrams et al.       (Mark II Collab.)         ALBAJAR       89       ZPHY C44 15       C. Albajar et al.       (UA1 Collab.)         BACALA       89       PL B218 112       A. Bacala et al.       (AMY Collab.)         BAND       89       PL B218 369       H.R. Band et al.       (MAC Collab.)         GREENSHAW       89       PPHY C42 1       T. Greenshaw et al.       (JADE Collab.)         OULD-SAADA       89       ZPHY C44 567       F. Ould-Saada et al.       (JADE Collab.)         SAGAWA       89       PRL 63 2341       H. Sagawa et al.       (AMY Collab.)         ADACHI       88C       PL B208 319       I. Adachi et al.       (TOPAZ Collab.)         ADEVA       88       PR D38 2665       B. Adeva et al.       (Mark-J Collab.)         BRAUNSCH       88D       ZPHY C40 163       W. Braunschweig et al.       (TASSO Collab.)         ANSARI       87       PL B186 440       R. Ansari et al.       (UA2 Collab.)         BEHREND       87C       PL B191 209       H.J. Behrend et al.       (CELLO Collab.)         Also       ZPHY C26 507       W. Bartel et al.       (JADE Collab.)         ASH       85       PRL 55 1831       W.W.V. As	ABE	89L			(VENUS Collab.)
ALBAJAR       89       ZPHY C44 15       C. Albajar et al.       (UA1 Collab.)         BACALA       89       PL B218 112       A. Bacala et al.       (AMY Collab.)         BAND       89       PL B218 369       H.R. Band et al.       (MAC Collab.)         GREENSHAW       89       PL B218 369       H.R. Band et al.       (MAC Collab.)         OULD-SAADA       89       ZPHY C42 1       T. Greenshaw et al.       (JADE Collab.)         SAGAWA       89       PRL 63 2341       H. Sagawa et al.       (AMY Collab.)         ADACHI       88C       PL B208 319       I. Adachi et al.       (TOPAZ Collab.)         ADEVA       88       PR D38 2665       B. Adeva et al.       (Mark-J Collab.)         BRAUNSCH       88D       ZPHY C40 163       W. Braunschweig et al.       (TASSO Collab.)         ANSARI       87       PL B186 440       R. Ansari et al.       (UA2 Collab.)         BEHREND       87C       PL B191 209       H.J. Behrend et al.       (CELLO Collab.)         BARTEL       86C       ZPHY C30 371       W. Bartel et al.       (JADE Collab.)         Also       ZPHY C26 507       W. Bartel et al.       (JADE Collab.)         ASH       85       PRL 55 1831       W.W. Ash et al. <td>ABRAMS</td> <td>89B</td> <td>PRL 63 2173</td> <td></td> <td>(Mark II Collab.)</td>	ABRAMS	89B	PRL 63 2173		(Mark II Collab.)
BACALA         89         PL B218 112         A. Bacala et al.         (AMY Collab.)           BAND         89         PL B218 369         H.R. Band et al.         (MAC Collab.)           GREENSHAW         89         ZPHY C42 1         T. Greenshaw et al.         (JADE Collab.)           OULD-SAADA         89         ZPHY C44 567         F. Ould-Saada et al.         (JADE Collab.)           SAGAWA         89         PRL 63 2341         H. Sagawa et al.         (AMY Collab.)           ADCHI         88C         PL B208 319         I. Adachi et al.         (TOPAZ Collab.)           ADEVA         88         PR D38 2665         B. Adeva et al.         (Mark-J Collab.)           BRAUNSCH         88D         ZPHY C40 163         W. Braunschweig et al.         (TASSO Collab.)           ANSARI         87         PL B186 440         R. Ansari et al.         (UA2 Collab.)           BEHREND         87C         PL B191 209         H.J. Behrend et al.         (CELLO Collab.)           BARTEL         86C         ZPHY C30 371         W. Bartel et al.         (JADE Collab.)           Also         ZPHY C26 507         W. Bartel et al.         (JADE Collab.)           ASH         85         PRL 55 1831         W.W. Ash et al.         (MAC Collab.)<	ABRAMS	89D	PRL 63 2780	G.S. Abrams et al.	(Mark II Collab.)
BACALA         89         PL B218 112         A. Bacala et al.         (AMY Collab.)           BAND         89         PL B218 369         H.R. Band et al.         (MAC Collab.)           GREENSHAW         89         ZPHY C42 1         T. Greenshaw et al.         (JADE Collab.)           OULD-SAADA         89         ZPHY C44 567         F. Ould-Saada et al.         (JADE Collab.)           SAGAWA         89         PRL 63 2341         H. Sagawa et al.         (AMY Collab.)           ADCHI         88C         PL B208 319         I. Adachi et al.         (TOPAZ Collab.)           ADEVA         88         PR D38 2665         B. Adeva et al.         (Mark-J Collab.)           BRAUNSCH         88D         ZPHY C40 163         W. Braunschweig et al.         (TASSO Collab.)           ANSARI         87         PL B186 440         R. Ansari et al.         (UA2 Collab.)           BEHREND         87C         PL B191 209         H.J. Behrend et al.         (CELLO Collab.)           BARTEL         86C         ZPHY C30 371         W. Bartel et al.         (JADE Collab.)           Also         ZPHY C26 507         W. Bartel et al.         (JADE Collab.)           ASH         85         PRL 55 1831         W.W. Ash et al.         (MAC Collab.)<	ALBAJAR	89	ZPHY C44 15	C. Albajar <i>et al.</i>	(UA1 Collab.)
BAND         89         PL         B218         369         H.R. Band et al.         (MAC Collab.)           GREENSHAW         89         ZPHY C42         1         T. Greenshaw et al.         (JADE Collab.)           OULD-SAADA         89         ZPHY C44         567         F. Ould-Saada et al.         (JADE Collab.)           SAGAWA         89         PRL 63         2341         H. Sagawa et al.         (AMY Collab.)           ADACHI         88C         PL B208         319         I. Adachi et al.         (TOPAZ Collab.)           ADEVA         88         PR D38         2665         B. Adeva et al.         (Mark-J Collab.)           BRAUNSCH         88D         ZPHY C40         163         W. Braunschweig et al.         (TASSO Collab.)           ANSARI         87         PL B186         440         R. Ansari et al.         (UA2 Collab.)           BEHREND         87C         PL B191         209         H.J. Behrend et al.         (CELLO Collab.)           BARTEL         86C         ZPHY C30         371         W. Bartel et al.         (JADE Collab.)           Also         ZPHY C26         507         W. Bartel et al.         (JADE Collab.)           ASH         85         PRL 55         1831	BACALA	89	PL B218 112		
GREENSHAW         89         ZPHY C42 1         T. Greenshaw et al.         (JADE Collab.)           OULD-SAADA         89         ZPHY C44 567         F. Ould-Saada et al.         (JADE Collab.)           SAGAWA         89         PRL 63 2341         H. Sagawa et al.         (AMY Collab.)           ADACHI         88C         PL B208 319         I. Adachi et al.         (TOPAZ Collab.)           ADEVA         88         PR D38 2665         B. Adeva et al.         (Mark-J Collab.)           BRAUNSCH         88D         ZPHY C40 163         W. Braunschweig et al.         (TASSO Collab.)           ANSARI         87         PL B186 440         R. Ansari et al.         (UA2 Collab.)           BEHREND         87C         PL B191 209         H.J. Behrend et al.         (CELLO Collab.)           BARTEL         86C         ZPHY C30 371         W. Bartel et al.         (JADE Collab.)           Also         ZPHY C26 507         W. Bartel et al.         (JADE Collab.)           ASH         85         PRL 55 1831         W.W. Ash et al.         (MAC Collab.)           BARTEL         85F         PL 161B 188         W. Bartel et al.         (JADE Collab.)           BARTEL         85F         PR D31 2352         M. Derrick et al.         (HRS Co					
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ADACHI         88C         PL B208 319         I. Adachi et al.         (TOPAZ Collab.)           ADEVA         88         PR D38 2665         B. Adeva et al.         (Mark-J Collab.)           BRAUNSCH         88D         ZPHY C40 163         W. Braunschweig et al.         (TASSO Collab.)           ANSARI         87         PL B186 440         R. Ansari et al.         (UA2 Collab.)           BEHREND         87C         PL B191 209         H.J. Behrend et al.         (CELLO Collab.)           BARTEL         86C         ZPHY C30 371         W. Bartel et al.         (JADE Collab.)           Also         ZPHY C26 507         W. Bartel et al.         (JADE Collab.)           ASH         85         PRL 55 1831         W.W. Ash et al.         (MAC Collab.)           ASH         85         PRL 55 1831         W. Bartel et al.         (JADE Collab.)           DERRICK         85         PR D31 2352         M. Derrick et al.         (HRS Collab.)           FERNANDEZ         85         PRL 54 1624         E. Fernandez et al.         (MAC Collab.)           LEVI         83         PRL 51 1941         M.E. Levi et al.         (Mark II Collab.)           BEHREND         82         PL 114B 282         H.J. Behrend et al.         (CELLO Collab.)					`. · · · · · · · · · · · · · · · · · · ·
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BRAUNSCH         88D         ZPHY C40 163         W. Braunschweig et al.         (TASSO Collab.)           ANSARI         87         PL B186 440         R. Ansari et al.         (UA2 Collab.)           BEHREND         87C         PL B191 209         H.J. Behrend et al.         (CELLO Collab.)           BARTEL         86C         ZPHY C30 371         W. Bartel et al.         (JADE Collab.)           Also         ZPHY C26 507         W. Bartel et al.         (JADE Collab.)           ASH         85         PRL 55 1831         W.W. Ash et al.         (MAC Collab.)           BARTEL         85F         PL 161B 188         W. Bartel et al.         (JADE Collab.)           DERRICK         85         PR D31 2352         M. Derrick et al.         (HRS Collab.)           FERNANDEZ         85         PRL 54 1624         E. Fernandez et al.         (MAC Collab.)           LEVI         83         PRL 51 1941         M.E. Levi et al.         (Mark II Collab.)           BEHREND         82         PL 114B 282         H.J. Behrend et al.         (CELLO Collab.)					`
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BARTEL       86C       ZPHY C30 371       W. Bartel et al.       (JADE Collab.)         Also       ZPHY C26 507       W. Bartel et al.       (JADE Collab.)         Also       PL 108B 140       W. Bartel et al.       (JADE Collab.)         ASH       85       PRL 55 1831       W.W. Ash et al.       (MAC Collab.)         BARTEL       85F       PL 161B 188       W. Bartel et al.       (JADE Collab.)         DERRICK       85       PR D31 2352       M. Derrick et al.       (HRS Collab.)         FERNANDEZ       85       PRL 54 1624       E. Fernandez et al.       (MAC Collab.)         LEVI       83       PRL 51 1941       M.E. Levi et al.       (Mark II Collab.)         BEHREND       82       PL 114B 282       H.J. Behrend et al.       (CELLO Collab.)					
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Also         PL 108B 140         W. Bartel et al.         (JADE Collab.)           ASH         85         PRL 55 1831         W.W. Ash et al.         (MAC Collab.)           BARTEL         85F         PL 161B 188         W. Bartel et al.         (JADE Collab.)           DERRICK         85         PR D31 2352         M. Derrick et al.         (HRS Collab.)           FERNANDEZ         85         PRL 54 1624         E. Fernandez et al.         (MAC Collab.)           LEVI         83         PRL 51 1941         M.E. Levi et al.         (Mark II Collab.)           BEHREND         82         PL 114B 282         H.J. Behrend et al.         (CELLO Collab.)	Also		ZPHY C26 507	W. Bartel <i>et al.</i>	(JADE Collab.)
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DERRICK         85         PR D31 2352         M. Derrick et al.         (HRS Collab.)           FERNANDEZ         85         PRL 54 1624         E. Fernandez et al.         (MAC Collab.)           LEVI         83         PRL 51 1941         M.E. Levi et al.         (Mark II Collab.)           BEHREND         82         PL 114B 282         H.J. Behrend et al.         (CELLO Collab.)					
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LEVI         83         PRL 51 1941         M.E. Levi et al.         (Mark II Collab.)           BEHREND         82         PL 114B 282         H.J. Behrend et al.         (CELLO Collab.)					
BEHREND 82 PL 114B 282 H.J. Behrend et al. (CELLO Collab.)					
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