

THE MUON ANOMALOUS MAGNETIC MOMENT

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The Dirac equation predicts a muon magnetic moment, $\vec{M} = g_\mu \frac{e}{2m_\mu} \vec{S}$, with gyromagnetic ratio $g_\mu = 2$. Quantum loop effects lead to a small calculable deviation from $g_\mu = 2$, parameterized by the anomalous magnetic moment

$$a_\mu \equiv \frac{g_\mu - 2}{2} . \quad (1)$$

That quantity can be accurately measured and, within the Standard Model (SM) framework, precisely predicted. Hence, comparison of experiment and theory tests the SM at its quantum loop level. A deviation in a_μ^{exp} from the SM expectation would signal effects of new physics, with current sensitivity reaching up to mass scales of $\mathcal{O}(\text{TeV})$ [1,2]. For recent and very thorough muon $g - 2$ reviews, see Refs. [3,4].

The E821 experiment at Brookhaven National Lab (BNL) studied the precession of μ^+ and μ^- in a constant external magnetic field as they circulated in a confining storage ring. It found [6]¹

$$\begin{aligned} a_{\mu^+}^{\text{exp}} &= 11\,659\,204(6)(5) \times 10^{-10} , \\ a_{\mu^-}^{\text{exp}} &= 11\,659\,215(8)(3) \times 10^{-10} , \end{aligned} \quad (2)$$

where the first errors are statistical and the second systematic. Assuming CPT invariance and taking into account correlations between systematic errors, one finds for their average [6]

$$a_\mu^{\text{exp}} = 11\,659\,208.9(5.4)(3.3) \times 10^{-10} . \quad (3)$$

These results represent about a factor of 14 improvement over the classic CERN experiments of the 1970's [7]. Improvement

¹ The original results reported by the experiment have been updated in Eqs. (2) and (3) to the newest value for the absolute muon-to-proton magnetic ratio $\lambda = 3.183345137(85)$ [5]. The change induced in a_μ^{exp} with respect to the value of $\lambda = 3.18334539(10)$ used in Ref. 6 amounts to $+0.92 \times 10^{-10}$.

of the measurement in Eq. (3) by a factor of four by moving the E821 storage ring to Fermilab, and utilizing a cleaner and more intense muon beam has been proposed.

The SM prediction for a_μ^{SM} is generally divided into three parts (see Fig. 1 for representative Feynman diagrams)

$$a_\mu^{\text{SM}} = a_\mu^{\text{QED}} + a_\mu^{\text{EW}} + a_\mu^{\text{Had}} . \quad (4)$$

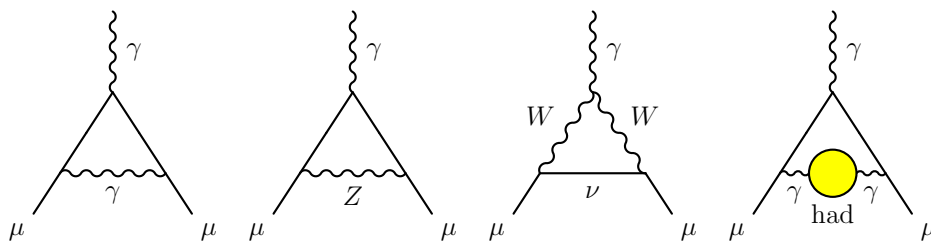


Figure 1: Representative diagrams contributing to a_μ^{SM} . From left to right: first order QED (Schwinger term), lowest-order weak, lowest-order hadronic.

The QED part includes all photonic and leptonic (e, μ, τ) loops starting with the classic $\alpha/2\pi$ Schwinger contribution. It has been computed through 4 loops and estimated at the 5-loop level [8]

$$a_\mu^{\text{QED}} = \frac{\alpha}{2\pi} + 0.765857410(27) \left(\frac{\alpha}{\pi}\right)^2 + 24.05050964(43) \left(\frac{\alpha}{\pi}\right)^3 + 130.8055(80) \left(\frac{\alpha}{\pi}\right)^4 + 663(20) \left(\frac{\alpha}{\pi}\right)^5 + \dots \quad (5)$$

Employing $\alpha^{-1} = 137.035999084(51)$, determined [8,9] from the electron a_e measurement, leads to

$$a_\mu^{\text{QED}} = 116\,584\,718.09(0.15) \times 10^{-11}, \quad (6)$$

where the error results from uncertainties in the coefficients of Eq. (5) and in α .

Loop contributions involving heavy W^\pm, Z or Higgs particles are collectively labeled as a_μ^{EW} . They are suppressed by at

least a factor of $\frac{\alpha}{\pi} \frac{m_\mu^2}{m_W^2} \simeq 4 \times 10^{-9}$. At 1-loop order [10]

$$\begin{aligned} a_\mu^{\text{EW}}[\text{1-loop}] &= \frac{G_\mu m_\mu^2}{8\sqrt{2}\pi^2} \left[\frac{5}{3} + \frac{1}{3} (1 - 4 \sin^2 \theta_W)^2 \right. \\ &\quad \left. + \mathcal{O}\left(\frac{m_\mu^2}{M_W^2}\right) + \mathcal{O}\left(\frac{m_\mu^2}{m_H^2}\right) \right], \\ &= 194.8 \times 10^{-11}, \end{aligned} \quad (7)$$

for $\sin^2 \theta_W \equiv 1 - M_W^2/M_Z^2 \simeq 0.223$, and where $G_\mu \simeq 1.166 \times 10^{-5} \text{ GeV}^{-2}$ is the Fermi coupling constant. Two-loop corrections are relatively large and negative [11]

$$a_\mu^{\text{EW}}[\text{2-loop}] = -40.7(1.0)(1.8) \times 10^{-11}, \quad (8)$$

where the errors stem from quark triangle loops and the assumed Higgs mass range between 100 and 500 GeV. The 3-loop leading logarithms are negligible [11,12], $\mathcal{O}(10^{-12})$, implying in total

$$a_\mu^{\text{EW}} = 154(1)(2) \times 10^{-11}. \quad (9)$$

Hadronic (quark and gluon) loop contributions to a_μ^{SM} give rise to its main theoretical uncertainties. At present, those effects are not calculable from first principles, but such an approach, at least partially, may become possible as lattice QCD matures. Instead, one currently relies on a dispersion relation approach to evaluate the lowest-order (*i.e.*, $\mathcal{O}(\alpha^2)$) hadronic vacuum polarization contribution $a_\mu^{\text{Had}}[\text{LO}]$ from corresponding cross section measurements [13]

$$a_\mu^{\text{Had}}[\text{LO}] = \frac{1}{3} \left(\frac{\alpha}{\pi} \right)^2 \int_{m_\pi^2}^{\infty} ds \frac{K(s)}{s} R^{(0)}(s), \quad (10)$$

where $K(s)$ is a QED kernel function [14], and where $R^{(0)}(s)$ denotes the ratio of the bare² cross section for e^+e^- annihilation

² The bare cross section is defined as the measured cross section corrected for initial-state radiation, electron-vertex loop contributions and vacuum-polarization effects in the photon propagator. However, QED effects in the hadron vertex and final state, as photon radiation, are included.

into hadrons to the pointlike muon-pair cross section at center-of-mass energy \sqrt{s} . The function $K(s) \sim 1/s$ in Eq. (10) gives a strong weight to the low-energy part of the integral. Hence, $a_\mu^{\text{Had}}[\text{LO}]$ is dominated by the $\rho(770)$ resonance.

Currently, the available $\sigma(e^+e^- \rightarrow \text{hadrons})$ data give a leading-order hadronic vacuum polarization (representative) contribution of [15]

$$a_\mu^{\text{Had}}[\text{LO}] = 6\,923(42)(3) \times 10^{-11}, \quad (11)$$

where the first error is experimental (dominated by systematic uncertainties), and the second due to perturbative QCD, which is used at intermediate and large energies to predict the contribution from the quark-antiquark continuum. New multi-hadron data from the BABAR experiment have increased the constraints on unmeasured exclusive final states and led to a small reduction in the hadronic contribution compared to the 2009 PDG value.

Alternatively, one can use precise vector spectral functions from $\tau \rightarrow \nu_\tau + \text{hadrons}$ decays [16] that can be related to isovector $e^+e^- \rightarrow \text{hadrons}$ cross sections by isospin symmetry. Replacing e^+e^- data in the two-pion and four-pion channels by the corresponding isospin-transformed τ data, and applying isospin-violating corrections (from QED and $m_d - m_u \neq 0$), one finds [15]

$$a_\mu^{\text{Had}}[\text{LO}] = 7\,015(42)(19)(3) \times 10^{-11} (\tau), \quad (12)$$

where the first error is experimental, the second estimates the uncertainty in the isospin-breaking corrections applied to the τ data, and the third error is due to perturbative QCD. The current discrepancy between the e^+e^- and τ -based determinations of $a_\mu^{\text{Had}}[\text{LO}]$ has been reduced to 1.8σ with respect to earlier evaluations. New e^+e^- and τ data from the B -factory experiments BABAR and Belle have increased the experimental information. Reevaluated isospin-breaking corrections have also contributed to this improvement [17]. BABAR recently reported good agreement with the τ data in the most important two-pion channel [18]. The remaining discrepancy with the older e^+e^- and τ datasets may be indicative of problems

with one or both data sets. It may also suggest the need for additional isospin-violating corrections to the τ data.

Higher order, $\mathcal{O}(\alpha^3)$, hadronic contributions are obtained from dispersion relations using the same $e^+e^- \rightarrow$ hadrons data [16,19,22], giving $a_\mu^{\text{Had,Disp}}[\text{NLO}] = (-98.4 \pm 0.6) \times 10^{-11}$, along with model-dependent estimates of the hadronic light-by-light scattering contribution, $a_\mu^{\text{Had,LBL}}[\text{NLO}]$, motivated by large- N_C QCD [23–29].³ Following [27], one finds for the sum of the two terms

$$a_\mu^{\text{Had}}[\text{NLO}] = 7(26) \times 10^{-11}, \quad (13)$$

where the error is dominated by hadronic light-by-light uncertainties.

Adding Eqs. (6), (9), (11) and (13) gives the representative e^+e^- data based SM prediction

$$a_\mu^{\text{SM}} = 116\,591\,802(2)(42)(26) \times 10^{-11}, \quad (14)$$

where the errors are due to the electroweak, lowest-order hadronic, and higher-order hadronic contributions, respectively. The difference between experiment and theory

$$\Delta a_\mu = a_\mu^{\text{exp}} - a_\mu^{\text{SM}} = 287(63)(49) \times 10^{-11}, \quad (15)$$

(with all errors combined in quadrature) represents an interesting but not yet conclusive discrepancy of 3.6 times the estimated 1σ error. All the recent estimates for the hadronic contribution compiled in Fig. 2 exhibit similar discrepancies. Switching to τ data reduces the discrepancy to 2.4σ , assuming the isospin-violating corrections are under control within the estimated uncertainties.

An alternate interpretation is that Δa_μ may be a new physics signal with supersymmetric particle loops as the leading candidate explanation. Such a scenario is quite natural, since

³ Some representative recent estimates of the hadronic light-by-light scattering contribution, $a_\mu^{\text{Had,LBL}}[\text{NLO}]$, that followed after the sign correction of [25], are: $105(26) \times 10^{-11}$ [27], $110(40) \times 10^{-11}$ [23], $136(25) \times 10^{-11}$ [24].

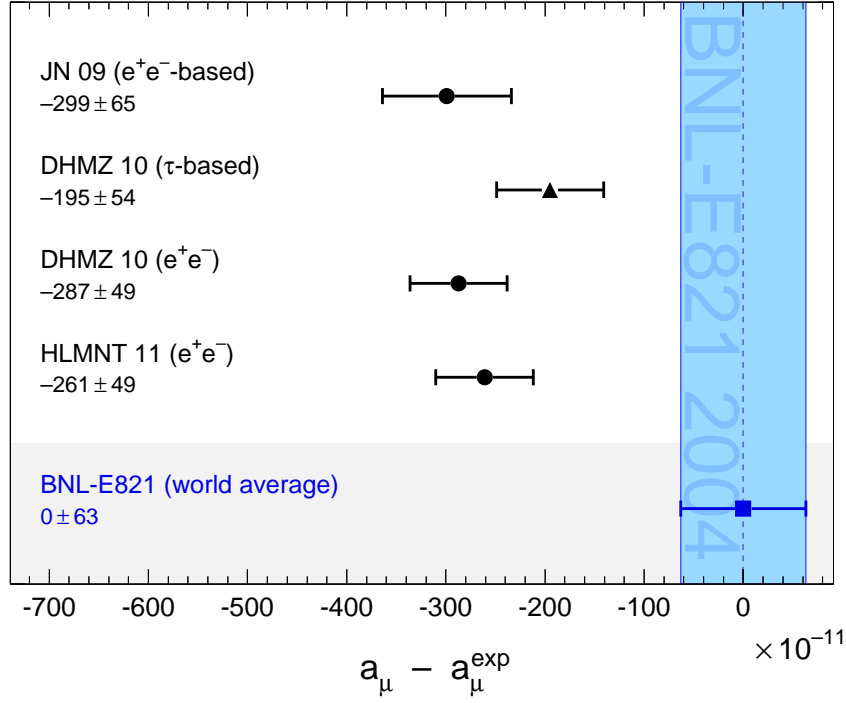


Figure 2: Compilation of recently published results for a_μ (in units of 10^{-11}), subtracted by the central value of the experimental average (3). The shaded band indicates the experimental error. The SM predictions are taken from: JN [4], DHMZ [15], and HMNT [19]. Note that the quoted errors do not include the uncertainty on the subtracted experimental value. To obtain for each theory calculation a result equivalent to Eq. (15), the errors from theory and experiment must be added in quadrature.

generically, supersymmetric models predict [1] an additional contribution to a_μ^{SM}

$$a_\mu^{\text{SUSY}} \simeq \pm 130 \times 10^{-11} \cdot \left(\frac{100 \text{ GeV}}{m_{\text{SUSY}}} \right)^2 \tan\beta, \quad (16)$$

where m_{SUSY} is a representative supersymmetric mass scale, and $\tan\beta \simeq 3\text{--}40$ is a potential enhancement factor. Supersymmetric particles in the mass range 100–500 GeV could be the source of the deviation Δa_μ . If so, those particles could be directly observed at the next generation of high energy colliders.

New physics effects [1] other than supersymmetry could also explain a non-vanishing Δa_μ . A recent popular scenario involves the “dark photon”, a relatively light hypothetical vector boson from the dark matter sector that couples to our world of particle physics through mixing with the ordinary photon [30,31]. As a result, it couples to ordinary charged particles with strength $\varepsilon \cdot e$ and gives rise to an additional muon anomalous magnetic moment contribution

$$a_\mu^{\text{dark photon}} = \frac{\alpha}{2\pi} \varepsilon^2 F(m_V/m_\mu), \quad (17)$$

where $F(x) = \int_0^1 2z(1-z)^2/[(1-z)^2 + x^2z] dz$. For values of $\varepsilon \sim 1-2 \cdot 10^{-3}$ and $m_V \sim 10-100$ MeV, the dark photon, which was originally motivated by cosmology, can provide a viable solution to the muon $g - 2$ discrepancy. Searches for the dark photon in that mass range are currently underway at Jefferson Lab, USA, and MAMI in Mainz, Germany.

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