



$$I(J^P) = \frac{1}{2}(0^-)$$

## $D^0$ MASS

The fit includes  $D^\pm$ ,  $D^0$ ,  $D_s^\pm$ ,  $D^{*\pm}$ ,  $D^{*0}$ , and  $D_s^{*\pm}$  mass and mass difference measurements.

VALUE (MeV)	EVTS	DOCUMENT ID	TECN	COMMENT
<b>1864.5 ± 0.5 OUR FIT</b>	Error includes scale factor of 1.1.			
<b>1864.1 ± 1.0 OUR AVERAGE</b>				
1864.6 ± 0.3 ± 1.0	641	BARLAG	90C ACCM	$\pi^-$ Cu 230 GeV
1852 ± 7	16	ADAMOVICH	87 EMUL	Photoproduction
1861 ± 4		DERRICK	84 HRS	$e^+ e^-$ 29 GeV
● ● ● We do not use the following data for averages, fits, limits, etc. ● ● ●				
1856 ± 36	22	ADAMOVICH	84B EMUL	Photoproduction
1847 ± 7	1	FIORINO	81 EMUL	$\gamma N \rightarrow \bar{D}^0 +$
1863.8 ± 0.5		<sup>1</sup> SCHINDLER	81 MRK2	$e^+ e^-$ 3.77 GeV
1864.7 ± 0.6		<sup>1</sup> TRILLING	81 RVUE	$e^+ e^-$ 3.77 GeV
1863.0 ± 2.5	238	ASTON	80E OMEG	$\gamma p \rightarrow \bar{D}^0$
1860 ± 2	143	<sup>2</sup> AVERY	80 SPEC	$\gamma N \rightarrow D^{*+}$
1869 ± 4	35	<sup>2</sup> AVERY	80 SPEC	$\gamma N \rightarrow D^{*+}$
1854 ± 6	94	<sup>2</sup> ATIYA	79 SPEC	$\gamma N \rightarrow D^0 \bar{D}^0$
1850 ± 15	64	BALTAY	78C HBC	$\nu N \rightarrow K^0 \pi \pi$
1863 ± 3		GOLDHABER	77 MRK1	$D^0, D^+$ recoil spectra
1863.3 ± 0.9		<sup>1</sup> PERUZZI	77 MRK1	$e^+ e^-$ 3.77 GeV
1868 ± 11		PICCOLO	77 MRK1	$e^+ e^-$ 4.03, 4.41 GeV
1865 ± 15	234	GOLDHABER	76 MRK1	$K \pi$ and $K 3\pi$

<sup>1</sup>PERUZZI 77 and SCHINDLER 81 errors do not include the 0.13% uncertainty in the absolute SPEAR energy calibration. TRILLING 81 uses the high precision  $J/\psi(1S)$  and  $\psi(2S)$  measurements of ZHOLENTZ 80 to determine this uncertainty and combines the PERUZZI 77 and SCHINDLER 81 results to obtain the value quoted. TRILLING 81 enters the fit in the  $D^\pm$  mass, and PERUZZI 77 and SCHINDLER 81 enter in the  $m_{D^\pm} - m_{D^0}$ , below.

<sup>2</sup>Error does not include possible systematic mass scale shift, estimated to be less than 5 MeV.

## $m_{D^\pm} - m_{D^0}$

The fit includes  $D^\pm$ ,  $D^0$ ,  $D_s^\pm$ ,  $D^{*\pm}$ ,  $D^{*0}$ , and  $D_s^{*\pm}$  mass and mass difference measurements.

VALUE (MeV)	DOCUMENT ID	TECN	COMMENT
<b>4.78 ± 0.10 OUR FIT</b>	Error includes scale factor of 1.1.		
<b>4.74 ± 0.28 OUR AVERAGE</b>			
4.7 ± 0.3	<sup>3</sup> SCHINDLER	81 MRK2	$e^+ e^-$ 3.77 GeV
5.0 ± 0.8	<sup>3</sup> PERUZZI	77 MRK1	$e^+ e^-$ 3.77 GeV

<sup>3</sup>See the footnote on TRILLING 81 in the  $D^0$  and  $D^\pm$  sections on the mass.

## $D^0$ MEAN LIFE

Measurements with an error  $> 20 \times 10^{-15}$  s have been omitted from the average.

VALUE ( $10^{-15}$ s)	EVTS	DOCUMENT ID	TECN	COMMENT
<b>411.7 ± 2.7 OUR AVERAGE</b>				
407.9 ± 6.0 ± 4.3	10k	KUSHNIR...	01 SELX	$D^0 \rightarrow K^- \pi^+$ , $K^- \pi^+ \pi^+ \pi^-$
413 ± 3 ± 4	35k	AITALA	99E E791	$K^- \pi^+$
408.5 ± 4.1 <sup>+</sup> <sub>3.5</sub> 3.4	25k	BONVICINI	99 CLE2	$e^+ e^- \approx \Upsilon(4S)$
413 ± 4 ± 3	16k	FRABETTI	94D E687	$K^- \pi^+$ , $K^- \pi^+ \pi^+ \pi^-$
424 ± 11 ± 7	5118	FRABETTI	91 E687	$K^- \pi^+$ , $K^- \pi^+ \pi^+ \pi^-$
• • • We do not use the following data for averages, fits, limits, etc. • • •				
417 ± 18 ± 15	890	ALVAREZ	90 NA14	$K^- \pi^+$ , $K^- \pi^+ \pi^+ \pi^-$
388 <sup>+23</sup> -21	641	<sup>4</sup> BARLAG	90C ACCM	$\pi^-$ Cu 230 GeV
480 ± 40 ± 30	776	ALBRECHT	88I ARG	$e^+ e^-$ 10 GeV
422 ± 8 ± 10	4212	RAAB	88 E691	Photoproduction
420 ± 50	90	BARLAG	87B ACCM	$K^-$ and $\pi^-$ 200 GeV

<sup>4</sup>BARLAG 90C estimate systematic error to be negligible.

---

## $D^0-\bar{D}^0$ MIXING

Written April 2002 by D. Asner (LLNL).

Standard Model contributions to  $D^0-\bar{D}^0$  mixing are strongly suppressed by CKM and GIM factors. Thus the observation of  $D^0-\bar{D}^0$  mixing might be evidence for physics beyond the Standard Model. See Bigi [1] for a review of  $D^0-\bar{D}^0$  mixing, and see Nelson [2] for a recent compilation of mixing predictions.

**Formalism:** The time evolution of the  $D^0-\bar{D}^0$  system is described by the Schrödinger equation

$$i \frac{\partial}{\partial t} \begin{pmatrix} D^0(t) \\ \bar{D}^0(t) \end{pmatrix} = \left( \mathbf{M} - \frac{i}{2} \mathbf{\Gamma} \right) \begin{pmatrix} D^0(t) \\ \bar{D}^0(t) \end{pmatrix}, \quad (1)$$

where the  $\mathbf{M}$  and  $\mathbf{\Gamma}$  matrices are Hermitian, and  $CPT$  invariance requires  $M_{11} = M_{22} \equiv M$  and  $\Gamma_{11} = \Gamma_{22} \equiv \Gamma$ . The off-diagonal elements of these matrices describe the dispersive and absorptive parts of  $D^0-\bar{D}^0$  mixing.

The two eigenstates  $D_1$  and  $D_2$  of the effective Hamiltonian matrix  $(\mathbf{M} - \frac{i}{2}\mathbf{\Gamma})$  are given by

$$|D_{1,2}\rangle = p|D^0\rangle \pm q|\bar{D}^0\rangle. \quad (2)$$

The corresponding eigenvalues are

$$\lambda_{1,2} \equiv m_{1,2} - \frac{i}{2}\Gamma_{1,2} = \left(M - \frac{i}{2}\Gamma\right) \pm \frac{q}{p} \left(M_{12} - \frac{i}{2}\Gamma_{12}\right), \quad (3)$$

where

$$\left|\frac{q}{p}\right|^2 = \frac{M_{12}^* - \frac{i}{2}\Gamma_{12}^*}{M_{12} - \frac{i}{2}\Gamma_{12}}. \quad (4)$$

We extend the formalism from our note on “ $B^0$ - $\bar{B}^0$  mixing” [3]. In addition to the ‘right-sign’ instantaneous decay amplitudes  $\bar{A}_f \equiv \langle f|H|\bar{D}^0\rangle$  and  $A_{\bar{f}} \equiv \langle \bar{f}|H|D^0\rangle$  for  $CP$  conjugate final states  $f$  and  $\bar{f}$ , we include the ‘wrong-sign’ amplitudes  $\bar{A}_{\bar{f}} \equiv \langle \bar{f}|H|\bar{D}^0\rangle$  and  $A_f \equiv \langle f|H|D^0\rangle$ .

It is usual to normalize the wrong-sign decay distributions to the integrated rate of right-sign decays and to express time in units of the precisely measured  $D^0$  mean life,  $\bar{\tau}_{D^0} = 1/\Gamma = 2/(\Gamma_1 + \Gamma_2)$ . We denote the resulting decay distributions as  $r(t)$  and  $\bar{r}(t)$ . Starting from a pure  $|D^0\rangle$  or  $|\bar{D}^0\rangle$  state at  $t = 0$ , the time dependence to the wrong-sign final states is then

$$r(t) = \frac{|\langle f|H|D^0(t)\rangle|^2}{|\bar{A}_f|^2} = \left|\frac{q}{p}\right|^2 \left|g_+(t)\chi_f^{-1} + g_-(t)\right|^2, \quad (5)$$

$$\bar{r}(t) = \frac{|\langle \bar{f}|H|\bar{D}^0(t)\rangle|^2}{|A_{\bar{f}}|^2} = \left|\frac{p}{q}\right|^2 \left|g_+(t)\chi_{\bar{f}} + g_-(t)\right|^2, \quad (6)$$

where

$$\chi_f = \frac{q\bar{A}_f}{pA_f} \quad (7)$$

and

$$g_{\pm}(t) = \frac{1}{2} (e^{-iz_1 t} \pm e^{-iz_2 t}) , \quad z_{1,2} = \frac{\lambda_{1,2}}{\Gamma} . \quad (8)$$

Note that a change of the relative phase of  $D^0$  and  $\bar{D}^0$  cancels between  $q/p$  and  $\bar{A}_f/A_f$  and leaves  $\chi_f$  invariant.

Since  $D^0$ - $\bar{D}^0$  mixing is a small effect, the identification tag of the initial particle as a  $D^0$  or a  $\bar{D}^0$  must be extremely accurate. The usual tag is the charge of the distinctive slow pion in the decay sequence  $D^{*+} \rightarrow D^0 \pi^+$  or  $D^{*-} \rightarrow \bar{D}^0 \pi^-$ . In current experiments, mis-tags occur at a rate of about one per thousand events. Another tag of sufficient accuracy is identification of one of the  $D$ 's from  $\psi(3770) \rightarrow D^0 \bar{D}^0$ .

**Semileptonic decays:** We expand  $r(t)$  and  $\bar{r}(t)$  to second order in time for modes where the ratio of decay amplitudes  $R_D = |A_f/\bar{A}_f|^2$  is very small. In semileptonic decays,  $A_f = \bar{A}_f = 0$  in the Standard Model. We define reduced mixing amplitudes  $x$  and  $y$  by

$$x \equiv \frac{2M_{12}}{\Gamma} = \frac{m_1 - m_2}{\Gamma} = \frac{\Delta m}{\Gamma} \quad (9)$$

and

$$y \equiv \frac{\Gamma_{12}}{\Gamma} = \frac{\Gamma_1 - \Gamma_2}{2\Gamma} = \frac{\Delta\Gamma}{2\Gamma} . \quad (10)$$

In these, the first equality holds in the limit of  $CP$  conservation; and now the subscripts 1 and 2 indicate the  $CP$ -even and  $CP$ -odd eigenstates, respectively. Then, in the limit of weak mixing, where  $|ix + y| \ll 1$ ,  $r(t)$  is given by

$$r(t) = |g_-(t)|^2 \left| \frac{q}{p} \right|^2 \approx \frac{e^{-t}}{4} (x^2 + y^2) t^2 \left| \frac{q}{p} \right|^2 . \quad (11)$$

For  $\bar{r}(t)$  one replaces  $q/p$  by  $p/q$ , and in the limit of  $CP$  conservation,  $r(t) = \bar{r}(t)$  and

$$R_M = \int_0^\infty r(t) dt \approx \frac{1}{2} (x^2 + y^2) . \quad (12)$$

***Wrong-sign decays to hadronic non-CP eigenstates:***

Consider the final state  $f = K^+\pi^-$ , where  $A_f$  is doubly Cabibbo suppressed, and the ratio of decay amplitudes is

$$\frac{A_f}{\bar{A}_f} = -\sqrt{R_D}e^{-i\delta} \sim O(\tan^2 \theta_c). \quad (13)$$

The minus sign originates from the sign of  $V_{us}$  with respect to  $V_{cd}$ , and  $\delta$  is a strong phase difference between doubly Cabibbo suppressed and Cabibbo-favored decay amplitudes.

We characterize the violation of  $CP$  in the mixing amplitude, decay amplitude, and the interference between those two processes, by the real-valued parameters  $A_M$ ,  $A_D$ , and  $\phi$ . In general  $\chi_{\bar{f}}$  and  $\chi_f^{-1}$  are two independent complex numbers. We adopt a parameterization similar to that of Nir [4] and CLEO [5] and express these quantities in a way that is convenient to describe the three types of  $CP$  violation:

$$\left| \frac{q}{p} \right| = 1 + A_M, \quad (14)$$

$$\chi_f^{-1} \equiv \frac{pA_f}{q\bar{A}_f} = \frac{-\sqrt{R_D}(1 + A_D)}{(1 + A_M)} e^{-i(\delta+\phi)}, \quad (15)$$

$$\chi_{\bar{f}} \equiv \frac{q\bar{A}_{\bar{f}}}{pA_{\bar{f}}} = \frac{-\sqrt{R_D}(1 + A_M)}{(1 + A_D)} e^{-i(\delta-\phi)}. \quad (16)$$

To leading order,

$$r(t) = e^{-t} \times \left[ R_D(1 + A_D)^2 + \sqrt{R_D}(1 + A_M)(1 + A_D)y'_-t + \frac{(1 + A_M)^2 R_M}{2} t^2 \right], \quad (17)$$

$$\bar{r}(t) = e^{-t} \times \left[ \frac{R_D}{(1+A_D)^2} + \frac{\sqrt{R_D}}{(1+A_D)(1+A_M)} y'_+ t + \frac{R_M}{2(1+A_M)^2} t^2 \right], \quad (18)$$

where

$$y'_\pm \equiv y' \cos \phi \pm x' \sin \phi = y \cos(\delta \mp \phi) - x \sin(\delta \mp \phi) \quad (19)$$

$$y' \equiv y \cos \delta - x \sin \delta, \quad x' \equiv x \cos \delta + y \sin \delta, \quad (20)$$

and  $R_D$  and  $R_M$  are the doubly Cabibbo-suppressed decay and mixing rates, respectively, relative to the time-integrated right-sign rate. Comparing the terms in Eq. (17) and Eq. (18) probes the three fundamental types of  $CP$  violation. In the limit of  $CP$  conservation,  $A_M$ ,  $A_D$ , and  $\phi$  are all zero, and  $r(t) = \bar{r}(t)$ . Eq. (17) and Eq. (18) become

$$r(t) = e^{-t} \left( R_D + \sqrt{R_D} y' t + \frac{1}{2} R_M t^2 \right), \quad (21)$$

$$R = \int_0^\infty r(t) dt = R_D + \sqrt{R_D} y' + R_M, \quad (22)$$

where  $R$  is the time-integrated wrong-sign rate relative to the time-integrated right-sign rate.

For multibody final states, Eqs. (13)–(22) are applicable for one point in the Dalitz space. Although  $x$  and  $y$  do not vary across the Dalitz space, knowledge of the resonant substructure is needed to extrapolate the strong phase difference  $\delta$  to a different point. Both the sign and magnitude of  $x$  and  $y$  are believed to be experimentally accessible through the study of the time-dependent resonant substructure in decay modes such as  $D^0 \rightarrow K_S \pi^+ \pi^-$  [6].

**Decays to  $CP$  eigenstates:** When the final state  $f$  is a  $CP$  eigenstate, there is no distinction between  $f$  and  $\bar{f}$ , and  $A_f = A_{\bar{f}}$

and  $\overline{A}_f = \overline{A}_f$ . We denote final states with  $CP$  eigenvalues  $\pm 1$  by  $f_{\pm}$ . In analogy with Eqs. (5)–(6), the time dependence of decays to  $CP$  eigenstates is then

$$\begin{aligned} r_{\pm}(t) &= \frac{|\langle f_{\pm} | H | D^0(t) \rangle|^2}{|\overline{A}_{\pm}|^2} \\ &= \frac{1}{4} \left| h_{\pm}(t) \left( \frac{A_{\pm}}{\overline{A}_{\pm}} \pm \frac{q}{p} \right) + h_{\mp}(t) \left( \frac{A_{\pm}}{\overline{A}_{\pm}} \mp \frac{q}{p} \right) \right|^2, \\ &\propto \frac{1}{|p|^2} \left| h_{\pm}(t) + \eta_{\pm} h_{\mp}(t) \right|^2, \end{aligned} \quad (23)$$

$$\overline{r}_{\pm}(t) = \frac{|\langle f_{\pm} | H | \overline{D}^0(t) \rangle|^2}{|A_{\pm}|^2} \propto \frac{1}{|q|^2} \left| h_{\pm}(t) - \eta_{\pm} h_{\mp}(t) \right|^2, \quad (24)$$

where

$$h_{\pm}(t) = g_+(t) \pm g_-(t) = e^{-iz_{\pm}t}, \quad (25)$$

$$\eta_{\pm} \equiv \frac{pA_{\pm} \mp q\overline{A}_{\pm}}{pA_{\pm} \pm q\overline{A}_{\pm}} = \frac{1 \mp \chi_{\pm}}{1 \pm \chi_{\pm}}, \quad (26)$$

and the variable  $\eta_{\pm}$  describes  $CP$  violation. This  $\eta_{\pm}$  can receive contributions from each of the three fundamental types of  $CP$  violation.

The quantity  $y$  is accessible experimentally, by comparing the lifetime measured, for example, with decays to non- $CP$  eigenstates such as  $D^0 \rightarrow K^- \pi^+$ , with that measured with decays to a  $CP$  eigenstate such as  $D^0 \rightarrow K^+ K^-$ ; see Bergmann [7]. A positive  $y$  would make  $K^+ K^-$  decays appear to have a shorter lifetime than  $K^- \pi^+$  decays.

In the limit of weak mixing, where  $|ix + y| \ll 1$ , and small  $CP$  violation, where  $|A_M|$ ,  $|A_D|$ , and  $|\sin \phi| \ll 1$ , the time

dependence of decays to  $CP$  eigenstates is proportional to a single exponential:

$$r_{\pm}(t) \propto e^{-[1 \pm \frac{q}{p}(y \cos \phi - x \sin \phi)]t}, \quad (27)$$

$$\bar{r}_{\pm}(t) \propto e^{-[1 \pm \frac{p}{q}(y \cos \phi + x \sin \phi)]t}, \quad (28)$$

$$r_{\pm}(t) + \bar{r}_{\pm}(t) \propto e^{-(1 \pm y_{CP})t}, \quad (29)$$

where

$$y_{CP} \equiv y \cos \phi \left[ \frac{1}{2} \left( \frac{q}{p} + \frac{p}{q} \right) + \frac{A_{\text{prod}}}{2} \left( \frac{q}{p} - \frac{p}{q} \right) \right] \\ - x \sin \phi \left[ \frac{1}{2} \left( \frac{q}{p} - \frac{p}{q} \right) + \frac{A_{\text{prod}}}{2} \left( \frac{q}{p} + \frac{p}{q} \right) \right], \quad (30)$$

and

$$A_{\text{prod}} \equiv \frac{N(D^0) - N(\bar{D}^0)}{N(D^0) + N(\bar{D}^0)} \quad (31)$$

is defined as the production asymmetry of the  $D^0$  and  $\bar{D}^0$ . Note that deviations from the lifetime measured in non- $CP$  eigenstates do not require  $y \neq 0$  but can be due to  $x \sin \phi \neq 0$ . This possibility is distinguished by a relative sign difference between the  $D^0$  and  $\bar{D}^0$  samples.

In the limit of  $CP$  conservation,  $A_{\pm} = \pm \bar{A}_{\pm}$ ,  $\eta_{\pm} = 0$ ,  $y = y_{CP}$ , and

$$r_{\pm}(t) |\bar{A}_{\pm}|^2 = \bar{r}_{\pm}(t) |A_{\pm}|^2 = e^{-(1 \pm y_{CP})t}. \quad (32)$$

## References

1. I. Bigi, in *Proceedings of the Tau-Charm Factory Workshop*, Stanford, Calif., ed. L.V. Beers, SLAC (1989).
2. H.N. Nelson, in *Proceedings of the 19th Intl. Symp. on Lepton and Photon Interactions at High Energy LP99*, ed. J.A. Jaros and M.E. Peskin, SLAC (1999).



3. See the review on  $B^0-\bar{B}^0$  mixing by O. Schneider in this *Review*.
4. Y. Nir, Lectures given at 27th SLAC Summer Institute on Particle Physics:  $CP$  Violation in and Beyond the Standard Model (SSI 99), Stanford, California, 7-16 Jul 1999. Published in Trieste 1999, *Particle Physics* pp. 165-243.
5. R. Godang *et al.*, Phys. Rev. Lett. **84**, 5038 (2000).
6. D. Asner, UCRL-JC-147296.
7. S. Bergmann *et al.*, Phys. Lett. **B486**, 418 (2000).

$$|m_{D_1^0} - m_{D_2^0}|$$

The  $D_1^0$  and  $D_2^0$  are the mass eigenstates of the  $D^0$  meson, as described in the note on " $D^0-\bar{D}^0$  Mixing," above.

VALUE ( $10^{10} \hbar s^{-1}$ )	CL%	DOCUMENT ID	TECN	COMMENT
<b>&lt; 7</b>	95	<sup>5</sup> GODANG	00 CLE2	$e^+ e^-$
● ● ● We do not use the following data for averages, fits, limits, etc. ● ● ●				
<32	90	<sup>6,7</sup> AITALA	98 E791	$\pi^-$ nucleus, 500 GeV
<24	90	<sup>8</sup> AITALA	96C E791	$\pi^-$ nucleus, 500 GeV
<21	90	<sup>7,9</sup> ANJOS	88C E691	Photoproduction

<sup>5</sup>This GODANG 00 limit is inferred from the  $D^0-\bar{D}^0$  mixing ratio  $\Gamma(K^+\pi^-)$  (via  $\bar{D}^0$ )/ $\Gamma(K^-\pi^+)$  given near the end of this  $D^0$  Listings. Decay-time information is used to distinguish DCS decays from  $D^0-\bar{D}^0$  mixing. The limit allows interference between the DCS and mixing ratios, and also allows  $CP$  violation. The strong phase between  $D^0 \rightarrow K^+\pi^-$  and  $\bar{D}^0 \rightarrow K^+\pi^-$  is assumed to be small. If an arbitrary relative strong phase is allowed, the limit degrades by a factor of two.

<sup>6</sup>AITALA 98 allows interference between the doubly Cabibbo-suppressed and mixing amplitudes, and also allows  $CP$  violation in this term, but assumes that  $A_D=A_R=0$ . See the note on " $D^0-\bar{D}^0$  Mixing," above.

<sup>7</sup>This limit is inferred from  $R_M$  for  $f = K^+\pi^-$  and  $f = K^+\pi^-\pi^+\pi^-$ . See the note on " $D^0-\bar{D}^0$  Mixing," above. Decay-time information is used to distinguish doubly Cabibbo-suppressed decays from  $D^0-\bar{D}^0$  mixing.

<sup>8</sup>This limit is inferred from  $R_M$  for  $f = K^+\ell^-\bar{\nu}_\ell$ . See the note on " $D^0-\bar{D}^0$  Mixing," above.

<sup>9</sup>ANJOS 88C assumes that  $y = 0$ . See the note on " $D^0-\bar{D}^0$  Mixing," above. Without this assumption, the limit degrades by about a factor of two.

$$(\Gamma_{D_1^0} - \Gamma_{D_2^0})/\Gamma = 2y$$

The  $D_1^0$  and  $D_2^0$  are the mass eigenstates of the  $D^0$  meson, as described in the note on " $D^0$ - $\bar{D}^0$  Mixing," above.

VALUE	CL%	EVTS	DOCUMENT ID	TECN	COMMENT
<b>-0.003 ± 0.022</b>	<b>OUR AVERAGE</b>				Error includes scale factor of 1.4. See the ideogram below.
-0.010 ± 0.020	+0.014 -0.016	18k	10 ABE	02l BELL	$e^+e^- \approx \gamma(4S)$
-0.024 ± 0.050	± 0.028	3393	11 CSORNA	02 CLE2	$e^+e^- \approx \gamma(4S)$
-0.050	+0.028 -0.032		12 GODANG	00 CLE2	$e^+e^-$
0.0684 ± 0.0278	± 0.0148	10k	10 LINK	00 FOCS	$\gamma$ nucleus
+0.016 ± 0.058	± 0.021		10 AITALA	99E E791	$K^- \pi^+$ , $K^+ K^-$

• • • We do not use the following data for averages, fits, limits, etc. • • •

$ \Delta\Gamma /\Gamma < 0.26$	90	13,14	AITALA	98 E791	$\pi^-$ nucleus, 500 GeV
$ \Delta\Gamma /\Gamma < 0.20$	90	15	AITALA	96C E791	$\pi^-$ nucleus, 500 GeV
$ \Delta\Gamma /\Gamma < 0.17$	90	14,16	ANJOS	88C E691	Photoproduction

<sup>10</sup> LINK 00, AITALA 99E, and ABE 02l measure the lifetime difference between  $D^0 \rightarrow K^- K^+$  ( $CP$  even) decays and  $D^0 \rightarrow K^- \pi^+$  ( $CP$  mixed) decays, or  $y_{CP} = [\Gamma(CP+) - \Gamma(CP-)] / [\Gamma(CP+) + \Gamma(CP-)]$ . We list  $2y_{CP} = \Delta\Gamma/\Gamma$ .

<sup>11</sup> CSORNA 02 measures the lifetime difference between  $D^0 \rightarrow K^- K^+$  and  $\pi^- \pi^+$  ( $CP$  even) decays and  $D^0 \rightarrow K^- \pi^+$  ( $CP$  mixed) decays, or  $y_{CP} = [\Gamma(CP+) - \Gamma(CP-)] / [\Gamma(CP+) + \Gamma(CP-)]$ . We list  $2y_{CP} = \Delta\Gamma/\Gamma$ .

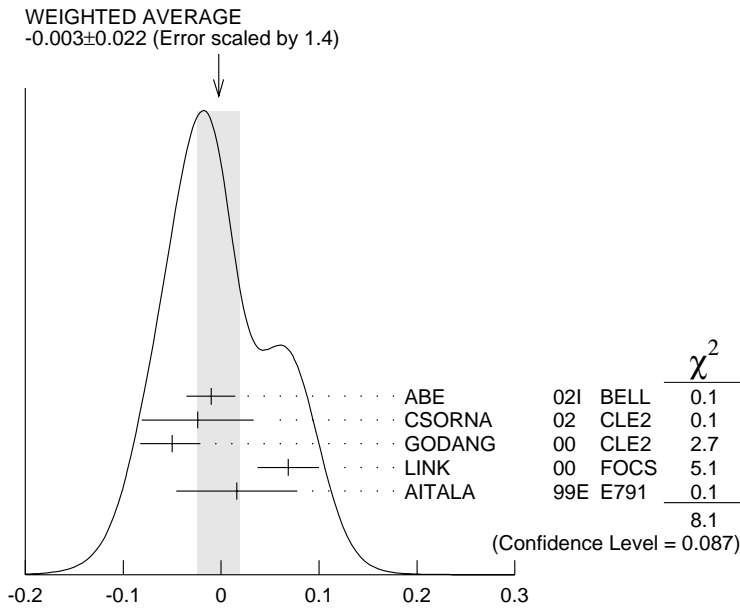
<sup>12</sup> This GODANG 00 limit is inferred from the  $D^0$ - $\bar{D}^0$  mixing ratio  $\Gamma(K^+ \pi^-)$  (via  $\bar{D}^0$ )/ $\Gamma(K^- \pi^+)$  given near the end of this  $D^0$  Listings. Decay-time information is used to distinguish DCS decays from  $D^0$ - $\bar{D}^0$  mixing. The limit allows interference between the DCS and mixing ratios, and also allows  $CP$  violation. The phase between  $D^0 \rightarrow K^+ \pi^-$  and  $\bar{D}^0 \rightarrow K^+ \pi^-$  is assumed to be small. This is a measurement of  $y'$  and is not the same as the  $y_{CP}$  of LINK 00.

<sup>13</sup> AITALA 98 allows interference between the doubly Cabibbo-suppressed and mixing amplitudes, and also allows  $CP$  violation in this term, but assumes that  $A_D = A_R = 0$ . See the note on " $D^0$ - $\bar{D}^0$  Mixing," above.

<sup>14</sup> This limit is inferred from  $R_M$  for  $f = K^+ \pi^-$  and  $f = K^+ \pi^- \pi^+ \pi^-$ . See the note on " $D^0$ - $\bar{D}^0$  Mixing," above. Decay-time information is used to distinguish doubly Cabibbo-suppressed decays from  $D^0$ - $\bar{D}^0$  mixing.

<sup>15</sup> This limit is inferred from  $R_M$  for  $f = K^+ \ell^- \bar{\nu}_\ell$ . See the note on " $D^0$ - $\bar{D}^0$  Mixing," above.

<sup>16</sup> ANJOS 88C assumes that  $y = 0$ . See the note on " $D^0$ - $\bar{D}^0$  Mixing," above. Without this assumption, the limit degrades by about a factor of two.



$$(\Gamma_1 - \Gamma_2)/\Gamma = 2y$$

## $D^0$ DECAY MODES

$\bar{D}^0$  modes are charge conjugates of the modes below.

Mode	Fraction ( $\Gamma_i/\Gamma$ )	Scale factor/ Confidence level
<b>Inclusive modes</b>		
$\Gamma_1$ $e^+$ anything	[a] ( 6.87 $\pm$ 0.28 ) %	
$\Gamma_2$ $\mu^+$ anything	( 6.5 $\pm$ 0.8 ) %	
$\Gamma_3$ $K^-$ anything	( 53 $\pm$ 4 ) %	S=1.3
$\Gamma_4$ $\bar{K}^0$ anything + $K^0$ anything	( 42 $\pm$ 5 ) %	
$\Gamma_5$ $K^+$ anything	( 3.4 $^{+0.6}_{-0.4}$ ) %	
$\Gamma_6$ $\eta$ anything	[b] < 13 %	CL=90%
$\Gamma_7$ $\phi$ anything	( 1.7 $\pm$ 0.8 ) %	
<b>Semileptonic modes</b>		
$\Gamma_8$ $K^- \ell^+ \nu_\ell$	[c] ( 3.43 $\pm$ 0.15 ) %	S=1.2
$\Gamma_9$ $K^- e^+ \nu_e$	( 3.58 $\pm$ 0.18 ) %	S=1.1
$\Gamma_{10}$ $K^- \mu^+ \nu_\mu$	( 3.19 $\pm$ 0.17 ) %	
$\Gamma_{11}$ $K^- \pi^0 e^+ \nu_e$	( 1.1 $^{+0.8}_{-0.6}$ ) %	S=1.6
$\Gamma_{12}$ $\bar{K}^0 \pi^- e^+ \nu_e$	( 1.8 $\pm$ 0.8 ) %	S=1.6

$\Gamma_{13}$	$\bar{K}^*(892)^- e^+ \nu_e$ $\times B(K^{*-} \rightarrow \bar{K}^0 \pi^-)$	( 1.43 ± 0.23 ) %	
$\Gamma_{14}$	$K^*(892)^- \ell^+ \nu_\ell$		
$\Gamma_{15}$	$\bar{K}^*(892)^0 \pi^- e^+ \nu_e$		
$\Gamma_{16}$	$K^- \pi^+ \pi^- \mu^+ \nu_\mu$	< 1.2	$\times 10^{-3}$ CL=90%
$\Gamma_{17}$	$(\bar{K}^*(892)\pi)^- \mu^+ \nu_\mu$	< 1.4	$\times 10^{-3}$ CL=90%
$\Gamma_{18}$	$\pi^- e^+ \nu_e$	( 3.6 ± 0.6 )	$\times 10^{-3}$

A fraction of the following resonance mode has already appeared above as a submode of a charged-particle mode.

$\Gamma_{19}$	$K^*(892)^- e^+ \nu_e$	( 2.14 ± 0.35 ) %
---------------	------------------------	-------------------

**Hadronic modes with a  $\bar{K}$  or  $\bar{K}K\bar{K}$**

$\Gamma_{20}$	$K^- \pi^+$	( 3.80 ± 0.09 ) %	
$\Gamma_{21}$	$\bar{K}^0 \pi^0$	( 2.28 ± 0.22 ) %	
$\Gamma_{22}$	$\bar{K}^0 \pi^+ \pi^-$	[d] ( 5.92 ± 0.35 ) %	S=1.1
$\Gamma_{23}$	$\bar{K}^0 \rho^0$	( 1.47 ± 0.29 ) %	
$\Gamma_{24}$	$\bar{K}^0 f_0(980)$ $\times B(f_0 \rightarrow \pi^+ \pi^-)$	( 3.2 ± 0.9 )	$\times 10^{-3}$
$\Gamma_{25}$	$\bar{K}^0 f_2(1270)$ $\times B(f_2 \rightarrow \pi^+ \pi^-)$	( 2.5 ± 1.0 )	$\times 10^{-3}$
$\Gamma_{26}$	$\bar{K}^0 f_0(1370)$ $\times B(f_0 \rightarrow \pi^+ \pi^-)$	( 4.7 ± 1.4 )	$\times 10^{-3}$
$\Gamma_{27}$	$K^*(892)^- \pi^+$ $\times B(K^{*-} \rightarrow \bar{K}^0 \pi^-)$	( 4.0 ± 0.4 ) %	
$\Gamma_{28}$	$K_0^*(1430)^- \pi^+$ $\times B(K_0^*(1430)^- \rightarrow \bar{K}^0 \pi^-)$	( 7.3 ± 1.6 )	$\times 10^{-3}$
$\Gamma_{29}$	$\bar{K}^0 \pi^+ \pi^-$ nonresonant		
$\Gamma_{30}$	$K^- \pi^+ \pi^0$	[d] (13.1 ± 0.9 ) %	S=1.3
$\Gamma_{31}$	$K^- \rho^+$	(10.2 ± 0.9 ) %	
$\Gamma_{32}$	$K^- \rho(1700)^+$ $\times B(\rho(1700)^+ \rightarrow \pi^+ \pi^0)$	( 7.5 ± 1.7 )	$\times 10^{-3}$
$\Gamma_{33}$	$K^*(892)^- \pi^+$ $\times B(K^{*-} \rightarrow K^- \pi^0)$	( 2.0 ± 0.2 ) %	
$\Gamma_{34}$	$\bar{K}^*(892)^0 \pi^0$ $\times B(\bar{K}^{*0} \rightarrow K^- \pi^+)$	( 1.87 ± 0.27 ) %	
$\Gamma_{35}$	$K_0^*(1430)^- \pi^+$ $\times B(K_0^*(1430)^- \rightarrow K^- \pi^0)$	( 3.6 ± 0.8 )	$\times 10^{-3}$
$\Gamma_{36}$	$\bar{K}_0^*(1430)^0 \pi^0$ $\times B(\bar{K}_0^*(1430)^0 \rightarrow K^- \pi^+)$	( 5.3 $\begin{smallmatrix} +4.2 \\ -1.4 \end{smallmatrix}$ )	$\times 10^{-3}$
$\Gamma_{37}$	$K^*(1680)^- \pi^+$ $\times B(K^*(1680)^- \rightarrow K^- \pi^0)$	( 1.7 ± 0.6 )	$\times 10^{-3}$
$\Gamma_{38}$	$K^- \pi^+ \pi^0$ nonresonant	( 1.05 $\begin{smallmatrix} +0.51 \\ -0.19 \end{smallmatrix}$ ) %	

Γ <sub>39</sub>	$\bar{K}^0 \pi^0 \pi^0$	—
Γ <sub>40</sub>	$\bar{K}^*(892)^0 \pi^0$ × B( $\bar{K}^{*0} \rightarrow \bar{K}^0 \pi^0$ )	( 9.3 ± 1.3 ) × 10 <sup>-3</sup>
Γ <sub>41</sub>	$\bar{K}^0 \pi^0 \pi^0$ nonresonant	( 8.4 ± 2.2 ) × 10 <sup>-3</sup>
Γ <sub>42</sub>	$K^- \pi^+ \pi^+ \pi^-$	[d] ( 7.46 ± 0.31 ) %
Γ <sub>43</sub>	$K^- \pi^+ \rho^0$ total	( 6.2 ± 0.4 ) %
Γ <sub>44</sub>	$K^- \pi^+ \rho^0$ 3-body	( 4.7 ± 2.1 ) × 10 <sup>-3</sup>
Γ <sub>45</sub>	$\bar{K}^*(892)^0 \rho^0$ × B( $\bar{K}^{*0} \rightarrow K^- \pi^+$ )	( 9.7 ± 2.1 ) × 10 <sup>-3</sup>
Γ <sub>46</sub>	$K^- a_1(1260)^+$ × B( $a_1(1260)^+ \rightarrow \pi^+ \pi^+ \pi^-$ )	( 3.6 ± 0.6 ) %
Γ <sub>47</sub>	$\bar{K}^*(892)^0 \pi^+ \pi^-$ total × B( $\bar{K}^{*0} \rightarrow K^- \pi^+$ )	( 1.5 ± 0.4 ) %
Γ <sub>48</sub>	$\bar{K}^*(892)^0 \pi^+ \pi^-$ 3-body × B( $\bar{K}^{*0} \rightarrow K^- \pi^+$ )	( 9.5 ± 2.1 ) × 10 <sup>-3</sup>
Γ <sub>49</sub>	$K_1(1270)^- \pi^+$ × B( $K_1(1270)^- \rightarrow K^- \pi^+ \pi^-$ )	[e] ( 3.7 ± 1.0 ) × 10 <sup>-3</sup>
Γ <sub>50</sub>	$K^- \pi^+ \pi^+ \pi^-$ nonresonant	( 1.74 ± 0.25 ) %
Γ <sub>51</sub>	$\bar{K}^0 \pi^+ \pi^- \pi^0$	[d] (10.8 ± 1.3 ) %
Γ <sub>52</sub>	$\bar{K}^0 \eta$ × B( $\eta \rightarrow \pi^+ \pi^- \pi^0$ )	( 1.7 ± 0.3 ) × 10 <sup>-3</sup>
Γ <sub>53</sub>	$\bar{K}^0 \omega$ × B( $\omega \rightarrow \pi^+ \pi^- \pi^0$ )	( 2.2 ± 0.4 ) %
Γ <sub>54</sub>	$K^*(892)^- \rho^+$ × B( $K^{*-} \rightarrow \bar{K}^0 \pi^-$ )	( 4.3 ± 1.7 ) %
Γ <sub>55</sub>	$\bar{K}^*(892)^0 \rho^0$ × B( $\bar{K}^{*0} \rightarrow \bar{K}^0 \pi^0$ )	( 4.8 ± 1.1 ) × 10 <sup>-3</sup>
Γ <sub>56</sub>	$K_1(1270)^- \pi^+$ × B( $K_1(1270)^- \rightarrow \bar{K}^0 \pi^- \pi^0$ )	[e] ( 5.3 ± 1.5 ) × 10 <sup>-3</sup>
Γ <sub>57</sub>	$\bar{K}^*(892)^0 \pi^+ \pi^-$ 3-body × B( $\bar{K}^{*0} \rightarrow \bar{K}^0 \pi^0$ )	( 4.7 ± 1.0 ) × 10 <sup>-3</sup>
Γ <sub>58</sub>	$\bar{K}^0 \pi^+ \pi^- \pi^0$ nonresonant	( 2.3 ± 2.3 ) %
Γ <sub>59</sub>	$K^- \pi^+ \pi^0 \pi^0$	
Γ <sub>60</sub>	$K^- \pi^+ \pi^+ \pi^- \pi^0$	( 4.0 ± 0.4 ) %
Γ <sub>61</sub>	$\bar{K}^*(892)^0 \pi^+ \pi^- \pi^0$ × B( $\bar{K}^{*0} \rightarrow K^- \pi^+$ )	( 1.2 ± 0.6 ) %
Γ <sub>62</sub>	$\bar{K}^*(892)^0 \eta$ × B( $\bar{K}^{*0} \rightarrow K^- \pi^+$ ) × B( $\eta \rightarrow \pi^+ \pi^- \pi^0$ )	( 2.8 ± 0.6 ) × 10 <sup>-3</sup>
Γ <sub>63</sub>	$K^- \pi^+ \omega$ × B( $\omega \rightarrow \pi^+ \pi^- \pi^0$ )	( 2.7 ± 0.5 ) %
Γ <sub>64</sub>	$\bar{K}^*(892)^0 \omega$ × B( $\bar{K}^{*0} \rightarrow K^- \pi^+$ ) × B( $\omega \rightarrow \pi^+ \pi^- \pi^0$ )	( 6.5 ± 2.4 ) × 10 <sup>-3</sup>
Γ <sub>65</sub>	$\bar{K}^0 \pi^+ \pi^+ \pi^- \pi^-$	( 6.3 ± 1.8 ) × 10 <sup>-3</sup>
Γ <sub>66</sub>	$\bar{K}^0 \pi^+ \pi^- \pi^0 \pi^0 (\pi^0)$	

$\Gamma_{67}$	$\bar{K}^0 K^+ K^-$	( 1.02 ± 0.10 ) %	
	In the fit as $\frac{1}{2}\Gamma_{79} + \Gamma_{69}$ , where $\frac{1}{2}\Gamma_{79} = \Gamma_{68}$ .		
$\Gamma_{68}$	$\bar{K}^0 \phi \times B(\phi \rightarrow K^+ K^-)$	( 4.7 ± 0.6 ) × 10 <sup>-3</sup>	
$\Gamma_{69}$	$\bar{K}^0 K^+ K^-$ non- $\phi$	( 5.5 ± 0.9 ) × 10 <sup>-3</sup>	
$\Gamma_{70}$	$K_S^0 K_S^0 K_S^0$	( 9.1 ± 1.6 ) × 10 <sup>-4</sup>	
$\Gamma_{71}$	$K^- \pi^+ \phi \times B(\phi \rightarrow K^+ K^-)$		
$\Gamma_{72}$	$K^+ K^- K^- \pi^+$	( 2.4 ± 0.7 ) × 10 <sup>-4</sup>	
$\Gamma_{73}$	$K^+ K^- \bar{K}^0 \pi^0$		

Fractions of many of the following modes with resonances have already appeared above as submodes of particular charged-particle modes. (Modes for which there are only upper limits and  $\bar{K}^*(892)\rho$  submodes only appear below.)

$\Gamma_{74}$	$\bar{K}^0 \eta$	( 7.6 ± 1.1 ) × 10 <sup>-3</sup>	
$\Gamma_{75}$	$\bar{K}^0 \rho^0$	( 1.47 ± 0.29 ) %	
$\Gamma_{76}$	$K^- \rho^+$	(10.2 ± 0.8 ) %	S=1.2
$\Gamma_{77}$	$\bar{K}^0 \omega$	( 2.2 ± 0.4 ) %	
$\Gamma_{78}$	$\bar{K}^0 \eta'(958)$	( 1.87 ± 0.28 ) %	
$\Gamma_{79}$	$\bar{K}^0 \phi$	( 9.4 ± 1.1 ) × 10 <sup>-3</sup>	
$\Gamma_{80}$	$K^- a_1(1260)^+$	( 7.2 ± 1.1 ) %	
$\Gamma_{81}$	$\bar{K}^0 a_1(1260)^0$	< 1.9 %	CL=90%
$\Gamma_{82}$	$\bar{K}^0 f_2(1270)$	( 4.5 ± 1.7 ) × 10 <sup>-3</sup>	
$\Gamma_{83}$	$K^- a_2(1320)^+$	< 2 × 10 <sup>-3</sup>	CL=90%
$\Gamma_{84}$	$K^*(892)^- \pi^+$	( 6.0 ± 0.5 ) %	S=1.2
$\Gamma_{85}$	$\bar{K}^*(892)^0 \pi^0$	( 2.8 ± 0.4 ) %	S=1.1
$\Gamma_{86}$	$\bar{K}^*(892)^0 \pi^+ \pi^-$ total	( 2.2 ± 0.5 ) %	
$\Gamma_{87}$	$\bar{K}^*(892)^0 \pi^+ \pi^-$ 3-body	( 1.42 ± 0.31 ) %	
$\Gamma_{88}$	$K^- \pi^+ \rho^0$ total	( 6.2 ± 0.4 ) %	
$\Gamma_{89}$	$K^- \pi^+ \rho^0$ 3-body	( 4.7 ± 2.1 ) × 10 <sup>-3</sup>	
$\Gamma_{90}$	$\bar{K}^*(892)^0 \rho^0$	( 1.45 ± 0.32 ) %	
$\Gamma_{91}$	$\bar{K}^*(892)^0 \rho^0$ transverse	( 1.5 ± 0.5 ) %	
$\Gamma_{92}$	$\bar{K}^*(892)^0 \rho^0$ S-wave	( 2.8 ± 0.6 ) %	
$\Gamma_{93}$	$\bar{K}^*(892)^0 \rho^0$ S-wave long.	< 3 × 10 <sup>-3</sup>	CL=90%
$\Gamma_{94}$	$\bar{K}^*(892)^0 \rho^0$ P-wave	< 3 × 10 <sup>-3</sup>	CL=90%
$\Gamma_{95}$	$\bar{K}^*(892)^0 \rho^0$ D-wave	( 1.9 ± 0.6 ) %	
$\Gamma_{96}$	$K^*(892)^- \rho^+$	( 6.5 ± 2.6 ) %	
$\Gamma_{97}$	$K^*(892)^- \rho^+$ longitudinal	( 3.1 ± 1.3 ) %	
$\Gamma_{98}$	$K^*(892)^- \rho^+$ transverse	( 3.4 ± 2.0 ) %	
$\Gamma_{99}$	$K^*(892)^- \rho^+$ P-wave	< 1.5 %	CL=90%
$\Gamma_{100}$	$K^- \pi^+ f_0(980)$		
$\Gamma_{101}$	$\bar{K}^*(892)^0 f_0(980)$		
$\Gamma_{102}$	$K_1(1270)^- \pi^+$	[e] ( 1.13 ± 0.31 ) %	
$\Gamma_{103}$	$K_1(1400)^- \pi^+$	< 1.2 %	CL=90%
$\Gamma_{104}$	$\bar{K}_1(1400)^0 \pi^0$	< 3.7 %	CL=90%

$\Gamma_{105}$	$K^*(1410)^- \pi^+$		
$\Gamma_{106}$	$K_0^*(1430)^- \pi^+$	( 1.18 ± 0.25 ) %	
$\Gamma_{107}$	$\bar{K}_0^*(1430)^0 \pi^0$	( 8.6 $^{+6.8}_{-2.3}$ ) × 10 <sup>-3</sup>	
$\Gamma_{108}$	$K_2^*(1430)^- \pi^+$	< 9 × 10 <sup>-3</sup>	CL=90%
$\Gamma_{109}$	$\bar{K}_2^*(1430)^0 \pi^0$	< 3.4 × 10 <sup>-3</sup>	CL=90%
$\Gamma_{110}$	$K^*(1680)^- \pi^+$	( 1.3 ± 0.5 ) %	
$\Gamma_{111}$	$\bar{K}^*(892)^0 \pi^+ \pi^- \pi^0$	( 1.8 ± 0.9 ) %	
$\Gamma_{112}$	$\bar{K}^*(892)^0 \eta$	( 1.8 ± 0.4 ) %	
$\Gamma_{113}$	$K^- \pi^+ \omega$	( 3.0 ± 0.6 ) %	
$\Gamma_{114}$	$\bar{K}^*(892)^0 \omega$	( 1.1 ± 0.4 ) %	
$\Gamma_{115}$	$K^- \pi^+ \eta'(958)$	( 6.9 ± 1.8 ) × 10 <sup>-3</sup>	
$\Gamma_{116}$	$\bar{K}^*(892)^0 \eta'(958)$	< 1.0 × 10 <sup>-3</sup>	CL=90%
$\Gamma_{117}$	$K^- \pi^+ \phi$	( 3.3 ± 1.7 ) × 10 <sup>-4</sup>	

### Pionic modes

$\Gamma_{118}$	$\pi^+ \pi^-$	( 1.43 ± 0.07 ) × 10 <sup>-3</sup>	
$\Gamma_{119}$	$\pi^0 \pi^0$	( 8.4 ± 2.2 ) × 10 <sup>-4</sup>	
$\Gamma_{120}$	$\pi^+ \pi^- \pi^0$	( 1.1 ± 0.4 ) %	
$\Gamma_{121}$	$\pi^+ \pi^+ \pi^- \pi^-$	( 7.3 ± 0.5 ) × 10 <sup>-3</sup>	
$\Gamma_{122}$	$\pi^+ \pi^+ \pi^- \pi^- \pi^0$		
$\Gamma_{123}$	$\pi^+ \pi^+ \pi^+ \pi^- \pi^- \pi^-$		

### Hadronic modes with a $K\bar{K}$ pair

$\Gamma_{124}$	$K^+ K^-$	( 4.12 ± 0.14 ) × 10 <sup>-3</sup>	
$\Gamma_{125}$	$K^0 \bar{K}^0$	( 7.1 ± 1.9 ) × 10 <sup>-4</sup>	S=1.2
$\Gamma_{126}$	$K^0 K^- \pi^+$	( 6.9 ± 1.0 ) × 10 <sup>-3</sup>	S=1.1
$\Gamma_{127}$	$\bar{K}^*(892)^0 K^0$ × B( $\bar{K}^{*0} \rightarrow K^- \pi^+$ )	< 1.1 × 10 <sup>-3</sup>	CL=90%
$\Gamma_{128}$	$K^*(892)^+ K^-$ × B( $K^{*+} \rightarrow K^0 \pi^+$ )	( 2.5 ± 0.5 ) × 10 <sup>-3</sup>	
$\Gamma_{129}$	$K^0 K^- \pi^+$ nonresonant	( 2.3 ± 2.3 ) × 10 <sup>-3</sup>	
$\Gamma_{130}$	$\bar{K}^0 K^+ \pi^-$	( 5.2 ± 1.0 ) × 10 <sup>-3</sup>	
$\Gamma_{131}$	$K^*(892)^0 \bar{K}^0$ × B( $K^{*0} \rightarrow K^+ \pi^-$ )	< 6 × 10 <sup>-4</sup>	CL=90%
$\Gamma_{132}$	$K^*(892)^- K^+$ × B( $K^{*-} \rightarrow \bar{K}^0 \pi^-$ )	( 1.3 ± 0.7 ) × 10 <sup>-3</sup>	
$\Gamma_{133}$	$\bar{K}^0 K^+ \pi^-$ nonresonant	( 3.8 $^{+2.3}_{-1.9}$ ) × 10 <sup>-3</sup>	
$\Gamma_{134}$	$K^+ K^- \pi^0$	( 1.24 ± 0.35 ) × 10 <sup>-3</sup>	
$\Gamma_{135}$	$K_S^0 K_S^0 \pi^0$	< 5.9 × 10 <sup>-4</sup>	
$\Gamma_{136}$	$K^+ K^- \pi^+ \pi^-$	[f] ( 2.49 ± 0.23 ) × 10 <sup>-3</sup>	
$\Gamma_{137}$	$\phi \pi^+ \pi^- \times B(\phi \rightarrow K^+ K^-)$	( 5.3 ± 1.4 ) × 10 <sup>-4</sup>	
$\Gamma_{138}$	$\phi \rho^0 \times B(\phi \rightarrow K^+ K^-)$	( 2.9 ± 1.5 ) × 10 <sup>-4</sup>	
$\Gamma_{139}$	$K^+ K^- \rho^0$ 3-body	( 9.0 ± 2.3 ) × 10 <sup>-4</sup>	

$\Gamma_{140}$	$K^*(892)^0 K^- \pi^+ + \text{c.c.}$ $\times B(K^{*0} \rightarrow K^+ \pi^-)$	$[g] < 5$	$\times 10^{-4}$	
$\Gamma_{141}$	$K^*(892)^0 \bar{K}^*(892)^0$ $\times B^2(K^{*0} \rightarrow K^+ \pi^-)$		$(6 \pm 2) \times 10^{-4}$	
$\Gamma_{142}$	$K^+ K^- \pi^+ \pi^- \text{ non-}\phi$			
$\Gamma_{143}$	$K^+ K^- \pi^+ \pi^- \text{ nonresonant}$	$< 8$	$\times 10^{-4}$	CL=90%
$\Gamma_{144}$	$K^0 \bar{K}^0 \pi^+ \pi^-$		$(7.5 \pm 2.9) \times 10^{-3}$	
$\Gamma_{145}$	$K^+ K^- \pi^+ \pi^- \pi^0$		$(3.1 \pm 2.0) \times 10^{-3}$	

Fractions of most of the following modes with resonances have already appeared above as submodes of particular charged-particle modes.

$\Gamma_{146}$	$\bar{K}^*(892)^0 K^0$	$< 1.7$	$\times 10^{-3}$	CL=90%
$\Gamma_{147}$	$K^*(892)^+ K^-$	$(3.8 \pm 0.8)$	$\times 10^{-3}$	
$\Gamma_{148}$	$K^*(892)^0 \bar{K}^0$	$< 9$	$\times 10^{-4}$	CL=90%
$\Gamma_{149}$	$K^*(892)^- K^+$	$(2.0 \pm 1.1)$	$\times 10^{-3}$	
$\Gamma_{150}$	$\phi \pi^0$	$< 1.4$	$\times 10^{-3}$	CL=90%
$\Gamma_{151}$	$\phi \eta$	$< 2.8$	$\times 10^{-3}$	CL=90%
$\Gamma_{152}$	$\phi \omega$	$< 2.1$	$\times 10^{-3}$	CL=90%
$\Gamma_{153}$	$\phi \pi^+ \pi^-$	$(1.07 \pm 0.28)$	$\times 10^{-3}$	
$\Gamma_{154}$	$\phi \rho^0$	$(5.7 \pm 3.0)$	$\times 10^{-4}$	
$\Gamma_{155}$	$\phi \pi^+ \pi^- \text{ 3-body}$	$(7 \pm 5)$	$\times 10^{-4}$	
$\Gamma_{156}$	$K^*(892)^0 K^- \pi^+ + \text{c.c.}$	$[g] < 7$	$\times 10^{-4}$	CL=90%
$\Gamma_{157}$	$K^*(892)^0 K^- \pi^+$			
$\Gamma_{158}$	$\bar{K}^*(892)^0 K^+ \pi^-$			
$\Gamma_{159}$	$K^*(892)^0 \bar{K}^*(892)^0$		$(1.4 \pm 0.5) \times 10^{-3}$	

#### Radiative modes

$\Gamma_{160}$	$\rho^0 \gamma$	$< 2.4$	$\times 10^{-4}$	CL=90%
$\Gamma_{161}$	$\omega \gamma$	$< 2.4$	$\times 10^{-4}$	CL=90%
$\Gamma_{162}$	$\phi \gamma$	$< 1.9$	$\times 10^{-4}$	CL=90%
$\Gamma_{163}$	$K^*(892)^0 \gamma$	$< 7.6$	$\times 10^{-4}$	CL=90%

**Doubly Cabibbo suppressed (DC) modes,  
 $\Delta C = 2$  forbidden via mixing (C2M) modes,  
 $\Delta C = 1$  weak neutral current (C1) modes,  
 Lepton Family number (LF) violating modes, or  
 Lepton number (L) violating modes**

$\Gamma_{164}$	$K^+ \ell^- \bar{\nu}_\ell \text{ (via } \bar{D}^0)$	C2M	$< 1.7$	$\times 10^{-4}$	CL=90%
$\Gamma_{165}$	$K^+ \pi^-$	DC	$(1.48 \pm 0.21)$	$\times 10^{-4}$	
$\Gamma_{166}$	$K^+ \pi^- \text{ (via } \bar{D}^0)$	C2M	$< 1.6$	$\times 10^{-5}$	CL=95%
$\Gamma_{167}$	$K^+ \pi^- \pi^0$		$(5.6 \pm 1.7)$	$\times 10^{-4}$	
$\Gamma_{168}$	$K^+ \pi^- \pi^+ \pi^-$	DC	$(3.1 \pm 1.0)$	$\times 10^{-4}$	
$\Gamma_{169}$	$K^+ \pi^- \pi^+ \pi^- \text{ (via } \bar{D}^0)$	C2M	$< 4$	$\times 10^{-4}$	CL=90%
$\Gamma_{170}$	$K^+ \pi^- \text{ or}$ $K^+ \pi^- \pi^+ \pi^- \text{ (via } \bar{D}^0)$		$< 1.0$	$\times 10^{-3}$	CL=90%



Γ <sub>171</sub>	$\mu^-$ anything (via $\bar{D}^0$ )	C2M	< 4	$\times 10^{-4}$	CL=90%
Γ <sub>172</sub>	$e^+ e^-$	C1	< 6.2	$\times 10^{-6}$	CL=90%
Γ <sub>173</sub>	$\mu^+ \mu^-$	C1	< 4.1	$\times 10^{-6}$	CL=90%
Γ <sub>174</sub>	$\pi^0 e^+ e^-$	C1	< 4.5	$\times 10^{-5}$	CL=90%
Γ <sub>175</sub>	$\pi^0 \mu^+ \mu^-$	C1	< 1.8	$\times 10^{-4}$	CL=90%
Γ <sub>176</sub>	$\eta e^+ e^-$	C1	< 1.1	$\times 10^{-4}$	CL=90%
Γ <sub>177</sub>	$\eta \mu^+ \mu^-$	C1	< 5.3	$\times 10^{-4}$	CL=90%
Γ <sub>178</sub>	$\pi^+ \pi^- e^+ e^-$	C1	< 3.73	$\times 10^{-4}$	CL=90%
Γ <sub>179</sub>	$\rho^0 e^+ e^-$	C1	< 1.0	$\times 10^{-4}$	CL=90%
Γ <sub>180</sub>	$\pi^+ \pi^- \mu^+ \mu^-$	C1	< 3.0	$\times 10^{-5}$	CL=90%
Γ <sub>181</sub>	$\rho^0 \mu^+ \mu^-$	C1	< 2.2	$\times 10^{-5}$	CL=90%
Γ <sub>182</sub>	$\omega e^+ e^-$	C1	< 1.8	$\times 10^{-4}$	CL=90%
Γ <sub>183</sub>	$\omega \mu^+ \mu^-$	C1	< 8.3	$\times 10^{-4}$	CL=90%
Γ <sub>184</sub>	$K^- K^+ e^+ e^-$	C1	< 3.15	$\times 10^{-4}$	CL=90%
Γ <sub>185</sub>	$\phi e^+ e^-$	C1	< 5.2	$\times 10^{-5}$	CL=90%
Γ <sub>186</sub>	$K^- K^+ \mu^+ \mu^-$	C1	< 3.3	$\times 10^{-5}$	CL=90%
Γ <sub>187</sub>	$\phi \mu^+ \mu^-$	C1	< 3.1	$\times 10^{-5}$	CL=90%
Γ <sub>188</sub>	$\bar{K}^0 e^+ e^-$	[h]	< 1.1	$\times 10^{-4}$	CL=90%
Γ <sub>189</sub>	$\bar{K}^0 \mu^+ \mu^-$	[h]	< 2.6	$\times 10^{-4}$	CL=90%
Γ <sub>190</sub>	$K^- \pi^+ e^+ e^-$	C1	< 3.85	$\times 10^{-4}$	CL=90%
Γ <sub>191</sub>	$\bar{K}^*(892)^0 e^+ e^-$	[h]	< 4.7	$\times 10^{-5}$	CL=90%
Γ <sub>192</sub>	$K^- \pi^+ \mu^+ \mu^-$	C1	< 3.59	$\times 10^{-4}$	CL=90%
Γ <sub>193</sub>	$\bar{K}^*(892)^0 \mu^+ \mu^-$	[h]	< 2.4	$\times 10^{-5}$	CL=90%
Γ <sub>194</sub>	$\pi^+ \pi^- \pi^0 \mu^+ \mu^-$	C1	< 8.1	$\times 10^{-4}$	CL=90%
Γ <sub>195</sub>	$\mu^\pm e^\mp$	LF	[i] < 8.1	$\times 10^{-6}$	CL=90%
Γ <sub>196</sub>	$\pi^0 e^\pm \mu^\mp$	LF	[i] < 8.6	$\times 10^{-5}$	CL=90%
Γ <sub>197</sub>	$\eta e^\pm \mu^\mp$	LF	[i] < 1.0	$\times 10^{-4}$	CL=90%
Γ <sub>198</sub>	$\pi^+ \pi^- e^\pm \mu^\mp$	LF	[i] < 1.5	$\times 10^{-5}$	CL=90%
Γ <sub>199</sub>	$\rho^0 e^\pm \mu^\mp$	LF	[i] < 4.9	$\times 10^{-5}$	CL=90%
Γ <sub>200</sub>	$\omega e^\pm \mu^\mp$	LF	[i] < 1.2	$\times 10^{-4}$	CL=90%
Γ <sub>201</sub>	$K^- K^+ e^\pm \mu^\mp$	LF	[i] < 1.8	$\times 10^{-4}$	CL=90%
Γ <sub>202</sub>	$\phi e^\pm \mu^\mp$	LF	[i] < 3.4	$\times 10^{-5}$	CL=90%
Γ <sub>203</sub>	$\bar{K}^0 e^\pm \mu^\mp$	LF	[i] < 1.0	$\times 10^{-4}$	CL=90%
Γ <sub>204</sub>	$K^- \pi^+ e^\pm \mu^\mp$	LF	[i] < 5.53	$\times 10^{-4}$	CL=90%
Γ <sub>205</sub>	$\bar{K}^*(892)^0 e^\pm \mu^\mp$	LF	[i] < 8.3	$\times 10^{-5}$	CL=90%
Γ <sub>206</sub>	$\pi^- \pi^- e^+ e^+ + \text{c.c.}$	L	< 1.12	$\times 10^{-4}$	CL=90%
Γ <sub>207</sub>	$\pi^- \pi^- \mu^+ \mu^+ + \text{c.c.}$	L	< 2.9	$\times 10^{-5}$	CL=90%
Γ <sub>208</sub>	$K^- \pi^- e^+ e^+ + \text{c.c.}$	L	< 2.06	$\times 10^{-4}$	CL=90%
Γ <sub>209</sub>	$K^- \pi^- \mu^+ \mu^+ + \text{c.c.}$	L	< 3.9	$\times 10^{-4}$	CL=90%
Γ <sub>210</sub>	$K^- K^- e^+ e^+ + \text{c.c.}$	L	< 1.52	$\times 10^{-4}$	CL=90%
Γ <sub>211</sub>	$K^- K^- \mu^+ \mu^+ + \text{c.c.}$	L	< 9.4	$\times 10^{-5}$	CL=90%
Γ <sub>212</sub>	$\pi^- \pi^- e^+ \mu^+ + \text{c.c.}$	L	< 7.9	$\times 10^{-5}$	CL=90%
Γ <sub>213</sub>	$K^- \pi^- e^+ \mu^+ + \text{c.c.}$	L	< 2.18	$\times 10^{-4}$	CL=90%
Γ <sub>214</sub>	$K^- K^- e^+ \mu^+ + \text{c.c.}$	L	< 5.7	$\times 10^{-5}$	CL=90%

$\Gamma_{215}$  A dummy mode used by the fit. (11.5  $\pm$  3.4 ) % S=1.1

- [a] The exclusive  $e^+$  modes  $K^- e^+ \nu_e$ ,  $K^- \pi^0 e^+ \nu_e$ ,  $\bar{K}^0 \pi^- e^+ \nu_e$  and  $\pi^- e^+ \nu_e$  are constrained to equal this (well-measured) inclusive fraction.
  - [b] This is a weighted average of  $D^\pm$  (44%) and  $D^0$  (56%) branching fractions. See " $D^+$  and  $D^0 \rightarrow (\eta \text{ anything}) / (\text{total } D^+ \text{ and } D^0)$ " under " $D^+$  Branching Ratios" in these Particle Listings.
  - [c] This value averages the  $e^+$  and  $\mu^+$  branching fractions, after making a small phase-space adjustment to the  $\mu^+$  fraction to be able to use it as an  $e^+$  fraction; hence our  $\ell^+$  here is really an  $e^+$ .
  - [d] The branching fraction for this mode may differ from the sum of the submodes that contribute to it, due to interference effects. See the relevant papers.
  - [e] The two experiments measuring this fraction are in serious disagreement. See the Particle Listings.
  - [f] The experiments on the division of this charge mode amongst its submodes disagree, and the submode branching fractions here add up to considerably more than the charged-mode fraction.
  - [g] However, these upper limits are in serious disagreement with values obtained in another experiment.
  - [h] This mode is not a useful test for a  $\Delta C=1$  weak neutral current because both quarks must change flavor in this decay.
  - [i] The value is for the sum of the charge states or particle/antiparticle states indicated.
-

## CONSTRAINED FIT INFORMATION

An overall fit to 56 branching ratios uses 125 measurements and one constraint to determine 31 parameters. The overall fit has a  $\chi^2 = 60.0$  for 95 degrees of freedom.

The following *off-diagonal* array elements are the correlation coefficients  $\langle \delta x_i \delta x_j \rangle / (\delta x_i \delta x_j)$ , in percent, from the fit to the branching fractions,  $x_i \equiv \Gamma_i / \Gamma_{\text{total}}$ . The fit constrains the  $x_i$  whose labels appear in this array to sum to one.

$x_9$	6									
$x_{10}$	31	19								
$x_{11}$	0	-8	-1							
$x_{12}$	-1	-15	-3	-91						
$x_{18}$	1	24	5	-4	-8					
$x_{19}$	1	10	3	-1	-1	2				
$x_{20}$	13	45	42	-3	-7	11	8			
$x_{21}$	1	6	5	0	-1	1	17	11		
$x_{22}$	2	10	7	-1	-1	2	28	18	58	
$x_{30}$	3	11	9	-1	-2	3	9	22	21	30
$x_{42}$	5	18	17	-1	-3	4	3	40	5	8
$x_{51}$	1	4	3	0	-1	1	12	8	25	42
$x_{60}$	3	9	8	-1	-1	2	2	19	2	3
$x_{69}$	1	4	3	0	-1	1	11	7	22	38
$x_{74}$	1	4	3	0	-1	1	11	7	54	40
$x_{77}$	1	3	2	0	0	1	9	5	18	30
$x_{79}$	1	5	4	0	-1	1	15	9	30	51
$x_{84}$	2	8	7	-1	-1	2	18	16	37	62
$x_{85}$	1	5	4	0	-1	1	5	10	19	18
$x_{87}$	1	3	3	0	0	1	1	7	1	1
$x_{91}$	1	2	2	0	0	0	1	4	2	4
$x_{102}$	0	2	2	0	0	0	5	4	9	16
$x_{106}$	1	3	2	0	0	1	7	5	15	25
$x_{112}$	1	3	3	0	0	1	2	6	5	7
$x_{124}$	9	31	29	-2	-5	7	5	68	7	12
$x_{125}$	0	2	2	0	0	0	6	4	12	21
$x_{126}$	1	4	4	0	-1	1	10	8	20	33
$x_{130}$	1	4	3	0	-1	1	7	7	14	23
$x_{147}$	1	3	2	0	0	1	7	5	15	25
$x_{215}$	-28	-17	-24	0	1	-4	-29	-34	-47	-67
	$x_2$	$x_9$	$x_{10}$	$x_{11}$	$x_{12}$	$x_{18}$	$x_{19}$	$x_{20}$	$x_{21}$	$x_{22}$

x <sub>42</sub>	9										
x <sub>51</sub>	13	4									
x <sub>60</sub>	4	28	2								
x <sub>69</sub>	11	3	16	1							
x <sub>74</sub>	14	3	17	1	15						
x <sub>77</sub>	9	3	39	1	11	12					
x <sub>79</sub>	15	4	21	2	-3	21	15				
x <sub>84</sub>	43	7	26	3	24	25	19	32			
x <sub>85</sub>	44	4	8	2	7	11	6	9	22		
x <sub>87</sub>	2	18	1	5	1	1	0	1	1	1	
x <sub>91</sub>	2	10	8	3	1	2	3	2	3	1	
x <sub>102</sub>	5	4	37	1	6	6	15	8	10	3	
x <sub>106</sub>	12	2	11	1	10	10	8	13	17	6	
x <sub>112</sub>	23	2	3	1	3	3	2	4	10	10	
x <sub>124</sub>	15	27	5	13	5	5	4	6	11	7	
x <sub>125</sub>	6	2	9	1	8	8	6	11	13	4	
x <sub>126</sub>	10	4	14	2	13	13	10	17	21	6	
x <sub>130</sub>	8	3	10	1	9	9	7	12	15	5	
x <sub>147</sub>	8	2	11	1	10	10	8	13	16	5	
x <sub>215</sub>	-56	-27	-66	-20	-27	-33	-42	-36	-59	-37	
	x <sub>30</sub>	x <sub>42</sub>	x <sub>51</sub>	x <sub>60</sub>	x <sub>69</sub>	x <sub>74</sub>	x <sub>77</sub>	x <sub>79</sub>	x <sub>84</sub>	x <sub>85</sub>	
x <sub>91</sub>	2										
x <sub>102</sub>	1	3									
x <sub>106</sub>	0	1	4								
x <sub>112</sub>	0	0	1	3							
x <sub>124</sub>	5	3	3	4	4						
x <sub>125</sub>	0	1	3	5	1	2					
x <sub>126</sub>	1	1	5	8	2	6	7				
x <sub>130</sub>	1	1	4	6	2	5	5	8			
x <sub>147</sub>	0	1	4	6	2	4	5	8	6		
x <sub>215</sub>	-14	-23	-33	-25	-26	-23	-14	-26	-19	-19	
	x <sub>87</sub>	x <sub>91</sub>	x <sub>102</sub>	x <sub>106</sub>	x <sub>112</sub>	x <sub>124</sub>	x <sub>125</sub>	x <sub>126</sub>	x <sub>130</sub>	x <sub>147</sub>	

## $D^0$ BRANCHING RATIOS

See the "Note on  $D$  Mesons" in the  $D^\pm$  Listings.

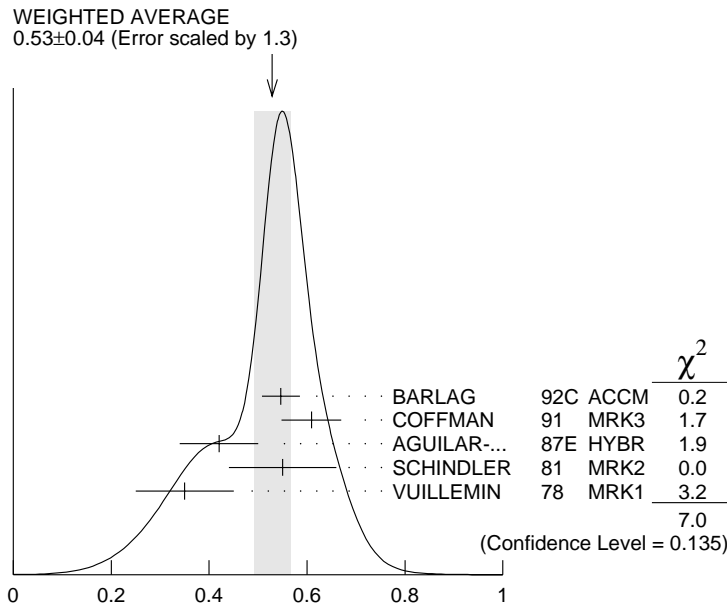
Some older now obsolete results have been omitted from these Listings.

### ———— Inclusive modes ————

$\Gamma(e^+ \text{ anything})/\Gamma_{\text{total}}$		$\Gamma_1/\Gamma = (\Gamma_9 + \Gamma_{11} + \Gamma_{12} + \Gamma_{18})/\Gamma$			
<u>VALUE</u>	<u>EVTS</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>	
<b>0.0687 ± 0.0028 OUR FIT</b>					
<b>0.0675 ± 0.0029 OUR AVERAGE</b>					
0.069 ± 0.003 ± 0.005	1670	ALBRECHT	96C ARG	$e^+ e^- \approx 10$ GeV	
0.0664 ± 0.0018 ± 0.0029	4609	<sup>17</sup> KUBOTA	96B CLE2	$e^+ e^- \approx \gamma(4S)$	
0.075 ± 0.011 ± 0.004	137	BALTRUSAIT..	85B MRK3	$e^+ e^- 3.77$ GeV	
• • • We do not use the following data for averages, fits, limits, etc. • • •					
0.15 ± 0.05		AGUILAR-...	87E HYBR	$\pi p, pp$ 360, 400 GeV	
0.055 ± 0.037	12	SCHINDLER	81 MRK2	$e^+ e^- 3.771$ GeV	
<sup>17</sup> KUBOTA 96B uses $D^{*+} \rightarrow D^0 \pi^+$ (and charge conjugate) events in which the $D^0$ subsequently decays to $X e^+ \nu_e$ .					

$\Gamma(\mu^+ \text{ anything})/\Gamma_{\text{total}}$						$\Gamma_2/\Gamma$
<u>VALUE</u>	<u>EVTS</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>		
<b>0.065 ± 0.008 OUR FIT</b>						
<b>0.060 ± 0.007 ± 0.012</b>	310	ALBRECHT	96C ARG	$e^+ e^- \approx 10$ GeV		

$\Gamma(K^- \text{ anything})/\Gamma_{\text{total}}$						$\Gamma_3/\Gamma$
<u>VALUE</u>	<u>EVTS</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>		
<b>0.53 ± 0.04 OUR AVERAGE</b>						
Error includes scale factor of 1.3. See the ideogram below.						
0.546 <sup>+0.039</sup> <sub>-0.038</sub>		<sup>18</sup> BARLAG	92C ACCM	$\pi^-$ Cu 230 GeV		
0.609 ± 0.032 ± 0.052		COFFMAN	91 MRK3	$e^+ e^- 3.77$ GeV		
0.42 ± 0.08		AGUILAR-...	87E HYBR	$\pi p, pp$ 360, 400 GeV		
0.55 ± 0.11	121	SCHINDLER	81 MRK2	$e^+ e^- 3.771$ GeV		
0.35 ± 0.10	19	VUILLEMIN	78 MRK1	$e^+ e^- 3.772$ GeV		
<sup>18</sup> BARLAG 92C computes the branching fraction using topological normalization.						



$$\Gamma(K^- \text{ anything})/\Gamma_{\text{total}}$$

$$\frac{[\Gamma(\bar{K}^0 \text{ anything}) + \Gamma(K^0 \text{ anything})]/\Gamma_{\text{total}}}{\Gamma_4/\Gamma}$$

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
<b>0.42 ± 0.05</b>				<b>OUR AVERAGE</b>
0.455 ± 0.050 ± 0.032		COFFMAN	91 MRK3	$e^+ e^-$ 3.77 GeV
0.29 ± 0.11	13	SCHINDLER	81 MRK2	$e^+ e^-$ 3.771 GeV
0.57 ± 0.26	6	VUILLEMIN	78 MRK1	$e^+ e^-$ 3.772 GeV

$$\frac{\Gamma(K^+ \text{ anything})/\Gamma_{\text{total}}}{\Gamma_5/\Gamma}$$

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
<b>0.034<sup>+0.006</sup><sub>-0.004</sub></b>				<b>OUR AVERAGE</b>
0.034 <sup>+0.007</sup> <sub>-0.005</sub>		<sup>19</sup> BARLAG	92C ACCM	$\pi^-$ Cu 230 GeV
0.028 ± 0.009 ± 0.004		COFFMAN	91 MRK3	$e^+ e^-$ 3.77 GeV
0.03 <sup>+0.05</sup> <sub>-0.02</sub>		AGUILAR-...	87E HYBR	$\pi p, pp$ 360, 400 GeV
0.08 ± 0.03	25	SCHINDLER	81 MRK2	$e^+ e^-$ 3.771 GeV

<sup>19</sup> BARLAG 92C computes the branching fraction using topological normalization.

$$\frac{\Gamma(\phi \text{ anything})/\Gamma_{\text{total}}}{\Gamma_7/\Gamma}$$

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
<b>0.0171<sup>+0.0076</sup><sub>-0.0071</sub> ± 0.0017</b>	9	<sup>20</sup> BAI	00C BES	$e^+ e^- \rightarrow D\bar{D}^*, D^*\bar{D}^*$

<sup>20</sup> BAI 00C finds the average ( $\phi$  anything) branching fraction for the 4.03-GeV mix of  $D^+$  and  $D^0$  mesons to be  $(1.34 \pm 0.52 \pm 0.12)\%$ .

————— **Semileptonic modes** —————

**$\Gamma(K^- \ell^+ \nu_\ell)/\Gamma_{\text{total}}$**

**$\Gamma_8/\Gamma$**

We average our  $K^- e^+ \nu_e$  and  $K^- \mu^+ \nu_\mu$  branching fractions, after multiplying the latter by a phase-space factor of 1.03 to be able to use it with the  $K^- e^+ \nu_e$  fraction.

Hence our  $\ell^+$  here is really an  $e^+$ .

<u>VALUE</u>	<u>DOCUMENT ID</u>	<u>COMMENT</u>
<b>0.0343 ± 0.0015 OUR AVERAGE</b>	Error includes scale factor of 1.2.	
0.0359 ± 0.0018	PDG	02 Our $\Gamma(K^- e^+ \nu_e)/\Gamma_{\text{total}}$
0.0329 ± 0.0017	PDG	02 1.03 × our $\Gamma(K^- \mu^+ \nu_\mu)/\Gamma_{\text{total}}$

**$\Gamma(K^- e^+ \nu_e)/\Gamma_{\text{total}}$**

**$\Gamma_9/\Gamma$**

<u>VALUE</u>	<u>EVTS</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
<b>0.0358 ± 0.0018 OUR FIT</b>		Error includes scale factor of 1.1.		
<b>0.034 ± 0.005 ± 0.004</b>	55	ADLER	89 MRK3	$e^+ e^-$ 3.77 GeV

**$\Gamma(K^- e^+ \nu_e)/\Gamma(K^- \pi^+)$**

**$\Gamma_9/\Gamma_{20}$**

<u>VALUE</u>	<u>EVTS</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
<b>0.94 ± 0.04 OUR FIT</b>				
<b>0.95 ± 0.04 OUR AVERAGE</b>				
0.978 ± 0.027 ± 0.044	2510	<sup>21</sup> BEAN	93C CLE2	$e^+ e^- \approx \gamma(4S)$
0.90 ± 0.06 ± 0.06	584	<sup>22</sup> CRAWFORD	91B CLEO	$e^+ e^- \approx 10.5$ GeV
0.91 ± 0.07 ± 0.11	250	<sup>23</sup> ANJOS	89F E691	Photoproduction

<sup>21</sup> BEAN 93C uses  $K^- \mu^+ \nu_\mu$  as well as  $K^- e^+ \nu_e$  events and makes a small phase-space adjustment to the number of the  $\mu^+$  events to use them as  $e^+$  events. A pole mass of  $2.00 \pm 0.12 \pm 0.18$  GeV/ $c^2$  is obtained from the  $q^2$  dependence of the decay rate.

<sup>22</sup> CRAWFORD 91B uses  $K^- e^+ \nu_e$  and  $K^- \mu^+ \nu_\mu$  candidates to measure a pole mass of  $2.1^{+0.4+0.3}_{-0.2-0.2}$  GeV/ $c^2$  from the  $q^2$  dependence of the decay rate.

<sup>23</sup> ANJOS 89F measures a pole mass of  $2.1^{+0.4}_{-0.2} \pm 0.2$  GeV/ $c^2$  from the  $q^2$  dependence of the decay rate.

**$\Gamma(K^- \mu^+ \nu_\mu)/\Gamma(K^- \pi^+)$**

**$\Gamma_{10}/\Gamma_{20}$**

<u>VALUE</u>	<u>EVTS</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
<b>0.84 ± 0.04 OUR FIT</b>				
<b>0.84 ± 0.04 OUR AVERAGE</b>				
0.852 ± 0.034 ± 0.028	1897	<sup>24</sup> FRABETTI	95G E687	$\gamma \text{Be } \bar{E}_\gamma = 220$ GeV
0.82 ± 0.13 ± 0.13	338	<sup>25</sup> FRABETTI	93I E687	$\gamma \text{Be } \bar{E}_\gamma = 221$ GeV
0.79 ± 0.08 ± 0.09	231	<sup>26</sup> CRAWFORD	91B CLEO	$e^+ e^- \approx 10.5$ GeV

<sup>24</sup> FRABETTI 95G extracts the ratio of form factors  $f_-(0)/f_+(0) = -1.3^{+3.6}_{-3.4} \pm 0.6$ , and measures a pole mass of  $1.87^{+0.11+0.07}_{-0.08-0.06}$  GeV/ $c^2$  from the  $q^2$  dependence of the decay rate.

<sup>25</sup> FRABETTI 93I measures a pole mass of  $2.1^{+0.7+0.7}_{-0.3-0.3}$  GeV/ $c^2$  from the  $q^2$  dependence of the decay rate.

<sup>26</sup> CRAWFORD 91B measures a pole mass of  $2.00 \pm 0.12 \pm 0.18$  GeV/ $c^2$  from the  $q^2$  dependence of the decay rate.

$\Gamma(K^- \mu^+ \nu_\mu) / \Gamma(\mu^+ \text{ anything})$   $\Gamma_{10} / \Gamma_2$

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
-------	------	-------------	------	---------

**0.49 ± 0.06 OUR FIT**

**0.472 ± 0.051 ± 0.040**    232    KODAMA    94    E653     $\pi^-$  emulsion 600 GeV

• • • We do not use the following data for averages, fits, limits, etc. • • •

0.32 ± 0.05 ± 0.05    124    KODAMA    91    EMUL     $pA$  800 GeV

$\Gamma(K^- \pi^0 e^+ \nu_e) / \Gamma_{\text{total}}$   $\Gamma_{11} / \Gamma$

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
-------	------	-------------	------	---------

**0.011<sup>+0.008</sup><sub>-0.006</sub> OUR FIT**    Error includes scale factor of 1.6.

**0.016<sup>+0.013</sup><sub>-0.005</sub> ± 0.002**    4    27 BAI    91    MRK3     $e^+ e^- \approx 3.77$  GeV

<sup>27</sup> BAI 91 finds that a fraction  $0.79^{+0.15+0.09}_{-0.17-0.03}$  of combined  $D^+$  and  $D^0$  decays to  $\bar{K} \pi e^+ \nu_e$  (24 events) are  $\bar{K}^*(892) e^+ \nu_e$ . BAI 91 uses 56  $K^- e^+ \nu_e$  events to measure a pole mass of  $1.8 \pm 0.3 \pm 0.2$  GeV/ $c^2$  from the  $q^2$  dependence of the decay rate.

$\Gamma(\bar{K}^0 \pi^- e^+ \nu_e) / \Gamma_{\text{total}}$   $\Gamma_{12} / \Gamma$

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
-------	------	-------------	------	---------

**0.018 ± 0.008 OUR FIT**    Error includes scale factor of 1.6.

**0.028<sup>+0.017</sup><sub>-0.008</sub> ± 0.003**    6    28 BAI    91    MRK3     $e^+ e^- \approx 3.77$  GeV

<sup>28</sup> BAI 91 finds that a fraction  $0.79^{+0.15+0.09}_{-0.17-0.03}$  of combined  $D^+$  and  $D^0$  decays to  $\bar{K} \pi e^+ \nu_e$  (24 events) are  $\bar{K}^*(892) e^+ \nu_e$ .

$\Gamma(K^*(892)^- e^+ \nu_e) / \Gamma(K^- e^+ \nu_e)$   $\Gamma_{19} / \Gamma_9$

Unseen decay modes of the  $K^*(892)^-$  are included.

VALUE	DOCUMENT ID	TECN	COMMENT
-------	-------------	------	---------

**0.60 ± 0.10 OUR FIT**

**0.51 ± 0.18 ± 0.06**    CRAWFORD    91B CLEO     $e^+ e^- \approx 10.5$  GeV

$\Gamma(K^*(892)^- e^+ \nu_e) / \Gamma(\bar{K}^0 \pi^+ \pi^-)$   $\Gamma_{19} / \Gamma_{22}$

Unseen decay modes of the  $\bar{K}^*(892)^-$  are included.

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
-------	------	-------------	------	---------

**0.36 ± 0.06 OUR FIT**

**0.38 ± 0.06 ± 0.03**    152    29 BEAN    93C CLE2     $e^+ e^- \approx \Upsilon(4S)$

<sup>29</sup> BEAN 93C uses  $K^{*-} \mu^+ \nu_\mu$  as well as  $K^{*-} e^+ \nu_e$  events and makes a small phase-space adjustment to the number of the  $\mu^+$  events to use them as  $e^+$  events.

$\Gamma(K^*(892)^- \ell^+ \nu_\ell) / \Gamma(\bar{K}^0 \pi^+ \pi^-)$   $\Gamma_{14} / \Gamma_{22}$

This an average of the  $K^*(892)^- e^+ \nu_e$  and  $K^*(892)^- \mu^+ \nu_\mu$  ratios. Unseen decay modes of the  $K^*(892)^-$  are included.

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
-------	------	-------------	------	---------

• • • We do not use the following data for averages, fits, limits, etc. • • •

0.24 ± 0.07 ± 0.06    137    30 ALEXANDER    90B CLEO     $e^+ e^-$  10.5–11 GeV

<sup>30</sup> ALEXANDER 90B cannot exclude extra  $\pi^0$ 's in the final state. See nearby data blocks for more detailed results.



$\Gamma(\bar{K}^*(892)^0 \pi^- e^+ \nu_e) / \Gamma(K^*(892)^- e^+ \nu_e)$   $\Gamma_{15} / \Gamma_{19}$

Unseen decay modes of the  $\bar{K}^*(892)^0$  are included.

VALUE	CL%	DOCUMENT ID	TECN	COMMENT
-------	-----	-------------	------	---------

• • • We do not use the following data for averages, fits, limits, etc. • • •

<0.64	90	<sup>31</sup> CRAWFORD	91B CLEO	$e^+ e^- \approx 10.5$ GeV
-------	----	------------------------	----------	----------------------------

<sup>31</sup> The limit on  $(\bar{K}^*(892)\pi)^- \mu^+ \nu_\mu$  below is much stronger.

$\Gamma(K^- \pi^+ \pi^- \mu^+ \nu_\mu) / \Gamma(K^- \mu^+ \nu_\mu)$   $\Gamma_{16} / \Gamma_{10}$

VALUE	CL%	DOCUMENT ID	TECN	COMMENT
-------	-----	-------------	------	---------

<0.037	90	KODAMA	93B E653	$\pi^-$ emulsion 600 GeV
--------	----	--------	----------	--------------------------

$\Gamma((\bar{K}^*(892)\pi)^- \mu^+ \nu_\mu) / \Gamma(K^- \mu^+ \nu_\mu)$   $\Gamma_{17} / \Gamma_{10}$

VALUE	CL%	DOCUMENT ID	TECN	COMMENT
-------	-----	-------------	------	---------

<0.043	90	<sup>32</sup> KODAMA	93B E653	$\pi^-$ emulsion 600 GeV
--------	----	----------------------	----------	--------------------------

<sup>32</sup> KODAMA 93B searched in  $K^- \pi^+ \pi^- \mu^+ \nu_\mu$ , but the limit includes other  $(\bar{K}^*(892)\pi)^-$  charge states.

$\Gamma(\pi^- e^+ \nu_e) / \Gamma_{\text{total}}$   $\Gamma_{18} / \Gamma$

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
-------	------	-------------	------	---------

**0.0036 ± 0.0006 OUR FIT**

<b>0.0039<sup>+0.0023</sup><sub>-0.0011</sub> ± 0.0004</b>	7	<sup>33</sup> ADLER	89 MRK3	$e^+ e^-$ 3.77 GeV
--	---	---------------------	---------	--------------------

<sup>33</sup> This result of ADLER 89 gives  $|\frac{V_{cd}}{V_{cs}} \cdot \frac{f_+^\pi(0)}{f_+^K(0)}|^2 = 0.057^{+0.038}_{-0.015} \pm 0.005$ .

$\Gamma(\pi^- e^+ \nu_e) / \Gamma(K^- e^+ \nu_e)$   $\Gamma_{18} / \Gamma_9$

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
-------	------	-------------	------	---------

**0.101 ± 0.017 OUR FIT**

**0.101 ± 0.018 OUR AVERAGE**

0.101 ± 0.020 ± 0.003	91	<sup>34</sup> FRABETTI	96B E687	$\gamma$ Be, $\bar{E}_\gamma \approx 200$ GeV
-----------------------	----	------------------------	----------	---

0.103 ± 0.039 ± 0.013	87	<sup>35</sup> BUTLER	95 CLE2	< 0.156 (90% CL)
-----------------------	----	----------------------	---------	------------------

<sup>34</sup> FRABETTI 96B uses both  $e$  and  $\mu$  events, and makes a small correction to the  $\mu$  events to make them effectively  $e$  events. This result gives  $|\frac{V_{cd}}{V_{cs}} \cdot \frac{f_+^\pi(0)}{f_+^K(0)}|^2 = 0.050 \pm 0.011 \pm 0.002$ .

<sup>35</sup> BUTLER 95 has  $87 \pm 33 \pi^- e^+ \nu_e$  events. The result gives  $|\frac{V_{cd}}{V_{cs}} \cdot \frac{f_+^\pi(0)}{f_+^K(0)}|^2 = 0.052 \pm 0.020 \pm 0.007$ .

————— Hadronic modes with a  $\bar{K}$  or  $\bar{K}K\bar{K}$  —————

$\Gamma(K^- \pi^+) / \Gamma_{\text{total}}$   $\Gamma_{20} / \Gamma$

We list measurements *before* radiative corrections are made.

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
-------	------	-------------	------	---------

**0.0380 ± 0.0009 OUR FIT**

**0.0385 ± 0.0009 OUR AVERAGE**

0.0382 ± 0.0007 ± 0.0012		<sup>36</sup> ARTUSO	98 CLE2	CLEO average
--------------------------	--	----------------------	---------	--------------

0.0390 ± 0.0009 ± 0.0012	5392	<sup>37</sup> BARATE	97C ALEP	From $Z$ decays
--------------------------	------	----------------------	----------	-----------------

0.045 ± 0.006 ± 0.004		<sup>38</sup> ALBRECHT	94 ARG	$e^+ e^- \approx \gamma(4S)$
-----------------------	--	------------------------	--------	------------------------------

0.0341 ± 0.0012 ± 0.0028	1173	<sup>37</sup> ALBRECHT	94F ARG	$e^+ e^- \approx \gamma(4S)$
--------------------------	------	------------------------	---------	------------------------------

0.0362 ± 0.0034 ± 0.0044		37 DECAMP	91J ALEP	From Z decays
0.045 ± 0.008 ± 0.005	56	37 ABACHI	88 HRS	$e^+e^-$ 29 GeV
0.042 ± 0.004 ± 0.004	930	ADLER	88C MRK3	$e^+e^-$ 3.77 GeV
0.041 ± 0.006	263	39 SCHINDLER	81 MRK2	$e^+e^-$ 3.771 GeV
0.043 ± 0.010	130	40 PERUZZI	77 MRK1	$e^+e^-$ 3.77 GeV

• • • We do not use the following data for averages, fits, limits, etc. • • •

0.0381 ± 0.0015 ± 0.0016	1165	41 ARTUSO	98 CLE2	$e^+e^-$ at $\Upsilon(4S)$
0.0369 ± 0.0011 ± 0.0016		42 COAN	98 CLE2	
0.0391 ± 0.0008 ± 0.0017	4208	37,43 AKERIB	93 CLE2	$e^+e^- \approx \Upsilon(4S)$

<sup>36</sup> This combines the CLEO results of ARTUSO 98, COAN 98, and AKERIB 93.

<sup>37</sup> ABACHI 88, DECAMP 91J, AKERIB 93, ALBRECHT 94F, and BARATE 97C use  $D^*(2010)^+ \rightarrow D^0\pi^+$  decays. The  $\pi^+$  is both slow and of low  $p_T$  with respect to the event thrust axis or nearest jet ( $\approx D^{*+}$  direction). The excess number of such  $\pi^+$ 's over background gives the number of  $D^*(2010)^+ \rightarrow D^0\pi^+$  events, and the fraction with  $D^0 \rightarrow K^-\pi^+$  gives the  $D^0 \rightarrow K^-\pi^+$  branching fraction.

<sup>38</sup> ALBRECHT 94 uses  $D^0$  mesons from  $\overline{B}^0 \rightarrow D^{*+}\ell^-\overline{\nu}_\ell$  decays. This is a different set of events than used by ALBRECHT 94F.

<sup>39</sup> SCHINDLER 81 (MARK-2) measures  $\sigma(e^+e^- \rightarrow \psi(3770)) \times$  branching fraction to be  $0.24 \pm 0.02$  nb. We use the MARK-3 (ADLER 88C) value of  $\sigma = 5.8 \pm 0.5 \pm 0.6$  nb.

<sup>40</sup> PERUZZI 77 (MARK-1) measures  $\sigma(e^+e^- \rightarrow \psi(3770)) \times$  branching fraction to be  $0.25 \pm 0.05$  nb. We use the MARK-3 (ADLER 88C) value of  $\sigma = 5.8 \pm 0.5 \pm 0.6$  nb.

<sup>41</sup> ARTUSO 98, following ALBRECHT 94, uses  $D^0$  mesons from  $\overline{B}^0 \rightarrow D^*(2010)^+X\ell^-\overline{\nu}_\ell$  decays. Our average uses the CLEO average of this value with the values of COAN 98 and AKERIB 93.

<sup>42</sup> COAN 98 assumes that  $\Gamma(B \rightarrow \overline{D}X\ell^+\nu)/\Gamma(B \rightarrow X\ell^+\nu) = 1.0 - 3|V_{ub}/V_{cb}|^2 - 0.010 \pm 0.005$ , the last term accounting for  $\overline{B} \rightarrow D_s^+KX\ell^-\overline{\nu}$ . COAN 98 is included in the CLEO average in ARTUSO 98.

<sup>43</sup> This AKERIB 93 value does not include radiative corrections; with them, the value is  $0.0395 \pm 0.0008 \pm 0.0017$ . AKERIB 93 is included in the CLEO average in ARTUSO 98.

$\Gamma(\overline{K}^0\pi^0)/\Gamma(K^-\pi^+)$   $\Gamma_{21}/\Gamma_{20}$

<u>VALUE</u>	<u>EVTS</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
<b>0.60 ± 0.06 OUR FIT</b>				
<b>1.36 ± 0.23 ± 0.22</b>	119	ANJOS	92B E691	$\gamma$ Be 80–240 GeV

$\Gamma(\overline{K}^0\pi^0)/\Gamma(\overline{K}^0\pi^+\pi^-)$   $\Gamma_{21}/\Gamma_{22}$

<u>VALUE</u>	<u>EVTS</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
<b>0.385 ± 0.031 OUR FIT</b>				
<b>0.378 ± 0.033 OUR AVERAGE</b>				
0.44 ± 0.02 ± 0.05	1942	PROCARIO	93B CLE2	$e^+e^-$ 10.36–10.7 GeV
0.34 ± 0.04 ± 0.02	92	44 ALBRECHT	92P ARG	$e^+e^- \approx 10$ GeV
0.36 ± 0.04 ± 0.08	104	KINOSHITA	91 CLEO	$e^+e^- \sim 10.7$ GeV

<sup>44</sup> This value is calculated from numbers in Table 1 of ALBRECHT 92P.

$\Gamma(\overline{K}^0 \pi^+ \pi^-) / \Gamma_{\text{total}}$   $\Gamma_{22} / \Gamma$

<u>VALUE</u>	<u>EVTS</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
<b>0.0592 ± 0.0035 OUR FIT</b>				Error includes scale factor of 1.1.
<b>0.055 ± 0.005 OUR AVERAGE</b>				
0.0503 ± 0.0039 ± 0.0049	284	<sup>45</sup> ALBRECHT	94F ARG	$e^+ e^- \approx \Upsilon(4S)$
0.064 ± 0.005 ± 0.010		ADLER	87 MRK3	$e^+ e^-$ 3.77 GeV
0.052 ± 0.016	32	<sup>46</sup> SCHINDLER	81 MRK2	$e^+ e^-$ 3.771 GeV
0.079 ± 0.023	28	<sup>47</sup> PERUZZI	77 MRK1	$e^+ e^-$ 3.77 GeV

<sup>45</sup> See the footnote on the ALBRECHT 94F measurement of  $\Gamma(K^- \pi^+) / \Gamma_{\text{total}}$  for the method used.

<sup>46</sup> SCHINDLER 81 (MARK-2) measures  $\sigma(e^+ e^- \rightarrow \psi(3770)) \times$  branching fraction to be  $0.30 \pm 0.08$  nb. We use the MARK-3 (ADLER 88C) value of  $\sigma = 5.8 \pm 0.5 \pm 0.6$  nb.

<sup>47</sup> PERUZZI 77 (MARK-1) measures  $\sigma(e^+ e^- \rightarrow \psi(3770)) \times$  branching fraction to be  $0.46 \pm 0.12$  nb. We use the MARK-3 (ADLER 88C) value of  $\sigma = 5.8 \pm 0.5 \pm 0.6$  nb.

$\Gamma(\overline{K}^0 \pi^+ \pi^-) / \Gamma(K^- \pi^+)$   $\Gamma_{22} / \Gamma_{20}$

<u>VALUE</u>	<u>EVTS</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
<b>1.56 ± 0.09 OUR FIT</b>				Error includes scale factor of 1.1.
<b>1.65 ± 0.17 OUR AVERAGE</b>				
1.61 ± 0.10 ± 0.15	856	FRABETTI	94J E687	$\gamma$ Be $\overline{E}_\gamma = 220$ GeV
1.7 ± 0.8	35	AVERY	80 SPEC	$\gamma N \rightarrow D^{*+}$
2.8 ± 1.0	116	PICCOLO	77 MRK1	$e^+ e^-$ 4.03, 4.41 GeV

$\Gamma(\overline{K}^0 \rho^0) / \Gamma(\overline{K}^0 \pi^+ \pi^-)$   $\Gamma_{23} / \Gamma_{22}$

<u>VALUE</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
<b>0.25 ± 0.05 OUR AVERAGE</b>			Error includes scale factor of 1.5.
0.350 ± 0.028 ± 0.067	FRABETTI	94G E687	$\gamma$ Be, $\overline{E}_\gamma \approx 220$ GeV
0.227 ± 0.032 ± 0.009	ALBRECHT	93D ARG	$e^+ e^- \approx 10$ GeV
• • • We do not use the following data for averages, fits, limits, etc. • • •			
0.215 ± 0.051 ± 0.037	ANJOS	93 E691	$\gamma$ Be 90–260 GeV
0.20 ± 0.06 ± 0.03	FRABETTI	92B E687	$\gamma$ Be $\overline{E}_\gamma = 221$ GeV
0.12 ± 0.01 ± 0.07	ADLER	87 MRK3	$e^+ e^-$ 3.77 GeV

$\Gamma(\overline{K}^0 f_0(980) \times B(f_0 \rightarrow \pi^+ \pi^-)) / \Gamma(\overline{K}^0 \pi^+ \pi^-)$   $\Gamma_{24} / \Gamma_{22}$

This includes only  $\pi^+ \pi^-$  decays of the  $f_0(980)$ , because branching fractions of this resonance are not known.

<u>VALUE</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
<b>0.054 ± 0.015 OUR AVERAGE</b>			
0.068 ± 0.016 ± 0.018	FRABETTI	94G E687	$\gamma$ Be, $\overline{E}_\gamma \approx 220$ GeV
0.046 ± 0.018 ± 0.006	ALBRECHT	93D ARG	$e^+ e^- \approx 10$ GeV

$\Gamma(\overline{K}^0 f_2(1270)) / \Gamma(\overline{K}^0 \pi^+ \pi^-)$   $\Gamma_{82} / \Gamma_{22}$

Unseen decay modes of the  $f_2(1270)$  are included.

<u>VALUE</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
<b>0.076 ± 0.028 OUR AVERAGE</b>			
0.065 ± 0.025 ± 0.030	FRABETTI	94G E687	$\gamma$ Be, $\overline{E}_\gamma \approx 220$ GeV
0.088 ± 0.037 ± 0.014	ALBRECHT	93D ARG	$e^+ e^- \approx 10$ GeV

$\Gamma(\bar{K}^0 f_0(1370) \times B(f_0 \rightarrow \pi^+ \pi^-))/\Gamma(\bar{K}^0 \pi^+ \pi^-)$   $\Gamma_{26}/\Gamma_{22}$

This includes only  $\pi^+ \pi^-$  decays of the  $f_0(1370)$ , because branching fractions of this resonance are not known.

VALUE	DOCUMENT ID	TECN	COMMENT
<b>0.080 ± 0.024 OUR AVERAGE</b>			
0.077 ± 0.022 ± 0.031	FRABETTI	94G E687	$\gamma$ Be, $\bar{E}_\gamma \approx 220$ GeV
0.082 ± 0.028 ± 0.013	ALBRECHT	93D ARG	$e^+ e^- \approx 10$ GeV

$\Gamma(K^*(892)^- \pi^+)/\Gamma(\bar{K}^0 \pi^+ \pi^-)$   $\Gamma_{84}/\Gamma_{22}$

Unseen decay modes of the  $K^*(892)^-$  are included.

VALUE	EVTs	DOCUMENT ID	TECN	COMMENT
<b>1.01 ± 0.06 OUR FIT</b>				Error includes scale factor of 1.4.
<b>1.00 ± 0.07 OUR AVERAGE</b>				Error includes scale factor of 1.4.
0.938 ± 0.054 ± 0.038		FRABETTI	94G E687	$\gamma$ Be, $\bar{E}_\gamma \approx 220$ GeV
1.08 ± 0.063 ± 0.045		ALBRECHT	93D ARG	$e^+ e^- \approx 10$ GeV
• • • We do not use the following data for averages, fits, limits, etc. • • •				
0.720 ± 0.145 ± 0.185		ANJOS	93 E691	$\gamma$ Be 90–260 GeV
0.96 ± 0.12 ± 0.075		FRABETTI	92B E687	$\gamma$ Be $\bar{E}_\gamma = 221$ GeV
0.84 ± 0.06 ± 0.08		ADLER	87 MRK3	$e^+ e^- 3.77$ GeV
1.05 <sup>+0.23</sup> <sub>-0.26</sub> <sup>+0.07</sup> <sub>-0.09</sub>	25	SCHINDLER	81 MRK2	$e^+ e^- 3.771$ GeV

$\Gamma(K_0^*(1430)^- \pi^+)/\Gamma(\bar{K}^0 \pi^+ \pi^-)$   $\Gamma_{106}/\Gamma_{22}$

Unseen decay modes of the  $\bar{K}_0^*(1430)^-$  are included.

VALUE	DOCUMENT ID	TECN	COMMENT
<b>0.20 ± 0.04 OUR FIT</b>			
<b>0.19 ± 0.05 OUR AVERAGE</b>			
0.176 ± 0.044 ± 0.047	FRABETTI	94G E687	$\gamma$ Be, $\bar{E}_\gamma \approx 220$ GeV
0.208 ± 0.055 ± 0.034	ALBRECHT	93D ARG	$e^+ e^- \approx 10$ GeV

$\Gamma(K_2^*(1430)^- \pi^+)/\Gamma(\bar{K}^0 \pi^+ \pi^-)$   $\Gamma_{108}/\Gamma_{22}$

Unseen decay modes of the  $\bar{K}_2^*(1430)^-$  are included.

VALUE	CL%	DOCUMENT ID	TECN	COMMENT
<b>&lt;0.15</b>	90	ALBRECHT	93D ARG	$e^+ e^- \approx 10$ GeV

$\Gamma(\bar{K}^0 \pi^+ \pi^- \text{ nonresonant})/\Gamma(\bar{K}^0 \pi^+ \pi^-)$   $\Gamma_{29}/\Gamma_{22}$

Neither FRABETTI 94G nor ALBRECHT 93D, the most detailed analyses so far, see evidence for a nonresonant component.

VALUE	DOCUMENT ID	TECN	COMMENT
• • • We do not use the following data for averages, fits, limits, etc. • • •			
0.263 ± 0.024 ± 0.041	ANJOS	93 E691	$\gamma$ Be 90–260 GeV
0.26 ± 0.08 ± 0.05	FRABETTI	92B E687	$\gamma$ Be $\bar{E}_\gamma = 221$ GeV
0.33 ± 0.05 ± 0.10	ADLER	87 MRK3	$e^+ e^- 3.77$ GeV

$\Gamma(K^- \pi^+ \pi^0)/\Gamma_{\text{total}}$   $\Gamma_{30}/\Gamma$

VALUE                      EVTS                      DOCUMENT ID                      TECN                      COMMENT

**0.131±0.009 OUR FIT**      Error includes scale factor of 1.3.

**0.131±0.016 OUR AVERAGE**

0.133±0.012±0.013	931	ADLER	88C MRK3	$e^+ e^-$ 3.77 GeV
0.117±0.043	37	<sup>48</sup> SCHINDLER	81 MRK2	$e^+ e^-$ 3.771 GeV

<sup>48</sup>SCHINDLER 81 (MARK-2) measures  $\sigma(e^+ e^- \rightarrow \psi(3770)) \times$  branching fraction to be  $0.68 \pm 0.23$  nb. We use the MARK-3 (ADLER 88C) value of  $\sigma = 5.8 \pm 0.5 \pm 0.6$  nb.

$\Gamma(K^- \pi^+ \pi^0)/\Gamma(K^- \pi^+)$   $\Gamma_{30}/\Gamma_{20}$

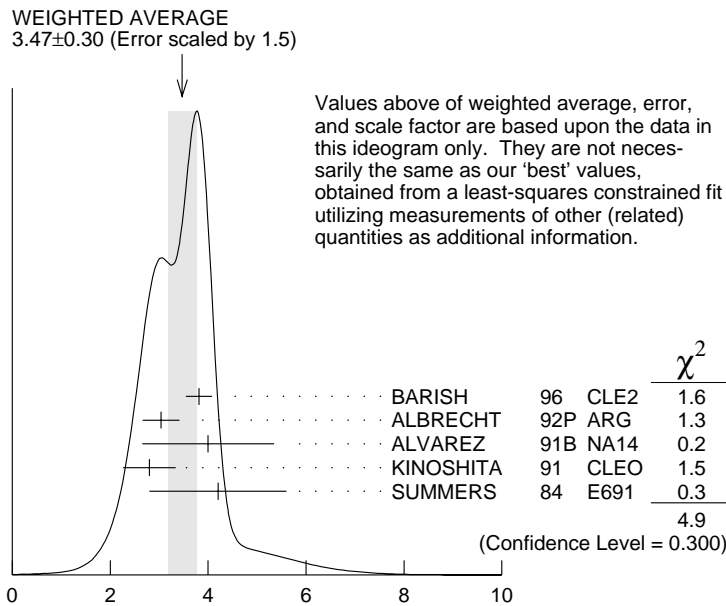
VALUE                      EVTS                      DOCUMENT ID                      TECN                      COMMENT

**3.44±0.22 OUR FIT**      Error includes scale factor of 1.3.

**3.47±0.30 OUR AVERAGE**      Error includes scale factor of 1.5. See the ideogram below.

3.81±0.07±0.26	10k	BARISH	96 CLE2	$e^+ e^- \approx \Upsilon(4S)$
3.04±0.16±0.34	931	<sup>49</sup> ALBRECHT	92P ARG	$e^+ e^- \approx 10$ GeV
4.0 ±0.9 ±1.0	69	ALVAREZ	91B NA14	Photoproduction
2.8 ±0.14±0.52	1050	KINOSHITA	91 CLEO	$e^+ e^- \sim 10.7$ GeV
4.2 ±1.4	41	SUMMERS	84 E691	Photoproduction

<sup>49</sup>This value is calculated from numbers in Table 1 of ALBRECHT 92P.



$\Gamma(K^- \pi^+ \pi^0)/\Gamma(K^- \pi^+)$

$\Gamma(K^- \rho^+)/\Gamma(K^- \pi^+ \pi^0)$   $\Gamma_{31}/\Gamma_{30}$

VALUE                      EVTS                      DOCUMENT ID                      TECN                      COMMENT

**0.78 ±0.04 OUR AVERAGE**

0.788±0.019±0.048		KOPP	01 CLE2	$e^+ e^- \approx 10.6$ GeV
0.765±0.041±0.054		FRABETTI	94G E687	$\gamma$ Be, $\bar{E}_\gamma \approx 220$ GeV

• • • We do not use the following data for averages, fits, limits, etc. • • •

$0.647 \pm 0.039 \pm 0.150$	ANJOS	93	E691	$\gamma$ Be 90–260 GeV
$0.81 \pm 0.03 \pm 0.06$	ADLER	87	MRK3	$e^+e^-$ 3.77 GeV
$0.31 \begin{smallmatrix} +0.20 \\ -0.14 \end{smallmatrix}$	13	SUMMERS	84	E691 Photoproduction
$0.85 \begin{smallmatrix} +0.11 & +0.09 \\ -0.15 & -0.10 \end{smallmatrix}$	31	SCHINDLER	81	MRK2 $e^+e^-$ 3.771 GeV

$\Gamma(K^-\rho(1700)^+ \times B(\rho(1700)^+ \rightarrow \pi^+\pi^0))/\Gamma(K^-\pi^+\pi^0)$   $\Gamma_{32}/\Gamma_{30}$

This only includes  $\pi^+\pi^0$  decays of the  $\rho(1700)^+$ , because branching fractions of this resonance are not known.

<u>VALUE</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
<b><math>0.057 \pm 0.008 \pm 0.009</math></b>	KOPP	01	CLE2 $e^+e^- \approx 10.6$ GeV

$\Gamma(K^*(892)^-\pi^+)/\Gamma(K^-\pi^+\pi^0)$   $\Gamma_{84}/\Gamma_{30}$

Unseen decay modes of the  $K^*(892)^-$  are included.

<u>VALUE</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
<b><math>0.46 \pm 0.04</math> OUR FIT</b>	Error includes scale factor of 1.2.		

**$0.48 \begin{smallmatrix} +0.08 \\ -0.04 \end{smallmatrix}$  OUR AVERAGE**

$0.483 \pm 0.021 \begin{smallmatrix} +0.081 \\ -0.032 \end{smallmatrix}$	KOPP	01	CLE2	$e^+e^- \approx 10.6$ GeV
$0.444 \pm 0.084 \pm 0.147$	FRABETTI	94G	E687	$\gamma$ Be, $\bar{E}_\gamma \approx 220$ GeV

• • • We do not use the following data for averages, fits, limits, etc. • • •

$0.252 \pm 0.033 \pm 0.035$	ANJOS	93	E691	$\gamma$ Be 90–260 GeV
$0.36 \pm 0.06 \pm 0.09$	ADLER	87	MRK3	$e^+e^-$ 3.77 GeV

$\Gamma(\bar{K}^*(892)^0\pi^0)/\Gamma(K^-\pi^+\pi^0)$   $\Gamma_{85}/\Gamma_{30}$

Unseen decay modes of the  $\bar{K}^*(892)^0$  are included.

<u>VALUE</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
<b><math>0.214 \pm 0.026</math> OUR FIT</b>	Error includes scale factor of 1.1.		

**$0.204 \pm 0.025$  OUR AVERAGE**

$0.191 \pm 0.014 \pm 0.024$	KOPP	01	CLE2	$e^+e^- \approx 10.6$ GeV
$0.248 \pm 0.047 \pm 0.023$	FRABETTI	94G	E687	$\gamma$ Be, $\bar{E}_\gamma \approx 220$ GeV

• • • We do not use the following data for averages, fits, limits, etc. • • •

$0.213 \pm 0.027 \pm 0.035$	ANJOS	93	E691	$\gamma$ Be 90–260 GeV
$0.20 \pm 0.03 \pm 0.05$	ADLER	87	MRK3	$e^+e^-$ 3.77 GeV

$\Gamma(K_0^*(1430)^-\pi^+)/\Gamma(K^-\pi^+\pi^0)$   $\Gamma_{106}/\Gamma_{30}$

Unseen decay modes of the  $K_0^*(1430)^-$  are included.

<u>VALUE</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
<b><math>0.090 \pm 0.020</math> OUR FIT</b>			

<b><math>0.107 \pm 0.019 \pm 0.045</math></b>	KOPP	01	CLE2	$e^+e^- \approx 10.6$ GeV
---	------	----	------	---------------------------

$\Gamma(\bar{K}_0^*(1430)^0\pi^0)/\Gamma(K^-\pi^+\pi^0)$   $\Gamma_{107}/\Gamma_{30}$

Unseen decay modes of the  $\bar{K}_0^*(1430)^0$  are included.

<u>VALUE</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>	
<b><math>0.066 \pm 0.010 \begin{smallmatrix} +0.051 \\ -0.014 \end{smallmatrix}</math></b>	KOPP	01	CLE2	$e^+e^- \approx 10.6$ GeV

$\Gamma(K^*(1680)^-\pi^+)/\Gamma(K^-\pi^+\pi^0)$   $\Gamma_{110}/\Gamma_{30}$

Unseen decay modes of the  $K^*(1680)^-$  are included.

VALUE	DOCUMENT ID	TECN	COMMENT
<b>0.101±0.023±0.033</b>	KOPP	01 CLE2	$e^+e^- \approx 10.6$ GeV

$\Gamma(K^-\pi^+\pi^0 \text{ nonresonant})/\Gamma(K^-\pi^+\pi^0)$   $\Gamma_{38}/\Gamma_{30}$

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
<b>0.080<sup>+0.038</sup><sub>-0.014</sub> OUR AVERAGE</b>				

0.075±0.009 <sup>+0.056</sup> <sub>-0.011</sub>		KOPP	01 CLE2	$e^+e^- \approx 10.6$ GeV
0.101±0.033±0.040		FRABETTI	94G E687	$\gamma\text{Be}, \bar{E}_\gamma \approx 220$ GeV

• • • We do not use the following data for averages, fits, limits, etc. • • •

0.036±0.004±0.018		ANJOS	93 E691	$\gamma\text{Be}$ 90–260 GeV
0.09 ±0.02 ±0.04		ADLER	87 MRK3	$e^+e^-$ 3.77 GeV
0.51 ±0.22	21	SUMMERS	84 E691	Photoproduction

$\Gamma(\bar{K}^*(892)^0\pi^0)/\Gamma(\bar{K}^0\pi^0)$   $\Gamma_{85}/\Gamma_{21}$

Unseen decay modes of the  $\bar{K}^*(892)^0$  are included.

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
<b>1.23±0.20 OUR FIT</b>				Error includes scale factor of 1.2.
<b>1.65<sup>+0.39</sup><sub>-0.31</sub>±0.20</b>	122	PROCARIO	93B CLE2	$\bar{K}^0\pi^0\pi^0$ Dalitz plot

$\Gamma(\bar{K}_2^*(1430)^0\pi^0)/\Gamma(\bar{K}^*(892)^0\pi^0)$   $\Gamma_{109}/\Gamma_{85}$

Unseen decay modes of the  $\bar{K}_2^*(1430)^0$  and  $\bar{K}^*(892)^0$  are included.

VALUE	CL%	DOCUMENT ID	TECN	COMMENT
<b>&lt;0.12</b>	90	PROCARIO	93B CLE2	$\bar{K}^0\pi^0\pi^0$ Dalitz plot

$\Gamma(\bar{K}^0\pi^0\pi^0 \text{ nonresonant})/\Gamma(\bar{K}^0\pi^0)$   $\Gamma_{41}/\Gamma_{21}$

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
<b>0.37±0.08±0.04</b>	76	PROCARIO	93B CLE2	$\bar{K}^0\pi^0\pi^0$ Dalitz plot

$\Gamma(K^-\pi^+\pi^+\pi^-)/\Gamma_{\text{total}}$   $\Gamma_{42}/\Gamma$

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
<b>0.0746±0.0031 OUR FIT</b>				

**0.075 ±0.006 OUR AVERAGE** Error includes scale factor of 1.3. See the ideogram below.

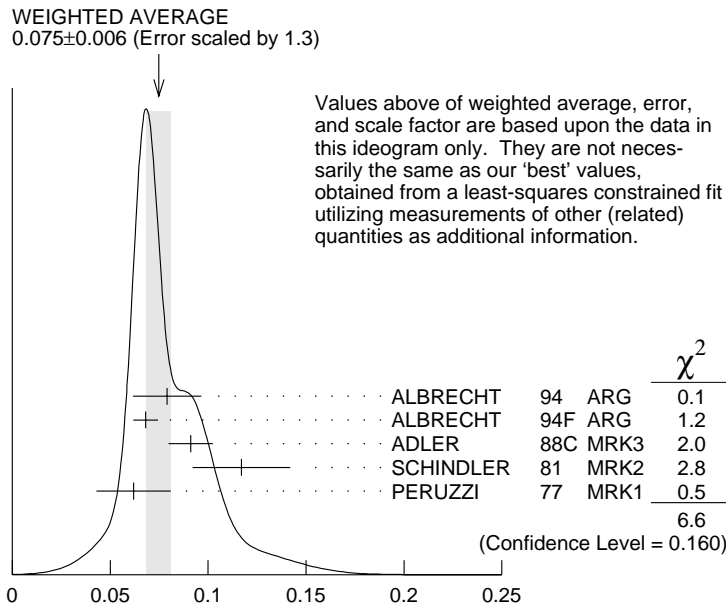
0.079 ±0.015 ±0.009		50 ALBRECHT	94 ARG	$e^+e^- \approx \Upsilon(4S)$
0.0680±0.0027±0.0057	1430	51 ALBRECHT	94F ARG	$e^+e^- \approx \Upsilon(4S)$
0.091 ±0.008 ±0.008	992	ADLER	88C MRK3	$e^+e^-$ 3.77 GeV
0.117 ±0.025	185	52 SCHINDLER	81 MRK2	$e^+e^-$ 3.771 GeV
0.062 ±0.019	44	53 PERUZZI	77 MRK1	$e^+e^-$ 3.77 GeV

<sup>50</sup> ALBRECHT 94 uses  $D^0$  mesons from  $\bar{B}^0 \rightarrow D^{*+}\ell^-\bar{\nu}_\ell$  decays. This is a different set of events than used by ALBRECHT 94F.

<sup>51</sup> See the footnote on the ALBRECHT 94F measurement of  $\Gamma(K^-\pi^+)/\Gamma_{\text{total}}$  for the method used.

<sup>52</sup> SCHINDLER 81 (MARK-2) measures  $\sigma(e^+e^- \rightarrow \psi(3770)) \times$  branching fraction to be  $0.68 \pm 0.11$  nb. We use the MARK-3 (ADLER 88C) value of  $\sigma = 5.8 \pm 0.5 \pm 0.6$  nb.

<sup>53</sup> PERUZZI 77 (MARK-1) measures  $\sigma(e^+e^- \rightarrow \psi(3770)) \times$  branching fraction to be  $0.36 \pm 0.10$  nb. We use the MARK-3 (ADLER 88C) value of  $\sigma = 5.8 \pm 0.5 \pm 0.6$  nb.



$$\Gamma(K^- \pi^+ \pi^+ \pi^-) / \Gamma_{\text{total}}$$

$$\Gamma(K^- \pi^+ \pi^+ \pi^-) / \Gamma(K^- \pi^+)$$

$$\Gamma_{42} / \Gamma_{20}$$

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
<b><math>1.96 \pm 0.08</math></b>				<b>OUR FIT</b>
<b><math>1.97 \pm 0.09</math></b>				<b>OUR AVERAGE</b>
$1.94 \pm 0.07^{+0.09}_{-0.11}$		JUN	00 SELX	$\Sigma^-$ nucleus, 600 GeV
$1.7 \pm 0.2 \pm 0.2$	1745	ANJOS	92C E691	$\gamma$ Be 90–260 GeV
$1.90 \pm 0.25 \pm 0.20$	337	ALVAREZ	91B NA14	Photoproduction
$2.12 \pm 0.16 \pm 0.09$		BORTOLETTO88	CLEO	$e^+ e^-$ 10.55 GeV
$2.0 \pm 0.9$	48	BAILEY	86 ACCM	$\pi^-$ Be fixed target
$2.17 \pm 0.28 \pm 0.23$		ALBRECHT	85F ARG	$e^+ e^-$ 10 GeV
$2.0 \pm 1.0$	10	BAILEY	83B SPEC	$\pi^-$ Be $\rightarrow D^0$
$2.2 \pm 0.8$	214	PICCOLO	77 MRK1	$e^+ e^-$ 4.03, 4.41 GeV

$$\Gamma(K^- \pi^+ \rho^0 \text{ total}) / \Gamma(K^- \pi^+ \pi^+ \pi^-)$$

$$\Gamma_{43} / \Gamma_{42}$$

This includes  $K^- a_1(1260)^+$ ,  $\bar{K}^*(892)^0 \rho^0$ , etc. The next entry gives the specifically 3-body fraction. We rely on the MARK III and E691 full amplitude analyses of the  $K^- \pi^+ \pi^+ \pi^-$  channel for values of the resonant substructure.

VALUE	DOCUMENT ID	TECN	COMMENT
<b><math>0.835 \pm 0.035</math></b>			<b>OUR AVERAGE</b>
$0.80 \pm 0.03 \pm 0.05$	ANJOS	92C E691	$\gamma$ Be 90–260 GeV
$0.855 \pm 0.032 \pm 0.030$	COFFMAN	92B MRK3	$e^+ e^-$ 3.77 GeV
• • • We do not use the following data for averages, fits, limits, etc. • • •			
$0.98 \pm 0.12 \pm 0.10$	ALVAREZ	91B NA14	Photoproduction



$\Gamma(K^- \pi^+ \rho^0 \text{ 3-body})/\Gamma(K^- \pi^+ \pi^+ \pi^-)$   $\Gamma_{44}/\Gamma_{42}$

We rely on the MARK III and E691 full amplitude analyses of the  $K^- \pi^+ \pi^+ \pi^-$  channel for values of the resonant substructure.

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
<b>0.063 ± 0.028 OUR AVERAGE</b>				
0.05 ± 0.03 ± 0.02		ANJOS	92C E691	$\gamma$ Be 90–260 GeV
0.084 ± 0.022 ± 0.04		COFFMAN	92B MRK3	$e^+ e^-$ 3.77 GeV
• • • We do not use the following data for averages, fits, limits, etc. • • •				
0.77 ± 0.06 ± 0.06		<sup>54</sup> ALVAREZ	91B NA14	Photoproduction
0.85 <sup>+0.11</sup> <sub>-0.22</sub>	180	PICCOLO	77 MRK1	$e^+ e^-$ 4.03, 4.41 GeV

<sup>54</sup>This value is for  $\rho^0 (K^- \pi^+)$ -nonresonant. ALVAREZ 91B cannot determine what fraction of this is  $K^- a_1(1260)^+$ .

$\Gamma(\bar{K}^*(892)^0 \rho^0)/\Gamma(K^- \pi^+ \pi^+ \pi^-)$   $\Gamma_{90}/\Gamma_{42}$

Unseen decay modes of the  $\bar{K}^*(892)^0$  are included. We rely on the MARK III and E691 full amplitude analyses of the  $K^- \pi^+ \pi^+ \pi^-$  channel for values of the resonant substructure.

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
<b>0.195 ± 0.03 ± 0.03</b>				
		ANJOS	92C E691	$\gamma$ Be 90–260 GeV
• • • We do not use the following data for averages, fits, limits, etc. • • •				
0.34 ± 0.09 ± 0.09		ALVAREZ	91B NA14	Photoproduction
0.75 ± 0.3	5	BAILEY	83B SPEC	$\pi$ Be $\rightarrow D^0$
0.15 <sup>+0.16</sup> <sub>-0.15</sub>	20	PICCOLO	77 MRK1	$e^+ e^-$ 4.03, 4.41 GeV

$\Gamma(\bar{K}^*(892)^0 \rho^0 \text{ transverse})/\Gamma(K^- \pi^+ \pi^+ \pi^-)$   $\Gamma_{91}/\Gamma_{42}$

Unseen decay modes of the  $\bar{K}^*(892)^0$  are included.

VALUE	DOCUMENT ID	TECN	COMMENT
<b>0.21 ± 0.07 OUR FIT</b>			
<b>0.213 ± 0.024 ± 0.075</b>	COFFMAN	92B MRK3	$e^+ e^-$ 3.77 GeV

$\Gamma(\bar{K}^*(892)^0 \rho^0 \text{ S-wave})/\Gamma(K^- \pi^+ \pi^+ \pi^-)$   $\Gamma_{92}/\Gamma_{42}$

Unseen decay modes of the  $\bar{K}^*(892)^0$  are included.

VALUE	DOCUMENT ID	TECN	COMMENT
<b>0.375 ± 0.045 ± 0.06</b>	ANJOS	92C E691	$\gamma$ Be 90–260 GeV

$\Gamma(\bar{K}^*(892)^0 \rho^0 \text{ S-wave long.})/\Gamma_{\text{total}}$   $\Gamma_{93}/\Gamma$

Unseen decay modes of the  $\bar{K}^*(892)^0$  are included.

VALUE	CL%	DOCUMENT ID	TECN	COMMENT
<b>&lt;0.003</b>	90	COFFMAN	92B MRK3	$e^+ e^-$ 3.77 GeV

$\Gamma(\bar{K}^*(892)^0 \rho^0 \text{ P-wave})/\Gamma_{\text{total}}$   $\Gamma_{94}/\Gamma$

Unseen decay modes of the  $\bar{K}^*(892)^0$  are included.

VALUE	CL%	DOCUMENT ID	TECN	COMMENT
<b>&lt;0.003</b>	90	COFFMAN	92B MRK3	$e^+ e^-$ 3.77 GeV
• • • We do not use the following data for averages, fits, limits, etc. • • •				
<0.009	90	ANJOS	92C E691	$\gamma$ Be 90–260 GeV

$\Gamma(\bar{K}^*(892)^0 \rho^0 D\text{-wave})/\Gamma(K^- \pi^+ \pi^+ \pi^-)$   $\Gamma_{95}/\Gamma_{42}$

Unseen decay modes of the  $\bar{K}^*(892)^0$  are included.

<u>VALUE</u>	<u>CL%</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
<b>0.255 ± 0.045 ± 0.06</b>		ANJOS	92C E691	$\gamma$ Be 90–260 GeV

$\Gamma(K^- \pi^+ f_0(980))/\Gamma_{\text{total}}$   $\Gamma_{100}/\Gamma$

<u>VALUE</u>	<u>CL%</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
--------------	------------	--------------------	-------------	----------------

• • • We do not use the following data for averages, fits, limits, etc. • • •

<0.011	90	ANJOS	92C E691	$\gamma$ Be 90–260 GeV
--------	----	-------	----------	------------------------

$\Gamma(\bar{K}^*(892)^0 f_0(980))/\Gamma_{\text{total}}$   $\Gamma_{101}/\Gamma$

<u>VALUE</u>	<u>CL%</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
--------------	------------	--------------------	-------------	----------------

• • • We do not use the following data for averages, fits, limits, etc. • • •

<0.007	90	ANJOS	92C E691	$\gamma$ Be 90–260 GeV
--------	----	-------	----------	------------------------

$\Gamma(K^- a_1(1260)^+)/\Gamma(K^- \pi^+ \pi^+ \pi^-)$   $\Gamma_{80}/\Gamma_{42}$

Unseen decay modes of the  $a_1(1260)^+$  are included, assuming that the  $a_1(1260)^+$  decays entirely to  $\rho\pi$  [or at least to  $(\pi\pi)_{I=1}\pi$ ].

<u>VALUE</u>	<u>CL%</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
--------------	------------	--------------------	-------------	----------------

**0.97 ± 0.14 OUR AVERAGE**

0.94 ± 0.13 ± 0.20		ANJOS	92C E691	$\gamma$ Be 90–260 GeV
0.984 ± 0.048 ± 0.16		COFFMAN	92B MRK3	$e^+ e^-$ 3.77 GeV

$\Gamma(K^- a_2(1320)^+)/\Gamma_{\text{total}}$   $\Gamma_{83}/\Gamma$

Unseen decay modes of the  $a_2(1320)^+$  are included.

<u>VALUE</u>	<u>CL%</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
--------------	------------	--------------------	-------------	----------------

<b>&lt;0.002</b>	90	ANJOS	92C E691	$\gamma$ Be 90–260 GeV
------------------	----	-------	----------	------------------------

• • • We do not use the following data for averages, fits, limits, etc. • • •

<0.006	90	COFFMAN	92B MRK3	$e^+ e^-$ 3.77 GeV
--------	----	---------	----------	--------------------

$\Gamma(K_1(1270)^- \pi^+)/\Gamma(K^- \pi^+ \pi^+ \pi^-)$   $\Gamma_{102}/\Gamma_{42}$

Unseen decay modes of the  $K_1(1270)^-$  are included. The MARK3 and E691 experiments disagree considerably here.

<u>VALUE</u>	<u>CL%</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
--------------	------------	--------------------	-------------	----------------

**0.15 ± 0.04 OUR FIT**

<b>0.194 ± 0.056 ± 0.088</b>		COFFMAN	92B MRK3	$e^+ e^-$ 3.77 GeV
------------------------------	--	---------	----------	--------------------

• • • We do not use the following data for averages, fits, limits, etc. • • •

<0.013	90	ANJOS	92C E691	$\gamma$ Be 90–260 GeV
--------	----	-------	----------	------------------------

$\Gamma(K_1(1400)^- \pi^+)/\Gamma_{\text{total}}$   $\Gamma_{103}/\Gamma$

<u>VALUE</u>	<u>CL%</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
--------------	------------	--------------------	-------------	----------------

<b>&lt;0.012</b>	90	COFFMAN	92B MRK3	$e^+ e^-$ 3.77 GeV
------------------	----	---------	----------	--------------------

$\Gamma(K^*(1410)^- \pi^+)/\Gamma_{\text{total}}$   $\Gamma_{105}/\Gamma$

<u>VALUE</u>	<u>CL%</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
--------------	------------	--------------------	-------------	----------------

• • • We do not use the following data for averages, fits, limits, etc. • • •

<0.012	90	COFFMAN	92B MRK3	$e^+ e^-$ 3.77 GeV
--------	----	---------	----------	--------------------

$\Gamma(\bar{K}^*(892)^0 \pi^+ \pi^- \text{ total}) / \Gamma(K^- \pi^+ \pi^+ \pi^-)$   $\Gamma_{86} / \Gamma_{42}$

This includes  $\bar{K}^*(892)^0 \rho^0$ , etc. The next entry gives the specifically 3-body fraction.  
Unseen decay modes of the  $\bar{K}^*(892)^0$  are included.

VALUE	DOCUMENT ID	TECN	COMMENT
<b>0.30 ± 0.06 ± 0.03</b>	ANJOS	92C E691	$\gamma$ Be 90–260 GeV

$\Gamma(\bar{K}^*(892)^0 \pi^+ \pi^- \text{ 3-body}) / \Gamma(K^- \pi^+ \pi^+ \pi^-)$   $\Gamma_{87} / \Gamma_{42}$

Unseen decay modes of the  $\bar{K}^*(892)^0$  are included.

VALUE	DOCUMENT ID	TECN	COMMENT
<b>0.19 ± 0.04 OUR FIT</b>			
<b>0.18 ± 0.04 OUR AVERAGE</b>			
0.165 ± 0.03 ± 0.045	ANJOS	92C E691	$\gamma$ Be 90–260 GeV
0.210 ± 0.027 ± 0.06	COFFMAN	92B MRK3	$e^+ e^-$ 3.77 GeV

$\Gamma(K^- \pi^+ \pi^+ \pi^- \text{ nonresonant}) / \Gamma(K^- \pi^+ \pi^+ \pi^-)$   $\Gamma_{50} / \Gamma_{42}$

VALUE	DOCUMENT ID	TECN	COMMENT
<b>0.233 ± 0.032 OUR AVERAGE</b>			
0.23 ± 0.02 ± 0.03	ANJOS	92C E691	$\gamma$ Be 90–260 GeV
0.242 ± 0.025 ± 0.06	COFFMAN	92B MRK3	$e^+ e^-$ 3.77 GeV

$\Gamma(\bar{K}^0 \pi^+ \pi^- \pi^0) / \Gamma_{\text{total}}$   $\Gamma_{51} / \Gamma$

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
<b>0.108 ± 0.013 OUR FIT</b>				
<b>0.103 ± 0.022 ± 0.025</b>	140	COFFMAN	92B MRK3	$e^+ e^-$ 3.77 GeV

• • • We do not use the following data for averages, fits, limits, etc. • • •

0.134 <sup>+0.032</sup> <sub>-0.033</sub>		<sup>55</sup> BARLAG	92C ACCM	$\pi^-$ Cu 230 GeV
---	--	----------------------	----------	--------------------

<sup>55</sup> BARLAG 92C computes the branching fraction using topological normalization.

$\Gamma(\bar{K}^0 \pi^+ \pi^- \pi^0) / \Gamma(\bar{K}^0 \pi^+ \pi^-)$   $\Gamma_{51} / \Gamma_{22}$

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
<b>1.82 ± 0.20 OUR FIT</b>				
<b>1.86 ± 0.23 OUR AVERAGE</b>				
1.80 ± 0.20 ± 0.21	190	<sup>56</sup> ALBRECHT	92P ARG	$e^+ e^- \approx 10$ GeV
2.8 ± 0.8 ± 0.8	46	ANJOS	92C E691	$\gamma$ Be 90–260 GeV
1.85 ± 0.26 ± 0.30	158	KINOSHITA	91 CLEO	$e^+ e^- \sim 10.7$ GeV

<sup>56</sup> This value is calculated from numbers in Table 1 of ALBRECHT 92P.

$\Gamma(\bar{K}^0 \eta) / \Gamma(K^- \pi^+)$   $\Gamma_{74} / \Gamma_{20}$

Unseen decay modes of the  $\eta$  are included.

VALUE	CL%	DOCUMENT ID	TECN	COMMENT
<b>• • • We do not use the following data for averages, fits, limits, etc. • • •</b>				
<0.64	90	ALBRECHT	89D ARG	$e^+ e^-$ 10 GeV

$\Gamma(\bar{K}^0 \eta) / \Gamma(\bar{K}^0 \pi^0)$   $\Gamma_{74} / \Gamma_{21}$

Unseen decay modes of the  $\eta$  are included.

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
<b>0.33 ± 0.04 OUR FIT</b>				
<b>0.32 ± 0.04 ± 0.03</b>	225	PROCARIO	93B CLE2	$\eta \rightarrow \gamma\gamma$

$\Gamma(\bar{K}^0 \eta) / \Gamma(\bar{K}^0 \pi^+ \pi^-)$   $\Gamma_{74} / \Gamma_{22}$

Unseen decay modes of the  $\eta$  are included.

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
<b>0.129 ± 0.017 OUR FIT</b>				
<b>0.14 ± 0.02 ± 0.02</b>	80	PROCARIO	93B CLE2	$\eta \rightarrow \pi^+ \pi^- \pi^0$

$\Gamma(\bar{K}^0 \omega) / \Gamma(K^- \pi^+)$   $\Gamma_{77} / \Gamma_{20}$

Unseen decay modes of the  $\omega$  are included.

VALUE	DOCUMENT ID	TECN	COMMENT
<b>0.59 ± 0.10 OUR FIT</b>			
<b>1.00 ± 0.36 ± 0.20</b>	ALBRECHT	89D ARG	$e^+ e^-$ 10 GeV

$\Gamma(\bar{K}^0 \omega) / \Gamma(\bar{K}^0 \pi^+ \pi^-)$   $\Gamma_{77} / \Gamma_{22}$

Unseen decay modes of the  $\omega$  are included.

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
<b>0.38 ± 0.06 OUR FIT</b>				
<b>0.33 ± 0.09 OUR AVERAGE</b>				Error includes scale factor of 1.1.
0.29 ± 0.08 ± 0.05	16	<sup>57</sup> ALBRECHT	92P ARG	$e^+ e^- \approx 10$ GeV
0.54 ± 0.14 ± 0.16	40	KINOSHITA	91 CLEO	$e^+ e^- \sim 10.7$ GeV

<sup>57</sup> This value is calculated from numbers in Table 1 of ALBRECHT 92P.

$\Gamma(\bar{K}^0 \omega) / \Gamma(\bar{K}^0 \pi^+ \pi^- \pi^0)$   $\Gamma_{77} / \Gamma_{51}$

Unseen decay modes of the  $\omega$  are included.

VALUE	DOCUMENT ID	TECN	COMMENT
<b>0.21 ± 0.04 OUR FIT</b>			
<b>0.220 ± 0.048 ± 0.0116</b>	COFFMAN	92B MRK3	$e^+ e^-$ 3.77 GeV

$\Gamma(\bar{K}^0 \eta'(958)) / \Gamma(\bar{K}^0 \pi^+ \pi^-)$   $\Gamma_{78} / \Gamma_{22}$

Unseen decay modes of the  $\eta'(958)$  are included.

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
<b>0.32 ± 0.04 OUR AVERAGE</b>				
0.31 ± 0.02 ± 0.04	594	PROCARIO	93B CLE2	$\eta' \rightarrow \eta \pi^+ \pi^-, \rho^0 \gamma$
0.37 ± 0.13 ± 0.06	18	<sup>58</sup> ALBRECHT	92P ARG	$e^+ e^- \approx 10$ GeV

<sup>58</sup> This value is calculated from numbers in Table 1 of ALBRECHT 92P.

$\Gamma(K^*(892)^- \rho^+) / \Gamma(\bar{K}^0 \pi^+ \pi^- \pi^0)$   $\Gamma_{96} / \Gamma_{51}$

Unseen decay modes of the  $K^*(892)^-$  are included.

VALUE	DOCUMENT ID	TECN	COMMENT
<b>0.606 ± 0.188 ± 0.126</b>	COFFMAN	92B MRK3	$e^+ e^-$ 3.77 GeV

$\Gamma(K^*(892)^- \rho^+ \text{longitudinal}) / \Gamma(\bar{K}^0 \pi^+ \pi^- \pi^0)$   $\Gamma_{97} / \Gamma_{51}$

Unseen decay modes of the  $K^*(892)^-$  are included.

VALUE	DOCUMENT ID	TECN	COMMENT
<b>0.290 ± 0.111</b>	COFFMAN	92B MRK3	$e^+ e^-$ 3.77 GeV

$\Gamma(K^*(892)^- \rho^+ \text{transverse}) / \Gamma(\bar{K}^0 \pi^+ \pi^- \pi^0)$   $\Gamma_{98} / \Gamma_{51}$

Unseen decay modes of the  $K^*(892)^-$  are included.

VALUE	DOCUMENT ID	TECN	COMMENT
<b>0.317 ± 0.180</b>	COFFMAN	92B MRK3	$e^+ e^-$ 3.77 GeV

$\Gamma(K^*(892)^- \rho^+ P\text{-wave})/\Gamma_{\text{total}}$   $\Gamma_{99}/\Gamma$

Unseen decay modes of the  $K^*(892)^-$  are included.

VALUE	CL%	DOCUMENT ID	TECN	COMMENT
<b>&lt;0.015</b>	90	<sup>59</sup> COFFMAN	92B MRK3	$e^+ e^-$ 3.77 GeV

<sup>59</sup> Obtained using other  $\bar{K}^*(892)\rho$   $P$ -wave limits and isospin relations.

$\Gamma(\bar{K}^*(892)^0 \rho^0 \text{transverse})/\Gamma(\bar{K}^0 \pi^+ \pi^- \pi^0)$   $\Gamma_{91}/\Gamma_{51}$

Unseen decay modes of the  $\bar{K}^*(892)^0$  are included.

VALUE	DOCUMENT ID	TECN	COMMENT
<b>0.14 ± 0.05 OUR FIT</b>			
<b>0.126 ± 0.111</b>	COFFMAN	92B MRK3	$e^+ e^-$ 3.77 GeV

$\Gamma(\bar{K}^0 a_1(1260)^0)/\Gamma_{\text{total}}$   $\Gamma_{81}/\Gamma$

Unseen decay modes of the  $a_1(1260)^+$  are included, assuming that the  $a_1(1260)^+$  decays entirely to  $\rho\pi$  [or at least to  $(\pi\pi)_{J=1}\pi$ ].

VALUE	CL%	DOCUMENT ID	TECN	COMMENT
<b>&lt;0.019</b>	90	COFFMAN	92B MRK3	$e^+ e^-$ 3.77 GeV

$\Gamma(K_1(1270)^- \pi^+)/\Gamma(\bar{K}^0 \pi^+ \pi^- \pi^0)$   $\Gamma_{102}/\Gamma_{51}$

Unseen decay modes of the  $K_1(1270)^-$  are included.

VALUE	DOCUMENT ID	TECN	COMMENT
<b>0.105 ± 0.028 OUR FIT</b>			
<b>0.10 ± 0.03</b>	COFFMAN	92B MRK3	$e^+ e^-$ 3.77 GeV

$\Gamma(\bar{K}_1(1400)^0 \pi^0)/\Gamma_{\text{total}}$   $\Gamma_{104}/\Gamma$

VALUE	CL%	DOCUMENT ID	TECN	COMMENT
<b>&lt;0.037</b>	90	COFFMAN	92B MRK3	$e^+ e^-$ 3.77 GeV

$\Gamma(\bar{K}^*(892)^0 \pi^+ \pi^- \text{3-body})/\Gamma(\bar{K}^0 \pi^+ \pi^- \pi^0)$   $\Gamma_{87}/\Gamma_{51}$

Unseen decay modes of the  $\bar{K}^*(892)^0$  are included.

VALUE	DOCUMENT ID	TECN	COMMENT
<b>0.131 ± 0.035 OUR FIT</b>			Error includes scale factor of 1.1.
<b>0.191 ± 0.105</b>	COFFMAN	92B MRK3	$e^+ e^-$ 3.77 GeV

$\Gamma(\bar{K}^0 \pi^+ \pi^- \pi^0 \text{nonresonant})/\Gamma(\bar{K}^0 \pi^+ \pi^- \pi^0)$   $\Gamma_{58}/\Gamma_{51}$

VALUE	DOCUMENT ID	TECN	COMMENT
<b>0.210 ± 0.147 ± 0.150</b>	COFFMAN	92B MRK3	$e^+ e^-$ 3.77 GeV

$\Gamma(K^- \pi^+ \pi^0 \pi^0)/\Gamma_{\text{total}}$   $\Gamma_{59}/\Gamma$

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
-------	------	-------------	------	---------

• • • We do not use the following data for averages, fits, limits, etc. • • •

0.177 ± 0.029		<sup>60</sup> BARLAG	92C ACCM	$\pi^-$ Cu 230 GeV
0.149 ± 0.037 ± 0.030	24	<sup>61</sup> ADLER	88C MRK3	$e^+ e^-$ 3.77 GeV
0.209 <sup>+0.074</sup> <sub>-0.043</sub> ± 0.012	9	<sup>60</sup> AGUILAR-...	87F HYBR	$\pi p, p p$ 360, 400 GeV

<sup>60</sup> AGUILAR-BENITEZ 87F and BARLAG 92C compute the branching fraction using topological normalization. They do not distinguish the presence of a third  $\pi^0$ , and thus are not included in the average.

<sup>61</sup> ADLER 88C uses an absolute normalization method finding this decay channel opposite a detected  $\bar{D}^0 \rightarrow K^+ \pi^-$  in pure  $D\bar{D}$  events.

$\Gamma(K^- \pi^+ \pi^+ \pi^- \pi^0)/\Gamma(K^- \pi^+)$   $\Gamma_{60}/\Gamma_{20}$

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
<b>1.05±0.10 OUR FIT</b>				
<b>0.98±0.11±0.11</b>	225	<sup>62</sup> ALBRECHT	92P ARG	$e^+ e^- \approx 10$ GeV

<sup>62</sup> This value is calculated from numbers in Table 1 of ALBRECHT 92P.

$\Gamma(K^- \pi^+ \pi^+ \pi^- \pi^0)/\Gamma(K^- \pi^+ \pi^+ \pi^-)$   $\Gamma_{60}/\Gamma_{42}$

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
<b>0.54±0.05 OUR FIT</b>				
<b>0.56±0.07 OUR AVERAGE</b>				

0.55±0.07 <sup>+0.12</sup> <sub>-0.09</sub>	167	KINOSHITA	91 CLEO	$e^+ e^- \sim 10.7$ GeV
0.57±0.06±0.05	180	ANJOS	90D E691	Photoproduction

$\Gamma(\bar{K}^*(892)^0 \pi^+ \pi^- \pi^0)/\Gamma(K^- \pi^+ \pi^+ \pi^- \pi^0)$   $\Gamma_{111}/\Gamma_{60}$

Unseen decay modes of the  $\bar{K}^*(892)^0$  are included.

VALUE	DOCUMENT ID	TECN	COMMENT
<b>0.45±0.15±0.15</b>	ANJOS	90D E691	Photoproduction

$\Gamma(\bar{K}^*(892)^0 \eta)/\Gamma(K^- \pi^+)$   $\Gamma_{112}/\Gamma_{20}$

Unseen decay modes of the  $\bar{K}^*(892)^0$  and  $\eta$  are included.

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
<b>0.46±0.12 OUR FIT</b>				
<b>0.58±0.19<sup>+0.24</sup></b> <b>-0.28</b>	46	KINOSHITA	91 CLEO	$e^+ e^- \sim 10.7$ GeV

$\Gamma(\bar{K}^*(892)^0 \eta)/\Gamma(K^- \pi^+ \pi^0)$   $\Gamma_{112}/\Gamma_{30}$

Unseen decay modes of the  $\bar{K}^*(892)^0$  and  $\eta$  are included.

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
<b>0.135±0.034 OUR FIT</b>				
<b>0.13 ±0.02 ±0.03</b>	214	PROCARIO	93B CLE2	$\bar{K}^{*0} \eta \rightarrow K^- \pi^+ / \gamma \gamma$

$\Gamma(K^- \pi^+ \omega)/\Gamma(K^- \pi^+)$   $\Gamma_{113}/\Gamma_{20}$

Unseen decay modes of the  $\omega$  are included.

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
<b>0.78±0.12±0.10</b>	99	<sup>63</sup> ALBRECHT	92P ARG	$e^+ e^- \approx 10$ GeV

<sup>63</sup> This value is calculated from numbers in Table 1 of ALBRECHT 92P.

$\Gamma(\bar{K}^*(892)^0 \omega)/\Gamma(K^- \pi^+)$   $\Gamma_{114}/\Gamma_{20}$

Unseen decay modes of the  $\bar{K}^*(892)^0$  and  $\omega$  are included.

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
<b>0.28±0.11±0.04</b>	17	<sup>64</sup> ALBRECHT	92P ARG	$e^+ e^- \approx 10$ GeV

<sup>64</sup> This value is calculated from numbers in Table 1 of ALBRECHT 92P.

$\Gamma(\bar{K}^*(892)^0 \omega)/\Gamma(K^- \pi^+ \pi^+ \pi^- \pi^0)$   $\Gamma_{114}/\Gamma_{60}$

Unseen decay modes of the  $\bar{K}^*(892)^0$  and  $\omega$  are included.

VALUE	CL%	DOCUMENT ID	TECN	COMMENT
-------	-----	-------------	------	---------

• • • We do not use the following data for averages, fits, limits, etc. • • •

<0.44	90	<sup>65</sup> ANJOS	90D E691	Photoproduction
-------	----	---------------------	----------	-----------------

<sup>65</sup> Recovered from the published limit,  $\Gamma(\bar{K}^*(892)^0 \omega)/\Gamma_{\text{total}}$ , in order to make our normalization consistent.

$\Gamma(K^- \pi^+ \eta'(958))/\Gamma(K^- \pi^+ \pi^+ \pi^-)$   $\Gamma_{115}/\Gamma_{42}$

Unseen decay modes of the  $\eta'(958)$  are included.

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
<b>0.093 ± 0.014 ± 0.019</b>	286	PROCARIO	93B CLE2	$\eta' \rightarrow \eta \pi^+ \pi^-, \rho^0 \gamma$

$\Gamma(\bar{K}^*(892)^0 \eta'(958))/\Gamma(K^- \pi^+ \eta'(958))$   $\Gamma_{116}/\Gamma_{115}$

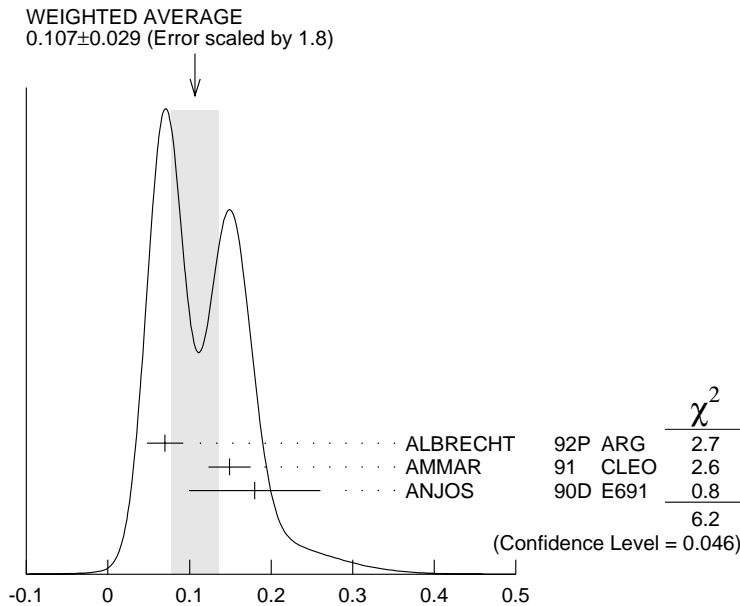
Unseen decay modes of the  $\bar{K}^*(892)^0$  are included.

VALUE	CL%	DOCUMENT ID	TECN	COMMENT
<b>&lt;0.15</b>	90	PROCARIO	93B CLE2	

$\Gamma(\bar{K}^0 \pi^+ \pi^+ \pi^- \pi^-)/\Gamma(\bar{K}^0 \pi^+ \pi^-)$   $\Gamma_{65}/\Gamma_{22}$

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
<b>0.107 ± 0.029 OUR AVERAGE</b>				Error includes scale factor of 1.8. See the ideogram below.
0.07 ± 0.02 ± 0.01	11	<sup>66</sup> ALBRECHT	92P ARG	$e^+ e^- \approx 10$ GeV
0.149 ± 0.026	56	AMMAR	91 CLEO	$e^+ e^- \approx 10.5$ GeV
0.18 ± 0.07 ± 0.04	6	ANJOS	90D E691	Photoproduction

<sup>66</sup> This value is calculated from numbers in Table 1 of ALBRECHT 92P.



$$\Gamma(\bar{K}^0 \pi^+ \pi^+ \pi^- \pi^-)/\Gamma(\bar{K}^0 \pi^+ \pi^-)$$

$\Gamma(\bar{K}^0 \pi^+ \pi^- \pi^0 \pi^0(\pi^0))/\Gamma_{\text{total}}$   $\Gamma_{66}/\Gamma$

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
<b>0.106<sup>+0.073</sup><sub>-0.029</sub> ± 0.006</b>	4	<sup>67</sup> AGUILAR-...	87F HYBR	$\pi p, p p$ 360, 400 GeV

<sup>67</sup> AGUILAR-BENITEZ 87F computes the branching fraction using topological normalization, and does not distinguish the presence of a third  $\pi^0$ .

$\Gamma(\overline{K}^0 K^+ K^-)/\Gamma(\overline{K}^0 \pi^+ \pi^-)$   $\Gamma_{67}/\Gamma_{22} = (\Gamma_{69} + \frac{1}{2}\Gamma_{79})/\Gamma_{22}$

<u>VALUE</u>	<u>EVTS</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
<b>0.172±0.014 OUR FIT</b>				
<b>0.178±0.019 OUR AVERAGE</b>				
0.20 ±0.05 ±0.04	47	FRABETTI	92B E687	$\gamma$ Be $\overline{E}_\gamma = 221$ GeV
0.170±0.022	136	AMMAR	91 CLEO	$e^+ e^- \approx 10.5$ GeV
0.24 ±0.08		BEBEK	86 CLEO	$e^+ e^-$ near $\Upsilon(4S)$
0.185±0.055	52	ALBRECHT	85B ARG	$e^+ e^- 10$ GeV

$\Gamma(\overline{K}^0 \phi)/\Gamma(\overline{K}^0 \pi^+ \pi^-)$   $\Gamma_{79}/\Gamma_{22}$

Unseen decay modes of the  $\phi$  are included.

<u>VALUE</u>	<u>EVTS</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
<b>0.158±0.016 OUR FIT</b>				
<b>0.156±0.017 OUR AVERAGE</b>				
0.13 ±0.06 ±0.02	13	FRABETTI	92B E687	$\gamma$ Be $\overline{E}_\gamma = 221$ GeV
0.163±0.023	63	AMMAR	91 CLEO	$e^+ e^- \approx 10.5$ GeV
0.155±0.033	56	ALBRECHT	87E ARG	$e^+ e^- 10$ GeV
0.14 ±0.05	29	BEBEK	86 CLEO	$e^+ e^-$ near $\Upsilon(4S)$
• • • We do not use the following data for averages, fits, limits, etc. • • •				
0.186±0.052	26	ALBRECHT	85B ARG	See ALBRECHT 87E

$\Gamma(\overline{K}^0 K^+ K^- \text{ non-}\phi)/\Gamma(\overline{K}^0 \pi^+ \pi^-)$   $\Gamma_{69}/\Gamma_{22}$

<u>VALUE</u>	<u>EVTS</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
<b>0.093±0.014 OUR FIT</b>				
<b>0.088±0.019 OUR AVERAGE</b>				
0.11 ±0.04 ±0.03	20	FRABETTI	92B E687	$\gamma$ Be $\overline{E}_\gamma = 221$ GeV
0.084±0.020		ALBRECHT	87E ARG	$e^+ e^- 10$ GeV

$\Gamma(K_S^0 K_S^0 K_S^0)/\Gamma(\overline{K}^0 \pi^+ \pi^-)$   $\Gamma_{70}/\Gamma_{22}$

<u>VALUE</u>	<u>EVTS</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
<b>0.0154±0.0025 OUR AVERAGE</b>				
0.0139±0.0019±0.0024	61	ASNER	96B CLE2	$e^+ e^- \approx \Upsilon(4S)$
0.035 ±0.012 ±0.006	10	FRABETTI	94J E687	$\gamma$ Be $\overline{E}_\gamma = 220$ GeV
0.016 ±0.005	22	AMMAR	91 CLEO	$e^+ e^- \approx 10.5$ GeV
0.017 ±0.007 ±0.005	5	ALBRECHT	90C ARG	$e^+ e^- \approx 10$ GeV

$\Gamma(K^+ K^- K^- \pi^+)/\Gamma(K^- \pi^+ \pi^+ \pi^-)$   $\Gamma_{72}/\Gamma_{42}$

<u>VALUE</u>	<u>EVTS</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
<b>0.0032±0.0009 OUR AVERAGE</b> Error includes scale factor of 1.4.				
0.0054±0.0016±0.0008	18	AITALA	01D E791	$\pi^-$ nucleus, 500 GeV
0.0028±0.0007±0.0001	20	FRABETTI	95C E687	$\gamma$ Be, $\overline{E}_\gamma \approx 200$ GeV

$\Gamma(K^- \pi^+ \phi)/\Gamma(K^+ K^- K^- \pi^+)$   $\Gamma_{117}/\Gamma_{72}$

Unseen decay modes of the  $\phi$  are included.

<u>VALUE</u>	<u>EVTS</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
<b>1.4±0.6</b>	13	<sup>68</sup> AITALA	01D E791	$\pi^-$ nucleus, 500 GeV

<sup>68</sup> This AITALA 01D result is from a projection fit, not a full amplitude analysis.



$\Gamma(K^+ K^- \bar{K}^0 \pi^0)/\Gamma_{\text{total}}$   $\Gamma_{73}/\Gamma$

VALUE DOCUMENT ID TECN COMMENT

• • • We do not use the following data for averages, fits, limits, etc. • • •

0.0072<sup>+0.0048</sup><sub>-0.0035</sub> <sup>69</sup> BARLAG 92C ACCM  $\pi^-$  Cu 230 GeV

<sup>69</sup> BARLAG 92C computes the branching fraction using topological normalization.

———— Pionic modes ————

$\Gamma(\pi^+ \pi^-)/\Gamma(K^- \pi^+)$   $\Gamma_{118}/\Gamma_{20}$

VALUE EVTS DOCUMENT ID TECN COMMENT

**0.0376 ± 0.0017 OUR AVERAGE** Error includes scale factor of 1.1.

0.0351 ± 0.0016 ± 0.0017 710 CSORNA 02 CLE2  $e^+ e^- \approx \Upsilon(4S)$

0.040 ± 0.002 ± 0.003 2043 AITALA 98C E791  $\pi^-$  nucleus, 500

0.043 ± 0.007 ± 0.003 177 FRABETTI 94C E687  $\gamma$ Be  $\bar{E}_\gamma = 220$

0.0348 ± 0.0030 ± 0.0023 227 SELEN 93 CLE2  $e^+ e^- \approx \Upsilon(4S)$

0.048 ± 0.013 ± 0.008 51 ADAMOVICH 92 OMEG  $\pi^-$  340 GeV

0.055 ± 0.008 ± 0.005 120 ANJOS 91D E691 Photoproduction

0.040 ± 0.007 ± 0.006 57 ALBRECHT 90C ARG  $e^+ e^- \approx 10$  GeV

0.050 ± 0.007 ± 0.005 110 ALEXANDER 90 CLEO  $e^+ e^-$  10.5–11

0.033 ± 0.010 ± 0.006 39 BALTRUSAIT..85E MRK3  $e^+ e^-$  3.77 GeV

0.033 ± 0.015 ABRAMS 79D MRK2  $e^+ e^-$  3.77 GeV

$\Gamma(\pi^0 \pi^0)/\Gamma(K^- \pi^+)$   $\Gamma_{119}/\Gamma_{20}$

VALUE EVTS DOCUMENT ID TECN COMMENT

**0.022 ± 0.004 ± 0.004** 40 SELEN 93 CLE2  $e^+ e^- \approx \Upsilon(4S)$

$\Gamma(\pi^+ \pi^- \pi^0)/\Gamma_{\text{total}}$   $\Gamma_{120}/\Gamma$

VALUE EVTS DOCUMENT ID TECN COMMENT

**0.011 ± 0.004 ± 0.002** 10 <sup>70</sup> BALTRUSAIT..85E MRK3  $e^+ e^-$  3.77 GeV

• • • We do not use the following data for averages, fits, limits, etc. • • •

0.0390<sup>+0.0100</sup><sub>-0.0095</sub> <sup>71</sup> BARLAG 92C ACCM  $\pi^-$  Cu 230 GeV

<sup>70</sup> All the BALTRUSAITIS 85E events are consistent with  $\rho^0 \pi^0$ .

<sup>71</sup> BARLAG 92C computes the branching fraction using topological normalization. Possible contamination by extra  $\pi^0$ 's may partly explain the unexpectedly large value.

$\Gamma(\pi^+ \pi^+ \pi^- \pi^-)/\Gamma(K^- \pi^+ \pi^+ \pi^-)$   $\Gamma_{121}/\Gamma_{42}$

VALUE EVTS DOCUMENT ID TECN COMMENT

**0.098 ± 0.006 OUR AVERAGE**

0.095 ± 0.007 ± 0.002 814 FRABETTI 95C E687  $\gamma$ Be,  $\bar{E}_\gamma \approx 200$  GeV

0.115 ± 0.023 ± 0.016 64 ADAMOVICH 92 OMEG  $\pi^-$  340 GeV

0.108 ± 0.024 ± 0.008 79 FRABETTI 92 E687  $\gamma$ Be

0.102 ± 0.013 345 <sup>72</sup> AMMAR 91 CLEO  $e^+ e^- \approx 10.5$  GeV

0.096 ± 0.018 ± 0.007 66 ANJOS 91 E691  $\gamma$ Be 80–240 GeV

<sup>72</sup> AMMAR 91 finds  $1.25 \pm 0.25 \pm 0.25 \rho^0$ 's per  $\pi^+ \pi^+ \pi^- \pi^-$  decay, but can't untangle the resonant substructure ( $\rho^0 \rho^0$ ,  $a_1^\pm \pi^\mp$ ,  $\rho^0 \pi^+ \pi^-$ ).

$\Gamma(\pi^+\pi^+\pi^-\pi^-\pi^0)/\Gamma_{\text{total}}$   $\Gamma_{122}/\Gamma$

VALUE DOCUMENT ID TECN COMMENT

• • • We do not use the following data for averages, fits, limits, etc. • • •

$0.0192^{+0.0041}_{-0.0038}$  73 BARLAG 92C ACCM  $\pi^-$  Cu 230 GeV

73 BARLAG 92C computes the branching fraction using topological normalization.

$\Gamma(\pi^+\pi^+\pi^+\pi^-\pi^-\pi^-)/\Gamma_{\text{total}}$   $\Gamma_{123}/\Gamma$

VALUE DOCUMENT ID TECN COMMENT

• • • We do not use the following data for averages, fits, limits, etc. • • •

$0.0004 \pm 0.0003$  74 BARLAG 92C ACCM  $\pi^-$  Cu 230 GeV

74 BARLAG 92C computes the branching fraction using topological normalization.

————— **Hadronic modes with a  $K\bar{K}$  pair** —————

$\Gamma(K^+K^-)/\Gamma(K^-\pi^+)$   $\Gamma_{124}/\Gamma_{20}$

VALUE EVTS DOCUMENT ID TECN COMMENT

**0.1083 ± 0.0027 OUR FIT**

**0.1084 ± 0.0026 OUR AVERAGE**

$0.1040 \pm 0.0033 \pm 0.0027$	1900	CSORNA	02	CLE2	$e^+e^- \approx \Upsilon(4S)$
$0.109 \pm 0.003 \pm 0.003$	3317	AITALA	98C	E791	$\pi^-$ nucleus, 500 GeV
$0.116 \pm 0.007 \pm 0.007$	1102	ASNER	96B	CLE2	$e^+e^- \approx \Upsilon(4S)$
$0.109 \pm 0.007 \pm 0.009$	581	FRABETTI	94C	E687	$\gamma\text{Be } \bar{E}_\gamma = 220$ GeV
$0.107 \pm 0.029 \pm 0.015$	103	ADAMOVICH	92	OMEG	$\pi^-$ 340 GeV
$0.138 \pm 0.027 \pm 0.010$	155	FRABETTI	92	E687	$\gamma\text{Be}$
$0.16 \pm 0.05$	34	ALVAREZ	91B	NA14	Photoproduction
$0.107 \pm 0.010 \pm 0.009$	193	ANJOS	91D	E691	Photoproduction
$0.10 \pm 0.02 \pm 0.01$	131	ALBRECHT	90C	ARG	$e^+e^- \approx 10$ GeV
$0.117 \pm 0.010 \pm 0.007$	249	ALEXANDER	90	CLEO	$e^+e^-$ 10.5–11 GeV
$0.122 \pm 0.018 \pm 0.012$	118	BALTRUSAIT.	85E	MRK3	$e^+e^-$ 3.77 GeV
$0.113 \pm 0.030$		ABRAMS	79D	MRK2	$e^+e^-$ 3.77 GeV

$\Gamma(K^+K^-)/\Gamma(\pi^+\pi^-)$   $\Gamma_{124}/\Gamma_{118}$

The unused results here are redundant with  $\Gamma(K^+K^-)/\Gamma(K^-\pi^+)$  and  $\Gamma(\pi^+\pi^-)/\Gamma(K^-\pi^+)$  measurements by the same experiments.

VALUE EVTS DOCUMENT ID TECN COMMENT

• • • We do not use the following data for averages, fits, limits, etc. • • •

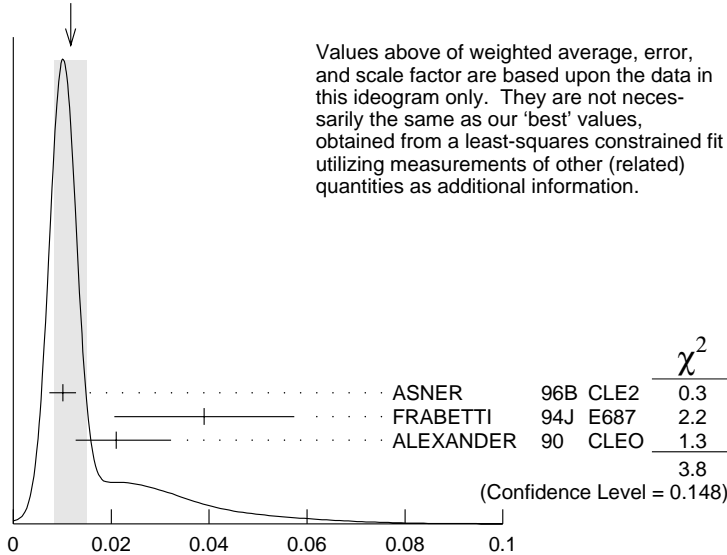
$2.96 \pm 0.16 \pm 0.15$	710	CSORNA	02	CLE2	$e^+e^- \approx \Upsilon(4S)$
$2.75 \pm 0.15 \pm 0.16$		AITALA	98C	E791	$\pi^-$ nucleus, 500 GeV
$2.53 \pm 0.46 \pm 0.19$		FRABETTI	94C	E687	$\gamma\text{Be } \bar{E}_\gamma = 220$ GeV
$2.23 \pm 0.81 \pm 0.46$		ADAMOVICH	92	OMEG	$\pi^-$ 340 GeV
$1.95 \pm 0.34 \pm 0.22$		ANJOS	91D	E691	Photoproduction
$2.5 \pm 0.7$		ALBRECHT	90C	ARG	$e^+e^- \approx 10$ GeV
$2.35 \pm 0.37 \pm 0.28$		ALEXANDER	90	CLEO	$e^+e^-$ 10.5–11 GeV

$\Gamma(K^0 \bar{K}^0) / \Gamma(\bar{K}^0 \pi^+ \pi^-)$

$\Gamma_{125} / \Gamma_{22}$

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
<b>0.0120 ± 0.0033 OUR FIT</b>				Error includes scale factor of 1.3.
<b>0.0117 ± 0.0033 OUR AVERAGE</b>				Error includes scale factor of 1.3. See the ideogram below.
0.0101 ± 0.0022 ± 0.0016	26	ASNER	96B CLE2	$e^+ e^- \approx \gamma(4S)$
0.039 ± 0.013 ± 0.013	20	FRABETTI	94J E687	$\gamma \text{Be } \bar{E}_\gamma = 220 \text{ GeV}$
0.021 $\begin{smallmatrix} +0.011 \\ -0.008 \end{smallmatrix}$ ± 0.002	5	ALEXANDER	90 CLEO	$e^+ e^- 10.5\text{--}11 \text{ GeV}$

WEIGHTED AVERAGE  
0.0117 ± 0.0033 (Error scaled by 1.3)



$\Gamma(K^0 \bar{K}^0) / \Gamma(\bar{K}^0 \pi^+ \pi^-)$

$\Gamma(K^0 \bar{K}^0) / \Gamma(K^+ K^-)$

$\Gamma_{125} / \Gamma_{124}$

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
<b>0.17 ± 0.05 OUR FIT</b>				Error includes scale factor of 1.2.
<b>0.24 ± 0.16</b>	4	<sup>75</sup> CUMALAT	88 SPEC	$nN 0\text{--}800 \text{ GeV}$

<sup>75</sup>Includes a correction communicated to us by the authors of CUMALAT 88.

$\Gamma(K^0 K^- \pi^+) / \Gamma(K^- \pi^+)$

$\Gamma_{126} / \Gamma_{20}$

VALUE	DOCUMENT ID	TECN	COMMENT
<b>0.182 ± 0.027 OUR FIT</b>			Error includes scale factor of 1.1.
<b>0.16 ± 0.06</b>	<sup>76</sup> ANJOS	91 E691	$\gamma \text{Be } 80\text{--}240 \text{ GeV}$

<sup>76</sup>The factor 100 at the top of column 2 of Table I of ANJOS 91 should be omitted.

$\Gamma(K^0 K^- \pi^+)/\Gamma(\bar{K}^0 \pi^+ \pi^-)$   $\Gamma_{126}/\Gamma_{22}$

<u>VALUE</u>	<u>EVTS</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
<b>0.117±0.017 OUR FIT</b>	Error includes scale factor of 1.1.			
<b>0.119±0.021 OUR AVERAGE</b>	Error includes scale factor of 1.3.			
0.108±0.019	61	AMMAR	91 CLEO	$e^+ e^- \approx 10.5$ GeV
0.16 ±0.03 ±0.02	39	ALBRECHT	90C ARG	$e^+ e^- \approx 10$ GeV

$\Gamma(\bar{K}^*(892)^0 K^0)/\Gamma(\bar{K}^0 \pi^+ \pi^-)$   $\Gamma_{146}/\Gamma_{22}$

Unseen decay modes of the  $\bar{K}^*(892)^0$  are included.

<u>VALUE</u>	<u>CL%</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
<b>&lt;0.029</b>	90	AMMAR	91 CLEO	$e^+ e^- \approx 10.5$ GeV
• • • We do not use the following data for averages, fits, limits, etc. • • •				
<0.03	90	ALBRECHT	90C ARG	$e^+ e^- \approx 10$ GeV

$\Gamma(K^*(892)^+ K^-)/\Gamma(K^- \pi^+)$   $\Gamma_{147}/\Gamma_{20}$

Unseen decay modes of the  $K^*(892)^+$  are included.

<u>VALUE</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
<b>0.099±0.021 OUR FIT</b>			
<b>0.16 <sup>+0.08</sup><sub>-0.06</sub></b>	<sup>77</sup> ANJOS	91 E691	$\gamma$ Be 80–240 GeV

<sup>77</sup> The factor 100 at the top of column 2 of Table I of ANJOS 91 should be omitted.

$\Gamma(K^*(892)^+ K^-)/\Gamma(\bar{K}^0 \pi^+ \pi^-)$   $\Gamma_{147}/\Gamma_{22}$

Unseen decay modes of the  $K^*(892)^+$  are included.

<u>VALUE</u>	<u>EVTS</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
<b>0.064±0.013 OUR FIT</b>				
<b>0.058±0.014 OUR AVERAGE</b>				
0.064±0.018	23	AMMAR	91 CLEO	$e^+ e^- \approx 10.5$ GeV
0.05 ±0.02 ±0.01	15	ALBRECHT	90C ARG	$e^+ e^- \approx 10$ GeV

$\Gamma(K^0 K^- \pi^+ \text{nonresonant})/\Gamma(K^- \pi^+)$   $\Gamma_{129}/\Gamma_{20}$

<u>VALUE</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
<b>0.06±0.06</b>	<sup>78</sup> ANJOS	91 E691	$\gamma$ Be 80–240 GeV

<sup>78</sup> The factor 100 at the top of column 2 of Table I of ANJOS 91 should be omitted.

$\Gamma(\bar{K}^0 K^+ \pi^-)/\Gamma(K^- \pi^+)$   $\Gamma_{130}/\Gamma_{20}$

<u>VALUE</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
<b>0.138±0.026 OUR FIT</b>			
<b>0.10 ±0.05</b>	<sup>79</sup> ANJOS	91 E691	$\gamma$ Be 80–240 GeV

<sup>79</sup> The factor 100 at the top of column 2 of Table I of ANJOS 91 should be omitted.

$\Gamma(\bar{K}^0 K^+ \pi^-)/\Gamma(\bar{K}^0 \pi^+ \pi^-)$   $\Gamma_{130}/\Gamma_{22}$

<u>VALUE</u>	<u>EVTS</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
<b>0.089±0.017 OUR FIT</b>				
<b>0.098±0.020</b>	55	AMMAR	91 CLEO	$e^+ e^- \approx 10.5$ GeV

$\Gamma(K^*(892)^0 \bar{K}^0)/\Gamma(\bar{K}^0 \pi^+ \pi^-)$   $\Gamma_{148}/\Gamma_{22}$

Unseen decay modes of the  $K^*(892)^0$  are included.

<u>VALUE</u>	<u>CL%</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
<b>&lt;0.015</b>	90	AMMAR	91 CLEO	$e^+ e^- \approx 10.5$ GeV

$\Gamma(K^*(892)^- K^+)/\Gamma(\bar{K}^0 \pi^+ \pi^-)$   $\Gamma_{149}/\Gamma_{22}$

Unseen decay modes of the  $K^*(892)^-$  are included.

<u>VALUE</u>	<u>EVTS</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
<b>0.034 ± 0.019</b>	12	AMMAR	91	CLEO $e^+ e^- \approx 10.5$ GeV

$\Gamma(\bar{K}^0 K^+ \pi^- \text{ nonresonant})/\Gamma(K^- \pi^+)$   $\Gamma_{133}/\Gamma_{20}$

<u>VALUE</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
<b>0.10<sup>+0.06</sup><sub>-0.05</sub></b>	<sup>80</sup> ANJOS	91	E691 $\gamma$ Be 80–240 GeV

<sup>80</sup> The factor 100 at the top of column 2 of Table I of ANJOS 91 should be omitted.

$\Gamma(K^+ K^- \pi^0)/\Gamma(K^- \pi^+ \pi^0)$   $\Gamma_{134}/\Gamma_{30}$

<u>VALUE</u>	<u>EVTS</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
<b>0.0095 ± 0.0026</b>	151	ASNER	96B	CLE2 $e^+ e^- \approx \Upsilon(4S)$

$\Gamma(K_S^0 K_S^0 \pi^0)/\Gamma_{\text{total}}$   $\Gamma_{135}/\Gamma$

<u>VALUE</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
<b>&lt;0.00059</b>	ASNER	96B	CLE2 $e^+ e^- \approx \Upsilon(4S)$

$\Gamma(\phi \pi^0)/\Gamma_{\text{total}}$   $\Gamma_{150}/\Gamma$

<u>VALUE</u>	<u>CL%</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
<b>&lt;0.0014</b>	90	ALBRECHT	94i	ARG $e^+ e^- \approx 10$ GeV

$\Gamma(\phi \eta)/\Gamma_{\text{total}}$   $\Gamma_{151}/\Gamma$

<u>VALUE</u>	<u>CL%</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
<b>&lt;0.0028</b>	90	ALBRECHT	94i	ARG $e^+ e^- \approx 10$ GeV

$\Gamma(\phi \omega)/\Gamma_{\text{total}}$   $\Gamma_{152}/\Gamma$

<u>VALUE</u>	<u>CL%</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
<b>&lt;0.0021</b>	90	ALBRECHT	94i	ARG $e^+ e^- \approx 10$ GeV

$\Gamma(K^+ K^- \pi^+ \pi^-)/\Gamma(K^- \pi^+ \pi^+ \pi^-)$   $\Gamma_{136}/\Gamma_{42}$

<u>VALUE</u>	<u>EVTS</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
<b>0.0334 ± 0.0028 OUR AVERAGE</b>				
0.0313 ± 0.0037 ± 0.0036	136	AITALA	98D	E791 $\pi^-$ nucleus, 500 GeV
0.035 ± 0.004 ± 0.002	244	FRABETTI	95C	E687 $\gamma$ Be, $\bar{E}_\gamma \approx 200$ GeV
0.041 ± 0.007 ± 0.005	114	ALBRECHT	94i	ARG $e^+ e^- \approx 10$ GeV
0.0314 ± 0.010	89	AMMAR	91	CLEO $e^+ e^- \approx 10.5$ GeV
0.028 <sup>+0.008</sup> <sub>-0.007</sub>		ANJOS	91	E691 $\gamma$ Be 80–240 GeV

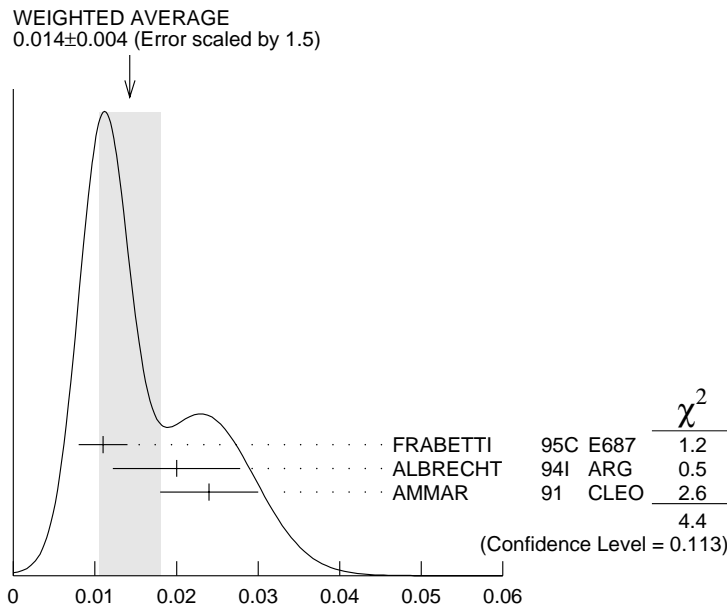
$\Gamma(\phi\pi^+\pi^-)/\Gamma(K^-\pi^+\pi^+\pi^-)$

$\Gamma_{153}/\Gamma_{42}$

Unseen decay modes of the  $\phi$  are included.

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
<b>0.014 ± 0.004 OUR AVERAGE</b>				Error includes scale factor of 1.5. See the ideogram below.
0.011 ± 0.003		FRABETTI	95C E687	$\gamma$ Be, $\bar{E}_\gamma \approx 200$ GeV
0.020 ± 0.006 ± 0.005	28	ALBRECHT	94I ARG	$e^+e^- \approx 10$ GeV
0.024 ± 0.006	34	<sup>81</sup> AMMAR	91 CLEO	$e^+e^- \approx 10.5$ GeV
• • • We do not use the following data for averages, fits, limits, etc. • • •				
0.0076 <sup>+0.0066</sup> <sub>-0.0049</sub>	3	ANJOS	91 E691	$\gamma$ Be 80–240 GeV

<sup>81</sup>AMMAR 91 measures  $\phi\rho^0$ , but notes that  $\phi\rho^0$  dominates  $\phi\pi^+\pi^-$ . We put the measurement here to keep from having more  $\phi\rho^0$  than  $\phi\pi^+\pi^-$ .



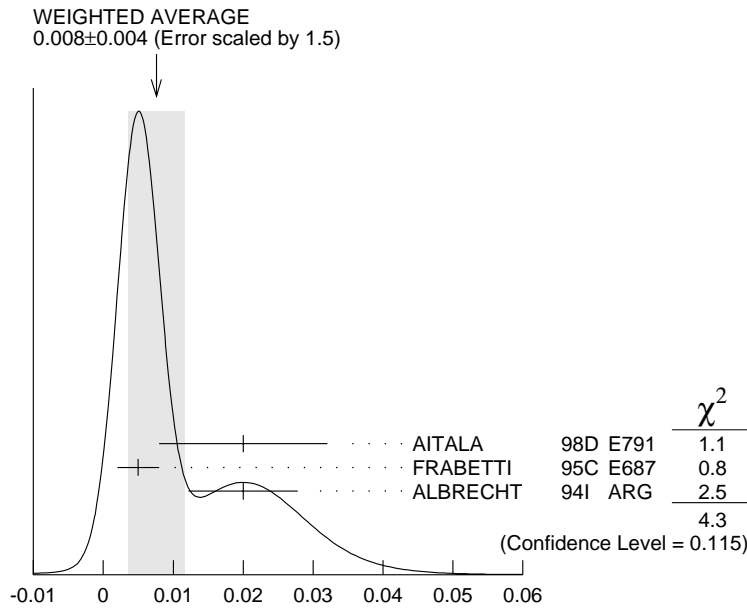
$\Gamma(\phi\pi^+\pi^-)/\Gamma(K^-\pi^+\pi^+\pi^-)$

$\Gamma(\phi\rho^0)/\Gamma(K^-\pi^+\pi^+\pi^-)$

$\Gamma_{154}/\Gamma_{42}$

Unseen decay modes of the  $\phi$  are included.

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
<b>0.008 ± 0.004 OUR AVERAGE</b>				Error includes scale factor of 1.5. See the ideogram below.
0.02 ± 0.009 ± 0.008		AITALA	98D E791	$\pi^-$ nucleus, 500 GeV
0.005 ± 0.003		FRABETTI	95C E687	$\gamma$ Be, $\bar{E}_\gamma \approx 200$ GeV
0.020 ± 0.006 ± 0.005	28	ALBRECHT	94I ARG	$e^+e^- \approx 10$ GeV



$$\Gamma(\phi \rho^0) / \Gamma(K^- \pi^+ \pi^+ \pi^-)$$

$$\Gamma(\phi \pi^+ \pi^- \text{ 3-body}) / \Gamma(K^- \pi^+ \pi^+ \pi^-) \quad \Gamma_{155} / \Gamma_{42}$$

Unseen decay modes of the  $\phi$  are included.

VALUE	CL%	DOCUMENT ID	TECN	COMMENT
<b><math>0.009 \pm 0.004 \pm 0.005</math></b>		AITALA	98D E791	$\pi^-$ nucleus, 500 GeV
• • • We do not use the following data for averages, fits, limits, etc. • • •				
<0.006	90	FRABETTI	95C E687	$\gamma$ Be, $\bar{E}_\gamma \approx 200$ GeV

$$\Gamma(K^+ K^- \rho^0 \text{ 3-body}) / \Gamma(K^- \pi^+ \pi^+ \pi^-) \quad \Gamma_{139} / \Gamma_{42}$$

VALUE	DOCUMENT ID	TECN	COMMENT
<b><math>0.012 \pm 0.003</math></b>	FRABETTI	95C E687	$\gamma$ Be, $\bar{E}_\gamma \approx 200$ GeV

$$\Gamma(K^*(892)^0 K^- \pi^+ + \text{c.c.}) / \Gamma(K^- \pi^+ \pi^+ \pi^-) \quad \Gamma_{156} / \Gamma_{42}$$

Unseen decay modes of the  $K^*(892)^0$  are included.

VALUE	CL%	DOCUMENT ID	TECN	COMMENT
<b>&lt;0.01</b>	90	<sup>82</sup> AITALA	98D E791	$\pi^-$ nucleus, 500 GeV
• • • We do not use the following data for averages, fits, limits, etc. • • •				
<0.017	90	<sup>82</sup> FRABETTI	95C E687	$\gamma$ Be, $\bar{E}_\gamma \approx 200$ GeV
$0.010^{+0.016}_{-0.010}$		ANJOS	91 E691	$\gamma$ Be 80–240 GeV

<sup>82</sup> These upper limits are in conflict with values in the next two data blocks.

$\Gamma(K^*(892)^0 K^- \pi^+)/\Gamma(K^- \pi^+ \pi^+ \pi^-)$   $\Gamma_{157}/\Gamma_{42}$

The  $K^{*0} K^- \pi^+$  and  $\bar{K}^{*0} K^+ \pi^-$  modes are distinguished by the charge of the pion in  $D^*(2010)^\pm \rightarrow D^0 \pi^\pm$  decays. Unseen decay modes of the  $K^*(892)^0$  are included.

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
-------	------	-------------	------	---------

• • • We do not use the following data for averages, fits, limits, etc. • • •

0.043±0.014±0.009      55      <sup>83</sup> ALBRECHT      94I ARG       $e^+ e^- \approx 10$  GeV

<sup>83</sup> This ALBRECHT 94I value is in conflict with upper limits given above.

$\Gamma(\bar{K}^*(892)^0 K^+ \pi^-)/\Gamma(K^- \pi^+ \pi^+ \pi^-)$   $\Gamma_{158}/\Gamma_{42}$

The  $K^{*0} K^- \pi^+$  and  $\bar{K}^{*0} K^+ \pi^-$  modes are distinguished by the charge of the pion in  $D^*(2010)^\pm \rightarrow D^0 \pi^\pm$  decays. Unseen decay modes of the  $\bar{K}^*(892)^0$  are included.

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
-------	------	-------------	------	---------

• • • We do not use the following data for averages, fits, limits, etc. • • •

0.023±0.013±0.009      30      <sup>84</sup> ALBRECHT      94I ARG       $e^+ e^- \approx 10$  GeV

<sup>84</sup> This ALBRECHT 94I value is in conflict with upper limits given above.

$\Gamma(K^*(892)^0 \bar{K}^*(892)^0)/\Gamma(K^- \pi^+ \pi^+ \pi^-)$   $\Gamma_{159}/\Gamma_{42}$

Unseen decay modes of the  $K^*(892)^0$  and  $\bar{K}^*(892)^0$  are included.

VALUE	CL%	EVTS	DOCUMENT ID	TECN	COMMENT
-------	-----	------	-------------	------	---------

**0.018±0.007 OUR AVERAGE**      Error includes scale factor of 1.2.

0.016±0.006      FRABETTI      95C E687       $\gamma$ Be,  $\bar{E}_\gamma \approx 200$  GeV

0.036<sup>+0.020</sup><sub>-0.016</sub>      11      ANJOS      91 E691       $\gamma$ Be 80–240 GeV

• • • We do not use the following data for averages, fits, limits, etc. • • •

<0.02      90      AITALA      98D E791       $\pi^-$  nucleus, 500 GeV

<0.033      90      <sup>85</sup> AMMAR      91 CLEO       $e^+ e^- \approx 10.5$  GeV

<sup>85</sup> A corrected value (G. Moneti, private communication).

$\Gamma(K^+ K^- \pi^+ \pi^- \text{ non-}\phi)/\Gamma_{\text{total}}$   $\Gamma_{142}/\Gamma$

VALUE	DOCUMENT ID	TECN	COMMENT
-------	-------------	------	---------

• • • We do not use the following data for averages, fits, limits, etc. • • •

0.0017±0.0005      <sup>86</sup> BARLAG      92C ACCM       $\pi^-$  Cu 230 GeV

<sup>86</sup> BARLAG 92C computes the branching fraction using topological normalization.

$\Gamma(K^+ K^- \pi^+ \pi^- \text{ nonresonant})/\Gamma(K^- \pi^+ \pi^+ \pi^-)$   $\Gamma_{143}/\Gamma_{42}$

VALUE	CL%	DOCUMENT ID	TECN	COMMENT
-------	-----	-------------	------	---------

**<0.011**      90      FRABETTI      95C E687       $\gamma$ Be,  $\bar{E}_\gamma \approx 200$  GeV

• • • We do not use the following data for averages, fits, limits, etc. • • •

0.001<sup>+0.011</sup><sub>-0.001</sub>      ANJOS      91 E691       $\gamma$ Be 80–240 GeV

$\Gamma(K^0 \bar{K}^0 \pi^+ \pi^-)/\Gamma(\bar{K}^0 \pi^+ \pi^-)$   $\Gamma_{144}/\Gamma_{22}$

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
-------	------	-------------	------	---------

**0.126±0.038±0.030**      25      ALBRECHT      94I ARG       $e^+ e^- \approx 10$  GeV



$\Gamma(K^+ K^- \pi^+ \pi^- \pi^0)/\Gamma_{\text{total}}$   $\Gamma_{145}/\Gamma$

<u>VALUE</u>	<u>CL%</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
<b>0.0031±0.0020</b>		<sup>87</sup> BARLAG	92C ACCM	$\pi^-$ Cu 230 GeV

<sup>87</sup> BARLAG 92C computes the branching fraction using topological normalization.

———— Radiative modes ————

$\Gamma(\rho^0 \gamma)/\Gamma_{\text{total}}$   $\Gamma_{160}/\Gamma$

<u>VALUE</u>	<u>CL%</u>	<u>DOCUMENT ID</u>	<u>TECN</u>
<b>&lt;2.4 × 10<sup>-4</sup></b>	90	ASNER	98 CLE2

$\Gamma(\omega \gamma)/\Gamma_{\text{total}}$   $\Gamma_{161}/\Gamma$

<u>VALUE</u>	<u>CL%</u>	<u>DOCUMENT ID</u>	<u>TECN</u>
<b>&lt;2.4 × 10<sup>-4</sup></b>	90	ASNER	98 CLE2

$\Gamma(\phi \gamma)/\Gamma_{\text{total}}$   $\Gamma_{162}/\Gamma$

<u>VALUE</u>	<u>CL%</u>	<u>DOCUMENT ID</u>	<u>TECN</u>
<b>&lt;1.9 × 10<sup>-4</sup></b>	90	ASNER	98 CLE2

$\Gamma(\bar{K}^*(892)^0 \gamma)/\Gamma_{\text{total}}$   $\Gamma_{163}/\Gamma$

<u>VALUE</u>	<u>CL%</u>	<u>DOCUMENT ID</u>	<u>TECN</u>
<b>&lt;7.6 × 10<sup>-4</sup></b>	90	ASNER	98 CLE2

———— Rare or forbidden modes ————

$\Gamma(K^+ \ell^- \bar{\nu}_\ell (\text{via } \bar{D}^0))/\Gamma(K^- \ell^+ \nu_\ell)$   $\Gamma_{164}/\Gamma_8$

This is a limit on  $R_M$  without the complications of possible doubly-Cabibbo-suppressed decays that occur when using hadronic modes. For the limits on  $|m_1 - m_2|$  and  $(\Gamma_1 - \Gamma_2)/\Gamma$  that come from the best mixing limit, see near the beginning of these  $D^0$  Listings.

<u>VALUE</u>	<u>CL%</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
<b>&lt;0.005</b>	90	<sup>88</sup> AITALA	96C E791	$\pi^-$ nucleus, 500 GeV

<sup>88</sup> AITALA 96C uses  $D^{*+} \rightarrow D^0 \pi^+$  (and charge conjugate) decays to identify the charm at production and  $D^0 \rightarrow K^- \ell^+ \nu_\ell$  (and charge conjugate) decays to identify the charm at decay.

$\Gamma(K^+ \pi^-)/\Gamma(K^- \pi^+)$

$\Gamma_{165}/\Gamma_{20}$

This is  $R_b$  in the note on “ $D^0$ - $\bar{D}^0$  Mixing,” near the start of the  $D^0$  Listings. The experiments here use the charge of the pion in  $D^*(2010)^\pm \rightarrow (D^0 \text{ or } \bar{D}^0) \pi^\pm$  decay to tell whether a  $D^0$  or a  $\bar{D}^0$  was born. The  $D^0 \rightarrow K^+ \pi^-$  decay can occur directly by doubly Cabibbo-suppressed (DCS) decay, or indirectly by  $D^0 \rightarrow \bar{D}^0$  mixing followed by  $\bar{D}^0 \rightarrow K^+ \pi^-$  decay. Some of the experiments can use the decay-time information to disentangle the two mechanisms. Here, we list the experimental branching ratio, which if there is no mixing is the DCS ratio; in the next data block we give the limits on the mixing ratio.

Some early limits have been omitted from this Listing; see our 1998 edition (EPJ **C3** 1).

VALUE	CL%	EVTS	DOCUMENT ID	TECN	COMMENT
<b>0.0039 ± 0.0006</b>		<b>OUR AVERAGE</b>			
0.0040 ± 0.00085 ± 0.00025		149	<sup>89</sup> LINK	01 FOCS	$\gamma$ nucleus
0.00332 <sup>+0.00063</sup> <sub>-0.00065</sub> ± 0.00040		45	<sup>90</sup> GODANG	00 CLE2	$e^+ e^-$
0.0068 <sup>+0.0034</sup> <sub>-0.0033</sub> ± 0.0007		34	<sup>91</sup> AITALA	98 E791	$\pi^-$ nucleus, 500 GeV
0.0184 ± 0.0059 ± 0.0034		19	<sup>92</sup> BARATE	98W ALEP	$e^+ e^-$ at $Z^0$
0.0077 ± 0.0025 ± 0.0025		19	<sup>93</sup> CINABRO	94 CLE2	$e^+ e^- \approx \Upsilon(4S)$

• • • We do not use the following data for averages, fits, limits, etc. • • •

<0.011	90		<sup>93</sup> AMMAR	91 CLEO	$e^+ e^- \approx 10.5$ GeV
<0.015	90	1 ± 6	<sup>94</sup> ANJOS	88C E691	Photoproduction
<0.014	90		<sup>93</sup> ALBRECHT	87K ARG	$e^+ e^-$ 10 GeV

<sup>89</sup>This LINK 01 result assumes no  $D^0$ - $\bar{D}^0$  mixing; see Fig. 4 of the paper for the DCS value as a function of the (unknown) mixing parameters  $x'$  and  $y'$ .

<sup>90</sup>This GODANG 00 result assumes no  $D^0$ - $\bar{D}^0$  mixing ( $R_M=0$  in the note on “ $D^0$ - $\bar{D}^0$  Mixing” near the start of the  $D^0$  Listings). The DCS ratio becomes  $0.0048 \pm 0.0012 \pm 0.0004$  when mixing is allowed.

<sup>91</sup>This AITALA 98 result assumes no  $D^0$ - $\bar{D}^0$  mixing ( $R_M=0$  in the note on “ $D^0$ - $\bar{D}^0$  Mixing” near the start of the  $D^0$  Listings). The DCS ratio becomes  $0.0090^{+0.0120}_{-0.0109} \pm 0.0044$  when mixing is allowed.

<sup>92</sup>BARATE 98W gets  $0.0177^{+0.0060}_{-0.0056} \pm 0.0031$  for the DCS ratio when mixing is allowed, assuming no interference between the DCS and mixing amplitudes ( $y' = 0$  in the note on “ $D^0$ - $\bar{D}^0$  Mixing” near the start of the  $D^0$  Listings).

<sup>93</sup>CINABRO 94, AMMAR 91, and ALBRECHT 87K cannot distinguish between doubly Cabibbo-suppressed decay and  $D^0$ - $\bar{D}^0$  mixing.

<sup>94</sup>ANJOS 88C allows mixing but assumes no interference between the DCS and mixing amplitudes ( $y' = 0$  in the note on “ $D^0$ - $\bar{D}^0$  Mixing” near the start of the  $D^0$  Listings). When interference is allowed, the limit degrades to 0.049.

$\Gamma(K^+ \pi^- (\text{via } \bar{D}^0))/\Gamma(K^- \pi^+)$

$\Gamma_{166}/\Gamma_{20}$

This is  $R_M$  in the note on “ $D^0$ - $\bar{D}^0$  Mixing” near the start of the  $D^0$  Listings. The experiments here (1) use the charge of the pion in  $D^*(2010)^\pm \rightarrow (D^0 \text{ or } \bar{D}^0) \pi^\pm$  decay to tell whether a  $D^0$  or a  $\bar{D}^0$  was born; and (2) use the decay-time distribution to disentangle doubly Cabibbo-suppressed decay and mixing. For the limits on  $|m_1 - m_2|$  and  $(\Gamma_1 - \Gamma_2)/\Gamma$  that come from the best mixing limit, see near the beginning of these  $D^0$  Listings.

VALUE	CL%	EVTS	DOCUMENT ID	TECN	COMMENT
<b>&lt;0.00041</b>	95		95 GODANG	00 CLE2	$e^+ e^-$
<0.0092	95		96 BARATE	98W ALEP	$e^+ e^-$ at $Z^0$
<0.005	90	$1 \pm 4$	97 ANJOS	88C E691	Photoproduction

• • • We do not use the following data for averages, fits, limits, etc. • • •

<sup>95</sup>This GODANG 00 result assumes that the strong phase between  $D^0 \rightarrow K^+ \pi^-$  and  $\bar{D}^0 \rightarrow K^+ \pi^-$  is small, and limits only  $D^0 \rightarrow \bar{D}^0$  transitions via off-shell intermediate states. The limit on transitions via on-shell intermediate states is 0.0017.

<sup>96</sup>This BARATE 98W result assumes no interference between the DCS and mixing amplitudes ( $y' = 0$  in the note on “ $D^0$ - $\bar{D}^0$  Mixing” near the start of the  $D^0$  Listings). When interference is allowed, the limit degrades to 0.036 (95%CL).

<sup>97</sup>This ANJOS 88C result assumes no interference between the DCS and mixing amplitudes ( $y' = 0$  in the note on “ $D^0$ - $\bar{D}^0$  Mixing” near the start of the  $D^0$  Listings). When interference is allowed, the limit degrades to 0.019. Combined with results on  $K^\pm \pi^\mp \pi^+ \pi^-$ , the limit is, assuming no interference, 0.0037.

$\Gamma(K^+ \pi^- \pi^0)/\Gamma(K^- \pi^+ \pi^0)$

$\Gamma_{167}/\Gamma_{30}$

The experiments here use the charge of the pion in  $D^*(2010)^\pm \rightarrow (D^0 \text{ or } \bar{D}^0) \pi^\pm$  decay to tell whether a  $D^0$  or a  $\bar{D}^0$  was born. The  $D^0 \rightarrow K^+ \pi^- \pi^0$  decay can occur directly by doubly Cabibbo-suppressed (DCS) decay, or indirectly by  $D^0 \rightarrow \bar{D}^0$  mixing followed by  $\bar{D}^0 \rightarrow K^+ \pi^- \pi^0$  decay.

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
<b><math>0.0043^{+0.0011}_{-0.0010} \pm 0.0007</math></b>	38	98 BRANDENB...	01 CLE2	$e^+ e^- \approx \Upsilon(4S)$

<sup>98</sup>BRANDENBURG 01 cannot distinguish between doubly Cabibbo-suppressed decay and  $D^0$ - $\bar{D}^0$  mixing.

$\Gamma(K^+ \pi^- \pi^+ \pi^-) / \Gamma(K^- \pi^+ \pi^+ \pi^-)$   $\Gamma_{168} / \Gamma_{42}$

The experiments here use the charge of the pion in  $D^*(2010)^\pm \rightarrow (D^0 \text{ or } \bar{D}^0) \pi^\pm$  decay to tell whether a  $D^0$  or a  $\bar{D}^0$  was born. The  $D^0 \rightarrow K^+ \pi^- \pi^+ \pi^-$  decay can occur directly by doubly Cabibbo-suppressed (DCS) decay, or indirectly by  $D^0 \rightarrow \bar{D}^0$  mixing followed by  $\bar{D}^0 \rightarrow K^+ \pi^- \pi^+ \pi^-$  decay. Some of the experiments can use the decay-time information to disentangle the two mechanisms. Here, we list the experimental branching ratio, which if there is no mixing is the DCS ratio; in the next data block we give the limits on the mixing ratio.

Some early limits have been omitted from this Listing; see our 1998 edition (EPJ **C3** 1).

VALUE	CL%	EVTS	DOCUMENT ID	TECN	COMMENT
<b>0.0042 ± 0.0013 OUR AVERAGE</b>					
$0.0044^{+0.0013}_{-0.0012} \pm 0.0006$		54	<sup>99</sup> DYTMAN	01 CLE2	$e^+ e^- \approx \Upsilon(4S)$
$0.0025^{+0.0036}_{-0.0034} \pm 0.0003$			<sup>100</sup> AITALA	98 E791	$\pi^-$ nucleus, 500 GeV
● ● ● We do not use the following data for averages, fits, limits, etc. ● ● ●					
<0.018	90		<sup>99</sup> AMMAR	91 CLEO	$e^+ e^- \approx 10.5$ GeV
<0.018	90	$5 \pm 12$	<sup>101</sup> ANJOS	88C E691	Photoproduction

- <sup>99</sup> AMMAR 91 and DYTMAN 01 cannot distinguish between doubly Cabibbo-suppressed decay and  $D^0$ - $\bar{D}^0$  mixing.
- <sup>100</sup> This AITALA 98 result assumes no  $D^0$ - $\bar{D}^0$  mixing ( $R_M$  in the note on “ $D^0$ - $\bar{D}^0$  Mixing”). It becomes  $-0.0020^{+0.0117}_{-0.0106} \pm 0.0035$  when mixing is allowed and decay-time information is used to distinguish doubly Cabibbo-suppressed decays from mixing.
- <sup>101</sup> ANJOS 88C uses decay-time information to distinguish doubly Cabibbo-suppressed (DCS) decays from  $D^0$ - $\bar{D}^0$  mixing. However, the result assumes no interference between the DCS and mixing amplitudes ( $y' = 0$  in the note on “ $D^0$ - $\bar{D}^0$  Mixing” near the start of the  $D^0$  Listings). When interference is allowed, the limit degrades to 0.033.

$\Gamma(K^+ \pi^- \pi^+ \pi^- \text{ (via } \bar{D}^0)) / \Gamma(K^- \pi^+ \pi^+ \pi^-)$   $\Gamma_{169} / \Gamma_{42}$

This is a  $D^0$ - $\bar{D}^0$  mixing limit. The experiments here (1) use the charge of the pion in  $D^*(2010)^\pm \rightarrow (D^0 \text{ or } \bar{D}^0) \pi^\pm$  decay to tell whether a  $D^0$  or a  $\bar{D}^0$  was born; and (2) use the decay-time distribution to disentangle doubly Cabibbo-suppressed decay and mixing. For the limits on  $|m_{D_1^0} - m_{D_2^0}|$  and  $(\Gamma_{D_1^0} - \Gamma_{D_2^0}) / \Gamma_{D^0}$  that come from the best mixing limit, see near the beginning of these  $D^0$  Listings.

VALUE	CL%	EVTS	DOCUMENT ID	TECN	COMMENT
<b>&lt;0.005</b>	90	$0 \pm 4$	<sup>102</sup> ANJOS	88C E691	Photoproduction

- <sup>102</sup> ANJOS 88C uses decay-time information to distinguish doubly Cabibbo-suppressed (DCS) decays from  $D^0$ - $\bar{D}^0$  mixing. However, the result assumes no interference between the DCS and mixing amplitudes ( $y' = 0$  in the note on “ $D^0$ - $\bar{D}^0$  Mixing” near the start of the  $D^0$  Listings). When interference is allowed, the limit degrades to 0.007. Combined with results on  $K^\pm \pi^\mp$ , the limit is, assuming no interference, 0.0037.

$\Gamma(K^+\pi^- \text{ or } K^+\pi^-\pi^+\pi^- \text{ (via } \bar{D}^0)) / \Gamma(K^-\pi^+ \text{ or } K^-\pi^+\pi^+\pi^-)$   $\Gamma_{170}/\Gamma_0$

This is a  $D^0-\bar{D}^0$  mixing limit. For the limits on  $|m_{D_1^0} - m_{D_2^0}|$  and  $(\Gamma_{D_1^0} - \Gamma_{D_2^0})/\Gamma_{D^0}$  that come from the best mixing limit, see near the beginning of these  $D^0$  Listings.

VALUE	CL%	DOCUMENT ID	TECN	COMMENT
<b>&lt;0.0085</b>	90	<sup>103</sup> AITALA	98 E791	$\pi^-$ nucleus, 500 GeV

<sup>103</sup>AITALA 98 uses decay-time information to distinguish doubly Cabibbo-suppressed decays from  $D^0-\bar{D}^0$  mixing. The fit allows interference between the two amplitudes, and also allows  $CP$  violation in this term. The central value obtained is  $0.0039^{+0.0036}_{-0.0032} \pm 0.0016$ . When interference is disallowed, the result becomes  $0.0021 \pm 0.0009 \pm 0.0002$ .

$\Gamma(\mu^- \text{ anything (via } \bar{D}^0)) / \Gamma(\mu^+ \text{ anything})$   $\Gamma_{171}/\Gamma_2$

This is a  $D^0-\bar{D}^0$  mixing limit. See the somewhat better limits above.

VALUE	CL%	DOCUMENT ID	TECN	COMMENT
<b>&lt;0.0056</b>	90	LOUIS	86 SPEC	$\pi^-$ W 225 GeV
● ● ● We do not use the following data for averages, fits, limits, etc. ● ● ●				
<0.012	90	BENVENUTI	85 CNTR	$\mu$ C, 200 GeV
<0.044	90	BODEK	82 SPEC	$\pi^-$ , $p$ Fe $\rightarrow D^0$

$\Gamma(e^+e^-) / \Gamma_{\text{total}}$   $\Gamma_{172}/\Gamma$

A test for the  $\Delta C = 1$  weak neutral current. Allowed by first-order weak interaction combined with electromagnetic interaction.

VALUE	CL%	EVTS	DOCUMENT ID	TECN	COMMENT
<b>&lt;6.2 <math>\times 10^{-6}</math></b>	90		AITALA	99G E791	$\pi^-$ N 500 GeV
● ● ● We do not use the following data for averages, fits, limits, etc. ● ● ●					
<8.19 $\times 10^{-6}$	90		PRIPSTEIN	00 E789	$p$ nucleus, 800 GeV
<1.3 $\times 10^{-5}$	90	0	FREYBERGER	96 CLE2	$e^+e^- \approx \gamma(4S)$
<1.3 $\times 10^{-4}$	90		ADLER	88 MRK3	$e^+e^-$ 3.77 GeV
<1.7 $\times 10^{-4}$	90	7	ALBRECHT	88G ARG	$e^+e^-$ 10 GeV
<2.2 $\times 10^{-4}$	90	8	HAAS	88 CLEO	$e^+e^-$ 10 GeV

$\Gamma(\mu^+\mu^-) / \Gamma_{\text{total}}$   $\Gamma_{173}/\Gamma$

A test for the  $\Delta C = 1$  weak neutral current. Allowed by first-order weak interaction combined with electromagnetic interaction.

VALUE	CL%	EVTS	DOCUMENT ID	TECN	COMMENT
<b>&lt;4.1 <math>\times 10^{-6}</math></b>	90		ADAMOVICH	97 BEAT	$\pi^-$ Cu, W 350 GeV
● ● ● We do not use the following data for averages, fits, limits, etc. ● ● ●					
<1.56 $\times 10^{-5}$	90		PRIPSTEIN	00 E789	$p$ nucleus, 800 GeV
<5.2 $\times 10^{-6}$	90		AITALA	99G E791	$\pi^-$ N 500 GeV
<4.2 $\times 10^{-6}$	90		ALEXOPOU...	96 E771	$p$ Si, 800 GeV
<3.4 $\times 10^{-5}$	90	1	FREYBERGER	96 CLE2	$e^+e^- \approx \gamma(4S)$
<7.6 $\times 10^{-6}$	90	0	ADAMOVICH	95 BEAT	See ADAMOVICH 97
<4.4 $\times 10^{-5}$	90	0	KODAMA	95 E653	$\pi^-$ emulsion 600 GeV
<3.1 $\times 10^{-5}$	90		<sup>104</sup> MISHRA	94 E789	$-4.1 \pm 4.8$ events
<7.0 $\times 10^{-5}$	90	3	ALBRECHT	88G ARG	$e^+e^-$ 10 GeV
<1.1 $\times 10^{-5}$	90		LOUIS	86 SPEC	$\pi^-$ W 225 GeV
<3.4 $\times 10^{-4}$	90		AUBERT	85 EMC	Deep inelast. $\mu^-$ N

<sup>104</sup>Here MISHRA 94 uses "the statistical approach advocated by the PDG." For an alternate approach, giving a limit of  $9 \times 10^{-6}$  at 90% confidence level, see the paper.

$\Gamma(\pi^0 e^+ e^-)/\Gamma_{\text{total}}$   $\Gamma_{174}/\Gamma$

A test for the  $\Delta C = 1$  weak neutral current. Allowed by higher-order electroweak interactions.

VALUE	CL%	EVTS	DOCUMENT ID	TECN	COMMENT
$<4.5 \times 10^{-5}$	90	0	FREYBERGER 96	CLE2	$e^+ e^- \approx \gamma(4S)$

$\Gamma(\pi^0 \mu^+ \mu^-)/\Gamma_{\text{total}}$   $\Gamma_{175}/\Gamma$

A test for the  $\Delta C=1$  weak neutral current. Allowed by higher-order electroweak interactions.

VALUE	CL%	EVTS	DOCUMENT ID	TECN	COMMENT
$<1.8 \times 10^{-4}$	90	2	KODAMA	95 E653	$\pi^-$ emulsion 600 GeV
• • • We do not use the following data for averages, fits, limits, etc. • • •					
$<5.4 \times 10^{-4}$	90	3	FREYBERGER 96	CLE2	$e^+ e^- \approx \gamma(4S)$

$\Gamma(\eta e^+ e^-)/\Gamma_{\text{total}}$   $\Gamma_{176}/\Gamma$

A test for the  $\Delta C = 1$  weak neutral current. Allowed by higher-order electroweak interactions.

VALUE	CL%	EVTS	DOCUMENT ID	TECN	COMMENT
$<1.1 \times 10^{-4}$	90	0	FREYBERGER 96	CLE2	$e^+ e^- \approx \gamma(4S)$

$\Gamma(\eta \mu^+ \mu^-)/\Gamma_{\text{total}}$   $\Gamma_{177}/\Gamma$

A test for the  $\Delta C = 1$  weak neutral current. Allowed by higher-order electroweak interactions.

VALUE	CL%	EVTS	DOCUMENT ID	TECN	COMMENT
$<5.3 \times 10^{-4}$	90	0	FREYBERGER 96	CLE2	$e^+ e^- \approx \gamma(4S)$

$\Gamma(\pi^+ \pi^- e^+ e^-)/\Gamma_{\text{total}}$   $\Gamma_{178}/\Gamma$

A test for the  $\Delta C = 1$  weak neutral current. Allowed by higher-order electroweak interactions.

VALUE	CL%	EVTS	DOCUMENT ID	TECN	COMMENT
$<3.73 \times 10^{-4}$	90	9	AITALA	01C E791	$\pi^-$ nucleus, 500 GeV

$\Gamma(\rho^0 e^+ e^-)/\Gamma_{\text{total}}$   $\Gamma_{179}/\Gamma$

A test for the  $\Delta C = 1$  weak neutral current. Allowed by higher-order electroweak interactions.

VALUE	CL%	EVTS	DOCUMENT ID	TECN	COMMENT
$<1.0 \times 10^{-4}$	90	2	<sup>105</sup> FREYBERGER 96	CLE2	$e^+ e^- \approx \gamma(4S)$

• • • We do not use the following data for averages, fits, limits, etc. • • •

$<1.24 \times 10^{-4}$	90	1	AITALA	01C E791	$\pi^-$ nucleus, 500 GeV
$<4.5 \times 10^{-4}$	90	2	HAAS	88 CLEO	$e^+ e^-$ 10 GeV

<sup>105</sup> This FREYBERGER 96 limit is obtained using a phase-space model. The limit changes to  $< 1.8 \times 10^{-4}$  using a photon pole amplitude model.

$\Gamma(\pi^+ \pi^- \mu^+ \mu^-)/\Gamma_{\text{total}}$   $\Gamma_{180}/\Gamma$

A test for the  $\Delta C = 1$  weak neutral current. Allowed by higher-order electroweak interactions.

VALUE	CL%	EVTS	DOCUMENT ID	TECN	COMMENT
$<3.0 \times 10^{-5}$	90	2	AITALA	01C E791	$\pi^-$ nucleus, 500 GeV

**$\Gamma(\rho^0 \mu^+ \mu^-)/\Gamma_{\text{total}}$**   **$\Gamma_{181}/\Gamma$**

A test for the  $\Delta C = 1$  weak neutral current. Allowed by higher-order electroweak interactions.

VALUE	CL%	EVTS	DOCUMENT ID	TECN	COMMENT
<b><math>&lt;2.2 \times 10^{-5}</math></b>	90	0	AITALA	01C E791	$\pi^-$ nucleus, 500 GeV
• • • We do not use the following data for averages, fits, limits, etc. • • •					
$<4.9 \times 10^{-4}$	90	1	<sup>106</sup> FREYBERGER 96	CLE2	$e^+ e^- \approx \gamma(4S)$
$<2.3 \times 10^{-4}$	90	0	KODAMA	95 E653	$\pi^-$ emulsion 600 GeV
$<8.1 \times 10^{-4}$	90	5	HAAS	88 CLEO	$e^+ e^-$ 10 GeV

<sup>106</sup> This FREYBERGER 96 limit is obtained using a phase-space model. The limit changes to  $< 4.5 \times 10^{-4}$  using a photon pole amplitude model.

**$\Gamma(\omega e^+ e^-)/\Gamma_{\text{total}}$**   **$\Gamma_{182}/\Gamma$**

A test for the  $\Delta C = 1$  weak neutral current. Allowed by higher-order electroweak interactions.

VALUE	CL%	EVTS	DOCUMENT ID	TECN	COMMENT
<b><math>&lt;1.8 \times 10^{-4}</math></b>	90	1	<sup>107</sup> FREYBERGER 96	CLE2	$e^+ e^- \approx \gamma(4S)$

<sup>107</sup> This FREYBERGER 96 limit is obtained using a phase-space model. The limit changes to  $< 2.7 \times 10^{-4}$  using a photon pole amplitude model.

**$\Gamma(\omega \mu^+ \mu^-)/\Gamma_{\text{total}}$**   **$\Gamma_{183}/\Gamma$**

A test for the  $\Delta C = 1$  weak neutral current. Allowed by higher-order electroweak interactions.

VALUE	CL%	EVTS	DOCUMENT ID	TECN	COMMENT
<b><math>&lt;8.3 \times 10^{-4}</math></b>	90	0	<sup>108</sup> FREYBERGER 96	CLE2	$e^+ e^- \approx \gamma(4S)$

<sup>108</sup> This FREYBERGER 96 limit is obtained using a phase-space model. The limit changes to  $< 6.5 \times 10^{-4}$  using a photon pole amplitude model.

**$\Gamma(K^- K^+ e^+ e^-)/\Gamma_{\text{total}}$**   **$\Gamma_{184}/\Gamma$**

A test for the  $\Delta C = 1$  weak neutral current. Allowed by higher-order electroweak interactions.

VALUE	CL%	EVTS	DOCUMENT ID	TECN	COMMENT
<b><math>&lt;3.15 \times 10^{-4}</math></b>	90	9	AITALA	01C E791	$\pi^-$ nucleus, 500 GeV

**$\Gamma(\phi e^+ e^-)/\Gamma_{\text{total}}$**   **$\Gamma_{185}/\Gamma$**

A test for the  $\Delta C = 1$  weak neutral current. Allowed by higher-order electroweak interactions.

VALUE	CL%	EVTS	DOCUMENT ID	TECN	COMMENT
<b><math>&lt;5.2 \times 10^{-5}</math></b>	90	2	<sup>109</sup> FREYBERGER 96	CLE2	$e^+ e^- \approx \gamma(4S)$

• • • We do not use the following data for averages, fits, limits, etc. • • •

$<5.9 \times 10^{-5}$	90	0	AITALA	01C E791	$\pi^-$ nucleus, 500 GeV
-----------------------	----	---	--------	----------	--------------------------

<sup>109</sup> This FREYBERGER 96 limit is obtained using a phase-space model. The limit changes to  $< 7.6 \times 10^{-5}$  using a photon pole amplitude model.

**$\Gamma(K^- K^+ \mu^+ \mu^-)/\Gamma_{\text{total}}$**   **$\Gamma_{186}/\Gamma$**

A test for the  $\Delta C = 1$  weak neutral current. Allowed by higher-order electroweak interactions.

VALUE	CL%	EVTS	DOCUMENT ID	TECN	COMMENT
<b><math>&lt;3.3 \times 10^{-5}</math></b>	90	0	AITALA	01C E791	$\pi^-$ nucleus, 500 GeV

$\Gamma(\phi\mu^+\mu^-)/\Gamma_{\text{total}}$   $\Gamma_{187}/\Gamma$

A test for the  $\Delta C = 1$  weak neutral current. Allowed by higher-order electroweak interactions.

VALUE	CL%	EVTS	DOCUMENT ID	TECN	COMMENT
$<3.1 \times 10^{-5}$	90	0	AITALA	01C E791	$\pi^-$ nucleus, 500 GeV

• • • We do not use the following data for averages, fits, limits, etc. • • •

$<4.1 \times 10^{-4}$  90 0 <sup>110</sup> FREYBERGER 96 CLE2  $e^+e^- \approx \gamma(4S)$

<sup>110</sup>This FREYBERGER 96 limit is obtained using a phase-space model. The limit changes to  $<2.4 \times 10^{-4}$  using a photon pole amplitude model.

$\Gamma(\bar{K}^0 e^+ e^-)/\Gamma_{\text{total}}$   $\Gamma_{188}/\Gamma$

Not a useful test for  $\Delta C = 1$  weak neutral current because both quarks must change flavor.

VALUE	CL%	EVTS	DOCUMENT ID	TECN	COMMENT
$<1.1 \times 10^{-4}$	90	0	FREYBERGER 96	CLE2	$e^+e^- \approx \gamma(4S)$

• • • We do not use the following data for averages, fits, limits, etc. • • •

$<1.7 \times 10^{-3}$  90 ADLER 89C MRK3  $e^+e^-$  3.77 GeV

$\Gamma(\bar{K}^0 \mu^+ \mu^-)/\Gamma_{\text{total}}$   $\Gamma_{189}/\Gamma$

Not a useful test for  $\Delta C = 1$  weak neutral current because both quarks must change flavor.

VALUE	CL%	EVTS	DOCUMENT ID	TECN	COMMENT
$<2.6 \times 10^{-4}$	90	2	KODAMA	95 E653	$\pi^-$ emulsion 600 GeV

• • • We do not use the following data for averages, fits, limits, etc. • • •

$<6.7 \times 10^{-4}$  90 1 FREYBERGER 96 CLE2  $e^+e^- \approx \gamma(4S)$

$\Gamma(K^- \pi^+ e^+ e^-)/\Gamma_{\text{total}}$   $\Gamma_{190}/\Gamma$

A test for the  $\Delta C = 1$  weak neutral current. Allowed by higher-order electroweak interactions.

VALUE	CL%	EVTS	DOCUMENT ID	TECN	COMMENT
$<3.85 \times 10^{-4}$	90	6	AITALA	01C E791	$\pi^-$ nucleus, 500 GeV

$\Gamma(\bar{K}^*(892)^0 e^+ e^-)/\Gamma_{\text{total}}$   $\Gamma_{191}/\Gamma$

Not a useful test for  $\Delta C = 1$  weak neutral current because both quarks must change flavor.

VALUE	CL%	EVTS	DOCUMENT ID	TECN	COMMENT
$<4.7 \times 10^{-5}$	90	2	AITALA	01C E791	$\pi^-$ nucleus, 500 GeV

• • • We do not use the following data for averages, fits, limits, etc. • • •

$<1.4 \times 10^{-4}$  90 1 <sup>111</sup> FREYBERGER 96 CLE2  $e^+e^- \approx \gamma(4S)$

<sup>111</sup>This FREYBERGER 96 limit is obtained using a phase-space model. The limit changes to  $<2.0 \times 10^{-4}$  using a photon pole amplitude model.

$\Gamma(K^- \pi^+ \mu^+ \mu^-)/\Gamma_{\text{total}}$   $\Gamma_{192}/\Gamma$

A test for the  $\Delta C = 1$  weak neutral current. Allowed by higher-order electroweak interactions.

VALUE	CL%	EVTS	DOCUMENT ID	TECN	COMMENT
$<3.59 \times 10^{-4}$	90	12	AITALA	01C E791	$\pi^-$ nucleus, 500 GeV



$\Gamma(\overline{K}^*(892)^0 \mu^+ \mu^-) / \Gamma_{\text{total}}$   $\Gamma_{193} / \Gamma$

Not a useful test for  $\Delta C = 1$  weak neutral current because both quarks must change flavor.

VALUE	CL%	EVTS	DOCUMENT ID	TECN	COMMENT
$< 2.4 \times 10^{-5}$	90	3	AITALA	01C E791	$\pi^-$ nucleus, 500 GeV

• • • We do not use the following data for averages, fits, limits, etc. • • •

$< 1.18 \times 10^{-3}$  90 1 <sup>112</sup> FREYBERGER 96 CLE2  $e^+ e^- \approx \gamma(4S)$

<sup>112</sup> This FREYBERGER 96 limit is obtained using a phase-space model. The limit changes to  $< 1.0 \times 10^{-3}$  using a photon pole amplitude model.

$\Gamma(\pi^+ \pi^- \pi^0 \mu^+ \mu^-) / \Gamma_{\text{total}}$   $\Gamma_{194} / \Gamma$

A test for the  $\Delta C = 1$  weak neutral current. Allowed by higher-order electroweak interactions.

VALUE	CL%	EVTS	DOCUMENT ID	TECN	COMMENT
$< 8.1 \times 10^{-4}$	90	1	KODAMA	95 E653	$\pi^-$ emulsion 600 GeV

$\Gamma(\mu^\pm e^\mp) / \Gamma_{\text{total}}$   $\Gamma_{195} / \Gamma$

A test of lepton family number conservation.

VALUE	CL%	EVTS	DOCUMENT ID	TECN	COMMENT
$< 8.1 \times 10^{-6}$	90		AITALA	99G E791	$\pi^- N$ 500 GeV

• • • We do not use the following data for averages, fits, limits, etc. • • •

$< 1.72 \times 10^{-5}$  90 PRIPSTEIN 00 E789  $p$  nucleus, 800 GeV

$< 1.9 \times 10^{-5}$  90 2 <sup>113</sup> FREYBERGER 96 CLE2  $e^+ e^- \approx \gamma(4S)$

$< 1.0 \times 10^{-4}$  90 4 ALBRECHT 88G ARG  $e^+ e^-$  10 GeV

$< 2.7 \times 10^{-4}$  90 9 HAAS 88 CLEO  $e^+ e^-$  10 GeV

$< 1.2 \times 10^{-4}$  90 BECKER 87C MRK3  $e^+ e^-$  3.77 GeV

$< 9 \times 10^{-4}$  90 PALKA 87 SILI 200 GeV  $\pi p$

$< 21 \times 10^{-4}$  90 0 <sup>114</sup> RILES 87 MRK2  $e^+ e^-$  29 GeV

<sup>113</sup> This is the corrected result given in the erratum to FREYBERGER 96.

<sup>114</sup> RILES 87 assumes  $B(D \rightarrow K\pi) = 3.0\%$  and has production model dependency.

$\Gamma(\pi^0 e^\pm \mu^\mp) / \Gamma_{\text{total}}$   $\Gamma_{196} / \Gamma$

A test of lepton family number conservation. The value is for the sum of the two charge states.

VALUE	CL%	EVTS	DOCUMENT ID	TECN	COMMENT
$< 8.6 \times 10^{-5}$	90	2	FREYBERGER 96	CLE2	$e^+ e^- \approx \gamma(4S)$

$\Gamma(\eta e^\pm \mu^\mp) / \Gamma_{\text{total}}$   $\Gamma_{197} / \Gamma$

A test of lepton family number conservation. The value is for the sum of the two charge states.

VALUE	CL%	EVTS	DOCUMENT ID	TECN	COMMENT
$< 1.0 \times 10^{-4}$	90	0	FREYBERGER 96	CLE2	$e^+ e^- \approx \gamma(4S)$

$\Gamma(\pi^+ \pi^- e^\pm \mu^\mp) / \Gamma_{\text{total}}$   $\Gamma_{198} / \Gamma$

A test of lepton family-number conservation. The value is for the sum of the two charge states.

VALUE	CL%	EVTS	DOCUMENT ID	TECN	COMMENT
$< 1.5 \times 10^{-5}$	90	1	AITALA	01C E791	$\pi^-$ nucleus, 500 GeV

**$\Gamma(\rho^0 e^\pm \mu^\mp)/\Gamma_{\text{total}}$**   **$\Gamma_{199}/\Gamma$**

A test of lepton family number conservation. The value is for the sum of the two charge states.

<u>VALUE</u>	<u>CL%</u>	<u>EVTS</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
$<4.9 \times 10^{-5}$	90	0	<sup>115</sup> FREYBERGER 96	CLE2	$e^+ e^- \approx \gamma(4S)$

• • • We do not use the following data for averages, fits, limits, etc. • • •

$<6.6 \times 10^{-5}$	90	1	AITALA	01C E791	$\pi^-$ nucleus, 500 GeV
-----------------------	----	---	--------	----------	--------------------------

<sup>115</sup> This FREYBERGER 96 limit is obtained using a phase-space model. The limit changes to  $< 5.0 \times 10^{-5}$  using a photon pole amplitude model.

**$\Gamma(\omega e^\pm \mu^\mp)/\Gamma_{\text{total}}$**   **$\Gamma_{200}/\Gamma$**

A test of lepton family number conservation. The value is for the sum of the two charge states.

<u>VALUE</u>	<u>CL%</u>	<u>EVTS</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
$<1.2 \times 10^{-4}$	90	0	<sup>116</sup> FREYBERGER 96	CLE2	$e^+ e^- \approx \gamma(4S)$

<sup>116</sup> This FREYBERGER 96 limit is obtained using a phase-space model. The same limit is obtained using a photon pole amplitude model.

**$\Gamma(K^- K^+ e^\pm \mu^\mp)/\Gamma_{\text{total}}$**   **$\Gamma_{201}/\Gamma$**

A test of lepton family-number conservation. The value is for the sum of the two charge states.

<u>VALUE</u>	<u>CL%</u>	<u>EVTS</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
$<1.8 \times 10^{-4}$	90	5	AITALA	01C E791	$\pi^-$ nucleus, 500 GeV

**$\Gamma(\phi e^\pm \mu^\mp)/\Gamma_{\text{total}}$**   **$\Gamma_{202}/\Gamma$**

A test of lepton family number conservation. The value is for the sum of the two charge states.

<u>VALUE</u>	<u>CL%</u>	<u>EVTS</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
$<3.4 \times 10^{-5}$	90	0	<sup>117</sup> FREYBERGER 96	CLE2	$e^+ e^- \approx \gamma(4S)$

• • • We do not use the following data for averages, fits, limits, etc. • • •

$<4.7 \times 10^{-5}$	90	0	AITALA	01C E791	$\pi^-$ nucleus, 500 GeV
-----------------------	----	---	--------	----------	--------------------------

<sup>117</sup> This FREYBERGER 96 limit is obtained using a phase-space model. The limit changes to  $< 3.3 \times 10^{-5}$  using a photon pole amplitude model.

**$\Gamma(\bar{K}^0 e^\pm \mu^\mp)/\Gamma_{\text{total}}$**   **$\Gamma_{203}/\Gamma$**

A test of lepton family number conservation. The value is for the sum of the two charge states.

<u>VALUE</u>	<u>CL%</u>	<u>EVTS</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
$<1.0 \times 10^{-4}$	90	0	FREYBERGER 96	CLE2	$e^+ e^- \approx \gamma(4S)$

**$\Gamma(K^- \pi^+ e^\pm \mu^\mp)/\Gamma_{\text{total}}$**   **$\Gamma_{204}/\Gamma$**

A test of lepton family-number conservation. The value is for the sum of the two charge states.

<u>VALUE</u>	<u>CL%</u>	<u>EVTS</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
$<5.53 \times 10^{-4}$	90	15	AITALA	01C E791	$\pi^-$ nucleus, 500 GeV

$\Gamma(\overline{K}^*(892)^0 e^\pm \mu^\mp)/\Gamma_{\text{total}}$   $\Gamma_{205}/\Gamma$

A test of lepton family number conservation. The value is for the sum of the two charge states.

VALUE	CL%	EVTS	DOCUMENT ID	TECN	COMMENT
$<8.3 \times 10^{-5}$	90	9	AITALA	01C E791	$\pi^-$ nucleus, 500 GeV

• • • We do not use the following data for averages, fits, limits, etc. • • •

$<1.0 \times 10^{-4}$	90	0	<sup>118</sup> FREYBERGER 96	CLE2	$e^+ e^- \approx \gamma(4S)$
-----------------------	----	---	------------------------------	------	------------------------------

<sup>118</sup>This FREYBERGER 96 limit is obtained using a phase-space model. The same limit is obtained using a photon pole amplitude model.

$\Gamma(\pi^- \pi^- e^+ e^+ + \text{c.c.})/\Gamma_{\text{total}}$   $\Gamma_{206}/\Gamma$

A test of lepton-number conservation. The value is for the sum of the two charge states.

VALUE	CL%	EVTS	DOCUMENT ID	TECN	COMMENT
$<1.12 \times 10^{-4}$	90	1	AITALA	01C E791	$\pi^-$ nucleus, 500 GeV

$\Gamma(\pi^- \pi^- \mu^+ \mu^+ + \text{c.c.})/\Gamma_{\text{total}}$   $\Gamma_{207}/\Gamma$

A test of lepton-number conservation. The value is for the sum of the two charge states.

VALUE	CL%	EVTS	DOCUMENT ID	TECN	COMMENT
$<2.9 \times 10^{-5}$	90	1	AITALA	01C E791	$\pi^-$ nucleus, 500 GeV

$\Gamma(K^- \pi^- e^+ e^+ + \text{c.c.})/\Gamma_{\text{total}}$   $\Gamma_{208}/\Gamma$

A test of lepton-number conservation. The value is for the sum of the two charge states.

VALUE	CL%	EVTS	DOCUMENT ID	TECN	COMMENT
$<2.06 \times 10^{-4}$	90	2	AITALA	01C E791	$\pi^-$ nucleus, 500 GeV

$\Gamma(K^- \pi^- \mu^+ \mu^+ + \text{c.c.})/\Gamma_{\text{total}}$   $\Gamma_{209}/\Gamma$

A test of lepton-number conservation. The value is for the sum of the two charge states.

VALUE	CL%	EVTS	DOCUMENT ID	TECN	COMMENT
$<3.9 \times 10^{-4}$	90	14	AITALA	01C E791	$\pi^-$ nucleus, 500 GeV

$\Gamma(K^- K^- e^+ e^+ + \text{c.c.})/\Gamma_{\text{total}}$   $\Gamma_{210}/\Gamma$

A test of lepton-number conservation. The value is for the sum of the two charge states.

VALUE	CL%	EVTS	DOCUMENT ID	TECN	COMMENT
$<1.52 \times 10^{-4}$	90	2	AITALA	01C E791	$\pi^-$ nucleus, 500 GeV

$\Gamma(K^- K^- \mu^+ \mu^+ + \text{c.c.})/\Gamma_{\text{total}}$   $\Gamma_{211}/\Gamma$

A test of lepton-number conservation. The value is for the sum of the two charge states.

VALUE	CL%	EVTS	DOCUMENT ID	TECN	COMMENT
$<9.4 \times 10^{-5}$	90	1	AITALA	01C E791	$\pi^-$ nucleus, 500 GeV

$\Gamma(\pi^- \pi^- e^+ \mu^+ + \text{c.c.})/\Gamma_{\text{total}}$   $\Gamma_{212}/\Gamma$

A test of lepton-number conservation. The value is for the sum of the two charge states.

VALUE	CL%	EVTS	DOCUMENT ID	TECN	COMMENT
$<7.9 \times 10^{-5}$	90	4	AITALA	01C E791	$\pi^-$ nucleus, 500 GeV

$\Gamma(K^- \pi^- e^+ \mu^+ + \text{c.c.})/\Gamma_{\text{total}}$   $\Gamma_{213}/\Gamma$

A test of lepton-number conservation. The value is for the sum of the two charge states.

VALUE	CL%	EVTS	DOCUMENT ID	TECN	COMMENT
$<2.18 \times 10^{-4}$	90	7	AITALA	01C E791	$\pi^-$ nucleus, 500 GeV

$\Gamma(K^- K^- e^+ \mu^+ + \text{c.c.})/\Gamma_{\text{total}}$   $\Gamma_{214}/\Gamma$

A test of lepton-number conservation. The value is for the sum of the two charge states.

VALUE	CL%	EVTS	DOCUMENT ID	TECN	COMMENT
$<5.7 \times 10^{-5}$	90	0	AITALA	01C E791	$\pi^-$ nucleus, 500 GeV

**$D^0$  CP-VIOLATING DECAY-RATE ASYMMETRIES**

**$A_{CP}(K^+ K^-)$  in  $D^0, \bar{D}^0 \rightarrow K^+ K^-$**

This is the difference between  $D^0$  and  $\bar{D}^0$  partial widths for these modes divided by the sum of the widths. The  $D^0$  and  $\bar{D}^0$  are distinguished by the charge of the parent  $D^*$ :  $D^{*+} \rightarrow D^0 \pi^+$  and  $D^{*-} \rightarrow \bar{D}^0 \pi^-$ .

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
<b>0.005 ± 0.016 OUR AVERAGE</b>				
$0.000 \pm 0.022 \pm 0.008$	3023	<sup>119</sup> CSORNA	02 CLE2	$e^+ e^- \approx \Upsilon(4S)$
$-0.001 \pm 0.022 \pm 0.015$	3330	<sup>119</sup> LINK	00B FOCS	
$-0.010 \pm 0.049 \pm 0.012$	609	<sup>119</sup> AITALA	98C E791	$-0.093 < A_{CP} < +0.073$ (90% CL)
$+0.080 \pm 0.061$		BARTELT	95 CLE2	$-0.022 < A_{CP} < +0.18$ (90%CL)
$+0.024 \pm 0.084$		<sup>119</sup> FRABETTI	94I E687	$-0.11 < A_{CP} < +0.16$ (90% CL)

<sup>119</sup> FRABETTI 94I, AITALA 98C, LINK 00B, and CSORNA 02 measure  $N(D^0 \rightarrow K^+ K^-)/N(D^0 \rightarrow K^- \pi^+)$ , the ratio of numbers of events observed, and similarly for the  $\bar{D}^0$ .

**$A_{CP}(K_S^0 K_S^0)$  in  $D^0, \bar{D}^0 \rightarrow K_S^0 K_S^0$**

This is the difference between  $D^0$  and  $\bar{D}^0$  partial widths for these modes divided by the sum of the widths. The  $D^0$  and  $\bar{D}^0$  are distinguished by the charge of the parent  $D^*$ :  $D^{*+} \rightarrow D^0 \pi^+$  and  $D^{*-} \rightarrow \bar{D}^0 \pi^-$ .

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
<b><math>-0.23 \pm 0.19</math></b>	65	BONVICINI	01 CLE2	$e^+ e^- \approx 10.6$ GeV

**$A_{CP}(\pi^+ \pi^-)$  in  $D^0, \bar{D}^0 \rightarrow \pi^+ \pi^-$**

This is the difference between  $D^0$  and  $\bar{D}^0$  partial widths for these modes divided by the sum of the widths. The  $D^0$  and  $\bar{D}^0$  are distinguished by the charge of the parent  $D^*$ :  $D^{*+} \rightarrow D^0 \pi^+$  and  $D^{*-} \rightarrow \bar{D}^0 \pi^-$ .

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
<b>0.021 ± 0.026 OUR AVERAGE</b>				
$0.019 \pm 0.032 \pm 0.008$	1136	<sup>120</sup> CSORNA	02 CLE2	$e^+ e^- \approx \Upsilon(4S)$
$+0.048 \pm 0.039 \pm 0.025$	1177	<sup>120</sup> LINK	00B FOCS	
$-0.049 \pm 0.078 \pm 0.030$	343	<sup>120</sup> AITALA	98C E791	$-0.186 < A_{CP} < +0.088$ (90% CL)

<sup>120</sup> AITALA 98C, LINK 00B, and CSORNA 02 measure  $N(D^0 \rightarrow \pi^+ \pi^-)/N(D^0 \rightarrow K^- \pi^+)$ , the ratio of numbers of events observed, and similarly for the  $\bar{D}^0$ .

### $A_{CP}(\pi^0\pi^0)$ in $D^0, \bar{D}^0 \rightarrow \pi^0\pi^0$

This is the difference between  $D^0$  and  $\bar{D}^0$  partial widths for these modes divided by the sum of the widths. The  $D^0$  and  $\bar{D}^0$  are distinguished by the charge of the parent  $D^*$ :  $D^{*+} \rightarrow D^0\pi^+$  and  $D^{*-} \rightarrow \bar{D}^0\pi^-$ .

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
$+0.001 \pm 0.048$	810	BONVICINI	01	CLE2 $e^+e^- \approx 10.6$ GeV

### $A_{CP}(K_S^0\phi)$ in $D^0, \bar{D}^0 \rightarrow K_S^0\phi$

This is the difference between  $D^0$  and  $\bar{D}^0$  partial widths for these modes divided by the sum of the widths. The  $D^0$  and  $\bar{D}^0$  are distinguished by the charge of the parent  $D^*$ :  $D^{*+} \rightarrow D^0\pi^+$  and  $D^{*-} \rightarrow \bar{D}^0\pi^-$ .

VALUE	DOCUMENT ID	TECN	COMMENT
$-0.028 \pm 0.094$	BARTELT 95	CLE2	$-0.182 < A_{CP} < +0.126$ (90%CL)

### $A_{CP}(K_S^0\pi^0)$ in $D^0, \bar{D}^0 \rightarrow K_S^0\pi^0$

This is the difference between  $D^0$  and  $\bar{D}^0$  partial widths for these modes divided by the sum of the widths. The  $D^0$  and  $\bar{D}^0$  are distinguished by the charge of the parent  $D^*$ :  $D^{*+} \rightarrow D^0\pi^+$  and  $D^{*-} \rightarrow \bar{D}^0\pi^-$ .

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
$+0.001 \pm 0.013$	9099	BONVICINI	01	CLE2 $e^+e^- \approx 10.6$ GeV

• • • We do not use the following data for averages, fits, limits, etc. • • •

$-0.018 \pm 0.030$	BARTELT 95	CLE2	See BONVICINI 01
--------------------	------------	------	------------------

### $A_{CP}(K^\pm\pi^\mp)$ in $D^0 \rightarrow K^+\pi^-, \bar{D}^0 \rightarrow K^-\pi^+$

This is the difference between  $D^0$  and  $\bar{D}^0$  partial widths for these modes divided by the sum of the widths. The  $D^0$  and  $\bar{D}^0$  are distinguished by the charge of the parent  $D^*$ :  $D^{*+} \rightarrow D^0\pi^+$  and  $D^{*-} \rightarrow \bar{D}^0\pi^-$ .

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
$+0.02^{+0.19}_{-0.20} \pm 0.01$	45	<sup>121</sup> GODANG	00	CLE2 $-0.43 < A_{CP} < +0.34$ (95%CL)

<sup>121</sup>This GODANG 00 result assumes no  $D^0$ - $\bar{D}^0$  mixing; it becomes  $-0.01^{+0.16}_{-0.17} \pm 0.01$  when mixing is allowed.

### $A_{CP}(K^\mp\pi^\pm\pi^0)$ in $D^0 \rightarrow K^-\pi^+\pi^0, \bar{D}^0 \rightarrow K^+\pi^-\pi^0$

This is the difference between  $D^0$  and  $\bar{D}^0$  partial widths for these modes divided by the sum of the widths. The  $D^0$  and  $\bar{D}^0$  are distinguished by the charge of the parent  $D^*$ :  $D^{*+} \rightarrow D^0\pi^+$  and  $D^{*-} \rightarrow \bar{D}^0\pi^-$ .

VALUE	DOCUMENT ID	TECN	COMMENT
$-0.031 \pm 0.086$	<sup>122</sup> KOPP	01	CLE2 $e^+e^- \approx 10.6$ GeV

<sup>122</sup>KOPP 01 fits separately the  $D^0$  and  $\bar{D}^0$  Dalitz plots and then calculates the integrated difference of normalized densities divided by the integrated sum.

### $A_{CP}(K^\pm\pi^\mp\pi^0)$ in $D^0 \rightarrow K^+\pi^-\pi^0, \bar{D}^0 \rightarrow K^-\pi^+\pi^0$

This is the difference between  $D^0$  and  $\bar{D}^0$  partial widths for these modes divided by the sum of the widths. The  $D^0$  and  $\bar{D}^0$  are distinguished by the charge of the parent  $D^*$ :  $D^{*+} \rightarrow D^0\pi^+$  and  $D^{*-} \rightarrow \bar{D}^0\pi^-$ .

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
$+0.09^{+0.25}_{-0.22}$	38	BRANDENB...	01	CLE2 $e^+e^- \approx \gamma(4S)$

## $D^0$ PRODUCTION CROSS SECTION AT $\psi(3770)$

A compilation of the cross sections for the direct production of  $D^0$  mesons at or near the  $\psi(3770)$  peak in  $e^+e^-$  production.

VALUE (nanobarns)	DOCUMENT ID	TECN	COMMENT
• • • We do not use the following data for averages, fits, limits, etc. • • •			
$5.8 \pm 0.5 \pm 0.6$	123 ADLER	88C MRK3	$e^+e^-$ 3.768 GeV
$7.3 \pm 1.3$	124 PARTRIDGE	84 CBAL	$e^+e^-$ 3.771 GeV
$8.00 \pm 0.95 \pm 1.21$	125 SCHINDLER	80 MRK2	$e^+e^-$ 3.771 GeV
$11.5 \pm 2.5$	126 PERUZZI	77 MRK1	$e^+e^-$ 3.774 GeV
123	This measurement compares events with one detected $D$ to those with two detected $D$ mesons, to determine the the absolute cross section. ADLER 88C find the ratio of cross sections (neutral to charged) to be $1.36 \pm 0.23 \pm 0.14$ .		
124	This measurement comes from a scan of the $\psi(3770)$ resonance and a fit to the cross section. PARTRIDGE 84 measures $6.4 \pm 1.15$ nb for the cross section. We take the phase space division of neutral and charged $D$ mesons in $\psi(3770)$ decay to be 1.33, and we assume that the $\psi(3770)$ is an isosinglet to evaluate the cross sections. The noncharm decays (e.g. radiative) of the $\psi(3770)$ are included in this measurement and may amount to a few percent correction.		
125	This measurement comes from a scan of the $\psi(3770)$ resonance and a fit to the cross section. SCHINDLER 80 assume the phase space division of neutral and charged $D$ mesons in $\psi(3770)$ decay to be 1.33, and that the $\psi(3770)$ is an isosinglet. The noncharm decays (e.g. radiative) of the $\psi(3770)$ are included in this measurement and may amount to a few percent correction.		
126	This measurement comes from a scan of the $\psi(3770)$ resonance and a fit to the cross section. The phase space division of neutral and charged $D$ mesons in $\psi(3770)$ decay is taken to be 1.33, and $\psi(3770)$ is assumed to be an isosinglet. The noncharm decays (e.g. radiative) of the $\psi(3770)$ are included in this measurement and may amount to a few percent correction. We exclude this measurement from the average because of uncertainties in the contamination from $\tau$ lepton pairs. Also see RAPIDIS 77.		

## $D^0$ REFERENCES

ABE	02I	PRL 88 162001	K. Abe <i>et al.</i>	(KEK BELLE Collab.)
CSORNA	02	PR D (to be publ.)	S.E. Csorna <i>et al.</i>	(CLEO Collab.)
		hep-ex/0111024		
PDG	02	PR D66 010001	K. Hagiwara <i>et al.</i>	
AITALA	01C	PRL 86 3969	E.M. Aitala <i>et al.</i>	(FNAL E791 Collab.)
AITALA	01D	PR D64 112003	E.M. Aitala <i>et al.</i>	(FNAL E791 Collab.)
BONVICINI	01	PR D63 071101R	G. Bonvicini <i>et al.</i>	(CLEO Collab.)
BRANDENB...	01	PRL 87 071802	G. Brandenburg <i>et al.</i>	(CLEO Collab.)
DYTMAN	01	PR D64 111101 (R)	S.A. Dytman <i>et al.</i>	(CLEO Collab.)
KOPP	01	PR D63 092001	S. Kopp <i>et al.</i>	(CLEO Collab.)
KUSHNIR...	01	PRL 86 5243	A. Kushnirenko <i>et al.</i>	(FNAL SELEX Collab.)
LINK	01	PRL 86 2955	J.M. Link <i>et al.</i>	(FNAL FOCUS Collab.)
BAI	00C	PR D62 052001	J.Z. Bai <i>et al.</i>	(BEPC BES Collab.)
GODANG	00	PRL 84 5038	R. Godang <i>et al.</i>	(CLEO Collab.)
JUN	00	PRL 84 1857	S.Y. Jun <i>et al.</i>	(FNAL SELEX Collab.)
LINK	00	PL B485 62	J.M. Link <i>et al.</i>	(FNAL FOCUS Collab.)
LINK	00B	PL B491 232	J.M. Link <i>et al.</i>	(FNAL FOCUS Collab.)
Also	00D	PL B495 443 (errata)	J.M. Link <i>et al.</i>	(FNAL FOCUS Collab.)
PRIPSTEIN	00	PR D61 032005	D. Pripstein <i>et al.</i>	(FNAL E789 Collab.)
AITALA	99E	PRL 83 32	E.M. Aitala <i>et al.</i>	(FNAL E791 Collab.)
AITALA	99G	PL B462 401	E.M. Aitala <i>et al.</i>	(FNAL E791 Collab.)
BONVICINI	99	PRL 82 4586	G. Bonvicini <i>et al.</i>	(CLEO Collab.)
AITALA	98	PR D57 13	E.M. Aitala <i>et al.</i>	(FNAL E791 Collab.)
AITALA	98C	PL B421 405	E.M. Aitala <i>et al.</i>	(FNAL E791 Collab.)
AITALA	98D	PL B423 185	E.M. Aitala <i>et al.</i>	(FNAL E791 Collab.)
ARTUSO	98	PRL 80 3193	M. Artuso <i>et al.</i>	(CLEO Collab.)

ASNER	98	PR D58 092001	D.M. Asner <i>et al.</i>	(CLEO Collab.)
BARATE	98W	PL B436 211	R. Barate <i>et al.</i>	(ALEPH Collab.)
COAN	98	PRL 80 1150	T.E. Coan <i>et al.</i>	(CLEO Collab.)
ADAMOVICH	97	PL B408 469	M.I. Adamovich <i>et al.</i>	(CERN BEATRICE Collab.)
BARATE	97C	PL B403 367	R. Barate <i>et al.</i>	(ALEPH Collab.)
AITALA	96C	PRL 77 2384	E.M. Aitala <i>et al.</i>	(FNAL E791 Collab.)
ALBRECHT	96C	PL B374 249	H. Albrecht <i>et al.</i>	(ARGUS Collab.)
ALEXOPOU...	96	PRL 77 2380	T. Alexopoulos <i>et al.</i>	(FNAL E771 Collab.)
ASNER	96B	PR D54 4211	D.M. Asner <i>et al.</i>	(CLEO Collab.)
BARISH	96	PL B373 334	B.C. Barish <i>et al.</i>	(CLEO Collab.)
FRABETTI	96B	PL B382 312	P.L. Frabetti <i>et al.</i>	(FNAL E687 Collab.)
FREYBERGER	96	PRL 76 3065	A. Freyberger <i>et al.</i>	(CLEO Collab.)
Also	96B	PRL 77 2147 (errata)	A. Freyberger <i>et al.</i>	(CLEO Collab.)
KUBOTA	96B	PR D54 2994	Y. Kubota <i>et al.</i>	(CLEO Collab.)
ADAMOVICH	95	PL B353 563	M.I. Adamovich <i>et al.</i>	(CERN BEATRICE Collab.)
BARTELT	95	PR D52 4860	J.E. Bartelt <i>et al.</i>	(CLEO Collab.)
BUTLER	95	PR D52 2656	F. Butler <i>et al.</i>	(CLEO Collab.)
FRABETTI	95C	PL B354 486	P.L. Frabetti <i>et al.</i>	(FNAL E687 Collab.)
FRABETTI	95G	PL B364 127	P.L. Frabetti <i>et al.</i>	(FNAL E687 Collab.)
KODAMA	95	PL B345 85	K. Kodama <i>et al.</i>	(FNAL E653 Collab.)
ALBRECHT	94	PL B324 249	H. Albrecht <i>et al.</i>	(ARGUS Collab.)
ALBRECHT	94F	PL B340 125	H. Albrecht <i>et al.</i>	(ARGUS Collab.)
ALBRECHT	94I	ZPHY C64 375	H. Albrecht <i>et al.</i>	(ARGUS Collab.)
CINABRO	94	PRL 72 1406	D. Cinabro <i>et al.</i>	(CLEO Collab.)
FRABETTI	94C	PL B321 295	P.L. Frabetti <i>et al.</i>	(FNAL E687 Collab.)
FRABETTI	94D	PL B323 459	P.L. Frabetti <i>et al.</i>	(FNAL E687 Collab.)
FRABETTI	94G	PL B331 217	P.L. Frabetti <i>et al.</i>	(FNAL E687 Collab.)
FRABETTI	94I	PR D50 R2953	P.L. Frabetti <i>et al.</i>	(FNAL E687 Collab.)
FRABETTI	94J	PL B340 254	P.L. Frabetti <i>et al.</i>	(FNAL E687 Collab.)
KODAMA	94	PL B336 605	K. Kodama <i>et al.</i>	(FNAL E653 Collab.)
MISHRA	94	PR D50 R9	C.S. Mishra <i>et al.</i>	(FNAL E789 Collab.)
AKERIB	93	PRL 71 3070	D.S. Akerib <i>et al.</i>	(CLEO Collab.)
ALBRECHT	93D	PL B308 435	H. Albrecht <i>et al.</i>	(ARGUS Collab.)
ANJOS	93	PR D48 56	J.C. Anjos <i>et al.</i>	(FNAL E691 Collab.)
BEAN	93C	PL B317 647	A. Bean <i>et al.</i>	(CLEO Collab.)
FRABETTI	93I	PL B315 203	P.L. Frabetti <i>et al.</i>	(FNAL E687 Collab.)
KODAMA	93B	PL B313 260	K. Kodama <i>et al.</i>	(FNAL E653 Collab.)
PROCARIO	93B	PR D48 4007	M. Procaro <i>et al.</i>	(CLEO Collab.)
SELEN	93	PRL 71 1973	M.A. Selen <i>et al.</i>	(CLEO Collab.)
ADAMOVICH	92	PL B280 163	M.I. Adamovich <i>et al.</i>	(CERN WA82 Collab.)
ALBRECHT	92P	ZPHY C56 7	H. Albrecht <i>et al.</i>	(ARGUS Collab.)
ANJOS	92B	PR D46 R1	J.C. Anjos <i>et al.</i>	(FNAL E691 Collab.)
ANJOS	92C	PR D46 1941	J.C. Anjos <i>et al.</i>	(FNAL E691 Collab.)
BARLAG	92C	ZPHY C55 383	S. Barlag <i>et al.</i>	(ACCMOR Collab.)
Also	90D	ZPHY C48 29	S. Barlag <i>et al.</i>	(ACCMOR Collab.)
COFFMAN	92B	PR D45 2196	D.M. Coffman <i>et al.</i>	(Mark III Collab.)
Also	90	PRL 64 2615	J. Adler <i>et al.</i>	(Mark III Collab.)
FRABETTI	92	PL B281 167	P.L. Frabetti <i>et al.</i>	(FNAL E687 Collab.)
FRABETTI	92B	PL B286 195	P.L. Frabetti <i>et al.</i>	(FNAL E687 Collab.)
ALVAREZ	91B	ZPHY C50 11	M.P. Alvarez <i>et al.</i>	(CERN NA14/2 Collab.)
AMMAR	91	PR D44 3383	R. Ammar <i>et al.</i>	(CLEO Collab.)
ANJOS	91	PR D43 R635	J.C. Anjos <i>et al.</i>	(FNAL-TPS Collab.)
ANJOS	91D	PR D44 R3371	J.C. Anjos <i>et al.</i>	(FNAL-TPS Collab.)
BAI	91	PRL 66 1011	Z. Bai <i>et al.</i>	(Mark III Collab.)
COFFMAN	91	PL B263 135	D.M. Coffman <i>et al.</i>	(Mark III Collab.)
CRAWFORD	91B	PR D44 3394	G. Crawford <i>et al.</i>	(CLEO Collab.)
DECAMP	91J	PL B266 218	D. Decamp <i>et al.</i>	(ALEPH Collab.)
FRABETTI	91	PL B263 584	P.L. Frabetti <i>et al.</i>	(FNAL E687 Collab.)
KINOSHITA	91	PR D43 2836	K. Kinoshita <i>et al.</i>	(CLEO Collab.)
KODAMA	91	PRL 66 1819	K. Kodama <i>et al.</i>	(FNAL E653 Collab.)
ALBRECHT	90C	ZPHY C46 9	H. Albrecht <i>et al.</i>	(ARGUS Collab.)
ALEXANDER	90	PRL 65 1184	J. Alexander <i>et al.</i>	(CLEO Collab.)
ALEXANDER	90B	PRL 65 1531	J. Alexander <i>et al.</i>	(CLEO Collab.)
ALVAREZ	90	ZPHY C47 539	M.P. Alvarez <i>et al.</i>	(CERN NA14/2 Collab.)
ANJOS	90D	PR D42 2414	J.C. Anjos <i>et al.</i>	(FNAL E691 Collab.)
BARLAG	90C	ZPHY C46 563	S. Barlag <i>et al.</i>	(ACCMOR Collab.)
ADLER	89	PRL 62 1821	J. Adler <i>et al.</i>	(Mark III Collab.)
ADLER	89C	PR D40 906	J. Adler <i>et al.</i>	(Mark III Collab.)
ALBRECHT	89D	ZPHY C43 181	H. Albrecht <i>et al.</i>	(ARGUS Collab.)
ANJOS	89F	PRL 62 1587	J.C. Anjos <i>et al.</i>	(FNAL E691 Collab.)
ABACHI	88	PL B205 411	S. Abachi <i>et al.</i>	(HRS Collab.)

ADLER	88	PR D37 2023	J. Adler <i>et al.</i>	(Mark III Collab.)
ADLER	88C	PRL 60 89	J. Adler <i>et al.</i>	(Mark III Collab.)
ALBRECHT	88G	PL B209 380	H. Albrecht <i>et al.</i>	(ARGUS Collab.)
ALBRECHT	88I	PL B210 267	H. Albrecht <i>et al.</i>	(ARGUS Collab.)
ANJOS	88C	PRL 60 1239	J.C. Anjos <i>et al.</i>	(FNAL E691 Collab.)
BORTOLETTO	88	PR D37 1719	D. Bortoletto <i>et al.</i>	(CLEO Collab.)
Also	89D	PR D39 1471 erratum	D. Bortoletto <i>et al.</i>	(CLEO Collab.)
CUMALAT	88	PL B210 253	J.P. Cumalat <i>et al.</i>	(E-400 Collab.)
HAAS	88	PRL 60 1614	P. Haas <i>et al.</i>	(CLEO Collab.)
RAAB	88	PR D37 2391	J.R. Raab <i>et al.</i>	(FNAL E691 Collab.)
ADAMOVIICH	87	EPL 4 887	M.I. Adamovich <i>et al.</i>	(Photon Emulsion Collab.)
ADLER	87	PL B196 107	J. Adler <i>et al.</i>	(Mark III Collab.)
AGUILAR-...	87E	ZPHY C36 551	M. Aguilar-Benitez <i>et al.</i>	(LEBC-EHS Collab.)
Also	88B	ZPHY C40 321	M. Aguilar-Benitez <i>et al.</i>	(LEBC-EHS Collab.)
AGUILAR-...	87F	ZPHY C36 559	M. Aguilar-Benitez <i>et al.</i>	(LEBC-EHS Collab.)
Also	88	ZPHY C38 520 erratum	M. Aguilar-Benitez <i>et al.</i>	(LEBC-EHS Collab.)
ALBRECHT	87E	ZPHY C33 359	H. Albrecht <i>et al.</i>	(ARGUS Collab.)
ALBRECHT	87K	PL B199 447	H. Albrecht <i>et al.</i>	(ARGUS Collab.)
BARLAG	87B	ZPHY C37 17	S. Barlag <i>et al.</i>	(ACCMOR Collab.)
BECKER	87C	PL B193 147	J.J. Becker <i>et al.</i>	(Mark III Collab.)
Also	87D	PL B198 590 erratum	J.J. Becker <i>et al.</i>	(Mark III Collab.)
PALKA	87	PL B189 238	H. Palka <i>et al.</i>	(ACCMOR Collab.)
RILES	87	PR D35 2914	K. Riles <i>et al.</i>	(Mark II Collab.)
BAILEY	86	ZPHY C30 51	R. Bailey <i>et al.</i>	(ACCMOR Collab.)
BEBEK	86	PRL 56 1893	C. Bebek <i>et al.</i>	(CLEO Collab.)
LOUIS	86	PRL 56 1027	W.C. Louis <i>et al.</i>	(PRIN, CHIC, ISU)
ALBRECHT	85B	PL 158B 525	H. Albrecht <i>et al.</i>	(ARGUS Collab.)
ALBRECHT	85F	PL 150B 235	H. Albrecht <i>et al.</i>	(ARGUS Collab.)
AUBERT	85	PL 155B 461	H. Aubert <i>et al.</i>	(EMC Collab.)
BALTRUSAIT-...	85B	PRL 54 1976	R.M. Baltrusaitis <i>et al.</i>	(Mark III Collab.)
BALTRUSAIT-...	85E	PRL 55 150	R.M. Baltrusaitis <i>et al.</i>	(Mark III Collab.)
BENVENUTI	85	PL 158B 531	A.C. Benvenuti <i>et al.</i>	(BCDMS Collab.)
ADAMOVIICH	84B	PL 140B 123	M.I. Adamovich <i>et al.</i>	(CERN WA58 Collab.)
DERRICK	84	PRL 53 1971	M. Derrick <i>et al.</i>	(HRS Collab.)
PARTRIDGE	84	Thesis CALT-68-1150	R.A. Partridge	(Crystal Ball Collab.)
SUMMERS	84	PRL 52 410	D.J. Summers <i>et al.</i>	(UCSB, CARL, COLO+)
BAILEY	83B	PL 132B 237	R. Bailey <i>et al.</i>	(ACCMOR Collab.)
BODEK	82	PL 113B 82	A. Bodek <i>et al.</i>	(ROCH, CIT, CHIC, FNAL+)
FIORINO	81	LNC 30 166	A. Fiorino <i>et al.</i>	
SCHINDLER	81	PR D24 78	R.H. Schindler <i>et al.</i>	(Mark II Collab.)
TRILLING	81	PRPL 75 57	G.H. Trilling	(LBL, UCB) J
ASTON	80E	PL 94B 113	D. Aston <i>et al.</i>	(BONN, CERN, EPOL, GLAS+)
VERY	80	PRL 44 1309	P. Avery <i>et al.</i>	(ILL, FNAL, COLU)
SCHINDLER	80	PR D21 2716	R.H. Schindler <i>et al.</i>	(Mark II Collab.)
ZHOLENTZ	80	PL 96B 214	A.A. Zholents <i>et al.</i>	(NOVO)
Also	81	SJNP 34 814	A.A. Zholents <i>et al.</i>	(NOVO)
		Translated from YAF 34	1471.	
ABRAMS	79D	PRL 43 481	G.S. Abrams <i>et al.</i>	(Mark II Collab.)
ATIYA	79	PRL 43 414	M.S. Atiya <i>et al.</i>	(COLU, ILL, FNAL)
BALTAY	78C	PRL 41 73	C. Baltay <i>et al.</i>	(COLU, BNL)
VUILLEMIN	78	PRL 41 1149	V. Vuillemin <i>et al.</i>	(Mark I Collab.)
GOLDHABER	77	PL 69B 503	G. Goldhaber <i>et al.</i>	(Mark I Collab.)
PERUZZI	77	PRL 39 1301	I. Peruzzi <i>et al.</i>	(Mark I Collab.)
PICCOLO	77	PL 70B 260	M. Piccolo <i>et al.</i>	(Mark I Collab.)
RAPIDIS	77	PRL 39 526	P.A. Rapidis <i>et al.</i>	(Mark I Collab.)
GOLDHABER	76	PRL 37 255	G. Goldhaber <i>et al.</i>	(Mark I Collab.)

## OTHER RELATED PAPERS

RICHMAN	95	RMP 67 893	J.D. Richman, P.R. Burchat	(UCSB, STAN)
ROSNER	95	CNPP 21 369	J. Rosner	(CHIC)