Double-β Decay

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Half-life Measurements and Limits for Double-β Decay

In all cases of double-beta decay, (Z,A) $\rightarrow$ (Z+2,A) $+2e^-$ + (0 or 2)$\overline{\nu}_e$. In the following Listings, only best or comparable limits or lifetimes for each isotope are reported. For 2ν decay, which is well established, only measured half-lives are reported.

$t_{1/2}(10^{21}$ yr)	CL%	ISOTOPE	TRANSITION	METHOD	DOCUMENT ID
• • • We do not use the following data for averages, fits, limits, etc. • • •					
> 550	90	^{130}Te	0ν	Cryog. det.	[1] ARNABOLDI 04
> 310	90	^{100}Mo	0ν	NEMO-3	[2] ARNOLD 04
> 140	90	^{82}Se	0ν	NEMO-3	[3] ARNOLD 04
$(7.68 \pm 0.02 \pm 0.54)$E-3		^{100}Mo	2ν	NEMO-3	[4] ARNOLD 04
$(10.3 \pm 0.3 \pm 0.7)$E-2		^{82}Se	2ν	NEMO-3	[5] ARNOLD 04
$14^{+4}_{-2} \pm 3$	68	^{150}Nd	0ν+2ν $0^+ \rightarrow 0^+_1$	γ in Ge det.	[6] BARABASH 04
11900^{+29900}_{-5000}	99.7	^{76}Ge	0ν	Enriched HPGe	[7] KLAPDOR-K... 04A
> 14	90	^{48}Ca	0ν	CaF_2 scint.	[8] OGAWA 04
> 210	90	^{130}Te	0ν	Cryog. det.	[9] ARNABOLDI 03
> 31	90	^{130}Te	0ν $0^+ \rightarrow 2^+$	Cryog. det.	[10] ARNABOLDI 03
$0.61 \pm 0.14^{+0.29}_{-0.35}$	90	^{130}Te	2ν	Cryog. det.	[11] ARNABOLDI 03
> 110	90	^{128}Te	0ν	Cryog. det.	[12] ARNABOLDI 03
$(0.029^{+0.004}_{-0.003})$		^{116}Cd	2ν	$^{116}CdWO_4$ scint.	[13] DANEVICH 03
> 170	90	^{116}Cd	0ν	$^{116}CdWO_4$ scint.	[14] DANEVICH 03
> 29	90	^{116}Cd	0ν $0^+ \rightarrow 2^+$	$^{116}CdWO_4$ scint.	[15] DANEVICH 03
> 14	90	^{116}Cd	0ν $0^+ \rightarrow 0^+_1$	$^{116}CdWO_4$ scint.	[16] DANEVICH 03
> 6	90	^{116}Cd	0ν $0^+ \rightarrow 0^+_2$	$^{116}CdWO_4$ scint.	[17] DANEVICH 03
$1.74 \pm 0.01^{+0.18}_{-0.16}$		^{76}Ge	2ν	Enriched HPGe	[18] DOERR 03
>15700	90	^{76}Ge	0ν	Enriched HPGe	[19] AALSETH 02B
> 58	90	^{134}Xe	0ν	Liquid Xe Scint.	[20] BERNABEI 02D
> 1200	90	^{136}Xe	0ν	Liquid Xe Scint.	[21] BERNABEI 02D
$15000^{+168000}_{-7500}$		^{76}Ge	0ν	Enriched HPGe	[22] KLAPDOR-K... 02D
$(7.2 \pm 0.9 \pm 1.8)$E-3		^{100}Mo	2ν	Liq. Ar ioniz.	[23] ASHITKOV 01
> 4.9	90	^{100}Mo	0ν	Liq. Ar ioniz.	[24] ASHITKOV 01
> 1.3	90	^{160}Gd	0ν	Gd_2SiO_5:Ce	[25] DANEVICH 01
> 1.3	90	^{160}Gd	0ν $0^+ \rightarrow 2^+$	Gd_2SiO_5:Ce	[26] DANEVICH 01
$0.59^{+0.17}_{-0.11} \pm 0.06$		^{100}Mo	0ν+2ν $0^+ \rightarrow 0^+_1$	Ge coinc.	[27] DEBRAECKEL.01
> 55	90	^{100}Mo	$0\nu,\langle m_\nu \rangle$	ELEGANT V	[28] EJIRI 01
> 42	90	^{100}Mo	$0\nu,\langle \lambda \rangle$	ELEGANT V	[28] EJIRI 01
> 49	90	^{100}Mo	$0\nu,\langle \eta \rangle$	ELEGANT V	[28] EJIRI 01
>19000	90	^{76}Ge	0ν	Enriched HPGe	[29] KLAPDOR-K... 01
$1.55 \pm 0.001^{+0.19}_{-0.15}$	90	^{76}Ge	2ν	Enriched HPGe	[30] KLAPDOR-K... 01

(9.4 ± 3.2)E-3	90	^{96}Zr	$0\nu+2\nu$		Geochem	[31] WIESER	01
$0.042^{+0.033}_{-0.013}$		^{48}Ca	2ν		Ge spectrometer	[32] BRUDANIN	00
$0.021^{+0.008}_{-0.004} \pm 0.002$		^{96}Zr	2ν		NEMO-2	[33] ARNOLD	99
> 1.0	90	^{96}Zr	0ν		NEMO-2	[33] ARNOLD	99
$(8.3 \pm 1.0 \pm 0.7)$E-2		^{82}Se	2ν		NEMO-2	[34] ARNOLD	98
$>$ 9.5	90	^{82}Se	0ν		NEMO-2	[35] ARNOLD	98
$>$ 2.8	90	^{82}Se	0ν	$0^+ \to 2^+$	NEMO-2	[36] ARNOLD	98
$(7.6^{+2.2}_{-1.4})$E-3		^{100}Mo	2ν		Si(Li)	[37] ALSTON-...	97
$(6.82^{+0.38}_{-0.53} \pm 0.68)$E-3		^{100}Mo	2ν		TPC	[38] DESILVA	97
$(6.75^{+0.37}_{-0.42} \pm 0.68)$E-3		^{150}Nd	2ν		TPC	[39] DESILVA	97
$>$ 1.2	90	^{150}Nd	0ν		TPC	[40] DESILVA	97
$(3.75 \pm 0.35 \pm 0.21)$E-2		^{116}Cd	2ν	$0^+ \to 0^+$	NEMO 2	[41] ARNOLD	96
$0.043^{+0.024}_{-0.011} \pm 0.014$		^{48}Ca	2ν		TPC	[42] BALYSH	96
0.79 ± 0.10		^{130}Te	$0\nu+2\nu$		Geochem	[43] TAKAOKA	96
$0.61^{+0.18}_{-0.11}$		^{100}Mo	$0\nu+2\nu$	$0^+ \to 0^+_1$	γ in HPGe	[44] BARABASH	95
$(9.5 \pm 0.4 \pm 0.9)$E-3		^{100}Mo	2ν		NEMO 2	DASSIE	95
$>$ 0.6	90	^{100}Mo	0ν	$0^+ \to 0^+_1$	NEMO 2	DASSIE	95
$0.026^{+0.009}_{-0.005}$		^{116}Cd	2ν	$0^+ \to 0^+$	ELEGANT IV	EJIRI	95
$0.017^{+0.010}_{-0.005} \pm 0.0035$		^{150}Nd	2ν	$0^+ \to 0^+$	TPC	ARTEMEV	93
0.039 ± 0.009		^{96}Zr	$0\nu+2\nu$		Geochem	KAWASHIMA	93
2.7 ± 0.1		^{130}Te	$0\nu+2\nu$		Geochem	BERNATOW...	92
7200 ± 400		^{128}Te	$0\nu+2\nu$		Geochem	[45] BERNATOW...	92
$>$ 27	68	^{82}Se	0ν	$0^+ \to 0^+$	TPC	ELLIOTT	92
$0.108^{+0.026}_{-0.006}$		^{82}Se	2ν	$0^+ \to 0^+$	TPC	ELLIOTT	92
2.0 ± 0.6		^{238}U	$0\nu+2\nu$		Radiochem	[46] TURKEVICH	91
$>$ 9.5	76	^{48}Ca	0ν		CaF_2 scint.	YOU	91
$0.12 \pm 0.01 \pm 0.04$	68	^{82}Se	$0\nu+2\nu$		Geochem.	[47] LIN	88
$0.75 \pm 0.03 \pm 0.23$	68	^{130}Te	$0\nu+2\nu$		Geochem.	[48] LIN	88
1800 ± 700	68	^{128}Te	$0\nu+2\nu$		Geochem.	[49] LIN	88B
2.60 ± 0.28		^{130}Te	$0\nu+2\nu$		Geochem	[50] KIRSTEN	83

[1] Supersedes ARNABOLDI 03. Bolometric TeO_2 detector array Cuoricino used for high resolution search for $0\nu\beta\beta$ decay.

[2] ARNOLD 04 use the NEMO-3 tracking detector to determine the limit for $0\nu\beta\beta$ halflife of ^{100}Mo. This represents an impovement, by a factor of ~ 6, when compared with EJIRI 01.

[3] ARNOLD 04 use the NEMO-3 tracking detector to determine the limit for $0\nu\beta\beta$ halflife of ^{82}Se. This represents an impovement, by a factor of ~ 10, when compared with ELLIOTT 92. It supersedes the limit of ARNOLD 98 for this decay using NEMO-2.

[4] ARNOLD 04 use the NEMO-3 tracking detector to determine the $2\nu\beta\beta$ halflife of chemical100Mo with high statistics and low background. The halflife is determined assuming the Single State Dominance. It is in agreement with, and more accurate than, previous determinations. Supersedes DASSIE 95 determination of this quantity with NEMO-2.

[5] ARNOLD 04 use the NEMO-3 tracking detector to determine the $2\nu\beta\beta$ halflife of ^{82}Se. The halflife is in agreement with ARNOLD 1998 with NEMO-2 which it supersedes.

[6] BARABASH 04 perform an inclusive measurement of the $\beta\beta$ decay of ^{150}Nd into the first excited (0^+_1) state of ^{150}Sm. Gamma radiation emitted in decay of the excited state is detected.

[7] Supersedes KLAPDOR-KLEINGROTHAUS 02D. Authors present new analysis of event excess seen in Heidelberg-Moscow experiment at $\beta\beta$-decay energy. Enhanced statistics leads to a 4.2 σ evidence for observation of $0\nu\beta\beta$-decay and a finite Majorana neutrino mass. Stated error is purely statistical. No systematic errors are mentioned in the paper. More details can be found in KLAPDOR-KLEINGROTHAUS 04C.

[8] CaF_2 scintillation calorimeter ELEGANT VI used to set limit on $0\nu\beta\beta$-decay rate of ^{48}Ca. The stated half-life limit benefits from a downward fluctuation on the number of background events. The experimental sensitivity is 5.9×10^{21} yr. at 90 % CL. Replaces YOU 91 as the most stringest experiment using ^{48}Ca.

[9] Supersedes ALESSANDRELLO 00. Array of TeO_2 crystals in high resolution cryogenic calorimeter. Some enriched in ^{130}Te. Ground state to ground state decay.

[10] Decay into first excited state of daughter nucleus.

[11] Two neutrino decay into ground state. Relatively large error mainly due to uncertainties in background determination. Reported value is shorter than the geochemical measurements of KIRSTEN 83 and BERNATOWICZ 92 but in agreement with LIN 88 and TAKAOKA 96.

[12] Supersedes ALESSANDRELLO 00. Array of TeO_2 crystals in high resolution cryogenic calorimeter. Some enriched in ^{128}Te. Ground state to ground state decay.

[13] Calorimetric measurement of 2ν ground state decay of ^{116}Cd using enriched $CdWO_4$ scintillators. Agrees with EJIRI 95 and ARNOLD 96. Supersedes DANEVICH 00.

[14] Limit on 0ν decay of ^{116}Cd using enriched $CdWO_4$ scintillators. Supersedes DANEVICH 00.

[15] Limit on 0ν decay of ^{116}Cd into first excited 2^+ state of daughter nucleus using enriched $CdWO_4$ scintillators. Supersedes DANEVICH 00.

[16] Limit on 0ν decay of ^{116}Cd into first excited 0^+ state of daughter nucleus using enriched $CdWO_4$ scintillators. Supersedes DANEVICH 00.

[17] Limit on 0ν decay of ^{116}Cd into second excited 0^+ state of daughter nucleus using enriched $CdWO_4$ scintillators. Supersedes DANEVICH 00.

[18] Results of the Heidelberg-Moscow experiment (KLAPDOR-KLEINGROTHAUS 01 and GUENTHER 97) are reanalyzed using a new simulation of the complete background spectrum. The $\beta\beta 2\nu$-decay rate is deduced from a 41.57 kg-y exposure. The result is in agreement and supersedes the above referenced halflives with similar statistical and systematic errors.

[19] AALSETH 02B limit is based on 117 mol·yr of data using enriched Ge detectors. Background reduction by means of pulse shape analysis is applied to part of the data set. Reported limit is slightly less restrictive than that in KLAPDOR-KLEINGROTHAUS 01 However, it excludes part of the allowed half-life range reported in KLAPDOR-KLEINGROTHAUS 01B for the same nuclide. The analysis has been criticized in KLAPDOR-KLEINGROTHAUS 04B. The criticism was addressed and disputed in AALSETH 04.

[20] BERNABEI 02D report a limit for the 0ν, $0^+ \rightarrow 0^+$ decay of ^{134}Xe, present in the source at 17%, by considering the maximum number of events for this mode compatible with the fitted smooth background.

[21] BERNABEI 02D report a limit for the 0ν, $0^+ \rightarrow 0^+$ decay of ^{136}Xe, by considering the maximum number of events for this mode compatible with the fitted smooth background. The quoted sensitivity is 450×10^{21} yr. The Feldman and Cousins method is used to obtain the quoted limit.

[22] KLAPDOR-KLEINGROTHAUS 02D is an expanded version of KLAPDOR-KLEINGROTHAUS 01B. The authors re-evaluate the data collected by the Heidelberg-Moscow experiment (KLAPDOR-KLEINGROTHAUS 01) and present a more detailed description of their analysis of an excess of counts at the energy expected for neutrinoless double-beta decay. They interpret this excess, which has a significance of 2.2 to 3.1 σ depending on the data analysis, as evidence for the observation of Lepton Number violation and violation of Baryon minus Lepton Number. The analysis has been criticized by AALSETH 02 and others. The criticisms have been addressed in KLAPDOR-KLEINGROTHAUS 02. See also KLAPDOR-KLEINGROTHAUS 02B.

[23] ASHITKOV 01 result for 2ν of ^{100}Mo is in agreement with other determinations of that halflife.
[24] ASHITKOV 01 result for 0ν of ^{100}Mo is less stringent than EJIRI 01.
[25] DANEVICH 01 place limit on 0ν decay of ^{160}Gd using Gd_2SiO_5:Ce crystal scintillators. The limit is more stringent than KOBAYASHI 95.
[26] DANEVICH 01 place limits on 0ν decay of ^{160}Gd into excited 2^+ state of daughter nucleus using Gd_2SiO_5:Ce crystal scintillators.
[27] DEBRAECKELEER 01 performed an inclusive measurement of the $\beta\beta$ decay into the second excited state of the daughter nucleus. A novel coincidence technique counting the de-excitation photons is employed. The result agrees with BARABASH 95.
[28] EJIRI 01 uses tracking calorimeter and isotopically enriched passive source. Efficiencies were calculated assuming $\langle m_\nu \rangle$, $\langle \lambda \rangle$, or $\langle \eta \rangle$ driven decay. This is a continuation of EJIRI 96 which it supersedes.
[29] KLAPDOR-KLEINGROTHAUS 01 is a continuation of the work published in BAUDIS 99. Isotopically enriched Ge detectors are used in calorimetric measurement. The most stringent bound is derived from the data set in which pulse-shape analysis has been used to reduce background. Exposure time is 35.5 kg y. Supersedes BAUDIS 99 as most stringent result.
[30] KLAPDOR-KLEINGROTHAUS 01 is a measurement of the $\beta\beta 2\nu$-decay rate with higher statistics than GUENTHER 97. The reported value has a larger systematic error than their previous result.
[31] WIESER 01 reports an inclusive geochemical measurement of ^{96}Zr $\beta\beta$ half life. Their result agrees within 2σ with ARNOLD 99 but only marginally, within 3σ, with KAWASHIMA 93.
[32] BRUDANIN 00 determine the 2ν halflife of ^{48}Ca. Their value is less accurate than BALYSH 96.
[33] ARNOLD 99 measure directly the 2ν decay of Zr for the first time, using the NEMO-2 tracking detector and an isotopically enriched source. The lifetime is more accurate than the geochemical result of KAWASHIMA 93.
[34] ARNOLD 98 measure the 2ν decay of ^{82}Se by comparing the spectra in an enriched and natural selenium source using the NEMO-2 tracking detector. The measured half-life is in agreement, perhaps slightly shorter, than ELLIOTT 92.
[35] ARNOLD 98 determine the limit for 0ν decay to the ground state of ^{82}Se using the NEMO-2 tracking detector. The half-life limit is in agreement, but less stringent, than ELLIOTT 92.
[36] ARNOLD 98 determine the limit for 0ν decay to the excited 2^+ state of ^{82}Se using the NEMO-2 tracking detector.
[37] ALSTON-GARNJOST 97 report evidence for 2ν decay of ^{100}Mo. This decay has been also observed by EJIRI 91, DASSIE 95, and DESILVA 97.
[38] DESILVA 97 result for 2ν decay of ^{100}Mo is in agreement with ALSTON-GARNJOST 97 and DASSIE 95. This measurement has the smallest errors.
[39] DESILVA 97 result for 2ν decay of ^{150}Nd is in marginal agreement with ARTEMEV 93. It has smaller errors.
[40] DESILVA 97 do not explain whether their efficiency for 0ν decay of ^{150}Nd was calculated under the assumption of a $\langle m_\nu \rangle$, $\langle \lambda \rangle$, or $\langle \eta \rangle$ driven decay.
[41] ARNOLD 96 measure the 2ν decay of ^{116}Cd. This result is in agreement with EJIRI 95, but has smaller errors. Supersedes ARNOLD 95.
[42] BALYSH 96 measure the 2ν decay of ^{48}Ca, using a passive source of enriched ^{48}Ca in a TPC.
[43] TAKAOKA 96 measure the geochemical half-life of ^{130}Te. Their value is in disagreement with the quoted values of BERNATOWICZ 92 and KIRSTEN 83; but agrees with several other unquoted determinations, e.g., MANUEL 91.
[44] BARABASH 95 cannot distinguish 0ν and 2ν, but it is inferred indirectly that the 0ν mode accounts for less than 0.026% of their event sample. They also note that their result disagrees with the previous experiment by the NEMO group (BLUM 92).
[45] BERNATOWICZ 92 finds ^{128}Te/^{130}Te activity ratio from slope of ^{128}Xe/^{132}Xe vs ^{130}Xe/^{132}Xe ratios during extraction, and normalizes to lead-dated ages for the ^{130}Te

lifetime. The authors state that their results imply that "(a) the double beta decay of ^{128}Te has been firmly established and its half-life has been determined ... without any ambiguity due to trapped Xe interferences... (b) Theoretical calculations ... underestimate the [long half-lives of ^{128}Te ^{130}Te] by 1 or 2 orders of magnitude, pointing to a real suppression in the 2ν decay rate of these isotopes. (c) Despite [this], most $\beta\beta$-models predict a *ratio* of 2ν decay widths ... in fair agreement with observation." Further details of the experiment are given in BERNATOWICZ 93. Our listed half-life has been revised downward from the published value by the authors, on the basis of reevaluated cosmic-ray ^{128}Xe production corrections.

[46] TURKEVICH 91 observes activity in old U sample. The authors compare their results with theoretical calculations. They state "Using the phase-space factors of Boehm and Vogel (BOEHM 87) leads to matrix element values for the ^{238}U transition in the same range as deduced for ^{130}Te and ^{76}Ge. On the other hand, the latest theoretical estimates (STAUDT 90) give an upper limit that is 10 times lower. This large discrepancy implies either a defect in the calculations or the presence of a faster path than the standard two-neutrino mode in this case." See BOEHM 87 and STAUDT 90.

[47] Result agrees with direct determination of ELLIOTT 92.

[48] Inclusive half life inferred from mass spectroscopic determination of abundance of $\beta\beta$-decay product ^{130}Te in mineral kitkaite (NiTeSe). Systematic uncertainty reflects variations in U-Xe gas-retention-age derived from different uranite samples. Agrees with geochemical determination of TAKAOKA 96 and direct measurement of ARNABOLDI 03. Inconsistent with results of KIRSTEN 83 and BERNATOWICZ 92.

[49] Ratio of inclusive double beta half lives of ^{128}Te and ^{130}Te determined from minerals melonite ($NiTe_2$) and altaite (PbTe) by means of mass spectroscopic measurement of abundance of $\beta\beta$-decay products. As gas-retention-age could not be determined the authors use half life of ^{130}Te (LIN 88) to infer the half life of ^{128}Te. No estimate of the systematic uncertainty of this method is given. The directly determined half life ratio agrees with BERNATOWICZ 92. However, the inferred ^{128}Te half life disagrees with KIRSTEN 83 and BERNATOWICZ 92.

[50] KIRSTEN 83 reports "2σ" error. References are given to earlier determinations of the ^{130}Te lifetime.

$\langle m_\nu \rangle$, The Effective Weighted Sum of Majorana Neutrino Masses Contributing to Neutrinoless Double-β Decay

$\langle m_\nu \rangle = |\Sigma\ U_{1j}^2 m_{\nu_j}|$, where the sum goes from 1 to n and where n = number of neutrino generations, and ν_j is a Majorana neutrino. Note that U_{ej}^2, not $|U_{ej}|^2$, occurs in the sum. The possibility of cancellations has been stressed. In the following Listings, only best or comparable limits or lifetimes for each isotope are reported.

VALUE (eV)	*CL%*	*ISOTOPE*	*TRANSITION*	*METHOD*	*DOCUMENT ID*	
• • • We do not use the following data for averages, fits, limits, etc. • • •						
< 0.37–1.9	90	^{130}Te		Cryog. det.	[51] ARNABOLDI	04
< 0.8–1.2	90	^{100}Mo	0ν	NEMO-3	[52] ARNOLD	04
< 1.5–3.1	90	^{82}Se	0ν	NEMO-3	[52] ARNOLD	04
0.1–0.9	99.7	^{76}Ge		Enriched HP Ge	[53] KLAPDOR-K...	04A
< 7.2–44.7	90	^{48}Ca		CaF_2 scint.	[54] OGAWA	04
$<$ 1.1–2.6	90	^{130}Te		Cryog. det.	[55] ARNABOLDI	03
< 1.5–1.7	90	^{116}Cd	0ν	$^{116}CdWO_4$ scint.	[56] DANEVICH	03
< 0.33–1.35	90			Enriched HPGe	[57] AALSETH	02B
<2.9	90	^{136}Xe	0ν	Liquid Xe Scint.	[58] BERNABEI	02D

$0.39^{+0.17}_{-0.28}$		^{76}Ge	0ν	Enriched HPGe	[59] KLAPDOR-K...	02D
< 2.1–4.8	90	^{100}Mo	0ν	ELEGANT V	[60] EJIRI	01
< 0.35	90	^{76}Ge		Enriched HPGe	[61] KLAPDOR-K...	01
<23	90	^{96}Zr		NEMO-2	[62] ARNOLD	99
< 1.1–1.5		^{128}Te		Geochem	[63] BERNATOW...	92
<5	68	^{82}Se		TPC	[64] ELLIOTT	92
<8.3	76	^{48}Ca	0ν	CaF_2 scint.	YOU	91

[51] Supersedes ARNABOLDI 03. Cryogenic calorimeter. Reported range of limits due to use of different nuclear matrix element calculations.

[52] ARNOLD 04 limit is based on the nuclear matrix elements of SIMKOVIC 99, STOICA 01 and CIVITARESE 03.

[53] Supersedes KLAPDOR-KLEINGROTHAUS 02D. Event excess at $\beta\beta$-decay energy is used to derive Majorana neutrino mass using the nuclear matrix elements of STAUDT 90. The mass range shown is based on the authors evaluation of the uncertainties of the STAUDT 90 matrix element calculation. If this uncertainty is neglected, and only statistical errors are considered, the range in $\langle m \rangle$ becomes (0.2–0.6) eV at the 3 σ level.

[54] Calorimetric CaF_2 scintillator. Range of limits reflects authors' estimate of the uncertainty of the nuclear matrix elements. Replaces YOU 91 as the most stringest limit based on ^{48}Ca.

[55] Supersedes ALESSANDRELLO 00. Cryogenic calorimeter search. Reported a range reflecting uncertainty in nuclear matrix element calculations.

[56] Limit for $\langle m_\nu \rangle$ is based on the nuclear matrix elements of STAUDT 90 and ARNOLD 96. Supersedes DANEVICH 00.

[57] AALSETH 02B reported range of limits on $\langle m_\nu \rangle$ reflects the spread of theoretical nuclear matrix elements. Excludes part of allowed mass range reported in KLAPDOR-KLEINGROTHAUS 01B.

[58] BERNABEI 02D limit is based on the matrix elements of SIMKOVIC 02. The range of neutrino masses based on a variety of matrix elements is 1.1–2.9 eV.

[59] KLAPDOR-KLEINGROTHAUS 02D is a detailed description of the analysis of the data collected by the Heidelberg-Moscow experiment, previously presented in KLAPDOR-KLEINGROTHAUS 01B. Matrix elements in STAUDT 90 have been used. See the footnote in the preceding table for further details. See also KLAPDOR-KLEINGROTHAUS 02B.

[60] The range of the reported $\langle m_\nu \rangle$ values reflects the spread of the nuclear matrix elements. On axis value assuming $\langle \lambda \rangle = \langle \eta \rangle = 0$.

[61] KLAPDOR-KLEINGROTHAUS 01 uses the calculation by STAUDT 90. Using several other models in the literature could worsen the limit up to 1.2 eV. This is the most stringent experimental bound on m_ν. It supersedes BAUDIS 99B.

[62] ARNOLD 99 limit based on the nuclear matrix elements of STAUDT 90.

[63] BERNATOWICZ 92 finds these majorana neutrino mass limits assuming that the measured geochemical decay width is a limit on the 0ν decay width. The range is the range found using matrix elements from HAXTON 84, TOMODA 87, and SUHONEN 91. Further details of the experiment are given in BERNATOWICZ 93.

[64] ELLIOTT 92 uses the matrix elements of HAXTON 84.

Limits on Lepton-Number Violating ($V+A$) Current Admixture

For reasons given in the discussion at the beginning of this section, we list only results from 1989 and later. $\langle \lambda \rangle = \lambda \sum U_{ej} V_{ej}$ and $\langle \eta \rangle = \eta \sum U_{ej} V_{ej}$, where the sum is over the number of neutrino generations. This sum vanishes for massless or unmixed neutrinos. In the following Listings, only best or comparable limits or lifetimes for each isotope are reported.

$\langle\lambda\rangle$ (10^{-6})	CL%	$\langle\eta\rangle$ (10^{-8})	CL%	ISOTOPE	METHOD	DOCUMENT ID	
• • • We do not use the following data for averages, fits, limits, etc. • • •							
< 1.5–2.0	90			^{100}Mo	0ν, NEMO-3	[65] ARNOLD	04
< 3.2–3.8	90			^{82}Se	0ν, NEMO-3	[66] ARNOLD	04
< 1.6–2.4	90	< 0.9–5.3	90	^{130}Te	Cryog. det.	[67] ARNABOLDI	03
<2.2	90	<2.5	90	^{116}Cd	$^{116}CdWO_4$ scint.	[68] DANEVICH	03
< 3.2–4.7	90	< 2.4–2.7	90	^{100}Mo	ELEGANT V	[69] EJIRI	01
<1.1	90	<0.64	90	^{76}Ge	Enriched HPGe	[70] GUENTHER	97
<4.4	90	<2.3	90	^{136}Xe	TPC	[71] VUILLEUMIER	93
		<5.3		^{128}Te	Geochem	[72] BERNATOW...	92

[65] ARNOLD 04 use the matrix elements of SUHONEN 94 to obtain a limit for $\langle\lambda\rangle$, no limit for $\langle\eta\rangle$ is given. This limit is more stringent than the limit in EJIRI 01 for the same nucleus.

[66] ARNOLD 04 use the matrix elements of TOMODA 91 and SUHONEN 91 to obtain a limit for $\langle\lambda\rangle$, no limit for $\langle\eta\rangle$ is given.

[67] Supersedes ALESSANDRELLO 00. Cryogenic calorimeter search. Reported a range reflecting uncertainty in nuclear matrix element calculations.

[68] Limits for $\langle\lambda\rangle$ and $\langle\eta\rangle$ are based on nuclear matrix elements of STAUDT 90. Supersedes DANEVICH 00.

[69] The range of the reported $\langle\lambda\rangle$ and $\langle\eta\rangle$ values reflects the spread of the nuclear matrix elements. On axis value assuming $\langle m_\nu\rangle$=0 and $\langle\lambda\rangle$=$\langle\eta\rangle$=0, respectively.

[70] GUENTHER 97 limits use the matrix elements of STAUDT 90. Supersedes BALYSH 95 and BALYSH 92.

[71] VUILLEUMIER 93 uses the matrix elements of MUTO 89. Based on a half-life limit 2.6×10^{23} y at 90%CL.

[72] BERNATOWICZ 92 takes the measured geochemical decay width as a limit on the 0ν width, and uses the SUHONEN 91 coefficients to obtain the least restrictive limit on η. Further details of the experiment are given in BERNATOWICZ 93.

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