## $b^{\prime}$ (4 ${ }^{\text {th }}$ Generation) Quark, Searches for

## $b^{\prime}$-quark/hadron mass limits in $p \bar{p}$ and $p p$ collisions

| VALUE (GeV) | CL\% | DOCUMENT ID | TECN | COMMENT |
| :---: | :---: | :---: | :---: | :---: |
| >372 | 95 | 1 AALTONEN 11J | CDF | $b^{\prime} \rightarrow t W$ |
| >361 | 95 | 2 CHATRCHYAN 11L | CMS | $b^{\prime} \rightarrow t W$ |
| >190 | 95 | 3 ABAZOV 08x | D0 | $\mathrm{c} \tau=200 \mathrm{~mm}$ |
| >268 | 95 | 4,5 AALTONEN 07C | CDF | $\mathrm{B}\left(b^{\prime} \rightarrow b Z\right)=1$ assumed |
| $>190$ | 95 | ${ }^{6}$ ACOSTA 03 | CDF | quasi-stable $b^{\prime}$ |
| $>128$ | 95 | $7 \mathrm{ABACHI} \quad 95 \mathrm{~F}$ | D0 | $\ell \ell+$ jets, $\ell+$ jets |
| - - We do not use the following data for averages, fits, limits, etc. - - |  |  |  |  |
| >338 | 95 | ${ }^{8}$ AALTONEN 10 H | CDF | $b^{\prime} \rightarrow t W$ |
| > 380-430 | 95 | 9 FLACCO 10 | RVUE | $m_{b^{\prime}}>m_{t^{\prime}}$ |
| >199 | 95 | 10 AFFOLDER 00 | CDF | NC: $b^{\prime} \rightarrow b Z$ |
| $>148$ | 95 | 11 ABE 98N | CDF | NC: $b^{\prime} \rightarrow b Z+$ decay vertex |
| $>96$ | 95 | 12 ABACHI 97 D | D0 | NC: $b^{\prime} \rightarrow b \gamma$ |
| $>75$ | 95 | 13 MUKHOPAD... 93 | RVUE | NC: $b^{\prime} \rightarrow$ ble |
| $>85$ | 95 | 14 ABE | CDF | CC: $\ell \ell$ |
| $>72$ | 95 | 15 ABE 90B | CDF | CC: $e+\mu$ |
| $>54$ | 95 | 16 AKESSON 90 | UA2 | CC: $e+$ jets + missing $E_{T}$ |
| $>43$ | 95 | 17 ALBAJAR 90B | UA1 | CC: $\mu+$ jets |
| $>34$ | 95 | 18 ALBAJAR 88 | UA1 | CC: e or $\mu+$ jets |

${ }^{1}$ Based on $4.8 \mathrm{fb}^{-1}$ of data in $p \bar{p}$ collisions at 1.96 TeV . AALTONEN 11 J looked for events with $\ell+E_{T}+\geq 5 \mathrm{j}(\geq 1 b$ or $c)$. No signal is observed and the bound $\sigma\left(b^{\prime} \bar{b}^{\prime}\right)$ $<30 \mathrm{fb}$ for $m_{b^{\prime}}>375 \mathrm{GeV}$ is found for $\mathrm{B}\left(b^{\prime} \rightarrow t W\right)=1$.
2 Based on $34 \mathrm{pb}^{-1}$ of data in $p p$ collisions at 7 TeV . CHATRCHYAN 11 L looked for multijet events with trileptons or same-sign dileptons. No excess above the SM background excludes $m_{b^{\prime}}$ between 255 and 361 GeV at $95 \% \mathrm{CL}$ for $\mathrm{B}\left(b^{\prime} \rightarrow t W\right)=1$.
${ }^{3}$ Result is based on $1.1 \mathrm{fb}^{-1}$ of data. No signal is found for the search of long-lived particles which decay into final states with two electrons or photons, and upper bound on the cross section times branching fraction is obtained for $2<\mathrm{c} \tau<7000 \mathrm{~mm}$; see Fig. 3. $95 \%$ CL excluded region of $b^{\prime}$ lifetime and mass is shown in Fig. 4.
${ }^{4}$ Result is based on $1.06 \mathrm{fb}^{-1}$ of data. No excess from the SM $Z+$ jet events is found when $Z$ decays into ee or $\mu \mu$. The $m_{b^{\prime}}$ bound is found by comparing the resulting upper bound on $\sigma\left(b^{\prime} \bar{b}^{\prime}\right)\left[1-\left(1-\mathrm{B}\left(b^{\prime} \rightarrow b Z\right)\right)^{2}\right]$ and the LO estimate of the $b^{\prime}$ pair production cross section shown in Fig. 38 of the article.
${ }^{5}$ HUANG 08 reexamined the $b^{\prime}$ mass lower bound of 268 GeV obtained in AALTONEN 07C that assumes $\mathrm{B}\left(b^{\prime} \rightarrow b Z\right)=1$, which does not hold for $m_{b^{\prime}}>255 \mathrm{GeV}$. The lower mass bound is given in the plane of $\sin ^{2}\left(\theta_{t b^{\prime}}\right)$ and $m_{b^{\prime}}$.
${ }^{6}$ ACOSTA 03 looked for long-lived fourth generation quarks in the data sample of 90 $\mathrm{pb}^{-1}$ of $\sqrt{s}=1.8 \mathrm{TeV} p \bar{p}$ collisions by using the muon-like penetration and anomalously high ionization energy loss signature. The corresponding lower mass bound for the charge $(2 / 3)$ e quark $\left(t^{\prime}\right)$ is 220 GeV . The $t^{\prime}$ bound is higher than the $b^{\prime}$ bound because $t^{\prime}$ is more likely to produce charged hadrons than $b^{\prime}$. The $95 \% \mathrm{CL}$ upper bounds for the production cross sections are given in their Fig. 3.
${ }^{7}$ ABACHI 95F bound on the top-quark also applies to $b^{\prime}$ and $t^{\prime}$ quarks that decay predominantly into $W$. See FROGGATT 97.
${ }^{8}$ Based on $2.7 \mathrm{fb}^{-1}$ of data in $p \bar{p}$ collisions at $\sqrt{s}=1.96 \mathrm{TeV}$. AALTONEN 10 H looked for pair production of heavy quarks which decay into $t W^{-}$or $t W^{+}$, in events with same sign dileptons ( $e$ or $\mu$ ), several jets and large missing $E_{T}$. The result is obtained for $b^{\prime}$ which decays into $t W^{-}$. For the charge $5 / 3$ quark ( $T_{5 / 3}$ ) which decays into $t W^{+}$, ${ }^{m} T_{5 / 3}>365 \mathrm{GeV}(95 \% \mathrm{CL})$ is found when it has the charge $-1 / 3$ partner B of the same mass.
${ }^{9}$ FLACCO 10 result is obtained from AALTONEN 10 H result of $m_{b^{\prime}}>338 \mathrm{GeV}$, by relaxing the condition $\mathrm{B}\left(b^{\prime} \rightarrow t W\right)=100 \%$ when $m_{b^{\prime}}>m_{t^{\prime}}$.
${ }^{10}$ AFFOLDER 00 looked for $b^{\prime}$ that decays in to $b+Z$. The signal searched for is $b b Z Z$ events where one $Z$ decays into $e^{+} e^{-}$or $\mu^{+} \mu^{-}$and the other $Z$ decays hadronically. The bound assumes $\mathrm{B}\left(b^{\prime} \rightarrow b Z\right)=100 \%$. Between 100 GeV and 199 GeV , the $95 \% \mathrm{CL}$ upper bound on $\sigma\left(b^{\prime} \rightarrow \bar{b}^{\prime}\right) \times \mathrm{B}^{2}\left(b^{\prime} \rightarrow b Z\right)$ is also given (see their Fig. 2).
${ }^{11}$ ABE 98N looked for $Z \rightarrow e^{+} e^{-}$decays with displaced vertices. Quoted limit assumes $\mathrm{B}\left(b^{\prime} \rightarrow b Z\right)=1$ and $c \tau_{b^{\prime}}=1 \mathrm{~cm}$. The limit is lower than $m_{Z}+m_{b}(\sim 96 \mathrm{GeV})$ if $c \tau>22 \mathrm{~cm}$ or $c \tau<0.009 \mathrm{~cm}$. See their Fig. 4.
12 ABACHI 97D searched for $b^{\prime}$ that decays mainly via FCNC. They obtained $95 \%$ CL upper bounds on $\mathrm{B}\left(b^{\prime} \bar{b}^{\prime} \rightarrow \gamma+3\right.$ jets) and $\mathrm{B}\left(b^{\prime} \bar{b}^{\prime} \rightarrow 2 \gamma+2\right.$ jets), which can be interpreted as the lower mass bound $m_{b^{\prime}}>m_{Z}+m_{b}$.
${ }^{13}$ MUKHOPADHYAYA 93 analyze CDF dilepton data of ABE 92 G in terms of a new quark decaying via flavor-changing neutral current. The above limit assumes $\mathrm{B}\left(b^{\prime} \rightarrow\right.$ $\left.b \ell^{+} \ell^{-}\right)=1 \%$. For an exotic quark decaying only via virtual $Z\left[\mathrm{~B}\left(b \ell^{+} \ell^{-}\right)=3 \%\right]$, the limit is 85 GeV .
${ }^{14} \mathrm{ABE} 92$ dilepton analysis limit of $>85 \mathrm{GeV}$ at $\mathrm{CL}=95 \%$ also applies to $b^{\prime}$ quarks, as discussed in ABE 90B.
${ }^{15}$ ABE 90B exclude the region $28-72 \mathrm{GeV}$.
${ }^{16}$ AKESSON 90 searched for events having an electron with $p_{T}>12 \mathrm{GeV}$, missing momentum $>15 \mathrm{GeV}$, and a jet with $E_{T}>10 \mathrm{GeV},|\eta|<2.2$, and excluded $m_{b^{\prime}}$ between 30 and 69 GeV .
17 For the reduction of the limit due to non-charged-current decay modes, see Fig. 19 of ALBAJAR 90b.
18 ALBAJAR 88 study events at $E_{\mathrm{cm}}=546$ and 630 GeV with a muon or isolated electron, accompanied by one or more jets and find agreement with Monte Carlo predictions for the production of charm and bottom, without the need for a new quark. The lower mass limit is obtained by using a conservative estimate for the $b^{\prime} \bar{b}^{\prime}$ production cross section and by assuming that it cannot be produced in $W$ decays. The value quoted here is revised using the full $O\left(\alpha_{s}^{3}\right)$ cross section of ALTARELLI 88.

## $\boldsymbol{b}^{\prime}$ mass limits from single production in $\boldsymbol{p} \overline{\boldsymbol{p}}$ and $p p$ collisions

| VALUE (GeV) | CL\% | DOCUMENT ID |  | TECN | COMMENT |
| :---: | :---: | :---: | :---: | :---: | :---: |
| >693 | 95 | 19 ABAZOV | 11F | D0 | $q u \rightarrow q^{\prime} b^{\prime} \rightarrow q^{\prime}(W u)$ |
|  |  |  |  |  | $\widetilde{\kappa}_{u b^{\prime}}=1, \mathrm{~B}\left(b^{\prime} \rightarrow W u\right)=1$ |
| >430 | 95 | 19 ABAZOV | 11F | D0 | $q d \rightarrow q b^{\prime} \rightarrow q(Z d)$ |
|  |  |  |  |  | $\widetilde{\kappa}_{d b^{\prime}}=\sqrt{2}, \mathrm{~B}\left(b^{\prime} \rightarrow Z d\right)=1$ |

${ }^{19}$ Based on $5.4 \mathrm{fb}^{-1}$ of data in ppbar collisions at 1.96 TeV . ABAZOV 11 F looked for single production of $b^{\prime}$ via the $W$ or $Z$ coupling to the first generation up or down quarks, respectively. Model independent cross section limits for the single production processes $p \bar{p} \rightarrow b^{\prime} q \rightarrow W u q$, and $p \bar{p} \rightarrow b^{\prime} q \rightarrow Z d q$ are given in Figs. 3 and 4, respectively, and the mass limits are obtained for the model of ATRE 09 with degenerate bi-doublets of vector-like quarks.

## MASS LIMITS for $\boldsymbol{b}^{\prime}$ (4 ${ }^{\text {th }}$ Generation) Quark or Hadron in $\boldsymbol{e}^{+} \boldsymbol{e}^{-}$Collisions

Search for hadrons containing a fourth-generation $-1 / 3$ quark denoted $b^{\prime}$.
The last column specifies the assumption for the decay mode ( $C$ C denotes the conventional charged-current decay) and the event signature which is looked for.

| VALUE (GeV) | CL\% | DOCUMENT ID |  | TECN | COMMENT |
| :---: | :---: | :---: | :---: | :---: | :---: |
| >46.0 | 95 | 20 DECAMP | 90F | ALEP | any decay |
| - - We do not use the following data for averages, fits, limits, etc. - - |  |  |  |  |  |
| none 96-103 | 95 | $21 \text { ABDALLAH } 07$ |  | DLPH | $b^{\prime} \rightarrow b Z, c W$ |
|  |  | 22 ADRIANI | 93G | L3 | Quarkonium |
| $>44.7$ | 95 | ADRIANI | 93M | L3 | $\Gamma(Z)$ |
| $>45$ | 95 | ABREU | 91F | DLPH | $\Gamma(Z)$ |
| none 19.4-28.2 | 95 | ABE | 90D | VNS | Any decay; event shape |
| $>45.0$ | 95 | ABREU | 90D | DLPH | $B(C C)=1$; event shape |
| $>44.5$ | 95 | 23 ABREU | 90D | DLPH | $\begin{aligned} & b^{\prime} \rightarrow c H^{-}, H^{-} \rightarrow \\ & \bar{c} s, \tau^{-} \nu \end{aligned}$ |
| $>40.5$ | 95 | 24 ABREU | 90D | DLPH | $\Gamma(Z \rightarrow$ hadrons $)$ |
| $>28.3$ | 95 | ADACHI | 90 | TOPZ | $\begin{aligned} & \mathrm{B}(\mathrm{FCNC})=100 \% \text {; isol. } \\ & \gamma \text { or } 4 \text { jets } \end{aligned}$ |
| $>41.4$ | 95 | 25 AKRAWY | 90B | OPAL | Any decay; acoplanarity |
| $>45.2$ | 95 | 25 AKRAWY | 90B | OPAL | $B(C C)=1$; acoplanarity |
| $>46$ | 95 | 26 AKRAWY | 90J | OPAL | $b^{\prime} \rightarrow \gamma+$ any |
| $>27.5$ | 95 | 27 ABE | 89E | VNS | $\mathrm{B}(\mathrm{CC})=1 ; \mu, e$ |
| none 11.4-27.3 | 95 | 28 ABE | 89G | VNS | $\underset{\underset{\text { isolated } \gamma}{\mathrm{B}}\left(b^{\prime} \rightarrow b \gamma\right)}{ }>10 \%$ |
| $>44.7$ | 95 | 29 ABRAMS | 89C | MRK2 | $B(C C)=100 \% \text {; isol. }$ <br> track |
| $>42.7$ | 95 | 29 ABRAMS | 89C | MRK2 | $\mathrm{B}(b g)=100 \% ; \text { event }$ shape |
| $>42.0$ | 95 | 29 ABRAMS | 89C | MRK2 | Any decay; event shape |
| $>28.4$ | 95 | 30,31 ADACHI | 89C | TOPZ | $\mathrm{B}(C C)=1 ; \mu$ |
| $>28.8$ | 95 | 32 ENO | 89 | AMY | $\mathrm{B}(\mathrm{CC}) \gtrsim 90 \%$; $\mu$, e |
| $>27.2$ | 95 | 32,33 ENO | 89 | AMY | any decay; event shape |
| $>29.0$ | 95 | 32 ENO | 89 | AMY | $\begin{gathered} \mathrm{B}\left(b^{\prime} \rightarrow b g\right) \gtrsim 85 \% ; \\ \text { event shape } \end{gathered}$ |
| $>24.4$ | 95 | 34 IGARASHI | 88 | AMY | $\mu, e$ |
| $>23.8$ | 95 | 35 SAGAWA | 88 | AMY | event shape |
| $>22.7$ | 95 | 36 ADEVA | 86 | MRKJ | $\mu$ |
| $>21$ |  | 37 ALTHOFF | 84C | TASS | $R$, event shape |
| >19 |  | 38 ALTHOFF | 841 | TASS | Aplanarity |

20 DECAMP 90F looked for isolated charged particles, for isolated photons, and for four-jet final states. The modes $b^{\prime} \rightarrow b g$ for $\mathrm{B}\left(b^{\prime} \rightarrow b g\right)>65 \% b^{\prime} \rightarrow b \gamma$ for $\mathrm{B}\left(b^{\prime} \rightarrow b \gamma\right)$ $>5 \%$ are excluded. Charged Higgs decay were not discussed.
${ }^{21}$ ABDALLAH 07 searched for $b^{\prime}$ pair production at $E_{\mathrm{cm}}=196-209 \mathrm{GeV}$, with $420 \mathrm{pb}^{-1}$. No signal leads to the $95 \% \mathrm{CL}$ upper limits on $\mathrm{B}\left(b^{\prime} \rightarrow b Z\right)$ and $\mathrm{B}\left(b^{\prime} \rightarrow c W\right)$ for $m_{b^{\prime}}$ $=96$ to 103 GeV .
22 ADRIANI 93G search for vector quarkonium states near $Z$ and give limit on quarkonium$Z$ mixing parameter $\delta m^{2}<(10-30) \mathrm{GeV}^{2}(95 \% \mathrm{CL})$ for the mass $88-94.5 \mathrm{GeV}$. Using Richardson potential, a $1 \mathrm{~S}\left(b^{\prime} \bar{b}^{\prime}\right)$ state is excluded for the mass range 87.7-94.7 GeV. This range depends on the potential choice.
23 ABREU 90D assumed $m_{H^{-}}<m_{b^{\prime}}-3 \mathrm{GeV}$.
24 Superseded by ABREU 91F.
${ }^{25}$ AKRAWY 90B search was restricted to data near the $Z$ peak at $E_{\mathrm{cm}}=91.26 \mathrm{GeV}$ at LEP. The excluded region is between 23.6 and 41.4 GeV if no $\mathrm{H}^{+}$decays exist. For charged Higgs decays the excluded regions are between ( $m_{H^{+}}+1.5 \mathrm{GeV}$ ) and 45.5 GeV.
26 AKRAWY 90J search for isolated photons in hadronic $Z$ decay and derive
$\mathrm{B}\left(Z \rightarrow b^{\prime} \bar{b}^{\prime}\right) \cdot \mathrm{B}\left(b^{\prime} \rightarrow \gamma \mathrm{X}\right) / \mathrm{B}(Z \rightarrow$ hadrons $)<2.2 \times 10^{-3}$. Mass limit assumes $\mathrm{B}\left(b^{\prime} \rightarrow \gamma X\right)>10 \%$.
${ }^{27}$ ABE 89E search at $E_{c m}=56-57 \mathrm{GeV}$ at TRISTAN for multihadron events with a spherical shape (using thrust and acoplanarity) or containing isolated leptons.
28 ABE 89 G search was at $E_{\mathrm{cm}}=55-60.8 \mathrm{GeV}$ at TRISTAN.
29 If the photonic decay mode is large $\left(\mathrm{B}\left(b^{\prime} \rightarrow b \gamma\right)>25 \%\right)$, the ABRAMS 89C limit is 45.4 GeV. The limit for for Higgs decay ( $b^{\prime} \rightarrow c H^{-}, H^{-} \rightarrow \bar{c} s$ ) is 45.2 GeV .
${ }^{30}$ ADACHI 89C search was at $E_{\mathrm{cm}}=56.5-60.8 \mathrm{GeV}$ at TRISTAN using multi-hadron events accompanying muons.
31 ADACHI 89C also gives limits for any mixture of CC and $b g$ decays.
32 ENO 89 search at $E_{\mathrm{cm}}=50-60.8$ at TRISTAN.
33 ENO 89 considers arbitrary mixture of the charged current, $b g$, and $b \gamma$ decays.
${ }^{34}$ IGARASHI 88 searches for leptons in low-thrust events and gives $\Delta R\left(b^{\prime}\right)<0.26(95 \%$ CL ) assuming charged current decay, which translates to $m_{b^{\prime}}>24.4 \mathrm{GeV}$.
35 SAGAWA 88 set limit $\sigma($ top $)<6.1 \mathrm{pb}$ at $\mathrm{CL}=95 \%$ for top-flavored hadron production from event shape analyses at $E_{\mathrm{cm}}=52 \mathrm{GeV}$. By using the quark parton model crosssection formula near threshold, the above limit leads to lower mass bounds of 23.8 GeV for charge $-1 / 3$ quarks.
36 ADEVA 86 give $95 \%$ CL upper bound on an excess of the normalized cross section, $\Delta R$, as a function of the minimum c.m. energy (see their figure 3). Production of a pair of $1 / 3$ charge quarks is excluded up to $E_{\mathrm{cm}}=45.4 \mathrm{GeV}$.
37 ALTHOFF 84C narrow state search sets limit $\Gamma\left(e^{+} e^{-}\right) \mathrm{B}$ (hadrons) $<2.4 \mathrm{keV} \mathrm{CL}=95 \%$ and heavy charge $1 / 3$ quark pair production $m>21 \mathrm{GeV}, \mathrm{CL}=95 \%$.
38 ALTHOFF 84I exclude heavy quark pair production for $7<m<19 \mathrm{GeV}$ ( $1 / 3$ charge) using aplanarity distributions ( $C L=95 \%$ ).

## REFERENCES FOR Searches for (Fourth Generation) $\boldsymbol{b}^{\prime}$ Quark

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