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J=1
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## THE MASS OF THE W BOSON

Revised November 2003 by C. Caso (University of Genova) and A. Gurtu (Tata Institute).

Till 1995 the production and study of the $W$ boson was the exclusive domain of the $\bar{p} p$ colliders at CERN and FNAL. $W$ production in these hadron colliders is tagged by a high $p_{T}$ lepton from $W$ decay. Owing to unknown parton-parton effective energy and missing energy in the longitudinal direction, the experiments reconstruct only the transverse mass of the $W$ and derive the $W$ mass from comparing the transverse mass distribution with Monte Carlo predictions as a function of $M_{W}$.

Beginning 1996 the energy of LEP increased to above 161 GeV , the threshold for $W$-pair production. A precise knowledge of the $e^{+} e^{-}$center-of-mass energy enables one to reconstruct the $W$ mass even if one of them decays leptonically. At LEP two methods have been used to obtain the $W$ mass. In the first method the measured $W$-pair production cross sections, $\sigma\left(e^{+} e^{-} \rightarrow W^{+} W^{-}\right)$, have been used to determine the $W$ mass using the predicted dependence of this cross section on $M_{W}$ (see Fig. 1). At 161 GeV , which is just above the $W$-pair production threshold, this dependence is a much more sensitive function of the $W$ mass than at the higher energies ( 172 to 208 GeV ) at which LEP has run during 1996-2000. In the second method, which is used at the higher energies, the $W$ mass has been determined by directly reconstructing the $W$ from its decay products.

Each LEP experiment has combined their own mass values properly taking into account the common systematic errors. In order to compute the LEP average $W$ mass each experiment


Figure 1: The $W$-pair cross section as a function of the center-of-mass energy. The data points are the LEP averages. The solid lines are predictions from different models of $W W$ production. For comparison the figure contains also the cross section if the ZWW coupling did not exist (dashed line), or if only the $t$-channel $\nu_{e}$ exchange diagram existed (dotted-dashed line). (Figure from http://lepewwg.web.cern.ch/ LEPEWWG/lepww/4f/Summer03/ wwxsec_nocouplings_2003.eps)
See full-color version on color pages at end of book.
has provided its measured $W$ mass for the $q q q q$ and $q q \ell \nu$ channels at each center-of-mass energy along with a detailed break-up of errors (statistical and uncorrelated, partially correlated and fully correlated systematics [1]). These have been properly combined to obtain a preliminary LEP $W$ mass $=$ $80.412 \pm 0.042 \mathrm{GeV}$ [2]. Errors due uncertainties in LEP energy ( 17 MeV ) and possible effect of color reconnection (CR) and Bose-Einstein (BE) correlations between quarks from different $W$ 's are included. The mass difference between $q q q q$ and $q q \ell \nu$ final states (due to possible CR and BE effects) is $+22 \pm 43$ MeV .

The two Tevatron experiments have also carried out the exercise of identifying common systematic errors and averaging with CERN UA2 data obtain an average $W$ mass [2]= $80.454 \pm 0.059 \mathrm{GeV}$.

Combining the above $W$ mass values from LEP and hadron colliders, which are based on all published and unpublished results, and assuming no common systematics between them, yields an average $W$ mass of $80.426 \pm 0.034 \mathrm{GeV}$.

Finally a fit to this directly determined $W$ mass together with measurements on the ratio of $W$ to $Z$ mass $\left(M_{W} / M_{Z}\right)$ and on their mass difference $\left(M_{Z}-M_{W}\right)$ yields a world average $W$-boson mass of $80.425 \pm 0.033 \mathrm{GeV}$.

The Standard Model prediction from the electroweak fit, using $Z$-pole data plus $m_{\text {top }}$ measurement, gives a $W$-boson mass of $80.378 \pm 0.023 \mathrm{GeV}[2]$.

OUR FIT in the listing below is obtained by combining only published LEP and $p-\bar{p}$ Collider results using the same procedure as above.

## References

1. The LEP Collaborations: ALEPH, DELPHI, L3, OPAL, the LEP Electroweak Working Group, and the SLD Heavy Flavour Group, CERN-EP-2002-091, hep-ex/0212036 (17 December 2002 ).
2. P. Wells, "Experimental Tests of the Standard Model," Int. Europhysics Conference on High-Energy Physics (Aachen, Germany, 17-23 July 2003).

## W MASS

To obtain the world average, common systematics between experiments are properly taken into account. The procedure for averaging the LEP data is given in the note LEPEWWG/MASS/2002-01 (March 11, 2002), accessible at http://lepewwg.web.cern.ch/LEPEWWG/lepww/mw/pdg_2002/. The LEP average $W$ mass based on published results is $80.400 \pm 0.056$ GeV . The combined $p \bar{p}$ collider data yields an average $W$ mass of $80.454 \pm 0.059 \mathrm{GeV}$ (KOTWAL 02).

OUR FIT uses these average LEP and $p \bar{p}$ collider $W$ mass values together with the $Z$ mass, the $W$ to $Z$ mass ratio, and mass difference measurements.

| VALUE (GeV) | EVTS | DOCUMENT ID | TECN | COMMENT |
| :---: | :---: | :---: | :---: | :---: |
| 80.425 $\pm$ 0.038 OUR FIT |  |  |  |  |
| $80.41 \pm 0.41 \pm 0.13$ | 1101 | 1 ABBIENDI | 03C OPAL | $E_{\mathrm{Cm}}^{e \mathrm{~m}}$ = $=183-207 \mathrm{GeV}$ |
| $80.483 \pm 0.084$ | 49247 | 2 ABAZOV | 02D D0 | $E_{\mathrm{cm}}^{p \bar{p}}=1.8 \mathrm{TeV}$ |
| $80.432 \pm 0.066 \pm 0.045$ | 2789 | 3 ABBIENDI | 01F OPAL | $\begin{gathered} E_{\mathrm{Cm}}^{e e}=161+172+183 \\ \quad+189 \mathrm{GeV} \end{gathered}$ |
| $80.359 \pm 0.074 \pm 0.049$ | 3077 | ${ }^{4}$ ABREU | 01K DLPH | $\begin{aligned} & E_{\mathrm{Cm}}^{\mathrm{ee}}=161+172+183 \\ & \pm 189 \mathrm{GeV} \end{aligned}$ |
| $80.433 \pm 0.079$ | 53841 | 5 AFFOLDER | 01E CDF | $E_{\mathrm{cm}}^{p \bar{p}}=1.8 \mathrm{TeV}$ |
| $80.418 \pm 0.061 \pm 0.047$ | 2977 | 6 BARATE | 00T ALEP | $\begin{aligned} & E_{\mathrm{Cm}}^{e e}=161+172+183 \\ & \quad+189 \mathrm{GeV} \end{aligned}$ |
| $80.61 \pm 0.15$ | 801 | 7 ACCIARRI | 99 L3 | $\begin{aligned} & E_{\mathrm{Cm}}^{\mathrm{ee}}=161+172+183 \\ & \mathrm{GeV} \end{aligned}$ |
| - - We do not use the following data for averages, fits, limits, etc. - . |  |  |  |  |
| $80.3 \pm 2.1 \pm 1.2 \pm 1.0$ | 645 | ${ }^{8}$ CHEKANOV | 02C ZEUS | $e^{-} \underset{\mathrm{GeV}}{p} \nu_{e} \mathrm{X}, \sqrt{s}=318$ |
| $79.9 \pm 2.2 \pm 2.3$ | 700 | 9 ADLOFF | 01A H1 | $e^{-} \underset{320 \mathrm{GeV}}{\nu_{e} X, \sqrt{s}} \approx$ |
| $80.482 \pm 0.091$ | 45394 | 10 ABBOTT | 00 D0 | Repl. by ABAZOV 02D |
| $80.9 \pm 3.7 \pm 3.7$ | 700 | 11 ADLOFF | 00b H1 | $\begin{aligned} & e^{+} p \rightarrow \bar{\nu}_{e} X, \sqrt{s} \approx \\ & 300 \mathrm{GeV} \end{aligned}$ |
| $81.4{ }_{-2.6}^{+2.7} \pm 2.0_{-3.0}^{+3.3}$ | 1086 | 12 BREITWEG | 00D ZEUS | $\begin{aligned} & e^{+}+\underset{\nu}{p} \bar{\nu}_{e} \mathrm{X}, \sqrt{s} \approx \\ & 300 \mathrm{GeV} \end{aligned}$ |
| $80.38 \pm 0.12 \pm 0.05$ | 701 | 13 ABBIENDI | 99C OPAL | Repl. by ABBIENDI 01F |
| $80.270 \pm 0.137 \pm 0.048$ | 809 | 14 ABREU | 99T DLPH | Repl. by ABREU 01k |
| $80.423 \pm 0.112 \pm 0.054$ | 812 | 15 BARATE | 99 ALEP | Repl. by BARATE 00t |

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| 80.80 | $\begin{aligned} & 0.48 \\ & +\quad 0.42 \end{aligned}$ | $\pm 0.03$ | 20 | 16 ACCIARRI | 97 | L3 | Repl. by ACCIARRI 99 |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 80.5 | $+\quad 1.4$ $+\quad 2.4$ | $\pm 0.3$ | 94 | 17 ACCIARRI | 97M | L3 | Repl. by ACCIARRI 99 |
| 80.71 | $\begin{array}{ll} + & 0.34 \\ - & 0.35 \end{array}$ | $\pm 0.09$ | 101 | 18 ACCIARRI | 97S | L3 | Repl. by ACCIARRI 99 |
| 80.41 | $\pm 0.18$ |  | 8986 | 19 ABE | 95P | CDF | Repl. by AFFOLDER 01E |
| 80.84 | $\pm 0.22$ | $\pm 0.83$ | 2065 | 20 ALITTI | 92B | UA2 | See $W / Z$ ratio below |
| 80.79 | $\pm 0.31$ | $\pm 0.84$ |  | 21 ALITTI | 90B | UA2 | $E_{\mathrm{cm}}^{p \bar{p}}=546,630 \mathrm{GeV}$ |
| 80.0 | $\pm 3.3$ | $\pm 2.4$ | 22 | 22 ABE | 891 | CDF | $E_{\mathrm{cm}}^{p \bar{p}}=1.8 \mathrm{TeV}$ |
| 82.7 | $\pm 1.0$ | $\pm 2.7$ | 149 | 23 ALBAJAR | 89 | UA1 | $E_{\mathrm{cm}}^{p \bar{p}}=546,630 \mathrm{GeV}$ |
| 81.8 | $\begin{array}{r} +6.0 \\ -\quad 5.3 \end{array}$ | $\pm 2.6$ | 46 | 24 ALBAJAR | 89 | UA1 | $E_{\mathrm{cm}}^{p \bar{p}}=546,630 \mathrm{GeV}$ |
| 89 | $\pm 3$ | $\pm 6$ | 32 | 25 ALBAJAR | 89 | UA1 | $E_{\mathrm{cm}}^{p \bar{p}}=546,630 \mathrm{GeV}$ |
| 81. | $\pm 5$. |  | 6 | ARNISON | 83 | UA1 | $E_{\mathrm{Cm}}^{e \mathrm{e}}=546 \mathrm{GeV}$ |
| 80. | $\begin{array}{r} +10 . \\ +\quad 6 . \end{array}$ |  | 4 | BANNER | 83B | UA2 | Repl. by ALITTI 90B |

${ }^{1}$ ABBIENDI 03C determine the mass of the $W$ boson using fully leptonic decays $W^{+} W^{-} \rightarrow \ell \nu_{\ell} \ell^{\prime} \nu_{\ell^{\prime}}$. They use the measured energies of the charged leptons and an approximate kinematic reconstruction of the event (both neutrinos are assumed in the same plane as the charged leptons) to get a $W$ pseudo-mass. All these variables are combined in a simultaneous maximum likelihood fit. The systematic error is dominated by the uncertainty on the lepton energy.
${ }^{2}$ ABAZOV 02D improve the measurement of the $W$-boson mass including $W \rightarrow e \nu_{e}$ events in which the electron is close to a boundary of a central electromagnetic calorimeter module. Properly combining the results obtained by fitting $m_{T}(W), p_{T}(e)$, and $p_{T}(\nu)$, this sample provides a mass value of $80.574 \pm 0.405 \mathrm{GeV}$. The value reported here is a combination of this measurement with all previous $D \varnothing W$-boson mass measurements.
${ }^{3}$ ABBIENDI 01F obtain this value properly combining results obtained from a direct $W$ mass reconstruction at 172,183 , and 189 GeV with that from measurement of the $W$-pair production cross section at 161 GeV . The systematic error includes $\pm 0.017 \mathrm{GeV}$ due to LEP energy uncertainty and $\pm 0.028 \mathrm{GeV}$ due to possible color reconnection and Bose-Einstein effects in the purely hadronic final state.
${ }^{4}$ ABREU 01 K obtain this value properly combining results obtained from a direct $W$ mass reconstruction at 172,183 , and 189 GeV with those from measurements of $W$ pair production cross sections at 161,172 , and 183 GeV . The systematic error includes $\pm 0.017 \mathrm{GeV}$ due to the beam energy uncertainty and $\pm 0.033 \mathrm{GeV}$ due to possible color reconnection and Bose-Einstein effects in the purely hadronic final state.
${ }^{5}$ AFFOLDER 01E fit the transverse mass spectrum of $30115 \mathrm{~W} \rightarrow e \nu_{e}$ events ( $M_{W}=$ $80.473 \pm 0.065 \pm 0.092 \mathrm{GeV})$ and of $14740 W \rightarrow \mu \nu_{\mu}$ events $\left(M_{W}=80.465 \pm 0.100 \pm\right.$ 0.103 GeV ) obtained in the run IB (1994-95). Combining the electron and muon results, accounting for correlated uncertainties, yields $M_{W}=80.470 \pm 0.089 \mathrm{GeV}$. They combine this value with their measurement of ABE 95P reported in run IA (1992-93) to obtain the quoted value.
${ }^{6}$ BARATE 00 T obtain this value properly combining results obtained from a direct $W$ mass reconstruction at 172,183 , and 189 GeV with those from measurements of $W$-pair production cross sections at 161 and 172 GeV . The systematic error includes $\pm 0.017$ GeV due to LEP energy uncertainty and $\pm 0.019 \mathrm{GeV}$ due to possible color reconnection and Bose-Einstein effects in the purely hadronic final state.
${ }^{7}$ ACCIARRI 99 obtain this value properly combining results obtained from a direct $W$ mass reconstruction at 172 and 183 GeV with those from the measurements of the total $W$ pair production cross sections at 161 and 172 GeV . The value of the mass obtained from the direct reconstruction at 172 and 183 GeV is $M(W)=80.58 \pm 0.14 \pm 0.08 \mathrm{GeV}$.
${ }^{8}$ CHEKANOV 02C fit the $Q^{2}$ dependence $\left(200<Q^{2}<60000 \mathrm{GeV}^{2}\right)$ of the charged-current differential cross sections with a propagator mass fit. The last error is due to the uncertainty on the probability density functions.
${ }^{9}$ ADLOFF 01A fit the $Q^{2}$ dependence ( $150<Q^{2}<30000 \mathrm{GeV}^{2}$ ) of the charged-current double-differential cross sections with a propagator mass fit. The second error includes 2.1 GeV due to the theoretical uncertainties.

10 ABBOTT 00 use $W \rightarrow e \nu_{e}$ events to measure the $W$ mass with a fit to the transverse mass distribution. The result quoted here corresponds to electrons detected both in the forward and in the central calorimeters for the data recorded in 1992-1995. For the large rapidity electrons recorded in 1994-1995, the analysis combines results obtained from $m_{T}, p_{T}(e)$, and $p_{T}(\nu)$.
11 ADLOFF 00B fit the $Q^{2}$ dependence ( $300<Q^{2}<15000 \mathrm{GeV}^{2}$ ) of the charged-current double-differential cross sections with a propagator mass fit. The second error is due to the theoretical uncertainties.
12 BREITWEG 00D fit the $Q^{2}$ dependence ( $200<Q^{2}<22500 \mathrm{GeV}^{2}$ ) of the chargedcurrent differential cross sections with a propagator mass fit. The last error is due to the uncertainty on the probability density functions.
13 ABBIENDI 99C obtain this value properly combining results from a direct $W$ mass reconstruction at 172 and 183 GeV with that from the measurement of the total $W$-pair production cross section at 161 GeV . The systematic error includes an uncertainty of $\pm 0.02 \mathrm{GeV}$ due to the possible color-reconnection and Bose-Einstein effects in the purely hadronic final states and an uncertainty of $\pm 0.02 \mathrm{GeV}$ due to the beam energy.
14 ABREU 99T obtain this value properly combining results obtained from a direct $W$ mass reconstruction at 172 and 183 GeV with those from measurement of $W$-pair production cross sections at 161,172 , and 183 GeV . The systematic error includes $\pm 0.021 \mathrm{GeV}$ due to the beam energy uncertainty and $\pm 0.030 \mathrm{GeV}$ due to possible color reconnection and Bose-Einstein effects in the purely hadronic final state.
15 BARATE 99 obtain this value properly combining results from a direct $W$ mass reconstruction at 172 and 183 GeV with those from the measurements of the total $W$-pair production cross sections at 161 and 172 GeV . The systematic error includes $\pm 0.023 \mathrm{GeV}$ due to LEP energy uncertainty and $\pm 0.021 \mathrm{GeV}$ due to theory uncertainty on account of possible color reconnection and Bose-Einstein correlations.
16 ACCIARRI 97 derive this value from their measured $W-W$ production cross section $\sigma_{W W}=2.89_{-0.70}^{+0.81} \pm 0.14 \mathrm{pb}$ using the Standard Model dependence of $\sigma_{W W}$ on $M_{W}$ at the given c.m. energy. Statistical and systematic errors are added in quadrature and the last error of $\pm 0.03 \mathrm{GeV}$ arises from the beam energy uncertainty. The same result is given by a fit of the production cross sections to the data.
17 ACCIARRI 97M derive this value from their measured $W W$ production cross section $\sigma_{W} W=12.27_{-1}^{+1.42} \pm 0.23 \mathrm{pb}$ using the Standard Model dependence of $\sigma_{W} W$ on $M_{W}$ at the given c.m. energy. Combining with ACCIARRI 97 authors find $M(W)=$ $80.78+0.41 \pm 0.03 \mathrm{GeV}$ where the last error is due to beam energy uncertainty.
18 ACCIARRI 975 obtain this value from a fit to the reconstructed $W$ mass distribution. The $W$ width was taken as its Standard Model value at the fitted $W$ mass. When both $W$ mass and width are varied they obtain $M(W)=80.72_{-0.31}^{+0.31} \pm 0.09 \mathrm{GeV}$. The systematic error includes $\pm 0.03 \mathrm{GeV}$ due to the beam energy uncertainty and $\pm 0.05 \mathrm{GeV}$ due to the possible color reconnection and Bose-Einstein effects in the purely hadronic final state. Combining with ACCIARRI 97 and ACCIARRI 97M authors find: $M(W)=$ $80.75_{-0.27}^{+0.26} \pm 0.03$ (LEP) GeV.
${ }^{19} \mathrm{ABE} 95 \mathrm{P}$ use $3268 \mathrm{~W} \rightarrow \mu \nu_{\mu}$ events to find $M=80.310 \pm 0.205 \pm 0.130 \mathrm{GeV}$ and $5718 W \rightarrow e \nu_{e}$ events to find $M=80.490 \pm 0.145 \pm 0.175 \mathrm{GeV}$. The result given here combines these while accounting for correlated uncertainties.
${ }^{20}$ ALITTI 92B result has two contributions to the systematic error $( \pm 0.83)$; one $( \pm 0.81)$ cancels in $m_{W} / m_{Z}$ and one $( \pm 0.17)$ is noncancelling. These were added in quadrature. We choose the ALITTI 92B value without using the LEP $m_{Z}$ value, because we perform our own combined fit.
${ }^{21}$ There are two contributions to the systematic error ( $\pm 0.84$ ): one ( $\pm 0.81$ ) which cancels in $m_{W} / m_{Z}$ and one ( $\pm 0.21$ ) which is non-cancelling. These were added in quadrature.
${ }^{22} \mathrm{ABE} 891$ systematic error dominated by the uncertainty in the absolute energy scale.
${ }^{23}$ ALBAJAR 89 result is from a total sample of $299 \mathrm{~W} \rightarrow e \nu$ events.
${ }^{24}$ ALBAJAR 89 result is from a total sample of $67 W \rightarrow \mu \nu$ events.
${ }^{25}$ ALBAJAR 89 result is from $W \rightarrow \tau \nu$ events.

## W/Z MASS RATIO

The fit uses the $W$ and $Z$ mass, mass difference, and mass ratio measurements.

| VALUE | EVTS | DOCUMENT ID | TECN | COMMENT |
| :---: | :---: | :---: | :---: | :---: |
| $0.88197 \pm 0.00042$ OUR FIT |  |  |  |  |
| $0.8821 \pm 0.0011 \pm 0.0008$ | 28323 | 26 ABBOTT | 98N D0 | $E_{\mathrm{cm}}^{p \bar{p}}=1.8 \mathrm{TeV}$ |
| $0.88114 \pm 0.00154 \pm 0.00252$ | 5982 | 27 ABBOTT | 98P D0 | $E_{\mathrm{cm}}^{p \bar{p}}=1.8 \mathrm{TeV}$ |
| $0.8813 \pm 0.0036 \pm 0.0019$ | 156 | 28 ALITTI | 92B UA2 | $E_{\mathrm{cm}}^{p \bar{p}}=630 \mathrm{GeV}$ |
| ${ }^{26}$ ABBOTT 98N obtain this from a study of $28323 W \rightarrow e \nu_{e}$ and $3294 Z \rightarrow e^{+} e^{-}$ decays. Of this latter sample, 2179 events are used to calibrate the electron energy scale. |  |  |  |  |
| 27 ABBOTT 98P obtain this from a study of $5982 \mathrm{~W} \rightarrow e \nu_{e}$ events. The systematic error includes an uncertainty of $\pm 0.00175$ due to the electron energy scale. |  |  |  |  |

$$
m_{Z}-m_{W}
$$

The fit uses the $W$ and $Z$ mass, mass difference, and mass ratio measurements.

VALUE $(\mathrm{GeV})$

## $10.763 \pm 0.038$ OUR FIT

$10.4 \quad \pm 1.4 \pm 0.8$

-     - We do not use the following data for averages, fits, limits, etc.
$11.3 \pm 1.3 \pm 0.9$
ANSARI
87 UA2 $E_{\mathrm{cm}}^{p \bar{p}}=546,630 \mathrm{GeV}$

$$
m_{w^{+}}-m_{W^{-}}
$$

Test of CPT invariance.

| $\frac{\operatorname{VALUE}(\mathrm{GeV})}{\mathbf{- 0 . 1 9 \pm \mathbf { 0 . 5 8 }}} \frac{\text { EVTS }}{1722}$ | $\frac{\text { DOCUMENT ID }}{\text { ABE }} \frac{\text { TECN }}{\text { COMMENT }}$ |
| :--- | :--- | :--- | :--- |

## W WIDTH

The CDF and $\mathrm{D} \emptyset$ widths labelled "extracted value" are obtained by measuring $R=[\sigma(W) / \sigma(Z)]\left[\Gamma\left(W \rightarrow \ell \nu_{\ell}\right)\right] /(\mathrm{B}(Z \rightarrow \ell \ell) \Gamma(W))$ where the bracketed quantities can be calculated with plausible reliability. $\Gamma(W)$ is then extracted by using a value of $\mathrm{B}(Z \rightarrow \ell \ell)$ measured at LEP. The UA1 and UA2 widths used $R=[\sigma(W) / \sigma(Z)]\left[\Gamma\left(W \rightarrow \quad \ell \nu_{\ell}\right) / \Gamma(Z \rightarrow\right.$ $\ell \ell)] \Gamma(Z) / \Gamma(W)$ and the measured value of $\Gamma(Z)$. The Standard Model prediction is $2.0921 \pm 0.0025 \mathrm{GeV}$ (see Review on "Electroweak model and constraints on new physics" in this Edition).

To obtain OUR FIT, the correlation between systematics for the Direct Measurements is properly taken into account. The following notes may be consulted for details as well as the respective average values: for the LEP experiments the note LEPEWWG/MASS/2002-01 (http://lepewwg.web.cern.ch/LEPEWWG/lepww/mw/pdg_2002/) of 11 March 2002 and for the Tevatron experiments the note FERMILAB-FN-716 of 1 July 2002 (KOTWAL 02). The respective average values ( $2.17 \pm 0.12 \mathrm{GeV}$ from LEP and $2.115 \pm 0.105 \mathrm{GeV}$ from Tevatron) yield an average $W$ width of $2.139 \pm 0.079 \mathrm{GeV}$ coming from direct measurements. Combined with the Extracted Values one obtains the quoted value.

| $\underline{V A L U E ~(G e V) ~ C L \% ~}$ | EVTS | DOCUMENT ID | TECN | COMMENT |
| :---: | :---: | :---: | :---: | :---: |
| $2.124 \pm 0.041$ OUR FIT |  |  |  |  |
| $2.23{ }^{+0.15} \pm 0.14{ }^{+0.10}$ | 294 | 29 ABAZOV | 02E D0 | Direct meas. |
| $2.04 \pm 0.16 \pm 0.09$ | 2756 | 30 ABBIENDI | 01F OPAL | $\begin{gathered} E_{\mathrm{Cm}}^{e e}=172+183 \\ +189 \mathrm{GeV} \end{gathered}$ |
| $2.266 \pm 0.176 \pm 0.076$ | 3005 | 31 ABREU | 01K DLPH | $\begin{aligned} & E_{\mathrm{Cm}}^{e e}=183+189 \\ & \mathrm{GeV} \end{aligned}$ |
| $2.152 \pm 0.066$ | 79176 | 32 ABBOTT | 00b D0 | Extracted value |
| $2.05 \pm 0.10 \pm 0.08$ | 662 | 33 AFFOLDER | 00m CDF | Direct meas. |
| $2.24 \pm 0.20 \pm 0.13$ | 1711 | 34 BARATE | 00T ALEP | $E_{\mathrm{cm}}^{e \mathrm{e}}=189 \mathrm{GeV}$ |
| $1.97 \pm 0.34 \pm 0.17$ | 687 | 35 ACCIARRI | 99 L3 | $\begin{aligned} & E_{\mathrm{Cm}}^{e \mathrm{ee}}=172+183 \\ & \mathrm{GeV} \end{aligned}$ |
| $2.064 \pm 0.060 \pm 0.059$ |  | 36 ABE | 95w CDF | Extracted value |
| $2.10{ }_{-0.13}^{+0.14} \pm 0.09$ | 3559 | 37 ALITTI | 92 UA2 | Extracted value |
| $2.18{ }_{-0.24}^{+0.26} \pm 0.04$ |  | 38 ALBAJAR | 91 UA1 | Extracted value |

-     - We do not use the following data for averages, fits, limits, etc. - -

| $1.84 \pm 0.32 \pm 0.20$ | 674 | 39 ABBIENDI | 99C OPAL | Repl. by ABBIENDI 01F |
| :---: | :---: | :---: | :---: | :---: |
| $2.044 \pm 0.097$ | 11858 | 40 ABBOTT | 99H D0 | Repl. by ABBOTT 00B |
| $2.48 \pm 0.40 \pm 0.10$ | 737 | 41 ABREU | 99T DLPH | Repl. by ABREU 01K |
| $2.126_{-0.048}^{+0.052} \pm 0.035$ |  | 42 BARATE | 99I ALEP | $\begin{aligned} & E_{\mathrm{Cm}}^{e e}= \\ & \quad 161+172+183 \\ & \quad \mathrm{GeV} \end{aligned}$ |
| $1.74{ }_{-0.78}^{+0.88} \pm 0.25$ | 101 | 43 ACCIARRI | 97S L3 | Repl. by ACCIARRI 99 |


${ }^{29}$ ABAZOV 02E obtain this result fitting the high-end tail ( $90-200 \mathrm{GeV}$ ) of the transversemass spectrum in semileptonic $W \rightarrow e \nu_{e}$ decays.
${ }^{30}$ ABBIENDI 01F obtain this value from a fit to the reconstructed $W$ mass distribution using data at 172,183 , and 189 GeV . The systematic error includes $\pm 0.010 \mathrm{GeV}$ due to LEP energy uncertainty and $\pm 0.078 \mathrm{GeV}$ due to possible color reconnection and Bose-Einstein effects in the purely hadronic final state.
31 ABREU 01K obtain this value properly combining results obtained at 183 and 189 GeV using $W W \rightarrow \ell \bar{\nu}_{\ell} q \bar{q}$ and $W W \rightarrow q \bar{q} q \bar{q}$ decays. The systematic error includes an uncertainty of $\pm 0.052 \mathrm{GeV}$ due to possible color reconnection and Bose-Einstein effects in the purely hadronic final state.
32 ABBOTT 00B measure $R=10.43 \pm 0.27$ for the $W \rightarrow e \nu_{e}$ decay channel. They use the SM theoretical predictions for $\sigma(W) / \sigma(Z)$ and $\Gamma\left(W \rightarrow e \nu_{e}\right)$ and the world average for $\mathrm{B}(Z \rightarrow e e)$. The value quoted here is obtained combining this result ( $2.169 \pm 0.070$ GeV ) with that of ABBOTT 99H.
33 AFFOLDER 00M fit the high transverse mass ( $100-200 \mathrm{GeV}$ ) $W \rightarrow e \nu_{e}$ and $W \rightarrow$ $\mu \nu_{\mu}$ events to obtain $\Gamma(W)=2.04 \pm 0.11$ (stat) $\pm 0.09$ (syst) GeV . This is combined with the earlier CDF measurement (ABE 95C) to obtain the quoted result.
${ }^{34}$ BARATE 00T obtain this value using $W W \rightarrow q \bar{q} q \bar{q}, W W \rightarrow e \nu_{e} q \bar{q}$, and $W W \rightarrow$ $\mu \nu_{\mu} q \bar{q}$ decays. The systematic error includes $\pm 0.015 \mathrm{GeV}$ due to LEP energy uncertainty and $\pm 0.080 \mathrm{GeV}$ due to possible color reconnection and Bose-Einstein effects in the purely hadronic final state.
${ }^{35}$ ACCIARRI 99 obtain this value from a fit to the reconstructed $W$ mass distribution using data at 172 and 183 GeV .
${ }^{36} \mathrm{ABE} 95 \mathrm{~W}$ measured $R=10.90 \pm 0.32 \pm 0.29$. They use $m_{W}=80.23 \pm 0.18 \mathrm{GeV}$, $\sigma(W) / \sigma(Z)=3.35 \pm 0.03, \Gamma(W \rightarrow e \nu)=225.9 \pm 0.9 \mathrm{MeV}, \Gamma\left(Z \rightarrow e^{+} e^{-}\right)=$ $83.98 \pm 0.18 \mathrm{MeV}$, and $\Gamma(Z)=2.4969 \pm 0.0038 \mathrm{GeV}$.
37 ALITTI 92 measured $R=10.4_{-0.6}^{+0.7} \pm 0.3$. The values of $\sigma(Z)$ and $\sigma(W)$ come from $O\left(\alpha_{s}^{2}\right)$ calculations using $m_{W}=80.14 \pm 0.27 \mathrm{GeV}$, and $m_{Z}=91.175 \pm 0.021 \mathrm{GeV}$ along with the corresponding value of $\sin ^{2} \theta_{W}=0.2274$. They use $\sigma(W) / \sigma(Z)=$ $3.26 \pm 0.07 \pm 0.05$ and $\Gamma(Z)=2.487 \pm 0.010 \mathrm{GeV}$.
38 ALBAJAR 91 measured $R=9.5_{-1.0}^{+1.1}$ (stat. + syst.). $\sigma(W) / \sigma(Z)$ is calculated in QCD at the parton level using $m_{W}=80.18 \pm 0.28 \mathrm{GeV}$ and $m_{Z}=91.172 \pm 0.031 \mathrm{GeV}$ along with $\sin ^{2} \theta_{W}=0.2322 \pm 0.0014$. They use $\sigma(W) / \sigma(Z)=3.23 \pm 0.05$ and $\Gamma(Z)$ $=2.498 \pm 0.020 \mathrm{GeV}$. This measurement is obtained combining both the electron and muon channels.
39 ABBIENDI 99C obtain this value from a fit to the reconstructed $W$ mass distribution using data at 172 and 183 GeV . The systematic error includes an uncertainty of $\pm 0.12 \mathrm{GeV}$ due to the possible color-reconnection and Bose-Einstein effects in the purely hadronic final states and an uncertainty of $\pm 0.01 \mathrm{GeV}$ due to the beam energy.
40 ABBOTT 99 H measure $R=10.90 \pm 0.52$ combining electron and muon channels. They use $M_{W}=80.39 \pm 0.06 \mathrm{GeV}$ and the SM theoretical predictions for $\sigma(W) / \sigma(Z), \mathrm{B}(Z \rightarrow$ $\ell \ell)$, and $\Gamma\left(W \rightarrow \ell \nu_{\ell}\right)$.
${ }^{41}$ ABREU 99T obtain this value using $W W \rightarrow \ell \bar{\nu}_{\ell} q \bar{q}$ and $W W \rightarrow q \bar{q} q \bar{q}$ events. The systematic error includes an uncertainty of $\pm 0.080 \mathrm{GeV}$ due to possible color reconnection and Bose-Einstein effects in the purely hadronic final state.

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42 BARATE 99 I obtain this result with a fit to the \(W W\) measured cross sections at 161 ,
    172 , and 183 GeV . The theoretical prediction takes into account the sensitivity to the
    \(W\) total width.
43 ACCIARRI 975 obtain this value from a fit to the reconstructed \(W\) mass distribution
\({ }^{44} \mathrm{ABE} 95 \mathrm{C}\) use the tail of the transverse mass distribution of \(W \rightarrow e \nu_{e}\) decays.
\({ }^{45}\) ALITTI 90C used the same technique as described for ABE 90. They measured \(R=\)
    \(9.38_{-0.72}^{+0.82} \pm 0.25\), obtained \(\Gamma(W) / \Gamma(Z)=0.902 \pm 0.074 \pm 0.024\). Using \(\Gamma(Z)=\)
    \(2.546 \pm 0.032 \mathrm{GeV}\), they obtained the \(\Gamma(W)\) value quoted above and the limits \(\Gamma(W)\)
    \(<2.56\) (2.64) GeV at the \(90 \%(95 \%) \mathrm{CL} . E_{\mathrm{cm}}^{p \bar{p}}=546,630 \mathrm{GeV}\).
46 ALBAJAR 89 result is from a total sample of \(299 \mathrm{~W} \rightarrow e \nu\) events.
47 If systematic error is neglected, result is \(2.7_{-1.5}^{+1.4} \mathrm{GeV}\). This is enhanced subsample of 172 total events.
```


## $W^{+}$DECAY MODES

$W^{-}$modes are charge conjugates of the modes below.

| Mode | Fraction $\left(\Gamma_{i} / \Gamma\right) \quad$ Con | Confidence level |
| :---: | :---: | :---: |
| $\Gamma_{1} \quad \ell^{+} \nu$ | [a] (10.68 $\pm 0.12) \%$ |  |
| $\Gamma_{2} e^{+} \nu$ | $(10.72 \pm 0.16) \%$ |  |
| $\Gamma 3 \mu^{+} \nu$ | $(10.57 \pm 0.22) \%$ |  |
| $\Gamma_{4} \tau^{+} \nu$ | $(10.74 \pm 0.27) \%$ |  |
| $\Gamma_{5}$ hadrons | (67.96 $\pm 0.35) \%$ |  |
| $\Gamma_{6} \pi^{+} \gamma$ | $<8 \times 10^{-5}$ | -5 95\% |
| $\Gamma_{7} \quad D_{s}^{+} \gamma$ | $<1.3 \times 10^{-3}$ | -3 95\% |
| $\Gamma_{8} c X$ | (33.6 $\pm 2.7$ ) \% |  |
| $\Gamma 9 \quad c \bar{s}$ | $\left(31 \begin{array}{l}+13 \\ -11\end{array}\right) \%$ |  |
| $\Gamma_{10}$ invisible | [b] ( $1.4 \pm 2.8) \%$ |  |

[a] $\ell$ indicates each type of lepton ( $e, \mu$, and $\tau$ ), not sum over them.
[b] This represents the width for the decay of the $W$ boson into a charged particle with momentum below detectability, $\mathrm{p}<200 \mathrm{MeV}$.

## W PARTIAL WIDTHS

This represents the width for the decay of the $W$ boson into a charged particle with momentum below detectability, $\mathrm{p}<200 \mathrm{MeV}$.

| VALUE ( MeV ) | DOCUMENT ID |  | TECN | COMMENT |
| :---: | :---: | :---: | :---: | :---: |
| $30_{-48}^{+52} \pm 33$ | 48 BARATE | 991 | ALEP | $E_{\mathrm{Cm}}^{e \mathrm{e}}=161+172+183$ |

-     - We do not use the following data for averages, fits, limits, etc.

49 BARATE 99 L ALEP $\quad E_{\mathrm{cm}}^{e e}=161+172+183$
GeV

48 BARATE 99I measure this quantity using the dependence of the total cross section $\sigma_{W} W$ upon a change in the total width. The fit is performed to the $W W$ measured cross sections at 161,172 , and 183 GeV . This partial width is $<139 \mathrm{MeV}$ at $95 \% \mathrm{CL}$.
49 BARATE 99L use $W$-pair production to search for effectively invisible $W$ decays, tagging with the decay of the other $W$ boson to Standard Model particles. The partial width for effectively invisible decay is $<27 \mathrm{MeV}$ at $95 \% \mathrm{CL}$.

## W BRANCHING RATIOS

Overall fits are performed to determine the branching ratios of the $W$. For each LEP experiment the correlation matrix of the leptonic branching ratios is used and the common systematic errors among LEP experiments are properly taken into account (see LEP Electroweak Working Group note LEPEWWG/XSEC/2001-02, 30 March 2001, accessible at http://lepewwg.web.cern.ch/LEPEWWG/lepww/4f/PDG01). A first fit determines three individual leptonic branching ratios, $\mathrm{B}\left(W \rightarrow e \nu_{e}\right)$, $\mathrm{B}\left(W \rightarrow \mu \nu_{\mu}\right)$, and $\mathrm{B}\left(W \rightarrow \tau \nu_{\tau}\right)$. This fit has a $\chi^{2}=11.0$ for 22 degrees of freedom. A second fit assumes lepton universality and determines the leptonic branching ratio $\mathrm{B}\left(W \rightarrow \ell \nu_{\ell}\right)$ and the hadronic branching ratio is derived as $\mathrm{B}(W \rightarrow$ hadrons $)=1-3 \mathrm{~B}(W \rightarrow \ell \nu)$. This fit has a $\chi^{2}=11.4$ for 24 degrees of freedom.

The LEP $W \rightarrow \ell \nu$ data are obtained by the Collaborations using individual leptonic channels and are, therefore, not included in the overall fits to avoid double counting.

## $\Gamma\left(\ell^{+} \nu\right) / \Gamma_{\text {total }}$

$\ell$ indicates average over $e, \mu$, and $\tau$ modes, not sum over modes.

| VALUE | EVTS | DOCUMENT ID |  | TECN | COMMENT |
| :---: | :---: | :---: | :---: | :---: | :---: |
| 0.1068 $\pm 0.0012$ OUR FIT |  |  |  |  |  |
| $0.1056 \pm 0.0020 \pm 0.0009$ | 5778 | ABBIENDI,G | 00 | OPAL | $\begin{gathered} E_{\mathrm{Cm}}^{e e}=161+172+183 \\ \quad+189 \mathrm{GeV} \end{gathered}$ |
| $0.1071 \pm 0.0024 \pm 0.0014$ | 4843 | ABREU | 00K | DLPH | $\begin{aligned} & E_{\mathrm{Cm}}^{\mathrm{ee}}=161+172+183 \\ & \quad+189 \mathrm{GeV} \end{aligned}$ |
| $0.1060 \pm 0.0023 \pm 0.0011$ | 5328 | ACCIARRI | 00V | L3 | $\begin{gathered} E_{\mathrm{Cm}}^{\mathrm{ee}}=161+172+183 \\ \quad+189 \mathrm{GeV} \end{gathered}$ |
| $0.1101 \pm 0.0022 \pm 0.0011$ | 5258 | BARATE | 00J | ALEP | $\begin{aligned} & E_{\mathrm{Cm}}^{\mathrm{ee}}=161+172+183 \\ & \quad \pm 189 \mathrm{GeV} \end{aligned}$ |
| $0.1102 \pm 0.0052$ | 11858 | 50 ABBOTT | 99 H | D0 | $E_{\mathrm{cm}}^{p \bar{p}}=1.8 \mathrm{TeV}$ |
| $0.104 \pm 0.008$ | 3642 | 51 ABE | 921 | CDF | $E_{\mathrm{cm}}^{p \bar{p}}=1.8 \mathrm{TeV}$ |

-     - We do not use the following data for averages, fits, limits, etc.

| $0.107 \pm 0.004 \pm 0.002$ | 1440 | ABBIENDI | 99D OPAL | Repl. by ABBI- <br> ENDI,G 00 |
| :--- | :--- | :--- | :--- | :--- |
| $0.1085 \pm 0.0048 \pm 0.0017$ | 1336 | ABREU | 99k DLPH | Repl. by ABREU 00K |
| $0.1036 \pm 0.0040 \pm 0.0017$ | 1322 | BARATE | 99। ALEP | Repl. by BARATE 00」 |
| $0.100 \pm 0.004 \pm 0.001$ | 1434 | ACCIARRI | 98P L3 | Repl. by ACCIA- |

${ }^{50}$ ABBOTT 99н measure $R \equiv\left[\sigma_{W} \mathrm{~B}\left(W \rightarrow \ell \nu_{\ell}\right)\right] /\left[\sigma_{Z} \mathrm{~B}(Z \rightarrow \ell \ell)\right]=10.90 \pm 0.52$ combining electron and muon channels. They use $M_{W}=80.39 \pm 0.06 \mathrm{GeV}$ and the SM theoretical predictions for $\sigma(W) / \sigma(Z)$ and $\mathrm{B}(Z \rightarrow \ell \ell)$.
$511216 \pm 38_{-31}^{+27} W \rightarrow \mu \nu$ events from $A B E 92$ and $2426 W \rightarrow e \nu$ events of ABE 91C. ABE 92I give the inverse quantity as $9.6 \pm 0.7$ and we have inverted.

| $\Gamma\left(e^{+} \nu\right) / \Gamma_{\text {total }}$VALUE | EVTS | DOCUMENT ID |  | TECN | COMMENT $\Gamma_{\mathbf{2} / \boldsymbol{~}}$ |
| :---: | :---: | :---: | :---: | :---: | :---: |
|  |  |  |  |  |  |
| $\mathbf{0 . 1 0 7 2} \pm \mathbf{0 . 0 0 1 6}$ OUR FIT |  |  |  |  |  |
| $0.1046 \pm 0.0042 \pm 0.0014$ | 801 | ABBIENDI,G | 00 | OPAL | $\begin{gathered} E_{\mathrm{Cm}}^{e e}=161+172+183 \\ \pm 189 \mathrm{GeV} \end{gathered}$ |
| $0.1044 \pm 0.0015 \pm 0.0028$ | 67318 | 52 ABBOTT | 00B | D0 | $E_{\mathrm{cm}}^{p \bar{p}}=1.8 \mathrm{TeV}$ |
| $0.1018 \pm 0.0054 \pm 0.0026$ | 527 | ABREU | 00K | DLPH | $\begin{gathered} E_{\mathrm{cm}}^{e e}=161+172+183 \\ \quad+189 \mathrm{GeV} \end{gathered}$ |
| $0.1077 \pm 0.0045 \pm 0.0016$ | 715 | ACCIARRI | OOV | L3 | $\begin{aligned} & E_{\mathrm{Cm}}^{\mathrm{ee}}=161+172+183 \\ & \quad+189 \mathrm{GeV} \end{aligned}$ |
| $0.1135 \pm 0.0046 \pm 0.0017$ | 720 | BARATE | 00J | ALEP | $\begin{gathered} E_{\mathrm{Cm}}^{e \mathrm{e}}=161+172+183 \\ +189 \mathrm{GeV} \end{gathered}$ |
| $0.1094 \pm 0.0033 \pm 0.0031$ |  | 53 ABE | 95w | CDF | $E_{\mathrm{cm}}^{p \bar{p}}=1.8 \mathrm{TeV}$ |
| $0.10 \pm 0.014{ }_{-0.03}^{+0.02}$ | 248 | 54 ANSARI | 87C | UA2 | $E_{\mathrm{cm}}^{p \bar{p}}=546,630 \mathrm{GeV}$ |

-     - We do not use the following data for averages, fits, limits, etc.

| $0.117 \pm 0.009 \pm 0.002$ | 224 | ABBIENDI | 99D OPAL | Repl. by ABBI- <br> ENDI,G 00 |
| :---: | :---: | :---: | :---: | :---: |
| $0.1012 \pm 0.0107 \pm 0.0028$ | 150 | ABREU | 99k DLPH | Repl. by ABREU 00k |
| $0.1115 \pm 0.0085 \pm 0.0024$ | 192 | BARATE | 991 ALEP | Repl. by BARATE 00J |
| $0.105 \pm 0.009 \pm 0.002$ | 173 | ACCIARRI | 98P L3 | Repl. by ACCIARRI 00V |
| seen | 119 | APPEL | 86 UA2 | $E_{\mathrm{cm}}^{p P}=546,630 \mathrm{GeV}$ |
| seen | 172 | ARNISON | 86 UA1 | $E_{\mathrm{cm}}^{p \bar{p}}=546,630 \mathrm{GeV}$ |

52 ABBOTT 00B measure $R \equiv\left[\sigma_{W} \mathrm{~B}\left(W \rightarrow e \nu_{e}\right)\right] /\left[\sigma_{Z} \mathrm{~B}(Z \rightarrow e e)\right]=10.43 \pm 0.27$ for the $W \rightarrow e \nu_{e}$ decay channel. They use the SM theoretical prediction for $\sigma(W) / \sigma(Z)$ and the world average for $\mathrm{B}(Z \rightarrow e e)$.
53 ABE 95 W result is from a measurement of $\sigma \mathrm{B}(W \rightarrow e \nu) / \sigma \mathrm{B}\left(Z \rightarrow e^{+} e^{-}\right)=$ $10.90 \pm 0.32 \pm 0.29$, the theoretical prediction for the cross section ratio, the experimental knowledge of $\Gamma\left(Z \rightarrow e^{+} e^{-}\right)=83.98 \pm 0.18 \mathrm{MeV}$, and $\Gamma(Z)=2.4969 \pm 0.0038$ GeV .
54 The first error was obtained by adding the statistical and systematic experimental uncertainties in quadrature. The second error reflects the dependence on theoretical prediction of total $W$ cross section: $\sigma(546 \mathrm{GeV})=4.7_{-0.7}^{+1.4} \mathrm{nb}$ and $\sigma(630 \mathrm{GeV})=5.8_{-1.0}^{+1.8} \mathrm{nb}$. See ALTARELLI 85B.

| $\Gamma\left(\mu^{+} \nu\right) / \Gamma_{\text {total }}$VALUE | EVTS | DOCUMENT ID |  | TECN | COMMENT |
| :---: | :---: | :---: | :---: | :---: | :---: |
|  |  |  |  |  |  |
| $\mathbf{0 . 1 0 5 7} \pm 0.0022$ OUR FIT |  |  |  |  |  |
| $0.1050 \pm 0.0041 \pm 0.0012$ | 803 | ABBIENDI,G | 00 | OPAL | $\begin{gathered} E_{\mathrm{Cm}}^{e e}=161+172+183 \\ +189 \mathrm{GeV} \end{gathered}$ |
| $0.1092 \pm 0.0048 \pm 0.0012$ | 649 | ABREU | 00K | DLPH | $\begin{gathered} E_{\mathrm{Cm}}^{\mathrm{ee}}=161+172+183 \\ \quad+189 \mathrm{GeV} \end{gathered}$ |
| $0.0990 \pm 0.0046 \pm 0.0015$ | 617 | ACCIARRI | 00V | L3 | $\begin{aligned} & E_{\mathrm{Cm}}^{e e}=161+172+183 \\ & \quad+189 \mathrm{GeV} \end{aligned}$ |
| $0.1110 \pm 0.0044 \pm 0.0016$ | 710 | BARATE | 00J | ALEP | $\begin{gathered} E_{\mathrm{Cm}}^{\mathrm{ee}}=161+172+183 \\ \pm 189 \mathrm{GeV} \end{gathered}$ |
| $0.10 \pm 0.01$ | 1216 | ABE | 921 | CDF | $E_{\mathrm{cm}}^{p \bar{p}}=1.8 \mathrm{TeV}$ |

-     - We do not use the following data for averages, fits, limits, etc.

| $0.102 \pm 0.008 \pm 0.002$ | 193 | ABBIENDI | 99D OPAL | Repl. by ABBI- |
| :--- | :--- | :--- | :--- | :--- |
|  |  |  | ENDI,G 00 |  |
| $0.1139 \pm 0.0096 \pm 0.0023$ | 186 | ABREU | 99k DLPH | Repl. by ABREU 00k |
| $0.1006 \pm 0.0078 \pm 0.0021$ | 179 | BARATE | 99। ALEP | RepI. by BARATE 00J |
| $0.102 \pm 0.009 \pm 0.002$ | 160 | ACCIARRI | 98P L3 | Repl. by ACCIA- |
|  |  |  |  | RRI 00V |

${ }^{55} \mathrm{ABE}$ 92। quote the inverse quantity as $9.9 \pm 1.2$ which we have inverted.


-     - We do not use the following data for averages, fits, limits, etc.

| $0.101 \pm 0.010 \pm 0.003$ | 183 | ABBIENDI | 99D OPAL | Repl. by ABBI- <br> ENDI,G 00 |
| :--- | :--- | :--- | :--- | :--- |
| $0.1095 \pm 0.0149 \pm 0.0041$ | 142 | ABREU | 99k DLPH | Repl. by ABREU 00K |
| $0.0976 \pm 0.0101 \pm 0.0033$ | 160 | BARATE | 99। ALEP | Repl. by BARATE 00J |
| $0.090 \pm 0.012 \pm 0.003$ | 123 | ACCIARRI | 98P L3 | Repl. by ACCIA- |

$\Gamma($ hadrons $) / \Gamma_{\text {total }}$
OUR FIT value is obtained by a fit to the lepton branching ratio data assuming lepton universality.

| VALUE | EVTS | DOCUMENT ID |  | TECN | COMMENT |
| :---: | :---: | :---: | :---: | :---: | :---: |
| $0.6796 \pm 0.0035$ OUR FIT |  |  |  |  |  |
| $0.679 \pm 0.004$ OUR AVERAGE |  |  |  |  |  |
| $0.6832 \pm 0.0061 \pm 0.0028$ | 5778 | ABBIENDI,G | 00 |  | OPAL | $\begin{gathered} E_{\mathrm{Cm}}^{e e}=161+172+183 \\ +189 \mathrm{GeV} \end{gathered}$ |
| $0.6789 \pm 0.0073 \pm 0.0043$ | 4843 | ABREU | 00k | DLPH | $\begin{aligned} & E_{\mathrm{Cm}}^{\mathrm{ee}}=161+172+183 \\ & \quad+189 \mathrm{GeV} \end{aligned}$ |
| $0.6820 \pm 0.0068 \pm 0.0033$ | 5328 | ACCIARRI | 00V | L3 | $\begin{gathered} E_{\mathrm{Cm}}^{e e}=161+172+183 \\ \quad+189 \mathrm{GeV} \end{gathered}$ |
| $0.6697 \pm 0.0065 \pm 0.0032$ | 5258 | BARATE | 00J | ALEP | $\begin{gathered} E_{\mathrm{Cm}}^{e e}=161+172+183 \\ +189 \mathrm{GeV} \end{gathered}$ |

-     - We do not use the following data for averages, fits, limits, etc.
$0.679 \pm 0.012 \pm 0.0051440$
$0.6746 \pm 0.0143 \pm 0.0052 \quad 1336$
$0.6893 \pm 0.0121 \pm 0.0051 \quad 1322$
$0.701 \pm 0.013 \pm 0.0041434$

ABBIENDI

ABREU
BARATE
ACCIARRI

99D OPAL Repl. by ABBIENDI,G 00 99k DLPH Repl. by ABREU 00k 99I ALEP Repl. by BARATE 00J 98P L3 Repl. by ACCIARRI 00v

| $\underset{\text { VALUE }}{\Gamma\left(\mu^{+} \nu\right) / \Gamma\left(e^{+} \nu\right)}$ | EVTS DOCUMENT ID |  | TECN | COMMENT ${ }^{3 / 1 / 2}$ |
| :---: | :---: | :---: | :---: | :---: |
|  |  |  |  |  |
| 0.986 $\pm 0.024$ OUR FIT |  |  |  |  |
| $0.89 \pm 0.10$ | 13k | 56 ABACHI | 95D D0 | $E_{\mathrm{cm}}^{p \bar{p}}=1.8 \mathrm{TeV}$ |
| $1.02 \pm 0.08$ | 1216 | 57 ABE | 92I CDF | $E_{\mathrm{cm}}^{p \bar{p}}=1.8 \mathrm{TeV}$ |
| $1.00 \pm 0.14 \pm 0.08$ | 67 | ALBAJAR | 89 UA1 | $E_{\mathrm{Cm}}^{p \bar{p}}=546,630 \mathrm{GeV}$ |

-     - We do not use the following data for averages, fits, limits, etc.
$1.24 \begin{gathered}+0.4 \\ +0.6\end{gathered} 14$ ARNISON 84D UA1 Repl. by ALBAJAR 89
${ }^{56} \mathrm{ABACHI} 95 \mathrm{D}$ obtain this result from the measured $\sigma_{W} \mathrm{~B}(W \rightarrow \mu \nu)=2.09 \pm 0.23 \pm$ 0.11 nb and $\sigma_{W} \mathrm{~B}(W \rightarrow e \nu)=2.36 \pm 0.07 \pm 0.13 \mathrm{nb}$ in which the first error is the combined statistical and systematic uncertainty, the second reflects the uncertainty in the luminosity.
${ }^{57} \mathrm{ABE} 92$ I obtain $\sigma_{W} \mathrm{~B}(W \rightarrow \mu \nu)=2.21 \pm 0.07 \pm 0.21$ and combine with ABE 91C $\sigma_{W}$ $\mathrm{B}((W \rightarrow e \nu))$ to give a ratio of the couplings from which we derive this measurement.
$\Gamma\left(\tau^{+} \nu\right) / \Gamma\left(e^{+} \nu\right)$
$\begin{array}{ll} & \Gamma_{4} / \Gamma_{2}\end{array}$
$\frac{\text { VALUE }}{1.002 \pm 0.029 \text { OUR FIT }}$
$0.961 \pm 0.061$
980
58 ABBOTT 00D D0 $E_{\mathrm{Cm}}^{p \bar{p}}=1.8 \mathrm{TeV}$
$0.94 \pm 0.14 \quad 179$
$59 \mathrm{ABE} \quad 92 \mathrm{E}$ CDF $\quad E_{\mathrm{cm}}^{p \bar{p}}=1.8 \mathrm{TeV}$
60 ALITTI 92F UA2 $E_{\mathrm{cm}}^{p \bar{p}}=630 \mathrm{GeV}$
$1.04 \pm 0.08 \pm 0.08 \quad 754$
ALBAJAR
89 UA1
$E_{\mathrm{cm}}^{p \bar{p}}=546,630 \mathrm{GeV}$
$1.02 \pm 0.20 \pm 0.12 \quad 32$
fits, limits, etc.

| $0.995 \pm 0.112 \pm 0.083$ | 198 | ALITTI | 91C UA2 | Repl. by ALITTI 92F |
| :--- | ---: | :--- | :--- | :--- |
| $1.02 \pm 0.20 \pm 0.10$ | 32 | ALBAJAR | 87 UA1 | Repl. by ALBAJAR 89 |

58 ABBOTT 00D measure $\sigma_{W} \times \mathrm{B}\left(W \rightarrow \tau \nu_{\tau}\right)=2.22 \pm 0.09 \pm 0.10 \pm 0.10 \mathrm{nb}$. Using the ABBOTT 00B result $\sigma_{W} \times \mathrm{B}\left(W \rightarrow e \nu_{e}\right)=2.31 \pm 0.01 \pm 0.05 \pm 0.10 \mathrm{nb}$, they quote the ratio of the couplings from which we derive this measurement.
59 ABE 92E use two procedures for selecting $W \rightarrow \tau \nu_{\tau}$ events. The missing $\mathrm{E}_{T}$ trigger leads to $132 \pm 14 \pm 8$ events and the $\tau$ trigger to $47 \pm 9 \pm 4$ events. Proper statistical and systematic correlations are taken into account to arrive at $\sigma \mathrm{B}(W \rightarrow \tau \nu)=2.05 \pm 0.27$ nb. Combined with ABE 91 C result on $\sigma \mathrm{B}(W \rightarrow e \nu)$, ABE 92 E quote a ratio of the couplings from which we derive this measurement.
60 This measurement is derived by us from the ratio of the couplings of ALITTI 92F.
$\Gamma\left(\pi^{+} \gamma\right) / \Gamma\left(e^{+} \nu\right) \quad \Gamma_{6} / \Gamma_{2}$

| VALUE | CL\% | DOCUMENT | TECN | COMMENT |
| :---: | :---: | :---: | :---: | :---: |
| $<7 \times 10^{-4}$ | 95 | ABE | 98H CDF | $E_{\mathrm{cm}}^{p \bar{p}}=1.8 \mathrm{TeV}$ |
| $<4.9 \times 10^{-3}$ | 95 | 61 ALITTI | 92D UA2 | $E_{\mathrm{cm}}^{p \bar{p}}=630 \mathrm{GeV}$ |
| $<58 \times 10^{-3}$ | 95 | 62 ALBAJAR | 90 UA1 | $E_{\mathrm{cm}}^{p \bar{P}}=546,630 \mathrm{GeV}$ |

61 ALITTI 92D limit is $3.8 \times 10^{-3}$ at $90 \%$ CL.
62 ALBAJAR 90 obtain $<0.048$ at $90 \%$ CL.


## $0.49 \pm 0.04$ OUR AVERAGE

$0.481 \pm 0.042 \pm 0.0323005$
$0.51 \pm 0.05 \pm 0.03 \quad 746$

| DOCUMENT ID | TECN | COMMENT |
| :---: | :---: | :---: |
| 63 ABBIENDI | 00V OPAL | $E_{\mathrm{Cm}}^{e \mathrm{e}}=183+189 \mathrm{GeV}$ |
| 64 BARATE | 99m ALEP | $E_{\mathrm{Cm}}^{e \mathrm{e}}=172+183 \mathrm{GeV}$ |

${ }^{63}$ ABBIENDI $00 v$ tag $W \rightarrow c X$ decays using measured jet properties, lifetime information, and leptons produced in charm decays. From this result, and using the additional measurements of $\Gamma(W)$ and $\mathrm{B}(W \rightarrow$ hadrons $),\left|V_{c s}\right|$ is determined to be $0.969 \pm 0.045 \pm 0.036$.
${ }^{64}$ BARATE 99M tag $c$ jets using a neural network algorithm. From this measurement $\left|V_{c s}\right|$ is determined to be $1.00 \pm 0.11 \pm 0.07$.
$R_{c s}=\Gamma(c \bar{s}) / \Gamma$ (hadrons)
VALUE
$0.46+0.18 \pm 0.07$
$\Gamma_{9} / \Gamma_{5}$
65 ABREU 98N DLPH $E_{\mathrm{cm}}^{e e}=161+172 \mathrm{GeV}$
${ }^{65}$ ABREU 98 N tag $c$ and $s$ jets by identifying a charged kaon as the highest momentum particle in a hadronic jet. They also use a lifetime tag to independently identify a c jet, based on the impact parameter distribution of charged particles in a jet. From this measurement $\left|V_{c s}\right|$ is determined to be $0.94_{-0.26}^{+0.32} \pm 0.13$.

## AVERAGE PARTICLE MULTIPLICITIES IN HADRONIC W DECAY

Summed over particle and antiparticle, when appropriate.
$\left\langle N_{\pi^{ \pm}}\right\rangle$
$\qquad$
$15.70 \pm 0.35$
$66 \frac{\text { DOCUMENT ID }}{\text { ABREU,P }}$ OOF $\frac{\text { TECN }}{\text { DLPH }} \frac{\text { COMMENT }}{E_{\mathrm{cm}}^{e e}=189 \mathrm{GeV}}$
${ }^{66}$ ABREU,P 00F measure $\left\langle N_{\pi^{ \pm}}\right\rangle=31.65 \pm 0.48 \pm 0.76$ and $15.51 \pm 0.38 \pm 0.40$ in the fully hadronic and semileptonic final states respectively. The value quoted is a weighted average without assuming any correlations.
$\left\langle\boldsymbol{N}_{\boldsymbol{K}^{ \pm}}\right\rangle$
$\frac{\text { VALUE }}{\mathbf{2 . 2 0} \mathbf{0 . 1 9}} \quad 67 \frac{\text { DOCUMENT ID }}{\text { ABREU,P }} \frac{\text { TECN }}{} \frac{\text { COMMENT }}{\text { DLPH }} \frac{\text { COM }}{E_{\mathrm{cm}}^{e e}=189 \mathrm{GeV}}$
${ }^{67}$ ABREU,P 00F measure $\left\langle N_{K^{ \pm}}\right\rangle=4.38 \pm 0.42 \pm 0.12$ and $2.23 \pm 0.32 \pm 0.17$ in the fully hadronic and semileptonic final states respectively. The value quoted is a weighted average without assuming any correlations.

## $\left\langle N_{p}\right\rangle$

VALUE
$0.92 \pm 0.14$
$68 \frac{\text { DOCUMENT ID }}{\text { ABREU,P }} \frac{\text { O0F }}{} \frac{\text { TECN }}{\text { DLPH }} \frac{\text { COMMENT }}{E_{\mathrm{Cm}}^{e e}=189 \mathrm{GeV}}$
${ }^{68} \mathrm{ABREU}, \mathrm{P} 00 \mathrm{~F}$ measure $\left\langle N_{p}\right\rangle=1.82 \pm 0.29 \pm 0.16$ and $0.94 \pm 0.23 \pm 0.06$ in the fully hadronic and semileptonic final states respectively. The value quoted is a weighted average without assuming any correlations.
$\frac{V A L U E}{19.41 \pm 0.15 \text { OUR AVERAGE }}$
$19.44 \pm 0.17$
$19.3 \pm 0.3 \pm 0.3$
$19.23 \pm 0.74$

| DOCUMENT | TECN | COMMENT |
| :---: | :---: | :---: |
| 69 ABREU,P | 00F DLPH | $E_{\text {cm }}^{e e}{ }_{\text {en }}=183+189 \mathrm{GeV}$ |
| 70 ABBIENDI | 99N OPAL | $E_{\mathrm{cm}}^{e e}=183 \mathrm{GeV}$ |
| 71 ABREU | 98C DLPH | $E_{\mathrm{cm}}^{e} \mathrm{e}=172 \mathrm{GeV}$ |

69 ABREU,P 00F measure $\left\langle N_{\text {charged }}\right\rangle=39.12 \pm 0.33 \pm 0.36$ and $38.11 \pm 0.57 \pm 0.44$ in the fully hadronic final states at 189 and 183 GeV respectively, and $\left\langle N_{\text {charged }}\right\rangle=$ $19.49 \pm 0.31 \pm 0.27$ and $19.78 \pm 0.49 \pm 0.43$ in the semileptonic final states. The value quoted is a weighted average without assuming any correlations.
${ }^{70}$ ABBIENDI 99N use the final states $W^{+} W^{-} \rightarrow q \bar{q} \ell \bar{\nu}_{\ell}$ to derive this value.
${ }^{71}$ ABREU 98C combine results from both the fully hadronic as well semileptonic $W W$ final states after demonstrating that the $W$ decay charged multiplicity is independent of the topology within errors.

## TRIPLE GAUGE COUPLINGS (TGC'S)

Revised February 2002 by C. Caso (University of Genova) and A. Gurtu (Tata Institute).

Fourteen independent couplings, 7 each for $Z W W$ and $\gamma W W$, completely describe the $V W W$ vertices within the most general framework of the electroweak Standard Model (SM) consistent with Lorentz invariance and $\mathrm{U}(1)$ gauge invariance. Of each of the 7 TGC's, 3 conserve $C$ and $P$ individually, 3 violate $C P$, and one TGC violates $C$ and $P$ individually while conserving $C P$. Assumption of $C$ and $P$ conservation and electromagnetic gauge invariance reduces the independent $V W W$ couplings to five: one common set $[1,2]$ is $\left(\Delta \kappa_{\gamma}, \Delta \kappa_{Z}, \lambda_{\gamma}, \lambda_{Z}, \Delta g_{1}^{Z}\right)$, where $\Delta \kappa_{\gamma}=\Delta \kappa_{Z}=\Delta g_{1}^{Z}=0$ and $\lambda_{\gamma}=\lambda_{Z}=0$ in the Standard Model at the tree level. The $W$ magnetic dipole moment, $\mu_{W}$, and the $W$ electric quadrupole moment, $q_{W}$, are expressed as $\mu_{W}=e\left(1+\kappa_{\gamma}+\lambda_{\gamma}\right) / 2 M_{W}$ and $q_{W}=-e\left(\kappa_{\gamma}-\lambda_{\gamma}\right) / M_{W}^{2}$.

Precision measurements of suitable observables at LEP1 has already led to an exploration of much of the TGC parameter space. For LEP2 data, the LEP Collaborations have agreed to express their results in terms of the parameters $\Delta g_{1}^{Z}, \Delta \kappa_{\gamma}$ and $\lambda_{\gamma}$ ( $\lambda_{Z}$ and $\Delta \kappa_{Z}$ are related to these by gauge invariance).

At LEP2 the $V W W$ coupling arises in $W$-pair production via $s$-channel exchange or in single $W$ production via the radiation of a virtual photon off the incident $e^{+}$or $e^{-}$. At the TEVATRON hard photon bremstrahlung off a produced $W$ or $Z$ signals the presence of a triple gauge vertex. In order to extract the value of one TGC the others are generally kept fixed to their SM values.

## References

1. K. Hagiwara et al., Nucl. Phys. B282, 253 (1987).
2. G. Gounaris et al., CERN 96-01 525.

## $\Delta g_{1}^{z}$

Combining published and unpublished LEP results (as of Summer 2003), a singleparameter fit yields $\Delta g_{1}^{Z}=-0.009_{-0.021}^{+0.022}$, where the other two parameters, $\Delta \kappa_{\gamma}$ and $\lambda_{\gamma}$, were kept fixed to their Standard Model values.
(See EP Preprint Summer 2003: CERN-EP/2003-091 and hep-ex/0312023, December 2003, on http://lepewwg.web.cern.ch/LEPEWWG/stanmod/)

| VALUE | EVTS | DOCUMENT II | TECN | COMMENT |
| :---: | :---: | :---: | :---: | :---: |
| $\begin{array}{r} -0.013+0.034 \\ -0.033 \end{array}$ | 9800 | 72 ABBIENDI | 04D OPAL | $E_{\mathrm{cm}}^{e \mathrm{e}}$ = $=183-209 \mathrm{GeV}$ |

-     - We do not use the following data for averages, fits, limits, etc.

| $-0.02 \pm 0.07 \pm 0.01$ | 2114 | 73 ABREU | 01I DLPH | $E_{\mathrm{cm}}^{e e}=183+189 \mathrm{GeV}$ |
| :---: | :---: | :---: | :---: | :---: |
| $\begin{array}{r} 0.023+0.059 \\ -0.055 \end{array}$ | 3586 | 74 HEISTER | 01C ALEP | $E_{\mathrm{Cm}}^{e \mathrm{e}}=161-189 \mathrm{GeV}$ |
|  | 331 | 75 ABBOTT | 991 D0 | $E_{\mathrm{cm}}^{p \bar{p}}=1.8 \mathrm{TeV}$ |
| $\mathrm{O}^{(11}{ }_{-0.18}^{+0.19} \pm 0.10$ | 1154 | 76 ACCIARRI | 99Q L3 | $E_{\mathrm{cm}}^{e e}=161+172+183$ |

72 ABBIENDI 04D combine results from $W^{+} W^{-}$in all decay channels. Only CP-conserving couplings are considered and each parameter is determined from a single-parameter fit in which the other parameters assume their Standard Model values. The $95 \%$ confidence interval is $-0.077<\Delta g_{1}^{Z}<0.054$.
73 ABREU 01I combine results from $e^{+} e^{-}$interactions at 189 GeV leading to $W^{+} W^{-}$ and $W e \nu_{e}$ final states with results from ABREU 99L at 183 GeV . The $95 \%$ confidence interval is $-0.16<\Delta g_{1}^{Z}<0.13$.
${ }^{74}$ HEISTER 01C study $W$-pair, single- $W$, and single photon events and combine with earlier results from BARATE,R 98, BARATE 98Y, and BARATE 99L to obtain the quoted value, fixing $\Delta \kappa_{\gamma}$ and $\lambda_{\gamma}$ to their Standard Model values. The $95 \%$ confidence interval is $-0.087<\Delta g_{1}^{Z}<0.141$. When all three couplings $\Delta g_{1}^{Z}, \Delta \kappa_{\gamma}$, and $\lambda_{\gamma}$ are floated freely in the fit, one obtains $\Delta g_{1}^{Z}=0.013_{-0.068}^{+0.066}$.
${ }^{75}$ ABBOTT 99ı perform a simultaneous fit to the $W \gamma, W W \rightarrow$ dilepton, $W W / W Z \rightarrow$ $e \nu j j, W W / W Z \rightarrow \mu \nu j j$, and $W Z \rightarrow$ trilepton data samples. For $\Lambda=2.0 \mathrm{TeV}$, the $95 \%$ CL limits are $-0.37<\Delta g_{1}^{Z}<0.57$, fixing $\lambda_{Z}=\Delta \kappa_{Z}=0$ and assuming Standard Model values for the $W W \gamma$ couplings.
76 ACCIARRI 99Q study $W$-pair, single- $W$, and single photon events.
$\Delta \kappa_{\gamma}$
Combining published and unpublished LEP results (as of Summer 2003), a singleparameter fit yields $\Delta \kappa_{\gamma}=-0.016_{-0.047}^{+0.042}$, where the other two parameters, $\Delta g_{1}^{Z}$ and $\lambda_{\gamma}$, were kept fixed to their Standard Model values.
(See EP Preprint Summer 2003: CERN-EP/2003-091 and hep-ex/0312023, December 2003, on http://lepewwg.web.cern.ch/LEPEWWG/stanmod/)

| VALUE | EVTS | DOCUMENT | TECN | COMMENT |
| :---: | :---: | :---: | :---: | :---: |
| $-0.12+0.09$ | 9800 | 77 ABBIENDI | 04D OPAL | $E_{\mathrm{Cm}}^{e \mathrm{e}}=183-209 \mathrm{GeV}$ |

-     - We do not use the following data for averages, fits, limits, etc. - -

| $0.116_{-0.086}^{+0.082} \pm 0.068$ | 315 | 78 ACHARD |  | L3 | $E_{\mathrm{Cm}}^{e \mathrm{e}}=161-209 \mathrm{GeV}$ |
| :---: | :---: | :---: | :---: | :---: | :---: |
| $0.25{ }_{-0.20}^{+0.21} \pm 0.06$ | 2298 | 79 ABREU | 011 | DLPH | $E_{\mathrm{Cm}}^{e e}=183+189 \mathrm{GeV}$ |
| $0.022+0.119$ | 3586 | 80 HEISTER | 01C | ALEP | $E_{\mathrm{cm}}^{e e}=161-189 \mathrm{GeV}$ |
|  |  | 81 BREITWEG | 00 | ZEUS | $\begin{gathered} e^{+} p \rightarrow e^{+} W^{ \pm} X \\ \sqrt{s} \approx 300 \mathrm{GeV} \end{gathered}$ |
| $-0.08 \pm 0.34$ | 331 | 82 ABBOTT | 991 | D0 | $E_{\mathrm{cm}}^{p \bar{p}}=1.8 \mathrm{TeV}$ |
| $0.11 \pm 0.25 \pm 0.17$ | 1154 | 83 ACCIARRI | 99Q | L3 | $E_{\mathrm{cm}}^{e \mathrm{e}}{ }^{e}=161+172+183$ |

77 ABBIENDI 04D combine results from $W^{+} W^{-}$in all decay channels. Only CP-conserving couplings are considered and each parameter is determined from a single-parameter fit in which the other parameters assume their Standard Model values. The 95\% confidence interval is $-0.27<\Delta \kappa_{\gamma}<0.07$.
78 ACHARD 02l study single $W$ production in $e^{+} e^{-}$interactions from 192 to 209 GeV . The result quoted here is obtained including data from 161 to 189 GeV , ACCIARRI 00N. The $95 \%$ C.L. limits are $-0.10<\Delta \kappa_{\gamma}<0.32$ (for $\lambda_{\gamma}=0$ ). When both couplings $\lambda_{\gamma}$ and $\kappa_{\gamma}$ are floated freely in the fit one obtains $\Delta \kappa_{\gamma}=0.07 \pm 0.10 \pm 0.07$.
${ }^{79}$ ABREU 01 l combine results from $e^{+} e^{-}$interactions at 189 GeV leading to $W^{+} W^{-}$, $W e \nu_{e}$, and $\nu \bar{\nu} \gamma$ final states with results from ABREU 99L at 183 GeV . The $95 \%$ confidence interval is $-0.13<\Delta \kappa_{\gamma}<0.68$.
80 HEISTER 01C study $W$-pair, single- $W$, and single photon events and combine with earlier results from BARATE,R 98, BARATE 98Y, and BARATE 99L to obtain the quoted value, fixing $\Delta g_{1}^{Z}$ and $\lambda_{\gamma}$ to their Standard Model values. The $95 \%$ confidence interval is $-0.200<\Delta \kappa_{\gamma}<0.258$. When all three couplings $\Delta g_{1}^{Z}, \Delta \kappa_{\gamma}$, and $\lambda_{\gamma}$ are floated freely in the fit, one obtains $\Delta \kappa_{\gamma}=0.043 \pm 0.110$.
81 BREITWEG 00 search for $W$ production in events with large hadronic $p_{T}$. For $p_{T}>20$ GeV , the upper limit on the cross section gives the $95 \% \mathrm{CL}$ limit $-4.7<\Delta \kappa_{\gamma}<1.5$ (for $\lambda_{\gamma}=0$ ).
${ }^{82}$ ABBOTT 99I perform a simultaneous fit to the $W \gamma, W W \rightarrow$ dilepton, $W W / W Z \rightarrow$ $e \nu j j, W W / W Z \rightarrow \mu \nu j j$, and $W Z \rightarrow$ trilepton data samples. For $\Lambda=2.0 \mathrm{TeV}$, the $95 \%$ CL limits are $-0.25<\Delta \kappa_{\gamma}<0.39$.
83 ACCIARRI 99Q study $W$-pair, single- $W$, and single photon events.
$\lambda_{\gamma}$
Combining published and unpublished LEP results (as of Summer 2003), a singleparameter fit yields $\lambda_{\gamma}=-0.016_{-0.023}^{+0.021}$, where the other two parameters, $\Delta g_{1}^{Z}$ and $\Delta \kappa_{\gamma}$, were kept fixed to their Standard Model values.
(See EP Preprint Summer 2003: CERN-EP/2003-091 and hep-ex/0312023, December 2003, on http://lepewwg.web.cern.ch/LEPEWWG/stanmod/)

| VALUE | $\frac{\text { EVTS }}{\text { DOCUMENT ID }}$ | $\frac{\text { TECN }}{\text { COMMENT }}$ |  |  |
| :--- | :--- | :--- | :--- | :--- |
| $\mathbf{- 0 . 0 6 0} \mathbf{+ 0 . 0 3 4}$ | 9800 | 84 ABBIENDI | 04D OPAL | $E_{\mathrm{cm}}^{e-0}=183-209 \mathrm{GeV}$ |

-     - We do not use the following data for averages, fits, limits, etc. - -

| $0.35{ }_{-0.13}^{+0.10} \pm 0.08$ | 315 | 85 ACHARD | 02I L3 | $E_{\mathrm{Cm}}^{e e}=161-209 \mathrm{GeV}$ |
| :---: | :---: | :---: | :---: | :---: |
| $0.05 \pm 0.09 \pm 0.01$ | 2298 | 86 ABREU | 01I DLPH | $E_{\mathrm{cm}}^{e e}=183+189 \mathrm{GeV}$ |
| $\begin{gathered} 0.040_{-0.052}^{+0.054} \end{gathered}$ | 3586 | 87 HEISTER | 01C ALEP | $E_{\mathrm{Cm}}^{e e}=161-189 \mathrm{GeV}$ |
|  |  | 88 BREITWEG | 00 ZEUS | $\begin{gathered} e^{+} p \rightarrow e^{+} W^{ \pm} X \\ \sqrt{s} \approx 300 \mathrm{GeV} \end{gathered}$ |
| $0.00 \begin{gathered} +0.10 \\ -0.09 \end{gathered}$ | 331 | 89 ABBOTT | 991 D0 | $E_{\mathrm{cm}}^{p \bar{p}}=1.8 \mathrm{TeV}$ |
| $0.10{ }_{-0.20}^{+0.22} \pm 0.10$ | 1154 | 90 ACCIARRI | 99Q L3 | $\begin{aligned} & E_{\mathrm{Cm}}^{e \mathrm{e}}=161+172+183 \\ & \mathrm{GeV} \end{aligned}$ |

84 ABBIENDI 04D combine results from $W^{+} W^{-}$in all decay channels. Only CP-conserving couplings are considered and each parameter is determined from a single-parameter fit in which the other parameters assume their Standard Model values. The 95\% confidence interval is $-0.13<\lambda_{\gamma}<0.01$.
85 ACHARD 02l study single $W$ production in $e^{+} e^{-}$interactions from 192 to 209 GeV . The result quoted here is obtained including data from 161 to 189 GeV , ACCIARRI 00N. The $95 \%$ C.L. limits are $-0.37<\lambda_{\gamma}<0.61$ (for $\kappa_{\gamma}=1$ ). When both couplings $\lambda_{\gamma}$ and $\kappa_{\gamma}$ are floated freely in the fit one obtains $\lambda_{\gamma}=0.31_{-0.20}^{+0.12} \pm 0.07$.
86 ABREU 01 combine results from $e^{+} e^{-}$interactions at 189 GeV leading to $W^{+} W^{-}$, $W e \nu_{e}$, and $\nu \bar{\nu} \gamma$ final states with results from ABREU 99L at 183 GeV . The $95 \%$ confidence interval is $-0.11<\lambda_{\gamma}<0.23$.
87 HEISTER 01C study $W$-pair, single- $W$, and single photon events and combine with earlier results from BARATE,R 98, BARATE 98Y, and BARATE 99L to obtain the quoted value, fixing $\Delta g_{1}^{Z}$ and $\Delta \kappa_{\gamma}$ to their Standard Model values. The $95 \%$ confidence interval is $-0.062<\lambda_{\gamma}<0.147$. When all three couplings $\Delta g_{1}^{Z}, \Delta \kappa_{\gamma}$, and $\lambda_{\gamma}$ are floated freely in the fit, one obtains $\lambda_{\gamma}=0.023_{-0.077}^{+0.074}$.
88 BREITWEG 00 search for $W$ production in events with large hadronic $p_{T}$. For $p_{T}>20$ GeV , the upper limit on the cross section gives the $95 \% \mathrm{CL}$ limit $-3.2<\lambda_{\gamma}<3.2$ (for $\Delta \kappa_{\gamma}=0$ ).
89 ABBOTT 99ı perform a simultaneous fit to the $W \gamma, W W \rightarrow$ dilepton, $W W / W Z \rightarrow$ $e \nu j j, W W / W Z \rightarrow \mu \nu j j$, and $W Z \rightarrow$ trilepton data samples. For $\Lambda=2.0 \mathrm{TeV}$, the $95 \% \mathrm{CL}$ limits are $-0.18<\lambda_{\gamma}<0.19$.
90 ACCIARRI 99Q study $W$-pair, single- $W$, and single photon events.

This coupling is $C P$-conserving but $C$ - and $P$-violating.

| VALUE | EVTS | DOCUMENT ID TECN |  | COMMENT |
| :---: | :---: | :---: | :---: | :---: |
| -0.11 $\pm 0.16$ OUR A | AGE | Error includes scale factor of 1.4. |  |  |
| $-0.04{ }_{-0.12}^{+0.13}$ | 9800 | ${ }^{91}$ ABBIENDI | 04D OPAL | $E_{\mathrm{cm}}^{e e}=183-209 \mathrm{GeV}$ |
| $-0.44-0.22 \pm 0.12$ | 1154 | 92 ACCIARRI | 99Q L3 | $\begin{aligned} & E_{\mathrm{Cm}}^{\mathrm{ee}}=161+172+183 \\ & \mathrm{GeV} \end{aligned}$ |

-     - We do not use the following data for averages, fits, limits, etc. - •
$-0.16 \pm 0.23 \quad 93$ EBOLI 00 THEO LEP1, SLC+ Tevatron
${ }^{91}$ ABBIENDI 04D combine results from $W^{+} W^{-}$in all decay channels. Only CP-conserving couplings are considered and each parameter is determined from a single-parameter fit in which the other parameters assume their Standard Model values. The $95 \%$ confidence interval is $-0.28<\Delta g_{5}^{Z}<0.21$.
${ }^{92}$ ACCIARRI 99Q study $W$-pair, single- $W$, and single photon events.
${ }^{93}$ EBOLI 00 extract this indirect value of the coupling studying the non-universal one-loop contributions to the experimental value of the $Z \rightarrow b \bar{b}$ width ( $\Lambda=1 \mathrm{TeV}$ is assumed).


## $g_{4}^{7}$

This coupling is $C P$-violating ( $C$-violating and $P$-conserving).

| VALUE | EVTS | DOCUMENT ID | TECN | COMMENT |
| :---: | :---: | :---: | :---: | :---: |
| $-0.02+0.32$ | 1065 | 94 ABBIENDI | 01H OPAL | $E_{\mathrm{cm}}^{e} \mathrm{e}=189 \mathrm{GeV}$ |

${ }^{94}$ ABBIENDI 01 H study $W$-pair events, with one leptonically and one hadronically decaying $W$. The coupling is extracted using information from the $W$ production angle together with decay angles from the leptonically decaying $W$.

## $\tilde{\kappa}_{Z}$

This coupling is $C P$-violating ( $C$-conserving and $P$-violating).

| VALUE | EVTS | DOCUMENT ID | TECN | COMMENT |
| :---: | :---: | :---: | :---: | :---: |
| $-0.20_{-0.07}^{+0.10}$ | 1065 | 95 ABBIENDI | 01H OPAL | $E_{\mathrm{cm}}^{e}{ }_{\text {en }}=189 \mathrm{GeV}$ |

${ }^{95}$ ABBIENDI 01H study $W$-pair events, with one leptonically and one hadronically decaying $W$. The coupling is extracted using information from the $W$ production angle together with decay angles from the leptonically decaying $W$.
$\tilde{\lambda}_{z}$
This coupling is $C P$-violating ( $C$-conserving and $P$-violating).

| VALUE | EVTS | DOCUMENT ID | TECN | COMMENT |
| :---: | :---: | :---: | :---: | :---: |
| $-0.18_{-0.16}^{+0.24}$ | 1065 | 96 ABBIENDI | OPAL | $E_{\mathrm{cm}}^{e e}=189 \mathrm{GeV}$ |

[^0]
## W ANOMALOUS MAGNETIC MOMENT

The full magnetic moment is given by $\mu_{W}=e(1+\kappa+\lambda) / 2 m_{W}$. In the Standard Model, at tree level, $\kappa=1$ and $\lambda=0$. Some papers have defined $\Delta \kappa=1-\kappa$ and assume that $\lambda=0$. Note that the electric quadrupole moment is given by $-e(\kappa-\lambda) / m_{W}^{2}$. A description of the parameterization of these moments and additional references can be found in HAGIWARA 87 and BAUR 88. The parameter $\Lambda$ appearing in the theoretical limits below is a regularization cutoff which roughly corresponds to the energy scale where the structure of the $W$ boson becomes manifest.


-     - We do not use the following data for averages, fits, limits, etc.

| 98 |  | 95 |
| :--- | :--- | :--- |
| ABE | CDF |  |
| 99 | ALITTI | $92 C$ |
| UA2 |  |  |
| 100 SAMUEL | 92 | THEO |
| 101 SAMUEL | 91 | THEO |
| 102 GRIFOLS | 88 | THEO |
| 103 GROTCH | 87 | THEO |
| 104 VANDERBIJ | 87 | THEO |
| 105 GRAU | 85 | THEO |
| 106 SUZUKI | 85 | THEO |
| 107 HERZOG | 84 | THEO |

97 ABREU 01I combine results from $e^{+} e^{-}$interactions at 189 GeV leading to $W^{+} W^{-}$, $W e \nu_{e}$, and $\nu \bar{\nu} \gamma$ final states with results from ABREU 99L at 183 GeV to determine $\Delta g_{1}^{Z}, \Delta \kappa_{\gamma}$, and $\lambda_{\gamma} . \Delta \kappa_{\gamma}$ and $\lambda_{\gamma}$ are simultaneously floated in the fit to determine $\mu_{W}$.
98 ABE 95G report $-1.3<\kappa<3.2$ for $\lambda=0$ and $-0.7<\lambda<0.7$ for $\kappa=1$ in $p \bar{p} \rightarrow e \nu_{e} \gamma X$ and $\mu \nu_{\mu} \gamma \mathrm{X}$ at $\sqrt{s}=1.8 \mathrm{TeV}$.
99 ALITTI 92C measure $\kappa=1_{-2.2}^{+2.6}$ and $\lambda=0_{-1.8}^{+1.7}$ in $p \bar{p} \rightarrow e \nu \gamma+X$ at $\sqrt{s}=630 \mathrm{GeV}$. At $95 \%$ CL they report $-3.5<\kappa<5.9$ and $-3.6<\lambda<3.5$.
100 SAMUEL 92 use preliminary CDF and UA2 data and find $-2.4<\kappa<3.7$ at $96 \% \mathrm{CL}$ and $-3.1<\kappa<4.2$ at $95 \%$ CL respectively. They use data for $W \gamma$ production and radiative $W$ decay.
101 SAMUEL 91 use preliminary CDF data for $p \bar{p} \rightarrow W \gamma X$ to obtain $-11.3 \leq \Delta \kappa \leq$ 10.9. Note that their $\kappa=1-\Delta \kappa$.

102 GRIFOLS 88 uses deviation from $\rho$ parameter to set limit $\Delta \kappa \lesssim 65\left(M_{W}^{2} / \Lambda^{2}\right)$.
103 GROTCH 87 finds the limit $-37<\Delta \kappa<73.5$ (90\% CL) from the experimental limits on $e^{+} e^{-} \rightarrow \nu \bar{\nu} \gamma$ assuming three neutrino generations and $-19.5<\Delta \kappa<56$ for four generations. Note their $\Delta \kappa$ has the opposite sign as our definition.
104 VANDERBIJ 87 uses existing limits to the photon structure to obtain $|\Delta \kappa|<33$ $\left(m_{W} / \Lambda\right)$. In addition VANDERBIJ 87 discusses problems with using the $\rho$ parameter of the Standard Model to determine $\Delta \kappa$.
105 GRAU 85 uses the muon anomaly to derive a coupled limit on the anomalous magnetic dipole and electric quadrupole $(\lambda)$ moments $1.05>\Delta \kappa \ln \left(\Lambda / m_{W}\right)+\lambda / 2>-2.77$. In the Standard Model $\lambda=0$.
106 SUZUKI 85 uses partial-wave unitarity at high energies to obtain $|\Delta \kappa| \lesssim 190$ $\left(m_{W} / \Lambda\right)^{2}$. From the anomalous magnetic moment of the muon, SUZUKI 85 obtains $|\Delta \kappa| \lesssim 2.2 / \ln \left(\Lambda / m_{W}\right)$. Finally SUZUKI 85 uses deviations from the $\rho$ parameter and
obtains a very qualitative, order-of-magnitude limit $|\Delta \kappa| \lesssim 150\left(m_{W} / \Lambda\right)^{4}$ if $|\Delta \kappa| \ll$ 107.

07 HERZOG 84 consider the contribution of $W$-boson to muon magnetic moment including anomalous coupling of $W W \gamma$. Obtain a limit $-1<\Delta \kappa<3$ for $\Lambda \gtrsim 1 \mathrm{TeV}$.

## ANOMALOUS $W / Z$ QUARTIC COUPLINGS

Revised November 2003 by C. Caso (University of Genova) and A. Gurtu (Tata Institute).

The Standard Model predictions for $W W W W, W W Z Z$, $W W Z \gamma, W W \gamma \gamma$, and $Z Z \gamma \gamma$ couplings are small at LEP, but expected to become important at a TeV Linear Collider. Outside the Standard Model framework such possible couplings, $a_{0}, a_{c}, a_{n}$, are expressed in terms of the following dimension-6 operators [1,2];

$$
\begin{aligned}
& L_{6}^{0}=-\frac{e^{2}}{16 \Lambda^{2}} a_{0} F^{\mu \nu} F_{\mu \nu} \overrightarrow{W^{\alpha}} \cdot \overrightarrow{W_{\alpha}} \\
& L_{6}^{c}=-\frac{e^{2}}{16 \Lambda^{2}} a_{c} F^{\mu \alpha} F_{\mu \beta} \overrightarrow{W^{\beta}} \cdot \vec{W}_{\alpha} \\
& L_{6}^{n}=-i \frac{e^{2}}{16 \Lambda^{2}} a_{n} \epsilon_{i j k} W_{\mu \alpha}^{(i)} W_{\nu}^{(j)} W^{(k) \alpha} F^{\mu \nu} \\
& \widetilde{L}_{6}^{0}=-\frac{e^{2}}{16 \Lambda^{2}} \widetilde{a}_{0} F^{\mu \nu} \widetilde{F}_{\mu \nu} \overrightarrow{W^{\alpha}} \cdot \vec{W}_{\alpha} \\
& \widetilde{L}_{6}^{n}=-i \frac{e^{2}}{16 \Lambda^{2}} \widetilde{a}_{n} \epsilon_{i j k} W_{\mu \alpha}^{(i)} W_{\nu}^{(j)} W^{(k) \alpha} \widetilde{F}^{\mu \nu}
\end{aligned}
$$

where $F, W$ are photon and $W$ fields, $L_{6}^{0}$ and $L_{6}^{c}$ conserve $C$, $P$ separately ( $\widetilde{L}_{6}^{0}$ conserves only $C$ ) and generate anomalous $W^{+} W^{-} \gamma \gamma$ and $Z Z \gamma \gamma$ couplings, $L_{6}^{n}$ violates $C P\left(\widetilde{L}_{6}^{n}\right.$ violates both $C$ and $P$ ) and generates an anomalous $W^{+} W^{-} Z \gamma$ coupling, and $\Lambda$ is a scale for new physics. For the $Z Z \gamma \gamma$ coupling the $C P$-violating term represented by $L_{6}^{n}$ does not contribute. These couplings are assumed to be real and to vanish at tree level in the Standard Model.

Within the same framework as above, a more recent description of the quartic couplings [3] treats the anomalous parts of the $W W \gamma \gamma$ and $Z Z \gamma \gamma$ couplings separately leading to two sets parameterized as $a_{0}^{V} / \Lambda^{2}$ and $a_{c}^{V} / \Lambda^{2}$, where $V=W$ or $Z$.

At LEP the processes studied in search of these quartic couplings are $e^{+} e^{-} \rightarrow W W \gamma, e^{+} e^{-} \rightarrow \gamma \gamma \nu \bar{\nu}$, and $e^{+} e^{-} \rightarrow$ $Z \gamma \gamma$ and limits are set on the quantities $a_{0}^{W} / \Lambda^{2}, a_{c}^{W} / \Lambda^{2}, a_{n} / \Lambda^{2}$. The characteristics of the first process depend on all the three couplings whereas those of the latter two depend only on the two $C P$-conserving couplings. The sensitive measured variables are the cross sections for these processes as well as the energy and angular distributions of the photon and recoil mass to the photon pair.

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$a_{0} / \Lambda^{2}, a_{c} / \Lambda^{2}, a_{n} / \Lambda^{2}$
Using the $W W \gamma$ final state, the LEP combined $95 \%$ CL limits on the anomalous contributions to the $W W \gamma \gamma$ and $W W Z_{\gamma}$ vertices (as of summer 2003) are given below:
(See P. Wells, "Experimental Tests of the Standard Model," Int. Europhysics Conference on High-Energy Physics, Aachen, Germany, 17-23 July 2003)

$$
\begin{aligned}
& -0.02<a_{0}^{W} / \Lambda^{2}<0.02 \mathrm{GeV}^{-2} \\
& -0.05<a_{c}^{W} / \Lambda^{2}<0.03 \mathrm{GeV}^{-2} \\
& -0.15<a_{n} / \Lambda^{2}<0.15 \mathrm{GeV}^{-2}
\end{aligned}
$$

VALUE DOCUMENTID TECN

-     - We do not use the following data for averages, fits, limits, etc. • - -

| 108 ABBIENDI | 04B OPAL |
| :--- | :--- |
| 109 ABDALLAH | 031 DLPH |
| 110 ACHARD | 02 F L3 |

${ }^{108}$ ABBIENDI 04B select $187 e^{+} e^{-} \rightarrow W^{+} W^{-} \gamma$ events in the C.M. energy range $180-209 \mathrm{GeV}$, where $E_{\gamma}>2.5 \mathrm{GeV}$, the photon has a polar angle $\left|\cos \theta_{\gamma}\right|<0.975$ and is well isolated from the nearest jet and charged lepton, and the effective masses of both fermion-antifermion systems agree with the $W$ mass within $3 \Gamma_{W}$. The measured differential cross section as a function of the photon energy and photon polar angle is used to extract the $95 \% \mathrm{CL}$ limits: $-0.020 \mathrm{GeV}^{-2}<a_{0} / \Lambda^{2}<0.020 \mathrm{GeV}^{-2}$, $-0.053 \mathrm{GeV}^{-2}<a_{c} / \Lambda^{2}<0.037 \mathrm{GeV}^{-2}$ and $-0.16 \mathrm{GeV}^{-2}<a_{n} / \Lambda^{2}<0.15 \mathrm{GeV}^{-2}$.
109 ABDALLAH 031 select $122 e^{+} e^{-} \rightarrow W^{+} W^{-} \gamma$ events in the C.M. energy range $189-209 \mathrm{GeV}$, where $E_{\gamma}>5 \mathrm{GeV}$, the photon has a polar angle $\left|\cos \theta_{\gamma}\right|<0.95$ and is well isolated from the nearest charged fermion. A fit to the photon energy spectra yields $a_{c} / \Lambda^{2}=0.000_{-0.040}^{+0.019} \mathrm{GeV}^{-2}, a_{0} / \Lambda^{2}=-0.004_{-0.010}^{+0.018} \mathrm{GeV}^{-2}, \widetilde{a}_{0} / \Lambda^{2}=$ $-0.007_{-0.008}^{+0.019} \mathrm{GeV}^{-2}, a_{n} / \Lambda^{2}=-0.09_{-0.05}^{+0.16} \mathrm{GeV}^{-2}$, and $\widetilde{a}_{n} / \Lambda^{2}=+0.05_{-0.15}^{+0.07}$ $\mathrm{GeV}^{-2}$, keeping the other parameters fixed to their Standard Model values ( 0 ). The $95 \% \mathrm{CL}$ limits are: $-0.063 \mathrm{GeV}^{-2}<\mathrm{a}_{\mathrm{C}} / \Lambda^{2}<+0.032 \mathrm{GeV}^{-2},-0.020$ $\mathrm{GeV}^{-2}<\mathrm{a}_{0} / \Lambda^{2}<+0.020 \mathrm{GeV}^{-2},-0.020 \mathrm{GeV}^{-2}<\widetilde{a}_{0} / \Lambda^{2}<+0.020 \mathrm{GeV}^{-2}$, $-0.18 \mathrm{GeV}^{-2}<a_{n} / \Lambda^{2}<+0.14 \mathrm{GeV}^{-2},-0.16 \mathrm{GeV}^{-2}<\widetilde{a}_{n} / \Lambda^{2}<+0.17 \mathrm{GeV}^{-2}$.
110 ACHARD 02 select $86 e^{+} e^{-} \rightarrow W^{+} W^{-} \gamma$ events at $192-207 \mathrm{GeV}$, where $E_{\gamma}>5$
GeV and the photon is well isolated. They also select 43 acoplanar $e^{+} e^{-} \rightarrow \nu \bar{\nu} \gamma \gamma$ events in this energy range, where the photon energies are $>5 \mathrm{GeV}$ and $>1 \mathrm{GeV}$ and the photon polar angles are between $14^{\circ}$ and $166^{\circ}$. All these 43 events are in the recoil mass region corresponding to the $Z(75-110 \mathrm{GeV})$. Using the shape and normalization of the photon spectra in the $W^{+} W^{-} \gamma$ events, and combining with the 42 event sample from 189 GeV data (ACCIARRI 00T), they obtain: $a_{0} / \Lambda^{2}=0.000 \pm 0.010 \mathrm{GeV}^{-2}, a_{c} / \Lambda^{2}=$ $-0.013 \pm 0.023 \mathrm{GeV}^{-2}$, and $a_{n} / \Lambda^{2}=-0.002 \pm 0.076 \mathrm{GeV}^{-2}$. Further combining the analyses of $W^{+} W^{-} \gamma$ events with the low recoil mass region of $\nu \bar{\nu} \gamma \gamma$ events (including samples collected at $183+189 \mathrm{GeV}$ ), they obtain the following one-parameter $95 \% \mathrm{CL}$ limits: $-0.015 \mathrm{GeV}^{-2}<a_{0} / \Lambda^{2}<0.015 \mathrm{GeV}^{-2},-0.048 \mathrm{GeV}^{-2}<a_{c} / \Lambda^{2}<0.026$ $\mathrm{GeV}^{-2}$, and $-0.14 \mathrm{GeV}^{-2}<a_{n} / \Lambda^{2}<0.13 \mathrm{GeV}^{-2}$.

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[^0]:    ${ }^{96}$ ABBIENDI 01H study $W$-pair events, with one leptonically and one hadronically decaying $W$. The coupling is extracted using information from the $W$ production angle together with decay angles from the leptonically decaying $W$.

