Higgs Bosons — H^0 and H^{\pm} , Searches for

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STANDARD MODEL H⁰ (Higgs Boson) MASS LIMITS

These limits apply to the Higgs boson of the three-generation Standard Model with the minimal Higgs sector. For a review and a bibliography, see the Note above on "Searches for Higgs Bosons."

H⁰ Direct Search Limits

Limits on the Standard Model Higgs obtained from the study of Z^0 decays rule out conclusively its existence in the whole mass region $m_{H^0} \lesssim 60$ GeV. These limits,

as well as stronger limits obtained from e^+e^- collisions at LEP at energies up to 202 GeV, and weaker limits obtained from other sources, have been superseded by the more recent data of LEP. They have been removed from this compilation, and are documented in previous editions of this Review of Particle Physics. The same holds for limits obtained from $p\overline{p}$ collisions at the Tevatron that have been superseded by more recent results incorporating a larger integrated luminosity.

In this Section, unless otherwise stated, limits from the four LEP experiments (ALEPH, DELPHI, L3, and OPAL) are obtained from the study of the $e^+e^- \rightarrow H^0Z$ process, at center-of-mass energies reported in the comment lines.

VALUE (GeV)	CL%	DOCUMENT ID		TECN	COMMENT
>114.1	95	¹ ABDALLAH	04	DLPH	$E_{ m cm} \leq$ 209 GeV
>112.7	95	$^{ m 1}$ abbiendi			$E_{\rm cm} \le 209 \text{ GeV}$
>114.4	95	1,2 HEISTER			$E_{\rm cm} \leq 209 \; {\rm GeV}$
>111.5	95	1,3 HEISTER	02	ALEP	$E_{\rm cm} \leq$ 209 GeV
>112.0	95	¹ ACHARD	01 C	L3	$E_{\rm cm}^{\rm off} \le 209 \; {\rm GeV}$

• • • We do not use the following data for averages, fits, limits, etc. • • •

⁴ AALTONEN	09A CDF	$p\overline{p} \rightarrow H^0 X, H^0 \rightarrow WW^{(*)}$
⁵ ABAZOV	09C D0	$p\overline{p} \rightarrow H^0 W X$
⁶ AALTONEN	08AF CDF	$p\overline{p} \rightarrow H^0 Z X$
⁷ AALTONEN	08V CDF	$p\overline{p} \rightarrow H^0 W X$
⁸ AALTONEN	08x CDF	$p\overline{p} \rightarrow H^0 ZX, H^0 WX$
⁹ ABAZOV	08AO D0	$p\overline{p} \rightarrow H^0 ZX, H^0 WX$
¹⁰ ABAZOV	08Y D0	$p\overline{p} \rightarrow H^0 W X$
¹¹ ABAZOV	07X D0	$p\overline{p} \rightarrow H^0 Z X$
¹² ABAZOV	06 D0	$p\overline{p} \rightarrow H^0 X, H^0 \rightarrow WW^*$
¹³ ABAZOV	060 D0	$p\overline{p} \rightarrow H^0 W X, H^0 \rightarrow W W^*$

¹ Search for $e^+e^- \to H^0 Z$ in the final states $H^0 \to b\overline{b}$ with $Z \to \ell\overline{\ell}$, $\nu\overline{\nu}$, $q\overline{q}$, $\tau^+\tau^-$ and $H^0 \to \tau^+\tau^-$ with $Z \to q\overline{q}$.

²Combination of the results of all LEP experiments.

 $^{^3}$ A 3σ excess of candidate events compatible with m_{H^0} near 114 GeV is observed in the combined channels $q \bar{q} q \bar{q}$, $q \bar{q} \ell \bar{\ell}$, $q \bar{q} \tau^+ \tau^-$.

- ⁴ AALTONEN 09A search for H^0 production in $p\overline{p}$ collisions at $E_{\rm cm}=1.96$ TeV in the decay mode $H^0\to WW^{(*)}\to \ell^+\ell^-\nu\overline{\nu}$. A limit on $\sigma(H^0)\cdot {\rm B}(H^0\to WW^{(*)})$ between 0.7 and 2.5 pb (95% CL) is given for $m_{H^0}=110$ –200 GeV, which is 1.7–45 times larger than the expected Standard Model cross section. The best limit is obtained for $m_{H^0}=160$ GeV.
- 5 ABAZOV 09C search for associated $H^0\,W$ production in $p\overline{p}$ collisions at $E_{\rm cm}=1.96$ TeV in the decay mode $H^0\to b\overline{b},\,W\to\ell\nu.$ A limit $\sigma(H^0\,W)\cdot {\rm B}(H^0\to b\overline{b})<(2.1–0.95)$ pb (95% CL) is given for $m_{H^0}=100$ –150 GeV, which is 9.1–84 times larger than the expected Standard Model cross section.
- ⁶ AALTONEN 08AF search for associated H^0 Z production in $p\overline{p}$ collisions at $E_{\rm cm}=1.96$ TeV in the decay mode $H^0\to b\overline{b}$, $Z\to \ell^+\ell^-$. A limit $\sigma(H^0Z)\cdot {\rm B}(H^0\to b\overline{b})<(1.5-1.2)$ pb (95% CL) is given for $m_{H^0}=110$ –150 GeV, which is 15–160 times larger than the expected Standard Model cross section.
- 7 AALTONEN 08V search for associated H^0 W production in $p\overline{p}$ collisions at $E_{\rm cm}=1.96$ TeV in the decay mode $H^0\to b\overline{b},~W\to \ell\nu.$ A limit $\sigma(H^0W)\cdot {\rm B}(H^0\to b\overline{b})<(3.9–1.3)$ pb (95% CL) is given for $m_{H^0}=110$ –150 GeV, which is 25–120 times larger than the expected Standard Model cross section.
- ⁸ AALTONEN 08X search for associated $H^0 Z$ and $H^0 W$ production in $p\overline{p}$ collisions at $E_{\rm cm}=1.96$ TeV in the decay mode $H^0 \to b\overline{b}$, $Z \to \nu\overline{\nu}$ and $W \to (\ell)\nu$, where ℓ is not detected. A limit $\sigma \cdot {\rm B}(H^0 \to b\overline{b}) < (4.7–3.3)$ pb (95% CL) is given for $m_{H^0}=110$ –140 GeV, which is 18–66 times larger than the expected Standard Model cross section.
- ⁹ABAZOV 08AO search for associated H^0Z and H^0W production in $p\overline{p}$ collisions at $E_{\rm cm}=1.96$ TeV in the decay mode $H^0\to b\overline{b}, Z\to \nu\overline{\nu}$ and $W\to (\ell)\nu$, where ℓ is not detected. A limit $\sigma + {\rm B}(H^0\to b\overline{b}) < (2.6-2.3)$ pb (95% CL) is given for $m_{H^0}=105-135$ GeV, which is 8.7–34 times larger than the expected Standard Model cross section.
- 10 ABAZOV 08Y search for associated H^0 W production in $p\overline{p}$ collisions at $E_{\rm cm}=1.96$ TeV in the decay mode $H^0\to b\overline{b},~W\to\ell\nu.$ A limit $\sigma(H^0W)\cdot {\rm B}(H^0\to b\overline{b})<(1.9–1.6)$ pb (95% CL) is given for $m_{H^0}=105–145$ GeV, which is 10–93 times larger than the expected Standard Model cross section. These results are combined with ABAZOV 06, ABAZOV 060, ABAZOV 06Q, and ABAZOV 07X to give cross section limits for $m_{H^0}=100–200$ GeV which are 6–24 times larger than the Standard Model expectation.
- ¹¹ ABAZOV 07X search for associated H^0Z production in $p\overline{p}$ collisions at $E_{\rm cm}=1.96$ TeV in the final state $Z\to e^+e^-$ or $\mu^+\mu^-$; $H^0\to b\overline{b}$. A limit $\sigma(ZH^0)\cdot {\rm B}(H^0\to b\overline{b})<(4.4-3.1)$ pb (95%CL) is given for $m_{H^0}=105-145$ GeV, which is more than 40 times larger than the expected Standard Model cross section.
- ¹² ABAZOV 06 search for Higgs boson production in $p\overline{p}$ collisions at $E_{\rm cm}=1.96$ TeV with the decay chain $H^0\to WW^*\to \ell^\pm\nu\ell'^\mp\overline{\nu}$. A limit $\sigma(H^0)\cdot {\rm B}(H^0\to WW^*)<(5.6-3.2)$ pb (95 %CL) is given for $m_{H^0}=120-200$ GeV, which far exceeds the expected Standard Model cross section.
- 13 ABAZOV 060 search for associated H^0 W production in $p\overline{p}$ collisions at $E_{\rm cm}=1.96$ TeV with the decay $H^0\to WW^*$, in the final states $\ell^\pm\ell'^\mp\nu\nu'$ X where $\ell=e,~\mu$. A limit $\sigma(H^0W)\cdot {\rm B}(H^0\to WW^*)<(3.2-2.8)$ pb (95 %CL) is given for $m_{H^0}=115-175$ GeV, which far exceeds the expected Standard Model cross section.

H⁰ Indirect Mass Limits from Electroweak Analysis

For limits obtained before the direct measurement of the top quark mass, see the 1996 (Physical Review **D54** 1 (1996)) Edition of this Review. Other studies based on data available prior to 1996 can be found in the 1998 Edition (The European

Physical Journal **C3** 1 (1998)) of this Review. For indirect limits obtained from other considerations of theoretical nature, see the Note on "Searches for Higgs Bosons."

VALUE (GeV)	DOCUMENT ID	TECN	
129 ⁺⁷⁴	¹⁴ LEP-SLC	06	RVUE

¹⁴ LEP-SLC 06 make Standard Model fits to Z parameters from LEP/SLC and m_t , m_W , and Γ_W measurements available in 2005 with $\Delta\alpha^{(5)}_{\rm had}(m_Z)=0.02758\pm0.00035$. The 95% CL limit is 285 GeV.

MASS LIMITS FOR NEUTRAL HIGGS BOSONS IN SUPERSYMMETRIC MODELS

The minimal supersymmetric model has two complex doublets of Higgs bosons. The resulting physical states are two scalars $[H_1^0 \text{ and } H_2^0, \text{ where we define } m_{H_1^0} < m_{H_2^0}]$, a pseudoscalar (A^0) , and a charged Higgs pair (H^\pm) . H_1^0 and H_2^0 are also called h and H in the literature. There are two free parameters in the theory which can be chosen to be m_{A^0} and $\tan\beta = v_2/v_1$, the ratio of vacuum expectation values of the two Higgs doublets. Tree-level Higgs masses are constrained by the model to be $m_{H_1^0} \leq m_Z$, $m_{H_2^0} \geq m_Z$, $m_{A^0} \geq m_{H_1^0}$, and $m_{H^\pm} \geq m_W$. However, as described in the review on "Searches for Higgs Bosons" in this Volume these relations are violated by radiative corrections.

Unless otherwise noted, the experiments in e^+e^- collisions search for the processes $e^+e^- \to H_1^0Z^0$ in the channels used for the Standard Model Higgs searches and $e^+e^- \to H_1^0A^0$ in the final states $b\overline{b}b\overline{b}$ and $b\overline{b}\tau^+\tau^-$. In $p\overline{p}$ collisions the experiments search for a variety of processes, as explicitly specified for each entry. Limits on the A^0 mass arise from these direct searches, as well as from the relations valid in the minimal supersymmetric model between m_{A^0} and $m_{H_1^0}$. As discussed in the review on "Searches for Higgs Bosons" in this Volume, these relations depend, via potentially large radiative corrections, on the mass of the t quark and on the supersymmetric parameters, in particular those of the stop sector. The limits are weaker for larger t and t masses. To include the radiative corrections to the Higgs masses, unless otherwise stated, the listed papers use theoretical predictions incorporating two-loop corrections and examine the two scenarios of no scalar top mixing and the $m_h^{\rm max}$ benchmark scenario (which gives rise to the most conservative upper bound on the mass of H_1^0 for given values of m_{A^0} and $\tan\beta$), see CARENA 998 and CARENA 03.

Limits in the low-mass region of H_1^0 , as well as other by now obsolete limits from different techniques, have been removed from this compilation, and can be found in earlier editions of this Review. Unless otherwise stated, the following results assume no invisible H_1^0 or A^0 decays.

H₁⁰ (Higgs Boson) MASS LIMITS in Supersymmetric Models

VALUE (GeV)	CL%	DOCUMENT ID		TECN	COMMENT
>89.7		¹⁵ ABDALLAH	08 B	DLPH	$E_{ m cm} \leq 209 \; { m GeV}$
>92.8	95	¹⁶ SCHAEL	06 B	LEP	$E_{\rm cm} \leq 209 \; {\rm GeV}$
>84.5		^{17,18} ABBIENDI	04M	OPAL	$E_{\rm cm} \leq 209 \; {\rm GeV}$
>86.0	95	^{17,19} ACHARD	02H	L3	$E_{ m cm}^{ m cm} \leq$ 209 GeV, $ an eta > 0.4$

• • • We do not use the following data for averages, fits, limits, etc. • • •

¹⁵ ABDALLAH 08B give limits in eight *CP*-conserving benchmark scenarios and some *CP*-violating scenarios. See paper for excluded regions for each scenario. Supersedes AB-DALLAH 04.

¹⁶ SCHAEL 06B make a combined analysis of the LEP data. The quoted limit is for the $m_h^{\rm max}$ scenario with $m_t=174.3$ GeV. In the *CP*-violating CPX scenario no lower bound on $m_{H_1^0}$ can be set at 95% CL. See paper for excluded regions in various scenarios. See Figs. 2–6 and Tabs. 14–21 for limits on $\sigma(ZH^0)$ · B($H^0\to b\overline{b}, \tau^+\tau^-$) and $\sigma(H^0_1H^0_2)$ · B($H^0_1, H^0_2\to b\overline{b}, \tau^+\tau^-$).

¹⁷ Search for $e^+e^- \rightarrow H_1^0 A^0$ in the final states $b\overline{b}b\overline{b}$ and $b\overline{b}\tau^+\tau^-$, and $e^+e^- \rightarrow H_1^0 Z$. Universal scalar mass of 1 TeV, SU(2) gaugino mass of 200 GeV, and $\mu=-200$ GeV are assumed, and two-loop radiative corrections incorporated. The limits hold for m_t =175 GeV, and for the $m_h^{\rm max}$ scenario.

¹⁸ ABBIENDI 04M exclude $0.7 < \tan \beta < 1.9$, assuming $m_t = 174.3$ GeV. Limits for other MSSM benchmark scenarios, as well as for *CP* violating cases, are also given.

MSSM benchmark scenarios, as well as for CP violating cases, are also given. 19 ACHARD 02H also search for the final state $H_1^0Z \to 2A^0q\overline{q}$, $A^0 \to q\overline{q}$. In addition, the MSSM parameter set in the "large- μ " and "no-mixing" scenarios are examined.

²⁰ ABAZOV 09F search for associated production of a Higgs boson and a b quark in $p\overline{p}$ collisions at $E_{\rm cm}=1.96$ TeV with the decay $H_{1,2}^0/A^0\to \tau^+\tau^-$. See their Fig. 2 for the limit on $\sigma\cdot {\rm B}(H_{1,2}^0/A^0\to \tau^+\tau^-)$ (for different Higgs masses) and for the excluded region in the MSSM parameter space for $\mu=\pm 200$ GeV.

²¹ ABAZOV 08AJ search for associated production of a Higgs boson and a b quark in $p\overline{p}$ collisions at $E_{\rm cm}=1.96$ TeV with the decay $H_{1,2}^0/A^0\to b\overline{b}$. See their Tab. 3 for the limit on $\sigma\cdot {\rm B}(H_{1,2}^0/A^0\to b\overline{b})$ for different Higgs masses, and see their Fig. 3 for the excluded region in the MSSM parameter space for $\mu=\pm 200$ GeV.

²² ABAZOV 08W search for Higgs boson production in $p \overline{p}$ collisions at $E_{\rm cm}=1.96$ TeV with the decay $H_{1,2}^0/A^0 \to \tau^+\tau^-$. See their Fig. 3 for the limit on $\sigma \cdot {\rm B}(H_{1,2}^0/A^0 \to \tau^+\tau^-)$ for different Higgs masses, and see their Fig. 4 for the excluded region in the MSSM parameter space.

- ²³ This limit applies also in the no-mixing scenario. Furthermore, ABDALLAH 04 excludes the range $0.54 < \tan\beta < 2.36$. The limit improves in the region $\tan\beta < 6$ (see Fig. 28). Limits for $\mu = 1$ TeV are given in Fig. 30.
- ²⁴ ABBIENDI 03G search for $e^+e^- \rightarrow H_1^0Z$ followed by $H_1^0 \rightarrow A^0A^0$, $A^0 \rightarrow c\overline{c}$, gg, or $\tau^+\tau^-$. In the no-mixing scenario, the region $m_{H_1^0}=45$ -85 GeV and $m_{A^0}=2$ -9.5 GeV is excluded at 95% CL.
- ²⁵ HEISTER 02 excludes the range 0.7 <tan β < 2.3. A wider range is excluded with different stop mixing assumptions. Updates BARATE 01C.

A⁰ (Pseudoscalar Higgs Boson) MASS LIMITS in Supersymmetric Models

VALUE (GeV)	CL%	DOCUMENT ID		TECN	COMMENT
>90.4		²⁶ ABDALLAH	08 B	DLPH	$E_{ m cm} \leq 209 \; { m GeV}$
>93.4	95	²⁷ SCHAEL			$E_{\rm cm} \le 209 {\rm GeV}$
>85.0		^{28,29} ABBIENDI	04M	OPAL	$E_{\rm cm} \leq 209 \; {\rm GeV}$
>86.5		^{28,30} ACHARD	02H	L3	$E_{\rm cm} \leq$ 209 GeV, $\tan \beta > 0.4$
>90.1	95	^{28,31} HEISTER	02	ALEP	$E_{ m cm}^{ m T} \leq$ 209 GeV, $ an eta > 0.5$

• • • We do not use the following data for averages, fits, limits, etc. • • •

- ²⁶ ABDALLAH 08B give limits in eight *CP*-conserving benchmark scenarios and some *CP*-violating scenarios. See paper for excluded regions for each scenario. Supersedes AB-DALLAH 04.
- ²⁷ SCHAEL 06B make a combined analysis of the LEP data. The quoted limit is for the $m_h^{\rm max}$ scenario with $m_t=174.3$ GeV. In the *CP*-violating CPX scenario no lower bound on $m_{H_1^0}$ can be set at 95% CL. See paper for excluded regions in various scenarios. See Figs. 2–6 and Tabs. 14–21 for limits on $\sigma(ZH^0)\cdot {\rm B}(H^0\to b\overline{b}, \tau^+\tau^-)$ and $\sigma(H_1^0H_2^0)\cdot {\rm B}(H^0\to b\overline{b}, \tau^+\tau^-)$
- ²⁸ Search for $e^+e^- \rightarrow H_1^0 A^0$ in the final states $b \overline{b} b \overline{b}$ and $b \overline{b} \tau^+ \tau^-$, and $e^+e^- \rightarrow H_1^0 Z$. Universal scalar mass of 1 TeV, SU(2) gaugino mass of 200 GeV, and $\mu = -200$ GeV are assumed, and two-loop radiative corrections incorporated. The limits hold for $m_t = 175$ GeV, and for the m_h^{max} scenario.
- ²⁹ ABBIENDI 04M exclude 0.7 $< \tan \beta < 1.9$, assuming $m_t = 174.3$ GeV. Limits for other MSSM benchmark scenarios, as well as for *CP* violating cases, are also given.
- 30 ACHARD 02H also search for the final state $H_1^0 Z \to 2A^0 \, q \, \overline{q}, \, A^0 \to q \, \overline{q}$. In addition, the MSSM parameter set in the "large- μ " and "no-mixing" scenarios are examined.
- 31 HEISTER 02 excludes the range 0.7 <tan β < 2.3. A wider range is excluded with different stop mixing assumptions. Updates BARATE 01C.

 $B(H_1^0, H_2^0 \rightarrow b\overline{b}, \tau^+\tau^-).$

- 32 ABAZOV 09F search for associated production of a Higgs boson and a b quark in $p\overline{p}$ collisions at $E_{\rm cm}=1.96$ TeV with the decay $H_{1.2}^0/A^0 \to \tau^+\tau^-$. See their Fig. 2 for the limit on $\sigma \cdot \mathrm{B}(H_{1\ 2}^0/A^0 \to\ \tau^+\,\tau^-)$ (for different Higgs masses) and for the excluded region in the MSSM parameter space for $\mu=\pm 200$ GeV.
- 33 ABAZOV 08AJ search for associated production of a Higgs boson and a b quark in $p\overline{p}$ collisions at $E_{\rm cm}=1.96$ TeV with the decay $H_{1.2}^0/A^0\to b\overline{b}$. See their Tab. 3 for the limit on $\sigma \cdot \mathsf{B}(H_{1.2}^0/A^0 \to b\,\overline{b})$ for different Higgs masses, and see their Fig. 3 for the excluded region in the MSSM parameter space for $\mu=\pm 200$ GeV.
- ³⁴ ABAZOV 08W search for Higgs boson production in $p\bar{p}$ collisions at $E_{\rm cm}=1.96$ TeV with the decay $H^0_{1,2}/A^0 \to \ \tau^+\,\tau^-$. See their Fig. 3 for the limit on $\sigma\cdot {\rm B}(H^0_{1,2}/A^0 \to \pi^+\,\tau^-)$ $au^+ au^-$) for different Higgs masses, and see their Fig. 4 for the excluded region in the MSSM parameter space.
- ³⁵ ABULENCIA 06 search for $H_{1,2}^0/A^0$ production in $p\bar{p}$ collisions at $E_{\rm cm}=1.96$ TeV with $H_{1,2}^0/A^0 \to \tau^+\tau^-$. A region with $\tan\beta > 40$ (100) is excluded for $m_{A^0} = 90$
- ³⁶ ACOSTA 05Q search for $H_{1,2}^0/A^0$ production in $p\overline{p}$ collisions at $E_{\rm cm}=1.8$ TeV with $H_{1.2}^0/{\it A}^0 \rightarrow \ \tau^+ \tau^-$. At $m_{{\it A}^0}=$ 100 GeV, the obtained cross section upper limit is above theoretical expectation.
- 37 This limit applies also in the no-mixing scenario. Furthermore, ABDALLAH 04 excludes the range $0.54 < \tan \beta < 2.36$. The limit improves in the region $\tan \beta < 6$ (see Fig. 28). Limits for $\mu=1$ TeV are given in Fig. 30.
- ³⁸ ABBIENDI 03G search for $e^+e^- \rightarrow H_1^0Z$ followed by $H_1^0 \rightarrow A^0A^0$, $A^0 \rightarrow c\overline{c}$, gg, or $au^+ au^-$. In the no-mixing scenario, the region $m_{H_1^0} =$ 45-85 GeV and $m_{A^0} =$ 2-9.5 GeV is excluded at 95% CL.
- ³⁹ AKEROYD 02 examine the possibility of a light A^0 with $\tan \beta < 1$. Electroweak measurements are found to be inconsistent with such a scenario.

H⁰ (Higgs Boson) MASS LIMITS in Extended Higgs Models

This Section covers models which do not fit into either the Standard Model or its simplest minimal Supersymmetric extension (MSSM), leading to anomalous production rates, or nonstandard final states and branching ratios. In particular, this Section covers limits which may apply to generic two-Higgs-doublet models (2HDM), or to special regions of the MSSM parameter space where decays to invisible particles or to photon pairs are dominant (see the Note on 'Searches for Higgs Bosons' at the beginning of this Chapter). See the footnotes or the comment lines for details on the nature of the models to which the limits apply.

VALUE (GeV)	CL%	DOCUMENT ID		TECN	COMMENT		
		⁴⁰ ABAZOV	08 U	D0	$H^0 \rightarrow \gamma \gamma$		
		⁴¹ LOVE	80	CLEO	$\Upsilon(1S) \rightarrow A^0 \gamma$		
			07	OPAL	invisible H^0 , large width		
		⁴³ BESSON			$\Upsilon(1S) ightarrow \ \eta_{m{b}} \gamma$		
>105.8	95	⁴⁴ SCHAEL	07	ALEP	$e^{+}e^{-} ightarrow$ $H^{0}Z,\;H^{0} ightarrow$		
none 1–55 none 3–63 >110.6	95 95 95	⁴⁵ ABBIENDI ⁴⁵ ABBIENDI ⁴⁶ ABDALLAH	05A	OPAL	WW^* H_1^0 , Type II model A^0 , Type II model $H^0 \rightarrow 2$ jets		
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>112.3	95	⁴⁷ ACHARD	05	L3	invisible H^0
>104	95	⁴⁸ ABBIENDI	04K	OPAL	$H^0 \rightarrow 2$ jets
		⁴⁹ ABDALLAH	04	DLPH	H ⁰ V V couplings
>112.1	95	⁴⁷ ABDALLAH	04 B	DLPH	Invisible H ⁰
>104.1	95	^{50,51} ABDALLAH	04L	DLPH	$e^+e^- \rightarrow H^0Z$, $H^0 \rightarrow \gamma\gamma$
		⁵² ABDALLAH	040	DLPH	$Z \rightarrow f \overline{f} H$
		⁵³ ABDALLAH	040	DLPH	$e^+e^- \rightarrow H^0Z$, H^0A^0
>110.3	95	⁵⁴ ACHARD	04 B	L3	$H^0 \rightarrow 2$ jets
		⁵⁵ ACHARD	04F	L3	Anomalous coupling
		⁵⁶ ABBIENDI	03F	OPAL	$e^+e^- o H^0Z$, $H^0 o$ any
		⁵⁷ ABBIENDI	03 G	OPAL	$H_1^0 \rightarrow A^0 A^0$
>107	95	⁵⁸ ACHARD	03 C	L3	$H^{ar{f 0}} ightarrow ~WW^*$, ZZ^* , $\gamma\gamma$
		⁵⁹ ABBIENDI	02 D	OPAL	$e^+e^- \rightarrow b\overline{b}H$
>105.5	95	^{50,60} ABBIENDI	02F	OPAL	$H_1^0 \rightarrow \gamma \gamma$
>105.4	95	⁶¹ ACHARD	02 C	L3	$H_1^{\bar{0}} \rightarrow \gamma \gamma$
>114.1	95	⁴⁷ HEISTER	02	ALEP	Invisible H^0 , $E_{\rm cm} \le 209 \text{ GeV}$
>105.4	95	^{50,62} HEISTER	02L	ALEP	$H_1^0 \rightarrow \gamma \gamma$
>109.1	95	⁶³ HEISTER	02м	ALEP	$H^{\clip{0}} ightarrow 2$ jets or $ au^+ au^-$
none 1–44	95	⁶⁴ ABBIENDI	01E	OPAL	H_1^0 , Type-II model
none 12-56	95	⁶⁴ ABBIENDI	01E	OPAL	A^{0} , Type-II model
> 98	95	⁶⁵ AFFOLDER	01н	CDF	$p\overline{p} \rightarrow H^0 W/Z, H^0 \rightarrow \gamma\gamma$
>106.4	95	⁴⁷ BARATE	01 C	ALEP	Invisible H^0 , $E_{cm} \leq 202 \text{ GeV}$
> 89.2	95	⁶⁶ ACCIARRI	00м	L3	Invisible H ⁰
		⁶⁷ ACCIARRI	00 R	L3	${ m e^+ e^-} ightarrow ~H^0 \gamma { m \ and/or} ~H^0 ightarrow$
		68			77
		68 ACCIARRI	00R	L3	$e^+e^- \rightarrow e^+e^-H^0$
> 94.9	95	69 ACCIARRI	00s	L3	$e^+e^- \rightarrow H^0Z, H^0 \rightarrow \gamma\gamma$
>100.7	95	70 BARATE	00L	ALEP	$e^+e^- \rightarrow H^0Z, H^0 \rightarrow \gamma\gamma$
> 68.0	95	⁷¹ ABBIENDI ⁷² ABBIENDI	99E	OPAL	$\tan \beta > 1$ $e^+e^- \rightarrow H^0Z, H^0 \rightarrow \gamma \gamma$
> 96.2	95 05	73 ABBOTT	990	OPAL	$e^+e^- \rightarrow H^0Z, H^0 \rightarrow \gamma\gamma$ $p\overline{p} \rightarrow H^0W/Z, H^0 \rightarrow \gamma\gamma$
> 78.5	95	74 ABREU	99B	D0 DLPH	$e^+e^- \rightarrow H^0 V/Z$, $H^0 \rightarrow \gamma \gamma$
		· · ABREU	99P	DLPH	$e^+e^- \rightarrow H^*\gamma$ and/or $H^* \rightarrow \gamma\gamma$
		⁷⁵ GONZALEZ-G.	. 98 B	RVUE	Anomalous coupling
		⁷⁶ KRAWCZYK	97	RVUE	$(g-2)_{\mu}$
		77 ALEXANDER	96н	OPAL	$Z \rightarrow H^0 \gamma$
		⁷⁸ ABREU	95H	DLPH	$Z \to H^0 Z^*, H^0 A^0$
		⁷⁹ PICH	92	RVUE	Very light Higgs

⁴⁰ ABAZOV 08U search for $H^0 \to \gamma \gamma$ in $p\overline{p}$ collisions at $E_{\rm cm}=1.96$ TeV in the mass range $m_{H^0}=70$ –150 GeV. Associated H^0 W, H^0 Z production and W W, Z Z fusion are considered. See their Tab. 1 for the limit on $\sigma \cdot {\rm B}(H^0 \to \gamma \gamma)$, and see their Fig. 3 for the excluded region in the $m_{H^0} - {\rm B}(H^0 \to \gamma \gamma)$ plane.

for the excluded region in the m_{H^0} — B($H^0 \to \gamma \gamma$) plane. ⁴¹ LOVE 08 search for the process $\Upsilon(1S) \to A^0 \gamma$ with $A^0 \to \mu^+ \mu^-$ (for $m_{A^0} < 2m_{\tau}$) and $A^0 \to \tau^+ \tau^-$. Limits on B($\Upsilon(1S) \to A^0 \gamma$) · B($A^0 \to \ell^+ \ell^-$) in the range 10^{-6} – 10^{-4} (90% CL) are given.

⁴² ABBIENDI 07 search for $e^+e^- \to H^0 Z$ with $Z \to q \overline{q}$ and H^0 decaying to invisible final states. The H^0 width is varied between 1 GeV and 3 TeV. A limit $\sigma \cdot B(H^0 \to \text{invisible})$ < (0.07–0.57) pb (95%CL) is obtained at $E_{\text{cm}} = 206$ GeV for $m_{H^0} = 60$ –114 GeV.

- ⁴³ BESSON 07 give a limit B($\Upsilon(1S) \rightarrow \eta_b \gamma$) · B($\eta_b \rightarrow \tau^+ \tau^-$) < 0.27% (95% CL), which constrains a possible A^0 exchange contribution to the η_b decay.
- ⁴⁴ SCHAEL 07 search for Higgs bosons in association with a fermion pair and decaying to WW^* . The limit is from this search and HEISTER 02L for a H^0 with SM production cross section and $B(H^0 \to f\overline{f}) = 0$ for all fermions f.
- ⁴⁵ ABBIENDI 05A search for $e^+e^- \to H_1^0\,A^0$ in general Type-II two-doublet models, with decays $H_1^0,\,A^0 \to q\,\overline{q},\,g\,g,\,\tau^+\tau^-$, and $H_1^0 \to A^0\,A^0$.
- ⁴⁶ ABDALLAH 05D search for $e^+e^- \rightarrow H^0Z$ and H^0A^0 with H^0 , A^0 decaying to two jets of any flavor including gg. The limit is for SM H^0Z production cross section with $B(H^0 \rightarrow jj) = 1$.
- ⁴⁷ Search for $e^+e^- \rightarrow H^0 Z$ with H^0 decaying invisibly. The limit assumes SM production cross section and B($H^0 \rightarrow$ invisible) = 1.
- ⁴⁸ ABBIENDI 04K search for $e^+e^- \rightarrow H^0Z$ with H^0 decaying to two jets of any flavor including gg. The limit is for SM production cross section with $B(H^0 \rightarrow jj) = 1$.
- 49 ABDALLAH 04 consider the full combined LEP and LEP2 datasets to set limits on the Higgs coupling to W or Z bosons, assuming SM decays of the Higgs. Results in Fig. 26.
- ⁵⁰ Search for associated production of a $\gamma\gamma$ resonance with a Z boson, followed by $Z\to q\overline{q}$, $\ell^+\ell^-$, or $\nu\overline{\nu}$, at $E_{\rm cm}\leq$ 209 GeV. The limit is for a H^0 with SM production cross section and B($H^0\to f\overline{f}$)=0 for all fermions f.
- ⁵¹ Updates ABREU 01F.
- ⁵² ABDALLAH 040 search for $Z \to b\overline{b}H^0$, $b\overline{b}A^0$, $\tau^+\tau^-H^0$ and $\tau^+\tau^-A^0$ in the final states 4b, $b\overline{b}\tau^+\tau^-$, and 4τ . See paper for limits on Yukawa couplings.
- ⁵³ ABDALLAH 040 search for $e^+e^- \rightarrow H^0 Z$ and $H^0 A^0$, with H^0 , A^0 decaying to $b \overline{b}$, $\tau^+\tau^-$, or $H^0 \rightarrow A^0 A^0$ at $E_{\rm cm}=189$ –208 GeV. See paper for limits on couplings.
- ⁵⁴ ACHARD 04B search for $e^+e^- \to H^0 Z$ with H^0 decaying to $b\overline{b}$, $c\overline{c}$, or gg. The limit is for SM production cross section with B($H^0 \to jj$) = 1.
- ⁵⁵ ACHARD 04F search for H^0 with anomalous coupling to gauge boson pairs in the processes $e^+e^- \to H^0\gamma$, $e^+e^-H^0$, H^0Z with decays $H^0 \to f\overline{f}$, $\gamma\gamma$, $Z\gamma$, and W^*W at $E_{\rm cm}=189$ –209 GeV. See paper for limits.
- ⁵⁶ ABBIENDI 03F search for $H^0 \to \text{anything in } e^+e^- \to H^0 Z$, using the recoil mass spectrum of $Z \to e^+e^-$ or $\mu^+\mu^-$. In addition, it searched for $Z \to \nu \overline{\nu}$ and $H^0 \to e^+e^-$ or photons. Scenarios with large width or continuum H^0 mass distribution are considered. See their Figs. 11–14 for the results.
- ⁵⁷ ABBIENDI 03G search for $e^+e^- \rightarrow H_1^0 Z$ followed by $H_1^0 \rightarrow A^0 A^0$, $A^0 \rightarrow c \overline{c}$, gg, or $\tau^+\tau^-$ in the region $m_{H_1^0}=$ 45-86 GeV and $m_{A^0}=$ 2-11 GeV. See their Fig. 7 for the limits.
- the limits. 58 ACHARD 03C search for $e^+e^- \to ZH^0$ followed by $H^0 \to WW^*$ or ZZ^* at $E_{\rm cm}=$ 200-209 GeV and combine with the ACHARD 02C result. The limit is for a H^0 with SM production cross section and B($H^0 \to f\overline{f}$) = 0 for all f. For B($H^0 \to WW^*$) + B($H^0 \to ZZ^*$) = 1, m $_{H^0} >$ 108.1 GeV is obtained. See fig. 6 for the limits under different BR assumptions.
- ^59 ABBIENDI 02D search for $Z \to b \overline{b} H_1^0$ and $b \overline{b} A^0$ with $H_1^0/A^0 \to \tau^+ \tau^-$, in the range $4 < m_H <$ 12 GeV. See their Fig. 8 for limits on the Yukawa coupling.
- ⁶⁰ For B($H^0 \rightarrow \gamma \gamma$)=1, $m_{H^0} >$ 117 GeV is obtained.
- ⁶¹ ACHARD 02C search for associated production of a $\gamma\gamma$ resonance with a Z boson, followed by $Z \to q \overline{q}$, $\ell^+ \ell^-$, or $\nu \overline{\nu}$, at $E_{\rm CM} \le 209$ GeV. The limit is for a H^0 with SM production cross section and B($H^0 \to f \overline{f}$)=0 for all fermions f. For B($H^0 \to \gamma\gamma$)=1, $m_{H^0} > 114$ GeV is obtained.

- $^{62}\,\mathrm{For}\,\,\mathrm{B}(\mathrm{H}^0\to~\gamma\gamma){=}1,~m_{\mathrm{H}^0}>113.1~\mathrm{GeV}$ is obtained.
- ⁶³ HEISTER 02M search for $e^+e^- \rightarrow H^0 Z$, assuming that H^0 decays to $q \overline{q}$, g g, or $\tau^+\tau^-$ only. The limit assumes SM production cross section.
- ⁶⁴ ABBIENDI 01E search for neutral Higgs bosons in general Type-II two-doublet models, at $E_{\rm cm} \leq$ 189 GeV. In addition to usual final states, the decays H_1^0 , $A^0 \rightarrow q \overline{q}$, g g are searched for. See their Figs. 15,16 for excluded regions.
- ⁶⁵ AFFOLDER 01H search for associated production of a $\gamma\gamma$ resonance and a W or Z (tagged by two jets, an isolated lepton, or missing E_T). The limit assumes Standard Model values for the production cross section and for the couplings of the H^0 to W and Z bosons. See their Fig. 11 for limits with $B(H^0 \to \gamma\gamma) < 1$.
- ⁶⁶ ACCIARRI 00M search for $e^+e^- \to ZH^0$ with H^0 decaying invisibly at $E_{\rm cm}=183-189$ GeV. The limit assumes SM production cross section and B($H^0 \to {\rm invisible})=1$. See their Fig. 6 for limits for smaller branching ratios.
- ⁶⁷ ACCIARRI 00R search for $e^+e^- \to H^0\gamma$ with $H^0 \to b\overline{b}$, $Z\gamma$, or $\gamma\gamma$. See their Fig. 3 for limits on $\sigma \cdot B$. Explicit limits within an effective interaction framework are also given, for which the Standard Model Higgs search results are used in addition.
- ⁶⁸ ACCIARRI 00R search for the two-photon type processes $e^+e^- \rightarrow e^+e^-H^0$ with $H^0 \rightarrow b\overline{b}$ or $\gamma\gamma$. See their Fig. 4 for limits on $\Gamma(H^0 \rightarrow \gamma\gamma)\cdot B(H^0 \rightarrow \gamma\gamma)$ or $b\overline{b}$ for m_{H^0} =70–170 GeV.
- ⁶⁹ ACCIARRI 00S search for associated production of a $\gamma\gamma$ resonance with a $q\overline{q}$, $\nu\overline{\nu}$, or $\ell^+\ell^-$ pair in e^+e^- collisions at $E_{\rm cm}=$ 189 GeV. The limit is for a H^0 with SM production cross section and B($H^0\to f\overline{f}$)=0 for all fermions f. For B($H^0\to \gamma\gamma$)=1, $m_{H^0}>$ 98 GeV is obtained. See their Fig. 5 for limits on B($H\to \gamma\gamma$)· $\sigma(e^+e^-\to Hf\overline{f})/\sigma(e^+e^-\to Hf\overline{f})$ (SM).
- ⁷⁰ BARATE 00L search for associated production of a $\gamma\gamma$ resonance with a $q\overline{q}$, $\nu\overline{\nu}$, or $\ell^+\ell^-$ pair in e^+e^- collisions at $E_{\rm cm}=$ 88–202 GeV. The limit is for a H^0 with SM production cross section and B($H^0\to f\overline{f}$)=0 for all fermions f. For B($H^0\to \gamma\gamma$)=1, $m_{H^0}>$ 109 GeV is obtained. See their Fig. 3 for limits on B($H\to \gamma\gamma$)· σ ($e^+e^-\to Hf\overline{f}$)/ σ ($e^+e^-\to Hf\overline{f}$) (SM).
- ⁷¹ ABBIENDI 99E search for $e^+e^- \to H^0A^0$ and H^0Z at $E_{\rm cm}=183$ GeV. The limit is with $m_H=m_A$ in general two Higgs-doublet models. See their Fig. 18 for the exclusion limit in the m_H-m_A plane. Updates the results of ACKERSTAFF 98S.
- ⁷² ABBIENDI 990 search for associated production of a $\gamma\gamma$ resonance with a $q\overline{q}$, $\nu\overline{\nu}$, or $\ell^+\ell^-$ pair in e^+e^- collisions at 189 GeV. The limit is for a H^0 with SM production cross section and B($H^0 \to f\overline{f}$)=0, for all fermions f. See their Fig. 4 for limits on $\sigma(e^+e^- \to H^0Z^0)\times B(H^0 \to \gamma\gamma)\times B(X^0 \to f\overline{f})$ for various masses. Updates the results of ACKERSTAFF 98Y.
- 73 ABBOTT 99B search for associated production of a $\gamma\gamma$ resonance and a dijet pair. The limit assumes Standard Model values for the production cross section and for the couplings of the H^0 to W and Z bosons. Limits in the range of $\sigma(H^0+Z/W)\cdot B(H^0\to\gamma\gamma)=0.80$ –0.34 pb are obtained in the mass range $m_{H^0}=65$ –150 GeV.
- ⁷⁴ ABREU 99P search for $e^+e^- \rightarrow H^0\gamma$ with $H^0 \rightarrow b \, \overline{b}$ or $\gamma\gamma$, and $e^+e^- \rightarrow H^0 \, q \, \overline{q}$ with $H^0 \rightarrow \gamma\gamma$. See their Fig. 4 for limits on $\sigma\times B$. Explicit limits within an effective interaction framework are also given.
- 75 GONZALEZ-GARCIA 98B use DØ limit for $\gamma\gamma$ events with missing E_{T} in $p\overline{p}$ collisions (ABBOTT 98) to constrain possible ZH or WH production followed by unconventional $H\to\gamma\gamma$ decay which is induced by higher-dimensional operators. See their Figs. 1 and 2 for limits on the anomalous couplings.
- ⁷⁶ KRAWCZYK 97 analyse the muon anomalous magnetic moment in a two-doublet Higgs model (with type II Yukawa couplings) assuming no H_1^0 ZZ coupling and obtain $m_{H_1^0} \gtrsim$

5 GeV or $m_{A^0} \gtrsim$ 5 GeV for $\tan \beta >$ 50. Other Higgs bosons are assumed to be much heavier.

heavier. 77 ALEXANDER 96H give B($Z \to H^0 \gamma$)×B($H^0 \to q \overline{q}$) < 1–4 × 10⁻⁵ (95%CL) and B($Z \to H^0 \gamma$)×B($H^0 \to b \overline{b}$) < 0.7–2 × 10⁻⁵ (95%CL) in the range 20 < m_{H^0} <80 GeV.

⁷⁸ See Fig. 4 of ABREU 95H for the excluded region in the $m_{H^0}-m_{A^0}$ plane for general two-doublet models. For $\tan\beta>1$, the region $m_{H^0}+m_{A^0}\lesssim 87$ GeV, $m_{H^0}<47$ GeV is excluded at 95% CL

excluded at 95% CL. $^{79}\,\mathrm{PICH}$ 92 analyse H^0 with $m_{H^0}^{}<\!2m_{\mu}^{}$ in general two-doublet models. Excluded regions in the space of mass-mixing angles from LEP, beam dump, and π^{\pm} , η rare decays are shown in Figs. 3,4. The considered mass region is not totally excluded.

H[±] (Charged Higgs) MASS LIMITS

Unless otherwise stated, the limits below assume B($H^+ \to \tau^+ \nu$)+B($H^+ \to c \bar{s}$)=1, and hold for all values of B($H^+ \to \tau^+ \nu_{\tau}$), and assume H^+ weak isospin of T_3 =+1/2. In the following, $\tan\beta$ is the ratio of the two vacuum expectation values in two-doublet models (2HDM).

The limits are also applicable to point-like technipions. For a discussion of techniparticles, see the Review of Dynamical Electroweak Symmetry Breaking in this Review.

For limits obtained in hadronic collisions before the observation of the top quark, and based on the top mass values inconsistent with the current measurements, see the 1996 (Physical Review **D54** 1 (1996)) Edition of this Review.

Searches in e^+e^- collisions at and above the Z pole have conclusively ruled out the existence of a charged Higgs in the region $m_{H^+}\lesssim 45$ GeV, and are now superseded by the most recent searches in higher energy e^+e^- collisions at LEP. Results by now obsolete are therefore not included in this compilation, and can be found in the previous Edition (The European Physical Journal **C15** 1 (2000)) of this Review.

In the following, and unless otherwise stated, results from the LEP experiments (ALEPH, DELPHI, L3, and OPAL) are assumed to derive from the study of the $e^+e^- \rightarrow H^+H^-$ process. Limits from $b \rightarrow s \gamma$ decays are usually stronger in generic 2HDM models than in Supersymmetric models.

VALUE (GeV)	CL%	DOCUMENT ID		TECN	COMMENT
> 74.4	95	ABDALLAH	041	DLPH	$E_{ m cm} \leq 209 \; { m GeV}$
> 76.5	95	ACHARD	03E	L3	$E_{\rm cm} \leq$ 209 GeV
> 79.3	95	HEISTER	02 P	ALEP	$E_{\rm cm} \leq$ 209 GeV
• • • We do not	use the fo	lowing data for a	averag	ges, fits,	limits, etc. • •
> 92.0 > 76.7	95 95 8	⁰ ABULENCIA ABBIENDI ¹ ABDALLAH ² ABBIENDI	04 04I	OPAL DLPH	B(au u) = 1 Type I
	8 8 8	³ ABAZOV ⁴ BORZUMATI ⁵ ABBIENDI ⁶ BARATE	02B 02 01Q 01E	D0 RVUE OPAL ALEP	$ au ightarrow \mu \overline{ u} \nu, e \overline{ u} \nu$ $t ightarrow b H^+, H ightarrow au u$ $B ightarrow au u_{ au} X$ $B ightarrow au u_{ au}$
>315		⁷ GAMBINO ⁸ AFFOLDER			$b \rightarrow s \gamma$ $t \rightarrow bH^+, H \rightarrow \tau \nu$
> 59.5	95	ABBIENDI	99E	OPAL	$E_{cm} \leq 183 \; GeV$

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<sup>89</sup> ABBOTT
                                                                      99E D0
                                         ^{90} ACKERSTAFF 99D OPAL 	au	o e
u
u, \mu
u
u
                                         <sup>91</sup> ACCIARRI
                                                                      97F L3
                                                                                            B \rightarrow \tau \nu_{\tau}
                                         <sup>92</sup> AMMAR
                                                                      97B CLEO 	au 	o 	au 
u 
u
                                         <sup>93</sup> COARASA
                                                                      97 RVUE B \rightarrow \tau \nu_{\tau} X
97 RVUE t \rightarrow b H^{+}, H \rightarrow \tau \nu
                                         <sup>94</sup> GUCHAIT
                                         <sup>95</sup> MANGANO
                                                                      97 RVUE B_{u(c)} \rightarrow \tau \nu_{\tau}
                                         <sup>96</sup> STAHL
                                                                      97 RVUE \tau \rightarrow \mu \nu \nu
                                         <sup>97</sup> ALAM
                                                                      95 CLE2 b \rightarrow s \gamma
95 ALEP b \rightarrow \tau \nu_{\tau} X
                           95
>244
                                         <sup>98</sup> BUSKULIC
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- ABULENCIA 06E search for associated H^0 W production in $p\overline{p}$ collisions at $E_{\rm cm}=1.96$ TeV. A fit is made for $t\overline{t}$ production processes in dilepton, lepton + jets, and lepton + τ final states, with the decays $t\to W^+b$ and $t\to H^+b$ followed by $H^+\to \tau^+\nu$, $c\overline{s}$, $t^*\overline{b}$, or W^+H^0 . Within the MSSM the search is sensitive to the region $\tan\beta<1$ or > 30 in the mass range $m_{H^+}=80$ –160 GeV. See Fig. 2 for the excluded region in a certain MSSM scenario.
- 81 ABDALLAH 04I search for $e^+e^- \rightarrow H^+H^-$ with H^\pm decaying to $\tau\nu$, cs, or W^*A^0 in Type-I two-Higgs-doublet models.
- ⁸² ABBIENDI 03 give a limit $m_{H^+}>1.28{\rm tan}\beta$ GeV (95%CL) in Type II two-doublet models.
- ⁸³ ABAZOV 02B search for a charged Higgs boson in top decays with $H^+ \to \tau^+ \nu$ at $E_{\rm cm} = 1.8$ TeV. For $m_{H^+} = 75$ GeV, the region $\tan \beta > 32.0$ is excluded at 95%CL. The excluded mass region extends to over 140 GeV for $\tan \beta$ values above 100.
- ⁸⁴ BORZUMATI 02 point out that the decay modes such as $b\overline{b}W$, A^0W , and supersymmetric ones can have substantial branching fractions in the mass range explored at LEP II and Tevatron.
- 85 ABBIENDI 01Q give a limit $\tan\beta/m_{H^+} < 0.53~{\rm GeV}^{-1}$ (95%CL) in Type II two-doublet models.
- ⁸⁶ BARATE 01E give a limit $\tan\beta/m_{H^+} <$ 0.40 GeV $^{-1}$ (90% CL) in Type II two-doublet models. An independent measurement of $B \to \tau \nu_{\tau} X$ gives $\tan\beta/m_{H^+} <$ 0.49 GeV $^{-1}$ (90% CL).
- ⁸⁷ GAMBINO 01 use the world average data in the summer of 2001 B($b \rightarrow s\gamma$)= (3.23 \pm 0.42) \times 10⁻⁴. The limit applies for Type-II two-doublet models.
- ⁸⁸ AFFOLDER 00I search for a charged Higgs boson in top decays with $H^+ \to \tau^+ \nu$ in $p\overline{p}$ collisions at $E_{\rm cm}{=}1.8$ TeV. The excluded mass region extends to over 120 GeV for $\tan\beta$ values above 100 and B $(\tau\nu){=}1$. If B $(t\to bH^+)\gtrsim$ 0.6, m_{H^+} up to 160 GeV is excluded. Updates ABE 97L.
- ⁸⁹ ABBOTT 99E search for a charged Higgs boson in top decays in $p\overline{p}$ collisions at $E_{\rm cm}{=}1.8$ TeV, by comparing the observed $t\overline{t}$ cross section (extracted from the data assuming the dominant decay $t\to bW^+$) with theoretical expectation. The search is sensitive to regions of the domains $\tan\beta\lesssim 1$, $50< m_{H^+}({\rm GeV})\lesssim 120$ and $\tan\beta\gtrsim 40$, $50< m_{H^+}({\rm GeV})\lesssim 160$. See Fig. 3 for the details of the excluded region.
- ⁹⁰ ACKERSTAFF 99D measure the Michel parameters ρ , ξ , η , and $\xi\delta$ in leptonic τ decays from $Z \to \tau\tau$. Assuming e- μ universality, the limit $m_{H^+} > 0.97 \tan\beta$ GeV (95%CL) is obtained for two-doublet models in which only one doublet couples to leptons.
- ⁹¹ ACCIARRI 97F give a limit $m_{H^+}>2.6~{\rm tan}\beta$ GeV (90% CL) from their limit on the exclusive $B\to ~\tau \nu_{\tau}$ branching ratio.
- 92 AMMAR 97B measure the Michel parameter ρ from $\tau\to e\nu\nu$ decays and assumes e/μ universality to extract the Michel η parameter from $\tau\to \mu\nu\nu$ decays. The measurement is translated to a lower limit on $m_{\mbox{$H^+$}}$ in a two-doublet model $m_{\mbox{$H^+$}}>0.97$ tan β GeV (90% CL).

- ⁹³COARASA 97 reanalyzed the constraint on the $(m_{H^\pm}, \tan\beta)$ plane derived from the inclusive $B \to \tau \nu_{\tau} X$ branching ratio in GROSSMAN 95B and BUSKULIC 95. They show that the constraint is quite sensitive to supersymmetric one-loop effects.
- ⁹⁴ GUCHAIT 97 studies the constraints on m_{H^+} set by Tevatron data on $\ell \tau$ final states in $t \bar{t} \to (W \, b) (H \, b), \, W \to \ell \nu, \, H \to \tau \nu_{\tau}$. See Fig. 2 for the excluded region.
- ⁹⁵ MANGANO 97 reconsiders the limit in ACCIARRI 97F including the effect of the potentially large $B_C \to \tau \nu_{\tau}$ background to $B_U \to \tau \nu_{\tau}$ decays. Stronger limits are obtained.
- $^{96}\,\rm STAHL$ 97 fit τ lifetime, leptonic branching ratios, and the Michel parameters and derive limit $m_{H^+}>1.5\,\tan\!\beta$ GeV (90% CL) for a two-doublet model. See also STAHL 94.
- 97 ALAM 95 measure the inclusive $b\to s\gamma$ branching ratio at $\Upsilon(4S)$ and give B(b $\to s\gamma)<4.2\times 10^{-4}$ (95% CL), which translates to the limit $m_{H^+}>[244+63/(\tan\beta)^{1.3}]$ GeV in the Type II two-doublet model. Light supersymmetric particles can invalidate this bound.
- 98 BUSKULIC 95 give a limit $m_{H^+}>1.9~{\rm tan}\beta$ GeV (90% CL) for Type-II models from $b\to \tau \nu_{\tau} X$ branching ratio, as proposed in GROSSMAN 94.

MASS LIMITS for $H^{\pm\pm}$ (doubly-charged Higgs boson)

This section covers searches for a doubly-charged Higgs boson with couplings to lepton pairs. Its weak isospin T_3 is thus restricted to two possibilities depending on lepton chiralities: $T_3(H^{\pm\pm})=\pm 1$, with the coupling $g_{\ell\ell}$ to $\ell_L^-\ell_L^{\prime-}$ and $\ell_R^+\ell_R^{\prime+}$ ("left-handed") and $T_3(H^{\pm\pm})=0$, with the coupling to $\ell_R^-\ell_R^{\prime-}$ and $\ell_L^+\ell_L^{\prime+}$ ("right-handed"). These Higgs bosons appear in some left-right symmetric models based on the gauge group $\mathrm{SU}(2)_L \times \mathrm{SU}(2)_R \times \mathrm{U}(1)$. These two cases are listed separately in the following. Unless noted, one of the lepton flavor combinations is assumed to be dominant in the decay.

LIMITS for $H^{\pm\pm}$ with $T_3=\pm 1$

VALUE	CL%	DOCUMENT ID		TECN	COMMENT
VALUL	CL/0			TLCIV	COMMENT
>114	95	⁹⁹ AALTONEN	08AA	CDF	e au
>112	95	⁹⁹ AALTONEN	08AA	CDF	μau
>150	95	¹⁰⁰ ABAZOV	08V	D0	$\mu\mu$
>136			04 G	CDF	$\mu\mu$
> 98.1	95	¹⁰² ABDALLAH	03	DLPH	au au
> 99.0	95	¹⁰³ ABBIENDI	0 2C	OPAL	au au

• • We do not use the following data for averages, fits, limits, etc.

		¹⁰⁴ AKTAS	06A	H1	single $H^{\pm\pm}$
>133	95	¹⁰⁵ ACOSTA	05L	CDF	stable
>118.4	95	¹⁰⁶ ABAZOV	04E	-	$\mu\mu$
		¹⁰⁷ ABBIENDI	03 Q	OPAL	$E_{\rm cm} \leq$ 209 GeV, single $H^{\pm\pm}$
		¹⁰⁸ GORDEEV	97	SPEC	muonium conversion
		¹⁰⁹ ASAKA	95	THEO	
> 45.6	95	¹¹⁰ ACTON	92M	OPAL	
> 30.4	95	111 ACTON	92M	OPAL	
none 6.5-36.6	95	¹¹² SWARTZ	90	MRK2	

- ⁹⁹ AALTONEN 08AA search for $H^{++}H^{--}$ production in $p\,\overline{p}$ collisions at $E_{\rm cm}=1.96$ TeV. The limit assumes 100% branching ratio to the specified final state.
- ABAZOV 08V search for $H^{++}H^{--}$ production in $p\overline{p}$ collisions at $E_{\rm cm}=$ 1.96 TeV. The limit is for B $(H\to\mu\mu)=$ 1.
- 101 ACOSTA 04G search for $H^{++}H^{--}$ pair production in $p\overline{p}$ collisions with muon and electron final states. The limit holds for $\mu\mu$. For ee and $e\mu$ modes, the limits are 133 and 115 GeV, respectively. The limits are valid for $g_{\rho\rho'}\gtrsim 10^{-5}$.
- ¹⁰² ABDALLAH 03 search for $H^{++}H^{--}$ pair production either followed by $H^{++} \rightarrow \tau^+ \tau^+$, or decaying outside the detector.
- ¹⁰³ ABBIENDI 02C searches for pair production of $H^{++}H^{--}$, with $H^{\pm\pm}\to\ell^{\pm}\ell^{\pm}$ ($\ell,\ell'=e,\mu,\tau$). The limit holds for $\ell=\ell'=\tau$, and becomes stronger for other combinations of leptonic final states. To ensure the decay within the detector, the limit only applies for $g(H\ell\ell)\gtrsim 10^{-7}$.
- 104 AKTAS 06A search for single $H^{\pm\pm}$ production in ep collisions at HERA. Assuming that H^{++} only couples to $e^+\mu^+$ with $g_{e\mu}=0.3$ (electromagnetic strength), a limit $m_{H^{++}}>141$ GeV (95% CL) is derived. For the case where H^{++} couples to $e\tau$ only the limit is 112 GeV
- the limit is 112 GeV. 105 ACOSTA 05L search for $H^{++}H^{--}$ pair production in $p\overline{p}$ collisions. The limit is valid for $g_{\ell\,\ell'} < 10^{-8}$ so that the Higgs decays outside the detector.
- ¹⁰⁶ ABAZOV 04E search for $H^{++}H^{--}$ pair production in $H^{\pm\pm}\to\mu^\pm\mu^\pm$. The limit is valid for $g_{\mu\mu}\gtrsim 10^{-7}$.
- ABBIENDI 03Q searches for single $H^{\pm\pm}$ via direct production in $e^+e^- \rightarrow e^\mp e^\mp H^{\pm\pm}$, and via t-channel exchange in $e^+e^- \rightarrow e^+e^-$. In the direct case, and assuming B($H^{\pm\pm} \rightarrow \ell^\pm \ell^\pm$) = 1, a 95% CL limit on h_{ee} < 0.071 is set for $m_{H^{\pm\pm}}$ < 160 GeV (see Fig. 6). In the second case, indirect limits on h_{ee} are set for $m_{H^{\pm\pm}}$ < 2 TeV (see Fig. 8).
- 108 GORDEEV 97 search for muonium-antimuonium conversion and find $G_{M\overline{M}}/G_F < 0.14$ (90% CL), where $G_{M\overline{M}}$ is the lepton-flavor violating effective four-fermion coupling. This limit may be converted to $m_{H^{++}} > 210$ GeV if the Yukawa couplings of H^{++} to ee and $\mu\mu$ are as large as the weak gauge coupling. For similar limits on muonium-antimuonium conversion, see the muon Particle Listings.
- 109 ASAKA 95 point out that H^{++} decays dominantly to four fermions in a large region of parameter space where the limit of ACTON 92M from the search of dilepton modes does not apply.
- ¹¹⁰ ACTON 92M limit assumes $H^{\pm\pm} \rightarrow \ell^{\pm}\ell^{\pm}$ or $H^{\pm\pm}$ does not decay in the detector. Thus the region $g_{\ell\ell} \approx 10^{-7}$ is not excluded.
- 111 ACTON 92M from $\Delta\Gamma_{7}$ <40 MeV.
- ¹¹² SWARTZ 90 assume $H^{\pm\pm} \to \ell^{\pm}\ell^{\pm}$ (any flavor). The limits are valid for the Higgs-lepton coupling g($H\ell\ell$) $\gtrsim 7.4 \times 10^{-7}/[m_H/\text{GeV}]^{1/2}$. The limits improve somewhat for ee and $\mu\mu$ decay modes.

LIMITS for $H^{\pm\pm}$ with $T_3 = 0$

VALUE	CL%	DOCUMENT ID		TECN	COMMENT
>127	95	¹¹³ ABAZOV	08V	D0	$\mu\mu$
>113	95	¹¹⁴ ACOSTA	04 G	CDF	$\mu\mu$
> 97.3	95	¹¹⁵ ABDALLAH	03	DLPH	au au
> 97.3	95	¹¹⁶ ACHARD	03F	L3	au au
> 98.5	95	¹¹⁷ ABBIENDI	0 2C	OPAL	au au

• • • We do not use the following data for averages, fits, limits, etc. • • •

		¹¹⁸ AKTAS	06A H1	single $H^{\pm\pm}$
>109	95	¹¹⁹ ACOSTA	05L CDF	stable
> 98.2	95	¹²⁰ ABAZOV	04E D0	$\mu\mu$
		¹²¹ ABBIENDI	03Q OPA	$E_{ m cm} \leq$ 209 GeV, single $H^{\pm\pm}$
		¹²² GORDEEV	97 SPE	
> 45.6	95	¹²³ ACTON	92м ОРА	_
> 25.5	95	¹²⁴ ACTON	92м ОРА	_
none 7.3-34.3	95	¹²⁵ SWARTZ	90 MRK	2

- ¹¹³ ABAZOV 08V search for $H^{++}H^{--}$ production in $p\bar{p}$ collisions at $E_{\rm cm}=$ 1.96 TeV. The limit is for B($H\to \mu\mu$) = 1.
- ¹¹⁴ ACOSTA 04G search for $H^{++}H^{--}$ pair production in $p\bar{p}$ collisions with muon and electron final states. The limit holds for $\mu\mu$.
- ¹¹⁵ ABDALLAH 03 search for $H^{++}H^{--}$ pair production either followed by $H^{++} \rightarrow \tau^+ \tau^+$, or decaying outside the detector.
- 116 ACHARD 03F search for $e^+e^- \to H^{++}H^{--}$ with $H^{\pm\pm} \to \ell^\pm\ell'^\pm$. The limit holds for $\ell=\ell'=\tau$, and slightly different limits apply for other flavor combinations. The limit is valid for $g_{\ell\,\ell'}\gtrsim 10^{-7}$.
- ¹¹⁷ ABBIENDI 02C searches for pair production of $H^{++}H^{--}$, with $H^{\pm\pm}\to\ell^{\pm}\ell^{\pm}$ ($\ell,\ell'=e,\mu,\tau$). the limit holds for $\ell=\ell'=\tau$, and becomes stronger for other combinations of leptonic final states. To ensure the decay within the detector, the limit only applies for $g(H\ell\ell)\gtrsim 10^{-7}$.
- 118 AKTAS 06A search for single $H^{\pm\pm}$ production in ep collisions at HERA. Assuming that H^{++} only couples to $e^+\mu^+$ with $g_{e\,\mu}=0.3$ (electromagnetic strength), a limit $m_{H^{++}}>141$ GeV (95% CL) is derived. For the case where H^{++} couples to $e\tau$ only the limit is 112 GeV.
- the limit is 112 GeV.
 119 ACOSTA 05L search for $H^{++}H^{--}$ pair production in $p\overline{p}$ collisions. The limit is valid for $g_{\ell\,\ell'} < 10^{-8}$ so that the Higgs decays outside the detector.
- ¹²⁰ ABAZOV 04E search for $H^{++}H^{--}$ pair production in $H^{\pm\pm}\to\mu^\pm\mu^\pm$. The limit is valid for $g_{\mu\mu}\gtrsim 10^{-7}$.
- 121 ABBIENDI 03Q searches for single $H^{\pm\pm}$ via direct production in $e^+e^- \rightarrow e^\mp e^\mp H^{\pm\pm}$, and via t-channel exchange in $e^+e^- \rightarrow e^+e^-$. In the direct case, and assuming B($H^{\pm\pm} \rightarrow \ell^\pm \ell^\pm$) = 1, a 95% CL limit on h_{ee} < 0.071 is set for $m_{H^{\pm\pm}}$ < 160 GeV (see Fig. 6). In the second case, indirect limits on h_{ee} are set for $m_{H^{\pm\pm}}$ < 2 TeV (see Fig. 8).
- 122 GORDEEV 97 search for muonium-antimuonium conversion and find $G_{M\,\overline{M}}/G_F<0.14$ (90% CL), where $G_{M\,\overline{M}}$ is the lepton-flavor violating effective four-fermion coupling. This limit may be converted to $m_{H^{++}}>210$ GeV if the Yukawa couplings of H^{++} to ee and $\mu\mu$ are as large as the weak gauge coupling. For similar limits on muonium-antimuonium conversion, see the muon Particle Listings.
- ¹²³ ACTON 92M limit assumes $H^{\pm\pm} \rightarrow \ell^{\pm}\ell^{\pm}$ or $H^{\pm\pm}$ does not decay in the detector. Thus the region $g_{\ell\ell} \approx 10^{-7}$ is not excluded.
- $^{124}\,\text{ACTON}$ 92M from $\Delta\Gamma_{\textit{7}}<\!40$ MeV.
- ¹²⁵ SWARTZ 90 assume $H^{\pm\pm} \to \ell^{\pm}\ell^{\pm}$ (any flavor). The limits are valid for the Higgslepton coupling $g(H\ell\ell) \gtrsim 7.4 \times 10^{-7}/[m_H/\text{GeV}]^{1/2}$. The limits improve somewhat for ee and $\mu\mu$ decay modes.

H^0 and H^{\pm} REFERENCES

AALTONEN	09A	PRL 102 021802		T. Aaltonen <i>et al.</i>	(CDF Collab.)
ABAZOV	09C	PRL 102 051803	,	V.M. Abazov <i>et al.</i>	(D0 Collab.)
ABAZOV	09F	PRL 102 051804	,	V.M. Abazov <i>et al.</i>	(D0 Collab.)
AALTONEN	AA80	PRL 101 121801		T. Aaltonen et al.	(CDF Collab.)
AALTONEN	08AF	PRL 101 251803		T. Aaltonen et al.	(CDF Collab.)
AALTONEN	08V	PRL 100 041801		T. Aaltonen <i>et al.</i>	(CDF Collab.)
Also		PR D78 032008		T. Aaltonen <i>et al.</i>	(CDF Collab.)
AALTONEN	08X	PRL 100 211801		T. Aaltonen et al.	(CDF Collab.)
ABAZOV		PRL 101 221802		V.M. Abazov et al.	(D0 Collab.)
ABAZOV		PRL 101 251802		V.M. Abazov et al.	(D0 Collab.)
ABAZOV	08U	PRL 101 051801		V.M. Abazov et al.	(D0 Collab.)
ABAZOV	08V	PRL 101 071803		V.M. Abazov <i>et al.</i>	(D0 Collab.)
ABAZOV	08W	PRL 101 071804		V.M. Abazov <i>et al.</i>	(D0 Collab.)
ABAZOV	08Y	PL B663 26		V.M. Abazov et al.	(D0 Collab.)
ABDALLAH	08B	EPJ C54 1		J. Abdallah <i>et al.</i>	(DELPHI Collab.)
Also	OOD	EPJ C56 165 (errat			(DELPHI Collab.)
LOVE	08	PRL 101 151802		W. Love <i>et al.</i>	(CLEO Collab.)
ABAZOV	07X	PL B655 209		V.M. Abazov <i>et al.</i>	`
					(D0 Collab.)
ABBIENDI	07 07	EPJ C49 457		G. Abbiendi <i>et al.</i>	(OPAL Collab.)
BESSON		PRL 98 052002		D. Besson <i>et al.</i>	(CLEO Collab.)
SCHAEL	07	EPJ C49 439		S. Schael <i>et al.</i>	(ALEPH Collab.)
ABAZOV	06	PRL 96 011801		V.M. Abazov <i>et al.</i>	(D0 Collab.)
ABAZOV	060	PRL 97 151804		V.M. Abazov et al.	(D0 Collab.)
ABAZOV	06Q	PRL 97 161803		V.M. Abazov et al.	(D0 Collab.)
ABULENCIA	06	PRL 96 011802		A. Abulencia <i>et al.</i>	(CDF Collab.)
ABULENCIA	06E	PRL 96 042003		A. Abulencia <i>et al.</i>	(CDF Collab.)
AKTAS	06A	PL B638 432		A. Aktas <i>et al.</i>	(H1 Collab.)
LEP-SLC	06	PRPL 427 257			OPAL, SLD and working groups
SCHAEL	06B	EPJ C47 547		S. Schael <i>et al.</i>	(LEP Collabs.)
ABBIENDI	05A	EPJ C40 317		G. Abbiendi <i>et al.</i>	(OPAL Collab.)
ABDALLAH	05D	EPJ C44 147		J. Abdallah <i>et al.</i>	(DELPHI Collab.)
ACHARD	05	PL B609 35	ļ	P. Achard <i>et al.</i>	(L3 Collab.)
ACOSTA	05L	PRL 95 071801	ļ	D. Acosta et al.	(CDF Collab.)
ACOSTA	05Q	PR D72 072004	I	D. Acosta et al.	(CDF Collab.)
ABAZOV	04E	PRL 93 141801	,	V.M. Abazov <i>et al.</i>	(D0 Collab.)
ABBIENDI	04	EPJ C32 453	(G. Abbiendi <i>et al.</i>	(OPAL Collab.)
ABBIENDI	04K	PL B597 11	(G. Abbiendi <i>et al.</i>	(OPAL Collab.)
ABBIENDI	04M	EPJ C37 49		G. Abbiendi <i>et al.</i>	(OPAL Collab.)
ABDALLAH	04	EPJ C32 145		J. Abdallah <i>et al.</i>	(DELPHI Collab.)
ABDALLAH	04B	EPJ C32 475		J. Abdallah <i>et al.</i>	(DELPHI Collab.)
ABDALLAH	04I	EPJ C34 399		J. Abdallah <i>et al.</i>	(DELPHI Collab.)
ABDALLAH	04L	EPJ C35 313		J. Abdallah <i>et al.</i>	(DELPHI Collab.)
ABDALLAH	040	EPJ C38 1		J. Abdallah <i>et al.</i>	(DELPHI Collab.)
ACHARD	04B	PL B583 14		P. Achard et al.	` (L3 Collab.
ACHARD	04F	PL B589 89		P. Achard et al.	(L3 Collab.)
ACOSTA	04G	PRL 93 221802		D. Acosta et al.	(CDF Collab.)
ABBIENDI	03	PL B551 35		G. Abbiendi et al.	(OPAL Collab.)
ABBIENDI	03B	EPJ C26 479		G. Abbiendi et al.	(OPAL Collab.)
ABBIENDI	03F	EPJ C27 311		G. Abbiendi <i>et al.</i>	(OPAL Collab.)
ABBIENDI	03G	EPJ C27 483		G. Abbiendi <i>et al.</i>	(OPAL Collab.)
ABBIENDI	03Q	PL B577 93		G. Abbiendi <i>et al.</i>	(OPAL Collab.)
ABDALLAH	03	PL B552 127		J. Abdallah <i>et al.</i>	(DELPHI Collab.)
ACHARD	03C	PL B568 191		P. Achard <i>et al.</i>	(L3 Collab.)
ACHARD	03E	PL B575 208		P. Achard <i>et al.</i>	(L3 Collab.)
ACHARD	03F	PL B576 18		P. Achard <i>et al.</i>	(L3 Collab.)
CARENA	03	EPJ C26 601		M.S. Carena <i>et al.</i>	(E3 Collab.)
HEISTER	03D	PL B565 61		A. Heister <i>et al.</i>	(ALEPH, DELPHI, L3+)
		L3, OPAL, LEP Hig			(ALLI II, DELI III, EST)
ABAZOV	02B	PRL 88 151803		V.M. Abazov <i>et al.</i>	(D0 Collab.)
ABBIENDI	02C	PL B526 221		G. Abbiendi <i>et al.</i>	(OPAL Collab.)
ABBIENDI	02C 02D	EPJ C23 397		G. Abbiendi <i>et al.</i>	(OPAL Collab.)
ABBIENDI	02D 02F	PL B544 44		G. Abbiendi <i>et al.</i>	(OPAL Collab.)
ACHARD	02F 02C	PL B534 28		P. Achard <i>et al.</i>	(L3 Collab.)
ACHARD	02C 02H	PL B545 30		P. Achard <i>et al.</i>	(L3 Collab.
AKEROYD	0211	PR D66 037702		A.G. Akeroyd <i>et al.</i>	(LS Collab.)
BORZUMATI	02	PL B549 170		A.G. Akeroya <i>et al.</i> F.M. Borzumati, A. D	iouadi
HEISTER	02	PL B526 191		A. Heister <i>et al.</i>	(ALEPH Collab.)
HEISTER	02 02L	PL B544 16		A. Heister <i>et al.</i>	(ALEPH Collab.)
11515151	UZL	1 F DO44 IO	4	A. FICISICI CL dl.	(ALLETT COURD.)

HEISTER	02M	PL B544 25	A. Heister et al.	(ALEPH Collab.)				
HEISTER	02P	PL B543 1	A. Heister et al.	(ALEPH Collab.)				
ABBIENDI	01E	EPJ C18 425	G. Abbiendi et al.	(OPAL Collab.)				
ABBIENDI	01Q	PL B520 1	G. Abbiendi et al.	(OPAL Collab.)				
ABREU	01F	PL B507 89	P. Abreu et al.	(DELPHI Collab.)				
ACHARD	01C	PL B517 319	P. Achard et al.	` (L3 Collab.)				
AFFOLDER	01H	PR D64 092002	T. Affolder et al.	(CDF Collab.)				
BARATE	01C	PL B499 53	R. Barate et al.	(ALEPH Collab.)				
BARATE	01E	EPJ C19 213	R. Barate et al.	(ALEPH Collab.)				
GAMBINO	01	NP B611 338	P. Gambino, M. Misiak	,				
ACCIARRI	M00	PL B485 85	M. Acciarri et al.	(L3 Collab.)				
ACCIARRI	00R	PL B489 102	M. Acciarri et al.	(L3 Collab.)				
ACCIARRI	00S	PL B489 115	M. Acciarri et al.	(L3 Collab.)				
AFFOLDER	001	PR D62 012004	T. Affolder et al.	(CDF Collab.)				
BARATE	00L	PL B487 241	R. Barate et al.	(ALEPH Collab.)				
PDG	00	EPJ C15 1	D.E. Groom et al.	,				
ABBIENDI	99E	EPJ C7 407	G. Abbiendi et al.	(OPAL Collab.)				
ABBIENDI	990	PL B464 311	G. Abbiendi et al.	(OPAL Collab.)				
ABBOTT	99B	PRL 82 2244	B. Abbott et al.	(D0 Collab.)				
ABBOTT	99E	PRL 82 4975	B. Abbott et al.	(D0 Collab.)				
ABREU	99P	PL B458 431	P. Abreu et al.	(DELPHI Collab.)				
ACKERSTAFF	99D	EPJ C8 3	K. Ackerstaff et al.	(OPAL Collab.)				
CARENA	99B	hep-ph/9912223	M.S. Carena et al.	()				
CERN-TH/99-374								
ABBOTT	98	PRL 80 442	B. Abbott et al.	(D0 Collab.)				
ACKERSTAFF		EPJ C5 19	K. Ackerstaff et al.	(OPAL Collab.)				
ACKERSTAFF		PL B437 218	K. Ackerstaff et al.	(OPAL Collab.)				
GONZALEZ-G		PR D57 7045	M.C. Gonzalez-Garcia, S.M. Lietti, S.F.					
PDG	98	EPJ C3 1	C. Caso et al.					
ABE	97L	PRL 79 357	F. Abe <i>et al.</i>	(CDF Collab.)				
ACCIARRI	97F	PL B396 327	M. Acciarri et al.	(L3 Collab.)				
AMMAR	97B	PRL 78 4686	R. Ammar et al.	(CLÈO Collab.)				
COARASA	97	PL B406 337	J.A. Coarasa, R.A. Jimenez, J. Sola	()				
GORDEEV	97	PAN 60 1164	V.A. Gordeev et al.	(PNPI)				
		Translated from YAF 60		()				
GUCHAIT	97	PR D55 7263	M. Guchait, D.P. Roy	(TATA)				
KRAWCZYK	97	PR D55 6968	M. Krawczyk, J. Zochowski	(WARS)				
MANGANO	97	PL B410 299	M. Mangano, S. Slabospitsky					
STAHL	97	ZPHY C74 73	A. Stahl, H. Voss	(BONN)				
ALEXANDER	96H	ZPHY C71 1	G. Alexander et al.	(OPAL Collab.)				
PDG	96	PR D54 1	R. M. Barnett et al.					
ABREU	95H	ZPHY C67 69	P. Abreu et al.	(DELPHI Collab.)				
ALAM	95	PRL 74 2885	M.S. Alam <i>et al.</i>	(CLEO Collab.)				
ASAKA	95	PL B345 36	T. Asaka, K.I. Hikasa	(TOHOK)				
BUSKULIC	95	PL B343 444	D. Buskulic et al.	(ALEPH Collab.)				
GROSSMAN	95B	PL B357 630	Y. Grossman, H. Haber, Y. Nir					
GROSSMAN	94	PL B332 373	Y. Grossman, Z. Ligeti					
STAHL	94	PL B324 121	A. Stahl	(BONN)				
ACTON	92M	PL B295 347	P.D. Acton et al.	(OPAL Collab.)				
PICH	92	NP B388 31	A. Pich, J. Prades, P. Yepes	(CERN, CPPM)				
SWARTZ	90	PRL 64 2877	M.L. Swartz et al.	(Mark II Collab.)				