$$
J=1
$$

## THE MASS OF THE W BOSON

Revised March 2000 by C. Caso (Univ. of Genova) and A. Gurtu (Tata Inst.)

Till 1995 the production and study of the $W$ boson was the exclusive domain of the $\bar{p} p$ colliders at CERN and FNAL. $W$ production in these hadron colliders is tagged by a high $p_{T}$ lepton from $W$ decay. Owing to unknown parton-parton effective energy and missing energy in the longitudinal direction, the experiments reconstruct only the transverse mass of the $W$ and derive the $W$ mass from comparing the transverse mass distribution with Monte Carlo predictions as a function of $M_{W}$.

Beginning 1996 the energy of LEP increased to above 161 GeV , the threshold for $W$-pair production. A precise knowledge of the $e^{+} e^{-}$centre of mass energy enables one to reconstruct the $W$ mass even if one of them decays leptonically. At LEP two methods have been used to obtain the $W$ mass. In the first method the measured $W$-pair production cross sections, $\sigma\left(e^{+} e^{-} \rightarrow W^{+} W^{-}\right)$, have been used to determine the $W$ mass using the predicted dependence of this cross section on $M_{W}$ (see Fig. 1). At 161 GeV , which is just above the $W$-pair production threshold, this dependence is a much more sensitive function of the $W$ mass than at the higher energies ( 172 to 202 GeV ) at which LEP has run during 1996-99. In the second method, which is used at the higher energies, the $W$ mass has been determined by directly reconstructing the $W$ from its decay products.

Each LEP experiment has combined their own mass values properly taking into account the common systematic errors. In order to compute the LEP average $W$ mass each experiment


Figure 1: The $W$-pair cross section as a function of the center-of-mass energy. The data points are the LEP averages. The solid lines are predictions from different models of $W W$ production. For comparison the figure contains also the cross section if the ZWW coupling did not exist (dashed line), or if only the $t$-channel $\nu_{e}$ exchange diagram existed (dotted line).
has provided its measured $W$ mass for the $q q q q$ and $q q \ell \nu$ channels at each center-of-mass energy along with a detailed break-up of errors (statistical and uncorrelated, partially correlated and fully correlated systematics [1]). These have been properly combined to obtain a preliminary [2] LEP $W$ mass
$=80.401 \pm 0.048 \mathrm{GeV}$. Errors due uncertainties in LEP energy ( 17 MeV ) and possible effect of color reconnection (CR) and Bose-Einstein (BE) correlations between quarks from different $W^{\prime}$ 's ( 18 MeV ) are included. The mass difference between $q q q q$ and $q q \ell \nu$ final states (due to possible CR and BE effects) is $35 \pm 55 \mathrm{MeV}$.

The two Tevatron experiments have also carried out the exercise of identifying common systematic errors and averaging with CERN UA2 data obtain an average $W$ mass $=$ $80.448 \pm 0.062 \mathrm{GeV}$.

Combining all the published and unpublished $p-\bar{p}$ Collider and LEP results (as of mid-March 2000) yields an average $W$-boson mass of $80.419 \pm 0.038 \mathrm{GeV}$ assuming no common systematics between LEP and hadron collider measurements.

The Standard Model prediction from the electroweak fit, excluding the direct $W$ mass measurements from LEP and Tevatron, gives a $W$-boson mass of $80.382 \pm 0.026 \mathrm{GeV}$.

OUR EVALUATION in the listing below is obtained by combining only published LEP and $p-\bar{p}$ Collider results using the same procedure as above.

## References

1. The LEP Collaborations: ALEPH, DELPHI, L3, OPAL, the LEP Electroweak Working Group, and the SLD Heavy Flavour and Electroweak Groups, CERN-EP-2000016 (January 21, 2000).
2. A. Straessner and C. Sbarra, talks presented at the XXXV Rencontres de Moriond, "Electroweak Interactions and Unified Theories," (Les Arcs, France, 11-18 March 2000).

## W MASS

OUR FIT uses the $W$ and $Z$ mass, mass difference, and mass ratio measurements.

To obtain OUR EVALUATION the correlation between systematics is properly taken into account.

| VALUE (GeV) | EVTS | DOCUMENT ID |  | TECN | COMMENT |
| :---: | :---: | :---: | :---: | :---: | :---: |
| $80.419 \pm 0.056$ OUR EVALUATION $80.43 \pm 0.05$ OUR FIT |  |  |  |  |  |
|  |  |  |  |  |  |
| $80.482 \pm 0.091$ | 45394 | ${ }^{1}$ ABBOTT | 00 | D0 | $E_{\mathrm{cm}}^{p \bar{p}}=1.8 \mathrm{TeV}$ |
| $80.38 \pm 0.12 \pm 0.05$ | 701 | 2 ABBIENDI | 99C | OPAL | $E_{\mathrm{cm}}^{e e}{ }_{\mathrm{e}}^{\mathrm{e}}=161+172+183$ |
| $80.270 \pm 0.137 \pm 0.048$ | 809 | 3 ABREU | 99 T | DLPH | $E_{\mathrm{cm}}^{e} \mathrm{Cm}_{\mathrm{GeV}}^{=} 161+172+183$ |
| $80.61 \pm 0.15$ | 801 | 4 ACCIARRI | 99 | L3 | $E_{\mathrm{cm}}^{e \mathrm{em}}=161+172+183$ |
| $80.423 \pm 0.112 \pm 0.054$ | 812 | ${ }^{5}$ BARATE | 99 | ALEP | $E_{\mathrm{cm}}^{e \mathrm{GeVV}}=161+172+183$ |
| $80.41 \pm 0.18$ | 8986 | ${ }^{6} \mathrm{ABE}$ | 95P | CDF | $E_{\mathrm{cm}}^{p \bar{p}}=1.8 \mathrm{TeV}$ |
| $79.91 \pm 0.39$ | 1722 | 7 ABE | 90G | CDF | $E_{\mathrm{cm}}^{p \bar{p}}=1.8 \mathrm{TeV}$ |

-     - We do not use the following data for averages, fits, limits, etc.

| 80.49 | $\pm 0.43$ | $\pm 0.095$ | 871 | 8 ABREU | 99K | DLPH | Repl. by ABREU 99T |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 80.44 | $\pm 0.10$ | $\pm 0.07$ | 28323 | ${ }^{9}$ ABBOTT | 980 | D0 | Repl. by ABBOTT 00 |
| 80.22 | $\pm 0.41$ | $\pm 0.07$ | 72 | 10 ABREU | 98B | DLPH | Repl. by ABREU 99T |
| 80.32 | $\pm 0.30$ | $\pm 0.094$ | 96 | 11 ACKERSTAFF | 98D | OPAL | Repl. by ABBIENDI 99c |
| 80.5 | $+\quad 1.4$ $+\quad 2.2$ | $\begin{array}{r} +0.5 \\ -0.6 \end{array}$ | 104 | 12 ACKERSTAFF | 98D | OPAL | Repl. by ABBIENDI 99c |
| 80.80 | $\pm 0.32$ | $\pm 0.114$ | 95 | 13 BARATE | 98B | ALEP | Repl. by BARATE 99 |
| 80.80 | 1 $+\quad 0.48$ $-\quad 0.42$ | $\pm 0.03$ | 20 | 14 ACCIARRI | 97 | L3 | Repl. by ACCIARRI 99 |
| 80.5 | P $+\quad 1.4$ $-\quad 2.4$ | $\pm 0.3$ | 94 | 15 ACCIARRI | 97M | L3 | Repl. by ACCIARRI 99 |
| 80.71 | $+\quad 0.34$ $-\quad 0.35$ | $\pm 0.09$ | 101 | 16 ACCIARRI | 97S | L3 | Repl. by ACCIARRI 99 |
| 80.14 | $\pm 0.34$ | $\pm 0.095$ | 32 | 17 BARATE | 97 | ALEP | Repl. by BARATE 99 |
| 81.17 | 1 $+\quad 1.15$ $-\quad 1.62$ |  | 106 | 18 BARATE | 97S | ALEP | Repl. by BARATE 99 |
| 80.35 | $\pm 0.14$ | $\pm 0.23$ | 5982 | 19 ABACHI | 96E | D0 | Repl. by ABBOTT 00 |
| 80.40 | $\begin{array}{r} + \\ + \\ -\quad 0.44 \end{array}$ | $\begin{array}{r} +0.09 \\ -0.10 \end{array}$ | 23 | 20 ACKERSTAFF | 96B | OPAL | Repl. by ABBIENDI 99c |
| 84 | $\begin{array}{r} +10 \\ -\quad 7 \end{array}$ |  | 13 | 21 AID | 96D | H1 | $\begin{aligned} e^{ \pm} p & \rightarrow \nu_{e}\left(\bar{\nu}_{e}\right)+X \\ \sqrt{s} & \approx 300 \mathrm{GeV} \end{aligned}$ |
| 80.84 | $\pm 0.22$ | $\pm 0.83$ | 2065 | 22 ALITTI | 92B | UA2 | See $W / Z$ ratio below |
| 80.79 | $\pm 0.31$ | $\pm 0.84$ |  | 23 ALITTI | 90B | UA2 | $E_{\mathrm{cm}}^{p \bar{p}}=546,630 \mathrm{GeV}$ |
| 80.0 | $\pm 3.3$ | $\pm 2.4$ | 22 | 24 ABE | 891 | CDF | $E_{\mathrm{Cm}}^{p \bar{p}}=1.8 \mathrm{TeV}$ |
| 82.7 | $\pm 1.0$ | $\pm 2.7$ | 149 | 25 ALBAJAR | 89 | UA1 | $E_{\mathrm{cm}}^{p \bar{p}}=546,630 \mathrm{GeV}$ |
| 81.8 | 1.0 +6.0 $-\quad 5.3$ | $\pm 2.6$ | 46 | 26 ALBAJAR | 89 | UA1 | $E_{\mathrm{cm}}^{p \bar{p}}=546,630 \mathrm{GeV}$ |
| 89 | $\pm 3$ | $\pm 6$ | 32 | 27 ALBAJAR | 89 | UA1 | $E_{\mathrm{cm}}^{p \bar{p}}=546,630 \mathrm{GeV}$ |
| 81. | $\pm 5$. |  | 6 | ARNISON | 83 | UA1 | $E_{\mathrm{Cm}}^{e \mathrm{e}}=546 \mathrm{GeV}$ |
| 80. | $\begin{gathered} +10 . \\ -6 . \end{gathered}$ |  | 4 | BANNER | 83B | UA2 | Repl. by ALITTI 90B |

${ }^{1}$ ABBOTT 00 use $W \rightarrow e \nu_{e}$ events to measure the $W$ mass with a fit to the transverse mass distribution. The result quoted here corresponds to electrons detected both in the forward and in the central calorimeters for the data recorded in 1992-1995. For the large rapidity electrons recorded in 1994-1995, the analysis combines results obtained from $m_{T}, p_{T}(e)$, and $p_{T}(\nu)$.
${ }^{2}$ ABBIENDI 99C obtain this value properly combining results from a direct $W$ mass reconstruction at 172 and 183 GeV with that from the measurement of the total $W$-pair production cross section at 161 GeV . The systematic error includes an uncertainty of $\pm 0.02 \mathrm{GeV}$ due to the possible color-reconnection and Bose-Einstein effects in the purely hadronic final states and an uncertainty of $\pm 0.02 \mathrm{GeV}$ due to the beam energy.
${ }^{3}$ ABREU 99T obtain this value properly combining results obtained from a direct $W$ mass reconstruction at 172 and 183 GeV with those from measurement of $W$-pair production cross sections at 161,172 , and 183 GeV . The systematic error includes $\pm 0.021 \mathrm{GeV}$ due to the beam energy uncertainty and $\pm 0.030 \mathrm{GeV}$ due to possible color reconnection and Bose-Einstein effects in the purely hadronic final state.
${ }^{4}$ ACCIARRI 99 obtain this value properly combining results obtained from a direct $W$ mass reconstruction at 172 and 183 GeV with those from the measurements of the total $W$ pair production cross sections at 161 and 172 GeV . The value of the mass obtained from the direct reconstruction at 172 and 183 GeV is $M(W)=80.58 \pm 0.14 \pm 0.08 \mathrm{GeV}$.
${ }^{5}$ BARATE 99 obtain this value properly combining results from a direct $W$ mass reconstruction at 172 and 183 GeV with those from the measurements of the total $W$-pair production cross sections at 161 and 172 GeV . The systematic error includes $\pm 0.023 \mathrm{GeV}$ due to LEP energy uncertainty and $\pm 0.021 \mathrm{GeV}$ due to theory uncertainty on account of possible color reconnection and Bose-Einstein correlations.
${ }^{6} \mathrm{ABE} 95 \mathrm{P}$ use $3268 W \rightarrow \mu \nu_{\mu}$ events to find $M=80.310 \pm 0.205 \pm 0.130 \mathrm{GeV}$ and $5718 W \rightarrow e \nu_{e}$ events to find $M=80.490 \pm 0.145 \pm 0.175 \mathrm{GeV}$. The result given here combines these while accounting for correlated uncertainties.
${ }^{7} \mathrm{ABE} 90 \mathrm{G}$ result from $W \rightarrow e \nu$ is $79.91 \pm 0.35 \pm 0.24 \pm 0.19$ (scale) GeV and from $W \rightarrow \mu \nu$ is $79.90 \pm 0.53 \pm 0.32 \pm 0.08$ (scale) GeV .
${ }^{8}$ ABREU 99K derive this value using the Standard Model dependence on $M_{W}$ of the $W-W$ production cross sections measured at 161,172 , and 183 GeV . The systematics include an error of $\pm 0.03 \mathrm{GeV}$ arising from the beam energy uncertainty.
${ }^{9}$ ABBOTT 980 fit the transverse mass distribution of $28323 \mathrm{~W} \rightarrow e \nu_{e}$ events. The systematic error includes a detector related uncertainty of $\pm 60 \mathrm{MeV}$ and a model uncertainty of $\pm 30 \mathrm{MeV}$. Combining with $\mathrm{ABACHI} 96 \mathrm{E} D$ obtain a $W$ mass value of $80.43 \pm 0.11 \mathrm{GeV}$.
10 ABREU 98B obtain this value from a fit to the reconstructed $W$ mass distribution. The $W$ width was taken as its Standard Model value at the fitted $W$ mass. The systematic error includes $\pm 0.03 \mathrm{GeV}$ due to the beam energy uncertainty and $\pm 0.05 \mathrm{GeV}$ due to the possible color reconnection and Bose-Einstein effects in the purely hadronic final state.
11 ACKERSTAFF 98D obtain this value from a fit to the reconstructed $W$ mass distribution. The $W$ width was taken as its Standard Model value at the fitted $W$ mass. When both $W$ mass and width are varied they obtain $M(W)=80.30 \pm 0.27 \pm 0.095 \mathrm{GeV}$. The systematic error includes $\pm 0.03 \mathrm{GeV}$ due to the beam energy uncertainty and $\pm 0.05 \mathrm{GeV}$ due to the possible color reconnection and Bose-Einstein effects in the purely hadronic final state. Combining both values of ACKERSTAFF 98D with ACKERSTAFF 96B authors find: $M(W)=80.35 \pm 0.24 \pm 0.07 \pm 0.03$ (LEP) GeV .
12 ACKERSTAFF 98D derive this value from their measured $W W$ production cross section $\sigma_{W W}=12.3 \pm 1.3 \pm 0.4 \mathrm{pb}$ using the Standard Model dependence of $\sigma_{W} W$ on $M_{W}$ at the given c.m. energy.
13 BARATE 98B obtain this value from a fit to the reconstructed $W$ mass distribution. The $W$ width was taken as its Standard Model value at the fitted $W$ mass. The systematic error includes $\pm 0.03 \mathrm{GeV}$ due to the beam energy uncertainty and $\pm 0.032 \mathrm{GeV}$ due to the possible color reconnection and Bose-Einstein effects in the purely hadronic final state. Combining with the $M_{W}$ values from cross section measurements at 161 and 172 GeV (BARATE 97 and BARATE 97s) authors find: $M(W)=80.51 \pm 0.23 \pm 0.08 \mathrm{GeV}$.

14 ACCIARRI 97 derive this value from their measured $W$ - $W$ production cross section $\sigma_{W W}=2.89_{-0.70}^{+0.81} \pm 0.14 \mathrm{pb}$ using the Standard Model dependence of $\sigma_{W W}$ on $M_{W}$ at the given c.m. energy. Statistical and systematic errors are added in quadrature and the last error of $\pm 0.03 \mathrm{GeV}$ arises from the beam energy uncertainty. The same result is given by a fit of the production cross sections to the data.
${ }^{15}$ ACCIARRI 97M derive this value from their measured $W W$ production cross section $\sigma_{W W}=12.27_{-1.32}^{+1.41} \pm 0.23 \mathrm{pb}$ using the Standard Model dependence of $\sigma_{W} W$ on $M_{W}$ at the given c.m. energy. Combining with ACCIARRI 97 authors find $M(W)=$ $80.78+0.41 \pm 0.03 \mathrm{GeV}$ where the last error is due to beam energy uncertainty.
${ }^{16}$ ACCIARRI 975 obtain this value from a fit to the reconstructed $W$ mass distribution. The $W$ width was taken as its Standard Model value at the fitted $W$ mass. When both $W$ mass and width are varied they obtain $M(W)=80.72_{-0.33}^{+0.31} \pm 0.09 \mathrm{GeV}$. The systematic error includes $\pm 0.03 \mathrm{GeV}$ due to the beam energy uncertainty and $\pm 0.05 \mathrm{GeV}$ due to the possible color reconnection and Bose-Einstein effects in the purely hadronic final state. Combining with ACCIARRI 97 and ACCIARRI 97M authors find: $M(W)=$ $80.75_{-0.27}^{+0.26} \pm 0.03$ (LEP) GeV.
17 BARATE 97 derive this value from their measured $W-W$ production cross section $\sigma_{W} W$ $=4.23 \pm 0.73 \pm 0.19 \mathrm{pb}$ using the Standard Model dependence of $\sigma_{W W}$ on $M_{W}$ at the given c.m. energy. The systematics include an error of $\pm 0.03 \mathrm{GeV}$ arising from the beam energy uncertainty.
18 BARATE 97S derive this value from their measured $W W$ production cross section $\sigma W W$ $=11.71 \pm 1.23 \pm 0.28 \mathrm{pb}$ using the Standard Model dependence of $\sigma_{W} W$ on $M_{W}$ at the given c.m. energy. The errors quoted on the mass are statistical only. Combining with BARATE 97 authors find: $M(W)=80.20 \pm 0.33 \pm 0.09 \pm 0.03$ (LEP) GeV .
${ }^{19}$ ABACHI 96E fit the transverse mass distribution of $5982 \mathrm{~W} \rightarrow e \nu_{e}$ decays. An error of $\pm 160 \mathrm{MeV}$ due to the uncertainty in the absolute energy scale of the EM calorimeter is included in the total systematics.
20 ACKERSTAFF 96B derive this value from an analysis of the predicted $M_{W}$ dependence of their accepted four-fermion cross section, explicitly taking into account interference effects. The systematics include an error of $\pm 0.03 \mathrm{GeV}$ arising from the beam energy uncertainty.
21 AID 96D derive this value as a propagator mass using the $Q^{2}$ shape and magnitude of the $e^{ \pm}$charged-current cross sections. $Q^{2}>5000 \mathrm{GeV}^{2}$ events with $p_{T}$ of the outgoing lepton $>25 \mathrm{GeV} / \mathrm{c}$ are used.
${ }^{22}$ ALITTI 92B result has two contributions to the systematic error $( \pm 0.83)$; one $( \pm 0.81)$ cancels in $m_{W} / m_{Z}$ and one ( $\pm 0.17$ ) is noncancelling. These were added in quadrature. We choose the ALITTI 92B value without using the LEP $m_{Z}$ value, because we perform our own combined fit.
23 There are two contributions to the systematic error $( \pm 0.84)$ : one $( \pm 0.81)$ which cancels in $m_{W} / m_{Z}$ and one ( $\pm 0.21$ ) which is non-cancelling. These were added in quadrature.
24 ABE 891 systematic error dominated by the uncertainty in the absolute energy scale.
${ }^{25}$ ALBAJAR 89 result is from a total sample of $299 W \rightarrow e \nu$ events.
${ }^{26}$ ALBAJAR 89 result is from a total sample of $67 \mathrm{~W} \rightarrow \mu \nu$ events.
27 ALBAJAR 89 result is from $W \rightarrow \tau \nu$ events.

## W/Z MASS RATIO

The fit uses the $W$ and $Z$ mass, mass difference, and mass ratio measurements.


| $0.88114 \pm 0.00154 \pm 0.00252$ | 5982 | 29 | ABBOTT | 98P D0 |
| :--- | ---: | :--- | :--- | :--- |$E_{\mathrm{cm}}^{p \bar{p}}=1.8 \mathrm{TeV}$

28 ABBOTT 98N obtain this from a study of $28323 W \rightarrow e \nu_{e}$ and $3294 Z \rightarrow e^{+} e^{-}$ decays. Of this latter sample, 2179 events are used to calibrate the electron energy scale.
29 ABBOTT 98P obtain this from a study of $5982 \mathrm{~W} \rightarrow e \nu_{e}$ events. The systematic error includes an uncertainty of $\pm 0.00175$ due to the electron energy scale.
30 Scale error cancels in this ratio.

$$
m_{Z}-m_{W}
$$

The fit uses the $W$ and $Z$ mass, mass difference, and mass ratio measurements.


-     - We do not use the following data for averages, fits, limits, etc.
$11.3 \pm 1.3 \pm 0.9 \quad$ ANSARI 87 UA2 $E_{\mathrm{cm}}^{p \bar{p}}=546,630 \mathrm{GeV}$

$$
\boldsymbol{m}_{\boldsymbol{w}^{+}}-\boldsymbol{m}_{\boldsymbol{w}^{-}}
$$

Test of CPT invariance.

| $\frac{\operatorname{VALUE}(\mathrm{GeV})}{\mathbf{- 0 . 1 9 \pm 0 . 5 8}} \frac{\text { EVTS }}{1722}$ | $\frac{\text { DOCUMENT ID }}{\text { ABE }} \frac{\text { TECN }}{}$ | $\frac{\text { COMMENT }}{E_{\mathrm{cm}}^{p \bar{p}}=1.8 \mathrm{TeV}}$ |
| :--- | :--- | :--- | :--- |

## W WIDTH

The CDF and $D \varnothing$ widths labelled "extracted value" are obtained by measuring $R=[\sigma(W) / \sigma(Z)]\left[\Gamma\left(W \rightarrow \ell \nu_{\ell}\right)\right] /(\mathrm{B}(Z \rightarrow \ell \ell) \Gamma(W))$ where the bracketed quantities can be calculated with plausible reliability. $\Gamma(W)$ is then extracted by using a value of $\mathrm{B}(Z \rightarrow \ell \ell)$ measured at LEP. The UA1 and UA2 widths used $R=[\sigma(W) / \sigma(Z)]\left[\Gamma\left(W \rightarrow \ell \nu_{\ell}\right) / \Gamma(Z \rightarrow\right.$ $\ell \ell)] \Gamma(Z) / \Gamma(W)$ and the measured value of $\Gamma(Z)$. The Standard Model prediction is $2.067 \pm 0.021$ (ROSNER 94).


-     - We do not use the following data for averages, fits, limits, etc.

${ }^{31}$ ABBOTT 00B measure $R=10.43 \pm 0.27$ for the $W \rightarrow e \nu_{e}$ decay channel. They use the SM theoretical predictions for $\sigma(W) / \sigma(Z)$ and $\Gamma\left(W \rightarrow e \nu_{e}\right)$ and the world average for $\mathrm{B}(Z \rightarrow e e)$. The value quoted here is obtained combining this result ( $2.169 \pm 0.070$ GeV ) with that of ABBOTT 99H.
32 ABBIENDI 99C obtain this value from a fit to the reconstructed $W$ mass distribution using data at 172 and 183 GeV . The systematic error includes an uncertainty of $\pm 0.12 \mathrm{GeV}$ due to the possible color-reconnection and Bose-Einstein effects in the purely hadronic final states and an uncertainty of $\pm 0.01 \mathrm{GeV}$ due to the beam energy.
33 ABREU 99T obtain this value using $W W \rightarrow \ell \bar{\nu}_{\ell} q \bar{q}$ and $W W \rightarrow q \bar{q} q \bar{q}$ events. The systematic error includes an uncertainty of $\pm 0.080 \mathrm{GeV}$ due to possible color reconnection and Bose-Einstein effects in the purely hadronic final state.
34 ACCIARRI 99 obtain this value from a fit to the reconstruced $W$ mass distribution using data at 172 and 183 GeV .
${ }^{35} \mathrm{ABE} 95 \mathrm{C}$ use the tail of the transverse mass distribution of $W \rightarrow e \nu_{e}$ decays.
${ }^{36} \mathrm{ABE} 95 \mathrm{~W}$ measured $R=10.90 \pm 0.32 \pm 0.29$. They use $m_{W}=80.23 \pm 0.18 \mathrm{GeV}$, $\sigma(W) / \sigma(Z)=3.35 \pm 0.03, \Gamma(W \rightarrow e \nu)=225.9 \pm 0.9 \mathrm{MeV}, \Gamma\left(Z \rightarrow e^{+} e^{-}\right)=$ $83.98 \pm 0.18 \mathrm{MeV}$, and $\Gamma(Z)=2.4969 \pm 0.0038 \mathrm{GeV}$.
37 ALITTI 92 measured $R=10.4_{-0.6}^{+0.7} \pm 0.3$. The values of $\sigma(Z)$ and $\sigma(W)$ come from $O\left(\alpha_{s}^{2}\right)$ calculations using $m_{W}=80.14 \pm 0.27 \mathrm{GeV}$, and $m_{Z}=91.175 \pm 0.021 \mathrm{GeV}$ along with the corresponding value of $\sin ^{2} \theta_{W}=0.2274$. They use $\sigma(W) / \sigma(Z)=$ $3.26 \pm 0.07 \pm 0.05$ and $\Gamma(Z)=2.487 \pm 0.010 \mathrm{GeV}$.
38 ALBAJAR 91 measured $R=9.5_{-1.0}^{+1.1}$ (stat. + syst.). $\sigma(W) / \sigma(Z)$ is calculated in QCD at the parton level using $m_{W}=80.18 \pm 0.28 \mathrm{GeV}$ and $m_{Z}=91.172 \pm 0.031 \mathrm{GeV}$ along with $\sin ^{2} \theta_{W}=0.2322 \pm 0.0014$. They use $\sigma(W) / \sigma(Z)=3.23 \pm 0.05$ and $\Gamma(Z)$ $=2.498 \pm 0.020 \mathrm{GeV}$. This measurement is obtained combining both the electron and muon channels.
39 ABBOTT 99 H measure $R=10.90 \pm 0.52$ combining electron and muon channels. They use $M_{W}=80.39 \pm 0.06 \mathrm{GeV}$ and the SM theoretical predictions for $\sigma(W) / \sigma(Z), \mathrm{B}(Z \rightarrow$ $\ell \ell)$, and $\Gamma\left(W \rightarrow \ell \nu_{\ell}\right)$.
${ }^{40}$ BARATE 99I obtain this result with a fit to the $W W$ measured cross sections at 161 , 172 , and 183 GeV . The theoretical prediction takes into account the sensitivity to the $W$ total width.
${ }^{41}$ ACKERSTAFF 98D obtain this value from a fit to the reconstructed $W$ mass distribution.
${ }^{42}$ ACCIARRI 975 obtain this value from a fit to the reconstructed $W$ mass distribution.
43 ALITTI 90C used the same technique as described for ABE 90. They measured $R=$ $9.38_{-0.72}^{+0.82} \pm 0.25$, obtained $\Gamma(W) / \Gamma(Z)=0.902 \pm 0.074 \pm 0.024$. Using $\Gamma(Z)=$ $2.546 \pm 0.032 \mathrm{GeV}$, they obtained the $\Gamma(W)$ value quoted above and the limits $\Gamma(W)$ $<2.56$ (2.64) GeV at the $90 \%(95 \%) \mathrm{CL} . E_{\mathrm{cm}}^{p \bar{p}}=546,630 \mathrm{GeV}$.

44 ALBAJAR 89 result is from a total sample of $299 W \rightarrow e \nu$ events.
45 If systematic error is neglected, result is $2.7_{-1.5}^{+1.4} \mathrm{GeV}$. This is enhanced subsample of 172 total events.

## $\mathbf{W}^{+}$DECAY MODES

$W^{-}$modes are charge conjugates of the modes below.

[a] $\ell$ indicates each type of lepton ( $e, \mu$, and $\tau$ ), not sum over them.
[b] This represents the width for the decay of the $W$ boson into a charged particle with momentum below detectability, $\mathrm{p}<200 \mathrm{MeV}$.

## W PARTIAL WIDTHS

This represents the width for the decay of the $W$ boson into a charged particle with momentum below detectability, $\mathrm{p}<200 \mathrm{MeV}$.

| VALUE ( MeV ) | DOCUMEN | TECN | COMMENT |
| :---: | :---: | :---: | :---: |
| $3{ }_{-48}^{+52} \pm 33$ | 46 BARATE | ALEP | $\begin{gathered} E_{\mathrm{Cm}}^{e \mathrm{em}}=161+172+183 \\ \mathrm{GeV} \end{gathered}$ |


${ }^{46}$ BARATE 991 measure this quantity using the dependence of the total cross section $\sigma_{W} W$ upon a change in the total width. The fit is performed to the $W W$ measured cross sections at 161,172 , and 183 GeV . This partial width is $<139 \mathrm{MeV}$ at $95 \% \mathrm{CL}$.
47 BARATE 99L use $W$-pair production to search for effectively invisible $W$ decays, tagging with the decay of the other $W$ boson to Standard Model particles. The partial width for effectively invisible decay is $<27 \mathrm{MeV}$ at $95 \% \mathrm{CL}$.

## W BRANCHING RATIOS

Overall fits are performed to determine the branching ratios of the $W$. For each LEP experiment the correlation matrix of the leptonic branching ratios is used and the common systematic errors among LEP experiments are properly taken into account. A first fit determines three individual leptonic branching ratios, $\mathrm{B}\left(W \rightarrow e \nu_{e}\right), \mathrm{B}\left(W \rightarrow \mu \nu_{\mu}\right)$, and $\mathrm{B}(W \rightarrow$ $\left.\tau \nu_{\tau}\right)$. This fit has a $\chi^{2}=10.5$ for 20 degrees of freedom. A second fit assumes lepton universality and determines the leptonic branching ratio $\mathrm{B}\left(W \rightarrow \ell \nu_{\ell}\right)$. This fit has a $\chi^{2}=11.0$ for 22 degrees of freedom. A separate fit is performed only to hadronic branching ratio data taking into account the common systematic errors. This fit has a $\chi^{2}=2.3$ for 3 degrees of freedom.
$\ell$ indicates average over $e, \mu$, and $\tau$ modes, not sum over modes.

Data marked "fit" are used for the fit. The other data is highly correlated with data appearing elsewhere in the Listings and are therefore not used in the fit.


-     - We do not use the following data for averages, fits, limits, etc.

| 0.1085 | $\pm 0.0048$ | $\pm 0.0017$ | 170 | ABREU | 99K DLPH | Repl. by ABREU 00k |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 0.101 | $\begin{array}{r} +0.011 \\ -0.010 \end{array}$ | $\pm 0.002$ | 61 | ACKERSTAFF | 98D OPAL | Repl. by ABBIENDI 99D |
| 0.119 | $\begin{array}{r} +0.013 \\ -0.012 \end{array}$ | $\pm 0.002$ | 51 | ACCIARRI | 97M L3 | Repl. by ACCIARRI 98P |

48 ABBOTT 99н measure $R \equiv\left[\sigma_{W} \mathrm{~B}\left(W \rightarrow \ell \nu_{\ell}\right)\right] /\left[\sigma_{Z} \mathrm{~B}(Z \rightarrow \ell \ell)\right]=10.90 \pm 0.52$ combining electron and muon channels. They use $M_{W}=80.39 \pm 0.06 \mathrm{GeV}$ and the SM theoretical predictions for $\sigma(W) / \sigma(Z)$ and $\mathrm{B}(Z \rightarrow \ell \ell)$.
$491216 \pm 38_{-31}^{+27} W \rightarrow \mu \nu$ events from ABE 92I and $2426 \mathrm{~W} \rightarrow e \nu$ events of ABE 91C. ABE 92l give the inverse quantity as $9.6 \pm 0.7$ and we have inverted.


-     - We do not use the following data for averages, fits, limits, etc. - . GeV


50 ABBOTT 00B measure $R \equiv\left[\sigma_{W} \mathrm{~B}\left(W \rightarrow e \nu_{e}\right)\right] /\left[\sigma_{Z} \mathrm{~B}(Z \rightarrow e e)\right]=10.43 \pm 0.27$ for the $W \rightarrow e \nu_{e}$ decay channel. They use the SM theoretical prediction for $\sigma(W) / \sigma(Z)$ and the world average for $\mathrm{B}(Z \rightarrow e e)$.
${ }^{51} \mathrm{ABE} 95 \mathrm{~W}$ result is from a measurement of $\sigma \mathrm{B}(W \rightarrow e \nu) / \sigma \mathrm{B}\left(Z \rightarrow e^{+} e^{-}\right)=$ $10.90 \pm 0.32 \pm 0.29$, the theoretical prediction for the cross section ratio, the experimental knowledge of $\Gamma\left(Z \rightarrow e^{+} e^{-}\right)=83.98 \pm 0.18 \mathrm{MeV}$, and $\Gamma(Z)=2.4969 \pm 0.0038$ GeV .
52 The first error was obtained by adding the statistical and systematic experimental uncertainties in quadrature. The second error reflects the dependence on theoretical prediction of total $W$ cross section: $\sigma(546 \mathrm{GeV})=4.7_{-0.7}^{+1.4} \mathrm{nb}$ and $\sigma(630 \mathrm{GeV})=5.8_{-1.0}^{+1.8} \mathrm{nb}$. See ALTARELLI 85B.

| $\Gamma\left(\mu^{+} \nu\right) / \Gamma_{\text {total }}$ |  |  |  | 「3/Г |
| :---: | :---: | :---: | :---: | :---: |
| VALUE | EVTS | DOCUMENT ID | TECN | COMMENT |
| $\overline{\mathbf{0 . 1 0 4 9} \pm 0.0029 ~ O U R ~ F I T ~-~-~}$ |  |  |  |  |
| $0.1092 \pm 0.0048 \pm 0.0012$ | 461 | ABREU | 00k DLPH | $E_{\mathrm{cm}}^{e e}=$ |
|  |  |  |  | $\begin{aligned} & 161+172+183 \\ & +189 \mathrm{GeV} \end{aligned}$ |
| $0.102 \pm 0.008 \pm 0.002$ | 169 | ABBIENDI | 99D OPAL | $E_{\mathrm{cm}}^{e \mathrm{ee}}=161+172+$ |
| $0.1006 \pm 0.0078 \pm 0.0021$ | 179 | BARATE | 991 ALEP | $E_{\mathrm{Cm}}^{\text {e }}{ }^{183}=1 \mathrm{GeV}$ |
|  |  |  |  | $161+172+183$ |
| $0.102 \pm 0.009 \pm 0.002$ | 115 | ACCIARRI | 98P L3 | $E_{\mathrm{cm}}^{e \mathrm{e}} \mathrm{GeV}_{=161+172+}$ |
|  |  | 53 ABE |  | ${ }^{1} \frac{1}{\bar{p}}{ }^{3} \mathrm{GeV}$ |
| $0.10 \pm 0.01$ | 1216 | ${ }^{53} \mathrm{ABE}$ | 921 CDF | $E_{\mathrm{cm}}^{\mathrm{cm}}=1.8 \mathrm{TeV}$ |

HTTP://PDG.LBL.GOV
Page $11 \quad$ Created: 12/18/2000 15:07

-     - We do not use the following data for averages, fits, limits, etc.

| 0.1139 | $\pm 0.0096$ | $\pm 0.0023$ | 67 | ABREU | 99k DLPH | Repl. by ABREU 00k |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 0.073 | +0.019 -0.017 | $\pm 0.002$ | 16 | ACKERSTAFF | 98D OPAL | Repl. by ABBIENDI 99D |
| 0.084 | $\begin{array}{r} +0.028 \\ -0.024 \end{array}$ | $\pm 0.003$ | 13 | ACCIARRI | 97M L3 | Repl. by ACCIARRI 98P |
| 0.112 | $\pm 0.02$ | $\pm 0.006$ | 25 | BARATE | 97S ALEP | Repl. by BARATE 99\| |

53 ABE 92l quote the inverse quantity as $9.9 \pm 1.2$ which we have inverted.

| $\Gamma\left(\tau^{+} \nu\right) / \Gamma_{\text {total }}$ |  |  |  | $\Gamma 4 / \Gamma$ |
| :---: | :---: | :---: | :---: | :---: |
| VALUE | EVTS | DOCUMENT ID | TECN | COMMENT |
| $\mathbf{0 . 1 0 4 3} \pm \mathbf{0 . 0 0 4 1 ~ O U R ~ F I T}$ |  |  |  |  |
| $0.1105 \pm 0.0075 \pm 0.0032$ | 424 | ABREU | 00K DLPH | $E_{\mathrm{cm}}^{e \mathrm{e}}$ = |
|  |  |  |  | $\begin{aligned} & 161+172+183 \\ & +189 \mathrm{GeV} \end{aligned}$ |
| $0.101 \pm 0.010 \pm 0.003$ | 144 | ABBIENDI | 99D OPAL | $\begin{aligned} & E_{\mathrm{Cm}}^{e e e}=161+172+ \\ & 183 \mathrm{GeV} \end{aligned}$ |
| $0.0976 \pm 0.0101 \pm 0.0033$ | 160 | BARATE | 991 ALEP | $\begin{aligned} & E_{\mathrm{cm}}^{e e}=161+172+ \\ & 183 \mathrm{GeV} \end{aligned}$ |
| $0.090 \pm 0.012 \pm 0.003$ | 81 | ACCIARRI | 98P L3 | $E_{\mathrm{cm}}^{e e}=161+172+$ |

- . . We do not use the following data for averages, fits, limits, etc. • . 183 GeV

| $0.1095 \pm 0.0149 \pm 0.0041$ | 47 | ABREU | 99K DLPH | Repl. by ABREU 00k |
| :---: | :---: | :---: | :---: | :---: |
| $0.140 \begin{gathered}+0.030 \\ -0.028\end{gathered} \pm 0.005$ | 23 | ACKERSTAFF | 98D OPAL | Repl. by ABBIENDI 99D |
| $0.109{ }_{-0.039}^{+0.042} \pm 0.005$ | 15 | ACCIARRI | 97M L3 | Repl. by ACCIARRI 98P |
| $0.113 \pm 0.027 \pm 0.006$ | 37 | BARATE | 97S ALEP | Repl. by BARATE 99। |
| $\Gamma($ hadrons $) / \Gamma_{\text {total }}$ |  |  |  | $\Gamma_{5} / \Gamma$ |
| VALUE | EVTS | DOCUMENT ID | TECN | COMMENT |
| $\mathbf{0 . 6 8 4 8} \pm \mathbf{0 . 0 0 5 9}$ OUR FIT |  |  |  |  |
| $0.6789 \pm 0.0073 \pm 0.0043$ | 1773 | ABREU | 00k DLPH | $\begin{aligned} & E_{\mathrm{Cm}}^{e e}= \\ & \quad 161+172+183 \\ & \quad+189 \mathrm{GeV} \end{aligned}$ |
| $0.679 \pm 0.012 \pm 0.005$ | 395 | ABBIENDI | 99D OPAL | $\begin{aligned} & E_{\mathrm{Cm}}^{\mathrm{ee}}=161+172+ \\ & 183 \mathrm{GeV} \end{aligned}$ |
| $0.6893 \pm 0.0121 \pm 0.0051$ | 1255 | BARATE | 991 ALEP | $\begin{aligned} & E_{\mathrm{Cm}}^{\mathrm{ee}}=161+172+ \\ & 183 \mathrm{GeV} \end{aligned}$ |
| $0.701 \pm 0.013 \pm 0.004$ | 462 | ACCIARRI | 98P L3 | $E_{\mathrm{cm}}^{e \mathrm{ee}}=161+172+$ |

-     - We do not use the following data for averages, fits, limits, etc. - •

| 0.6746 | $\pm 0.0143$ | $\pm 0.0052$ | 465 | ABREU | 99k DLPH | Repl. by ABREU 00k |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 0.698 | $\begin{array}{r} +0.030 \\ -0.032 \end{array}$ | $\pm 0.007$ | 52 | ACKERSTAFF | 98D OPAL | Repl. by ABBIENDI 99D |
| 0.642 | $\begin{array}{r} +0.037 \\ -0.038 \end{array}$ | $\pm 0.005$ | 70 | ACCIARRI | 97M L3 | Repl. by ACCIARRI 98P |
| 0.677 | $\pm 0.031$ | $\pm 0.007$ | 65 | BARATE | 97S ALEP | Repl. by BARATE 99I |


| $\underset{\text { VALUE }}{\Gamma\left(\mu^{+} \nu\right) / \Gamma\left(e^{+} \nu\right)}$ |  |  |  | $\Gamma_{3} / \Gamma_{2}$ |
| :---: | :---: | :---: | :---: | :---: |
|  | EVTS | DOCUMENT ID | TECN | COMMENT |
| $0.984 \pm 0.032$ OUR FIT |  |  |  |  |
| $0.89 \pm 0.10$ | 13k | 54 ABACHI | 95D D0 | $E_{\mathrm{Cm}}^{p \bar{p}}=1.8 \mathrm{TeV}$ |
| $1.02 \pm 0.08$ | 1216 | 55 ABE | 92I CDF | $E_{\mathrm{cm}}^{p \bar{p}}=1.8 \mathrm{TeV}$ |

-     -         - We do not use the following data for averages, fits, limits, etc. • • •
$1.00 \pm 0.14 \pm 0.08$
67
$\begin{array}{ll}\text { ALBAJAR } & 89 \text { UA1 } \\ \text { ARNISON } & 84 D \text { UA1 }\end{array}$
$E_{\mathrm{cm}}^{p \bar{p}}=546,630$ GeV

1.24 \begin{tabular}{c}
+0.6 <br>
-0.4

$\quad 14 \quad$ ARNISON 84D UA1 

Repl. by ALBA- <br>
JAR 89
\end{tabular}

54 ABACHI 95D obtain this result from the measured $\sigma_{W} \mathrm{~B}(W \rightarrow \mu \nu)=2.09 \pm 0.23 \pm$ 0.11 nb and $\sigma_{W} \mathrm{~B}(W \rightarrow e \nu)=2.36 \pm 0.07 \pm 0.13 \mathrm{nb}$ in which the first error is the combined statistical and systematic uncertainty, the second reflects the uncertainty in the luminosity.
55 ABE 92 I obtain $\sigma_{W} \mathrm{~B}(W \rightarrow \mu \nu)=2.21 \pm 0.07 \pm 0.21$ and combine with ABE 91C $\sigma_{W}$ $\mathrm{B}((W \rightarrow e \nu))$ to give a ratio of the couplings from which we derive this measurement.

| $\Gamma\left(\tau^{+} \nu\right) / \Gamma\left(e^{+} \nu\right)$ | EVTS DOCUMENT ID |  | TECN | COMMENT $\Gamma_{\mathbf{4} / \Gamma_{\mathbf{2}}}$ |
| :---: | :---: | :---: | :---: | :---: |
| VALUE |  |  |  |  |
| $0.979 \pm 0.044$ OUR FIT |  |  |  |  |
| $0.94 \pm 0.14$ | 179 | 56 ABE | 92E CDF | $E_{\mathrm{cm}}^{p \bar{p}}=1.8 \mathrm{TeV}$ |
| $1.04 \pm 0.08 \pm 0.08$ | 754 | 57 ALITTI | 92F UA2 | $E_{\mathrm{cm}}^{p \bar{p}}=630 \mathrm{GeV}$ |
| $1.02 \pm 0.20 \pm 0.12$ | 32 | ALBAJAR | 89 UA1 | $\begin{gathered} E_{\mathrm{cm}}^{p \bar{p}}=546,630 \\ \mathrm{GeV} \end{gathered}$ |

-     - We do not use the following data for averages, fits, limits, etc. - -
$0.995 \pm 0.112 \pm 0.083$
198
$1.02 \pm 0.20 \pm 0.10$
32
ALITTI
ALBAJAR
91C UA2
Repl. by ALITTI 92F
87 UA1
Repl. by ALBAJAR 89
${ }^{56}$ ABE 92E use two procedures for selecting $W \rightarrow \tau \nu_{\tau}$ events. The missing $\mathrm{E}_{T}$ trigger leads to $132 \pm 14 \pm 8$ events and the $\tau$ trigger to $47 \pm 9 \pm 4$ events. Proper statistical and systematic correlations are taken into account to arrive at $\sigma \mathrm{B}(W \rightarrow \tau \nu)=2.05 \pm 0.27$ nb. Combined with ABE 91 C result on $\sigma \mathrm{B}(W \rightarrow e \nu)$, ABE 92 E quote a ratio of the couplings from which we derive this measurement.
57 This measurement is derived by us from the ratio of the couplings of ALITTI 92F.

| $\Gamma\left(\pi^{+} \gamma\right) / \Gamma\left(e^{+} \nu\right)$$\underline{V A L U E}$ |  |  |  | $\Gamma_{6} / \Gamma_{2}$ |
| :---: | :---: | :---: | :---: | :---: |
|  | CL\% | DOCUMENT ID | TECN | COMMENT |
| $<7 \times 10^{-4}(C L=95 \%)$ |  |  |  |  |
| $<7 \times 10^{-4}$ | 95 | ABE | 98H CDF | $E_{\mathrm{cm}}^{p \bar{p}}=1.8 \mathrm{TeV}$ |
| $<4.9 \times 10^{-3}$ | 95 | 58 ALITTI | 92D UA2 | $E_{\mathrm{cm}}^{p \bar{p}}=630 \mathrm{GeV}$ |
| $<58 \times 10^{-3}$ | 95 | 59 ALBAJAR | 90 UA1 | $E_{\mathrm{cm}}^{p \bar{p}}=546,630 \mathrm{GeV}$ |
| 58 ALITTI 92D limit is $3.8 \times 10^{-3}$ at $90 \%$ CL. <br> 59 ALBAJAR 90 obtain $<0.048$ at $90 \%$ CL. |  |  |  |  |


| $\Gamma\left(D_{s}^{+} \gamma\right) / \Gamma\left(e^{+} \nu\right)$ |  |  |  |  | $\Gamma_{7} / \Gamma_{2}$ |
| :---: | :---: | :---: | :---: | :---: | :---: |
| VALUE | CL\% | DOCUMENT ID | TECN | COMMENT |  |
| $<1.2 \times 10^{-2}$ | 95 | ABE | CDF | $E_{\mathrm{cm}}^{p \bar{p}}=1.8 \mathrm{TeV}$ |  |


| X)/ |  |  |  |  | $\Gamma_{8} / \Gamma_{5}$ |
| :---: | :---: | :---: | :---: | :---: | :---: |
| VALUE | EVTS | DOCUMENT ID TECN |  | COMMENT |  |
| $0.51 \pm 0.05 \pm 0.03$ | 746 | 60 BARATE | 99M ALEP | $E_{\mathrm{Cm}}^{e e}=172+183 \mathrm{GeV}$ |  |
| 60 BARATE 99M ta is determined to | $00 \pm$ | $\begin{aligned} & \text { neural net } \\ & 1 \pm 0.07 \text {. } \end{aligned}$ | rithm. | his me | $\mathrm{t} \mid V_{c}$ |

$\mathrm{R}_{\boldsymbol{c s}}=\Gamma(c \bar{s}) / \Gamma$ (hadrons)

VALUE
$\frac{\text { DOCUMENT ID }}{61 \text { ABREU }} \frac{\text { TECN }}{\text { COMMENT }} \frac{\text { 98N DLPH }}{E_{\mathrm{cm}}^{e e}=161+172 \mathrm{GeV}}$
61 ABREU 98 N tag $c$ and $s$ jets by identifying a charged kaon as the highest momentum particle in a hadronic jet. They also use a lifetime tag to independently identify a $c$ jet, based on the impact parameter distribution of charged particles in a jet. From this measurement $\left|V_{c s}\right|$ is determined to be $0.94_{-0.26}^{+0.32} \pm 0.13$.

## AVERAGE PARTICLE MULTIPLICITIES IN HADRONIC W DECAY

Summed over particle and antiparticle, when appropriate.

## $\left\langle N_{\text {charged }}\right\rangle$

$\frac{V A L U E}{19.3 \pm 0.4 \text { OUR AVERAGE }}$
$19.3 \pm 0.3 \pm 0.3$
$19.23 \pm 0.74$

DOCUMENT ID TECN COMMENT
62 ABBIENDI 99N OPAL $E_{\mathrm{Cm}}^{e e}=183 \mathrm{GeV}$
63 ABREU 98C DLPH $\quad E_{\mathrm{cm}}^{e e}=172 \mathrm{GeV}$
${ }^{62} \mathrm{ABBIENDI} 99 \mathrm{~N}$ use the final states $W^{+} W^{-} \rightarrow q \bar{q} \ell \bar{\nu}_{\ell}$ to derive this value.
63 ABREU 98C combine results from both the fully hadronic as well semileptonic $W W$ final states after demonstrating that the $W$ decay charged multiplicity is independent of the topology within errors.

## TRIPLE GAUGE COUPLINGS (TGC'S)

Revised March 2000 by C. Caso (Univ. of Genova) and A. Gurtu (Tata Inst.)

Fourteen independent couplings, 7 each for $Z W W$ and $\gamma W W$, completely describe the $V W W$ vertices within the most general framework of the electroweak Standard Model (SM) consistent with Lorentz invariance and $U(1)$ gauge invariance. Of each of the 7 TGC's, 3 conserve $C$ and $P$ individually, 3 violate $C P$, and one TGC violates $C$ and $P$ individually while conserving $C P$. Assumption of $C$ and $P$ conservation and electromagnetic gauge invariance reduces the independent $V W W$ couplings to five: one common set is $\left(\kappa_{\gamma}, \kappa_{Z}, \lambda_{\gamma}, \lambda_{Z}, g_{1}^{Z}\right)$, where $\kappa_{\gamma}=\kappa_{Z}=g_{1}^{Z}=1$ and $\lambda_{\gamma}=\lambda_{Z}=0$ in the Standard

Model at the tree level. The $W$ magnetic dipole moment, $\mu_{W}$, and the $W$ electric quadrupole moment, $q_{W}$, are expressed as $\mu_{W}=e\left(1+\kappa_{\gamma}+\lambda_{\gamma}\right) / 2 M_{W}$ and $q_{W}=-e\left(\kappa_{\gamma}-\lambda_{\gamma}\right) / M_{W}^{2}$.

Precision measurements of suitable observables at LEP1 has already led to an exploration of much of the TGC parameter space. Three linear combinations of the TGC's, $\alpha_{W \phi}, \alpha_{B \phi}$ and $\alpha_{W}$, have been proposed to investigate the leftover "blind" directions in the $C P$-conserving TGC parameter space, and two linear couplings, $\widetilde{\alpha}_{B W}$ and $\widetilde{\alpha}_{W}$ in the $C P$-violating TGC parameter space (see e.g., papers by Hagiwara [1], Bilenky [2], and Gounaris $[3,4]$ ). The relations between these parameters and those contained in the above set, expressed as deviations from the SM, are $\Delta g_{1}^{Z}=\alpha_{W \phi} / c_{w}^{2}, \Delta \kappa_{\gamma}=\alpha_{W \phi}+\alpha_{B \phi}, \Delta \kappa_{Z}=$ $\alpha_{W \phi}-t_{w}^{2} \alpha_{B \phi}$ and $\lambda_{\gamma}=\lambda_{Z}=\alpha_{W}$, where $c_{w}$ and $t_{w}$ are the cosine and tangent of the electroweak mixing angle. Similarly, $\widetilde{\kappa}_{\gamma}=\widetilde{\alpha}_{B W}, \widetilde{\kappa}_{Z}=t_{w}^{2} \widetilde{\alpha}_{B W}$ and $\widetilde{\lambda}_{\gamma}=\widetilde{\lambda}_{Z}=\widetilde{\alpha}_{W}$ within the $C P-$ violating sector. The LEP Collaborations have recently agreed to express their results directly in terms of the parameters $\Delta g_{1}^{Z}$, $\Delta \kappa_{\gamma}$ and $\lambda_{\gamma}$.

At LEP2 the $V W W$ coupling arises in $W$-pair production via $s$-channel exchange or in single $W$ production via the radiation of a virtual photon off the incident $e^{+}$or $e^{-}$. At the TEVATRON hard photon bremstrahlung off a produced $W$ or $Z$ signals the presence of a triple gauge vertex. In order to extract the value of one TGC the others are generally kept fixed to their SM values.

## References

1. K. Hagiwara et al., Nucl. Phys. B282, 253 (1987).
2. M. Bilenky et al., Nucl. Phys. B409, 22 (1993).
3. G. Gounaris et al., CERN 96-01 525.
4. G. Gounaris et al., Eur. Phys. J. C2, 365 (1998).

HTTP://PDG.LBL.GOV
Page $15 \quad$ Created: 12/18/2000 15:07

| VALUE |  |  | EVTS | DOCUMENT ID | TECN | COMMENT |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 0.01 | $\begin{array}{r} +0.09 \\ -0.08 \end{array}$ | OUR | VRAGE |  |  |  |
| 0.01 | $\begin{array}{r} +0.13 \\ -0.12 \end{array}$ |  | 853 | 64 ABBIENDI | 99D OPAL | $\begin{aligned} & E_{\mathrm{Cm}}^{e e}=161+172+183 \\ & G e V \end{aligned}$ |
| -0.04 | $\begin{array}{r} +0.14 \\ -0.12 \end{array}$ |  | 566 | 65 ABREU | 99L DLPH | $E_{\mathrm{cm}}^{e e}{ }_{\text {e }}=183 \mathrm{GeV}$ |
| 0.11 | $\begin{array}{r} +0.19 \\ -0.18 \end{array}$ | $\pm 0.10$ | 1154 | 66 ACCIARRI | 99Q L3 | $E_{\mathrm{cm}}^{e e}=161+172+183$ |

-     - We do not use the following data for averages, fits, limits, etc. - -

$-0.017 \pm 0.018_{-0.003}^{+0.018} \quad 68$ MOLNAR 99 THEO LEP1, SLAC+Tevatron
${ }^{64}$ ABBIENDI 99D combine results from $W^{+} W^{-}$production at different energies. The $95 \%$ confidence interval is $-0.23<\Delta g_{1}^{Z}<0.26$.
65 ABREU 99L use $W^{+} W^{-}, W e \nu_{e}$, and $\nu \bar{\nu} \gamma$ final states. The $95 \%$ confidence interval is $-0.28<\Delta g_{1}^{Z}<0.24$.
66 ACCIARRI 99Q study $W$-pair, single- $W$, and single photon events.
67 ABBOTT 99I perform a simultaneous fit to the $W \gamma, W W \rightarrow$ dilepton, $W W / W Z \rightarrow$ $e \nu j j, W W / W Z \rightarrow \mu \nu j j$, and $W Z \rightarrow$ trilepton data samples. For $\Lambda=2.0 \mathrm{TeV}$, the $95 \%$ CL limits are $-0.37<\Delta g_{1}^{Z}<0.57$, fixing $\lambda_{Z}=\Delta \kappa_{Z}=0$ and assuming Standard Model values for the $W W \gamma$ couplings.
68 MOLNAR 99 extract this value indirectly by fitting high energy electroweak data within the framework of the Standard Model. The central value of the Higgs mass used is 300 GeV and the quoted systematic error is due to its variation between 90 to 1000 GeV .
$\Delta \kappa_{\gamma}$

| VALUE |  |  | EVTS | DOCUMENT ID |  | TECN | COMMENT |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 0.08 | $\pm 0.17$ | OUR | ERAGE |  |  |  |  |
| 0.11 | $\begin{array}{r} +0.52 \\ -0.37 \end{array}$ |  | 853 | 69 ABBIENDI | 99D | OPAL | $\begin{aligned} & E_{\mathrm{Cm}}^{e e}=161+172+183 \\ & \underline{\mathrm{GeV}} \end{aligned}$ |
| -0.08 | $\pm 0.34$ |  | 331 | 70 ABBOTT | 991 | D0 | $E_{\mathrm{cm}}^{p \bar{p}}=1.8 \mathrm{TeV}$ |
| 0.19 | $\begin{array}{r} +0.32 \\ -0.34 \end{array}$ |  | 566 | 71 ABREU | 99L | DLPH | $E_{\mathrm{cm}}^{e e}=183 \mathrm{GeV}$ |
| 0.11 | $\pm 0.25$ | $\pm 0.17$ | 1154 | 72 ACCIARRI | 99Q | L3 | $\begin{aligned} & E_{\mathrm{Cm}}^{e e}=161+172+183 \\ & \mathrm{GeV} \end{aligned}$ |
| 0.05 | $\begin{array}{r} +1.15 \\ -1.10 \end{array}$ | $\pm 0.25$ | 207 | 73 BARATE,R | 98 | ALEP | $\begin{gathered} E_{\mathrm{Cm}}^{e \mathrm{em}}=161+172+183 \\ \mathrm{GeV} \end{gathered}$ |

-     - We do not use the following data for averages, fits, limits, etc.

|  | 15 | 74 BARATE | 99 ALEP | $E_{\mathrm{Cm}}^{e e}=161+172+183$ <br> GeV |
| :--- | :--- | :--- | :--- | :--- |
| $0.016 \pm 0.019+0.009$ |  |  |  |  |

69 ABBIENDI 99D combine results from $W^{+} W^{-}$production at different energies. The $95 \%$ confidence interval is $-0.55<\Delta \kappa_{\gamma}<1.28$.
${ }^{70}$ ABBOTT 99ı perform a simultaneous fit to the $W \gamma, W W \rightarrow$ dilepton, $W W / W Z \rightarrow$ $e \nu j j, W W / W Z \rightarrow \mu \nu j j$, and $W Z \rightarrow$ trilepton data samples. For $\Lambda=2.0 \mathrm{TeV}$, the $95 \% \mathrm{CL}$ limits are $-0.25<\Delta \kappa_{\gamma}<0.39$.
${ }^{71}$ ABREU 99L use $W^{+} W^{-}, W e \nu_{e}$, and $\nu \bar{\nu} \gamma$ final states. The $95 \%$ confidence interval is $-0.46<\Delta \kappa_{\gamma}<0.84$.
${ }^{72}$ ACCIARRI 99Q study $W$-pair, single- $W$, and single photon events.
73 BARATE, R 98 study single photon production in $e^{+} e^{-}$interactions from 161 to 183 GeV . A likelihood fit is performed to the cross section and to the photon energy and angular distributions, taking into account systematic uncertainties. The 95\%CL limits are $-2.2<\Delta \kappa_{\gamma}<2.3$.
${ }^{74}$ BARATE 99L study single $W$ production in $e^{+} e^{-}$interactions from 161 to 183 GeV . They obtain $95 \%$ CL limits of $-1.6<\kappa_{\gamma}<1.5$, which we convert to $-2.6<\Delta \kappa_{\gamma}<0.5$ for $\lambda_{\gamma}=0$.
${ }^{75}$ MOLNAR 99 extract this value indirectly by fitting high energy electroweak data within the framework of the Standard Model. The central value of the Higgs mass used is 300 GeV and the quoted systematic error is due to its variation between 90 to 1000 GeV .
76 ACCIARRI 98 N study single $W$ production in $e^{+} e^{-}$interactions from 130 to 183 GeV . The $95 \% C L$ limits are $-0.46<\Delta \kappa_{\gamma}<0.57$.
$\lambda_{\gamma}$

| VALUE | EVTS | DOCUMENT ID | TECN | COMMENT |
| :---: | :---: | :---: | :---: | :---: |
| $-0.04{ }_{=}^{+0.07}$ OUR AVERAGE |  |  |  |  |
| $-0.10{ }_{-0.12}^{+0.13}$ | 853 | 77 ABBIENDI | 99D OPAL | $\underset{\mathrm{GeV}}{E_{\mathrm{Gm}}^{e e}=161+172+183}$ |
| $\begin{gathered} 0.00_{-0.09}^{+0.10} \end{gathered}$ | 331 | 78 ABBOTT | 991 D0 | $E_{\mathrm{Cm}}^{p \bar{p}}=1.8 \mathrm{TeV}$ |
| $\begin{array}{r} -0.15_{-0.15}^{+0.19} \end{array}$ | 566 | 79 ABREU | 99L DLPH | $E_{\mathrm{Cm}}^{e e}=183 \mathrm{GeV}$ |
| $0.10{ }_{-0.20}^{+0.22} \pm 0.10$ | 1154 | 80 ACCIARRI | 99Q L3 | $\begin{aligned} & E_{\mathrm{Cm}}^{e e}=161+172+183 \\ & \mathrm{GeV} \end{aligned}$ |
| $-0.05{ }_{-1.45}^{+1.55} \pm 0.30$ | 207 | 81 BARATE,R | 98 ALEP | $\begin{aligned} & E_{\mathrm{Cm}}^{e e}=161+172+183 \\ & \quad \mathrm{GeV} \end{aligned}$ |

-     - We do not use the following data for averages, fits, limits, etc. - -
-0.48

-0.21 $\quad 86 \quad 83$ ACCIARRI $\quad 98 \mathrm{~N}$ L3 $\quad$| GeV |
| :---: |
| Repl. by ACCIARRI 99Q |

77 ABBIENDI 99D combine results from $W^{+} W^{-}$production at different energies. The $95 \%$ confidence interval is $-0.33<\lambda_{\gamma}<0.16$.
78 ABBOTT 99I perform a simultaneous fit to the $W \gamma, W W \rightarrow$ dilepton, $W W / W Z \rightarrow$ $e \nu j j, W W / W Z \rightarrow \mu \nu j j$, and $W Z \rightarrow$ trilepton data samples. For $\Lambda=2.0 \mathrm{TeV}$, the $95 \%$ CL limits are $-0.18<\lambda_{\gamma}<0.19$.
${ }^{79}$ ABREU 99L use $W^{+} W^{-}, W e \nu_{e}$, and $\nu \bar{\nu} \gamma$ final states. The $95 \%$ confidence interval is $-0.44<\lambda_{\gamma}<0.24$.
80 ACCIARRI 99Q study $W$-pair, single- $W$, and single photon events.
81 BARATE, R 98 study single photon production in $e^{+} e^{-}$interactions from 161 to 183 GeV . A likelihood fit is performed to the cross section and to the photon energy and angular distributions, taking into account systematic uncertainties. The 95\%CL limits are $-3.1<\lambda_{\gamma}<3.2$.
82 BARATE 99L study single $W$ production in $e^{+} e^{-}$interactions from 161 to 183 GeV . The $95 \% \mathrm{CL}$ limits are $-1.6<\lambda_{\gamma}<1.6$ for $\Delta \kappa_{\gamma}=0$.
83 ACCIARRI 98N study single $W$ production in $e^{+} e^{-}$interactions from 130 to 183 GeV . The $95 \%$ CL limits are $-0.86<\lambda_{\gamma}<0.75$.

## $\Delta g_{5}^{7}$

This coupling is $C P$ conserving but $C$ and $P$ violating.

| VALUE | EVTS | DOCUMENT | TECN | COMMENT |
| :---: | :---: | :---: | :---: | :---: |
| $-0.44_{-0.22}^{+0.23} \pm 0.12$ | 1154 | 84 ACCIARRI | 99Q L3 | $\underset{\mathrm{GeV}}{E_{\mathrm{Cm}}^{e e}=161+172+183}$ |


| $\alpha_{W}{ }_{\boldsymbol{\phi}}$ |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: |
| VALUE |  | DOCUMEN | TECN | COMMENT |
| 0.05 $\pm 0.20$ OUR AVERAGE |  |  |  |  |
| $0.22_{-0.28}^{+0.25} \pm 0.06$ | 89 | ${ }^{85}$ Abreu | 98K DLPH | $E_{\mathrm{Cm}}^{e}=161+172 \mathrm{GeV}$ |
| $-0.14{ }_{-0.25}^{+0.27}+0.14$ | 78 | 86 BARATE | 98Y ALEP | $E_{\mathrm{cm}}^{e e}=172 \mathrm{GeV}$ |

-     - We do not use the following data for averages, fits, limits, etc. - - -

$$
331 \quad 87 \text { ABBOTT } \quad 99 \text { D0 } \quad E_{\mathrm{cm}}^{p \bar{p}}=1.8 \mathrm{TeV}
$$

${ }^{85}$ ABREU 98 K obtain this result using both $W$ pair production and single $W\left(W e \nu_{e}\right)$ production.
${ }^{86}$ BARATE 98 Y obtain this value using semileptonic and hadronic decay modes in $W$ pair production.
${ }^{87}$ ABBOTT 991 perform a simultaneous fit to the $W \gamma, W W \rightarrow$ dilepton, $W W / W Z \rightarrow$ $e \nu j j, W W / W Z \rightarrow \mu \nu j j$, and $W Z \rightarrow$ trilepton data samples. For $\Lambda=2.0 \mathrm{TeV}$, the $95 \%$ CL limits are $-0.18<\alpha_{W}{ }_{\phi}<0.36$, fixing $\alpha_{B \phi}=\alpha_{W}=0$.
$\alpha_{W}$
$\frac{\text { VALUE }}{0.1} \pm \mathbf{0 . 4}$ OUR AVERAGE EVTS DOCUMENT ID TECN COMMENT

| $0.11_{-0.49}^{+0.48} \pm 0.09$ | 89 | 88 ABREU | $98 \kappa$ DLPH | $E_{\mathrm{cm}}^{e e}=161+172 \mathrm{GeV}$ |
| :--- | :--- | :--- | :--- | :--- |
| $0.06_{-0.50}^{+0.56+0.12}$ | 78 | 89 BARATE | 98 Y ALEP | $E_{\mathrm{cm}}^{e e}=172 \mathrm{GeV}$ |

-     - We do not use the following data for averages, fits, limits, etc. - • -

$$
33190 \text { ABBOTT } 99 \text { D0 } \quad E_{\mathrm{cm}}^{p \bar{p}}=1.8 \mathrm{TeV}
$$

${ }^{88}$ ABREU 98 K obtain this result using both $W$ pair production and single $W\left(W e \nu_{e}\right)$ production.
${ }^{89}$ BARATE 98 Y obtain this value using semileptonic and hadronic decay modes in $W$ pair production.
${ }^{90}$ ABBOTT 99ı perform a simultaneous fit to the $W \gamma, W W \rightarrow$ dilepton, $W W / W Z \rightarrow$ $e \nu j j, W W / W Z \rightarrow \mu \nu j j$, and $W Z \rightarrow$ trilepton data samples. For $\Lambda=2.0 \mathrm{TeV}$, the $95 \%$ CL limits are $-0.18<\alpha_{W}<0.19$, fixing $\alpha_{B \phi}=\alpha_{W}=0$.

## $\alpha_{B \phi}$

VALUE EVTS
DOCUMENT ID TECN COMMENT
$0.4{ }_{-0.8}^{+0.5}$ OUR AVERAGE

| $0.22_{-0.83}^{+0.66} \pm 0.24$ | 89 | 91 ABREU | 98 k DLPH | $E_{\mathrm{cm}}^{e e}=161+172 \mathrm{GeV}$ |
| :--- | :--- | :--- | :--- | :--- |
| $1.01_{-1.75}^{+0.71} \pm 0.33$ | 78 | 92 BARATE | $98 Y$ ALEP | $E_{\mathrm{cm}}^{e e}=172 \mathrm{GeV}$ |

-     - We do not use the following data for averages, fits, limits, etc.

$$
331 \quad 93 \text { ABBOTT } \quad 99 \text { D0 } E_{\mathrm{cm}}^{p \bar{p}}=1.8 \mathrm{TeV}
$$

91 ABREU 98K obtain this result using both $W$ pair production and single $W\left(W e \nu_{e}\right)$ production.
92 BARATE 98 Y obtain this value using semileptonic and hadronic decay modes in $W$ pair production.
93 ABBOTT 99ı perform a simultaneous fit to the $W \gamma, W W \rightarrow$ dilepton, $W W / W Z \rightarrow$ $e \nu j j, W W / W Z \rightarrow \mu \nu j j$, and $W Z \rightarrow$ trilepton data samples. For $\Lambda=2.0 \mathrm{TeV}$, the $95 \%$ CL limits are $-0.67<\alpha_{B \phi}<0.56$, fixing $\alpha_{W}=\alpha_{W}=0$.

## $\widetilde{\alpha}_{B} W$

| VALUE | EVTS | DOCUMEN | TECN | COMMENT |
| :---: | :---: | :---: | :---: | :---: |
| $0.11_{-0.88}^{+0.71} \pm 0.09$ | 89 | 94 ABREU | 98K DLPH | $E_{\mathrm{cm}}^{e \mathrm{em}}=161+172 \mathrm{GeV}$ |

${ }^{94}$ ABREU 98 K obtain this result using both $W$ pair production and single $W\left(W e \nu_{e}\right)$ production.

## $\widetilde{\alpha}_{W}$

| VALUE | EVTS | DOCUMEN | TECN | COMMENT |
| :---: | :---: | :---: | :---: | :---: |
| $0.19{ }_{=0.41}^{+0.28} \pm 0.11$ | 89 | 95 ABREU | 98K DLPH | $E_{\mathrm{Cm}}^{\mathrm{ee}} \mathrm{e}=161+172 \mathrm{GeV}$ |

${ }^{95}$ ABREU 98K obtain this result using both $W$ pair production and single $W\left(W e \nu_{e}\right)$ production.

## W ANOMALOUS MAGNETIC MOMENT $(\Delta \kappa)$

The full magnetic moment is given by $\mu_{W}=e(1+\kappa+\lambda) / 2 m_{W}$. In the Standard Model, at tree level, $\kappa=1$ and $\lambda=0$. Some papers have defined $\Delta \kappa=1-\kappa$ and assume that $\lambda=0$. Note that the electric quadrupole moment is given by $-e(\kappa-\lambda) / m_{W}^{2}$. A description of the parameterization of these moments and additional references can be found in HAGIWARA 87 and BAUR 88. The parameter $\Lambda$ appearing in the theoretical limits below is a regularization cutoff which roughly corresponds to the energy scale where the structure of the $W$ boson becomes manifest.
$\operatorname{VALUE}\left(e / 2 m_{W}\right)$
DOCUMENT ID TECN

-     - We do not use the following data for averages, fits, limits, etc. - - -

| 96 ABE | 95G CDF |
| :---: | :---: |
| 97 ALITTI | 92C UA2 |
| 98 SAMUEL | 92 THEO |
| 99 SAMUEL | 91 THEO |
| 100 GRIFOLS | 88 THEO |
| 101 GROTCH | 87 THEO |
| 102 VANDERBIJ | 87 THEO |
| 103 GRAU | 85 THEO |
| 104 SUZUKI | 85 THEO |
| 105 HERZOG | 84 THEO |

${ }^{96} \mathrm{ABE} 95 \mathrm{G}$ report $-1.3<\kappa<3.2$ for $\lambda=0$ and $-0.7<\lambda<0.7$ for $\kappa=1$ in $p \bar{p} \rightarrow e \nu_{e} \gamma \mathrm{X}$ and $\mu \nu_{\mu} \gamma \mathrm{X}$ at $\sqrt{s}=1.8 \mathrm{TeV}$.
${ }^{97}$ ALITTI 92C measure $\kappa=1_{-2.2}^{+2.6}$ and $\lambda=0_{-1.8}^{+1.7}$ in $p \bar{p} \rightarrow e \nu \gamma+X$ at $\sqrt{s}=630 \mathrm{GeV}$. At $95 \%$ CL they report $-3.5<\kappa<5.9$ and $-3.6<\lambda<3.5$.

98 SAMUEL 92 use preliminary CDF and UA2 data and find $-2.4<\kappa<3.7$ at $96 \% \mathrm{CL}$ and $-3.1<\kappa<4.2$ at $95 \% \mathrm{CL}$ respectively. They use data for $W \gamma$ production and radiative $W$ decay.
99 SAMUEL 91 use preliminary CDF data for $p \bar{p} \rightarrow W \gamma X$ to obtain $-11.3 \leq \Delta \kappa \leq$ 10.9. Note that their $\kappa=1-\Delta \kappa$.

100 GRIFOLS 88 uses deviation from $\rho$ parameter to set limit $\Delta \kappa \lesssim 65\left(M_{W}^{2} / \Lambda^{2}\right)$.
101 GROTCH 87 finds the limit $-37<\Delta \kappa<73.5$ ( $90 \%$ CL) from the experimental limits on $e^{+} e^{-} \rightarrow \nu \bar{\nu} \gamma$ assuming three neutrino generations and $-19.5<\Delta \kappa<56$ for four generations. Note their $\Delta \kappa$ has the opposite sign as our definition.
102 VANDERBIJ 87 uses existing limits to the photon structure to obtain $|\Delta \kappa|<33$ $\left(m_{W} / \Lambda\right)$. In addition VANDERBIJ 87 discusses problems with using the $\rho$ parameter of the Standard Model to determine $\Delta \kappa$.
103 GRAU 85 uses the muon anomaly to derive a coupled limit on the anomalous magnetic dipole and electric quadrupole $(\lambda)$ moments $1.05>\Delta \kappa \ln \left(\Lambda / m_{W}\right)+\lambda / 2>-2.77$. In the Standard Model $\lambda=0$.
104 SUZUKI 85 uses partial-wave unitarity at high energies to obtain $|\Delta \kappa| \lesssim 190$ $\left(m_{W} / \Lambda\right)^{2}$. From the anomalous magnetic moment of the muon, SUZUKI 85 obtains $|\Delta \kappa| \lesssim 2.2 / \ln \left(\Lambda / m_{W}\right)$. Finally SUZUKI 85 uses deviations from the $\rho$ parameter and obtains a very qualitative, order-of-magnitude limit $|\Delta \kappa| \lesssim 150\left(m_{W} / \Lambda\right)^{4}$ if $|\Delta \kappa| \ll$ 1.

105 HERZOG 84 consider the contribution of $W$-boson to muon magnetic moment including anomalous coupling of $W W \gamma$. Obtain a limit $-1<\Delta \kappa<3$ for $\Lambda \gtrsim 1 \mathrm{TeV}$.

## W REFERENCES

| ABBOTT | 00 | PRL 84222 | B. Abbott et al. | (D0 Collab.) |
| :---: | :---: | :---: | :---: | :---: |
| ABBOTT | 00B | PR D61 072001 | B. Abbott et al. | (D0 Collab.) |
| ABREU | 00K | PL B479 89 | P. Abreu et al. | (DELPHI Collab.) |
| ABBIENDI | 99C | PL B453 138 | G. Abbiendi et al. | (OPAL Collab.) |
| ABBIENDI | 99D | EPJ C8 191 | G. Abbiendi et al. | (OPAL Collab.) |
| ABBIENDI | 99N | PL B453 153 | G. Abbiendi et al. | (OPAL Collab.) |
| ABBOTT | 99H | PR D60 052003 | B. Abbott et al. | (D0 Collab.) |
| ABBOTT | 991 | PR D60 072002 | B. Abbott et al. | (D0 Collab.) |
| ABREU | 99K | PL B456 310 | P. Abreu et al. | (DELPHI Collab.) |
| ABREU | 99L | PL B459 382 | P. Abreu et al. | (DELPHI Collab.) |
| ABREU | 99T | PL B462 410 | P. Abreu et al. | (DELPHI Collab.) |
| ACCIARRI | 99 | PL B454 386 | M. Acciarri et al. | (L3 Collab.) |
| ACCIARRI | 99Q | PL B467 171 | M. Acciarri et al. | (L3 Collab.) |
| BARATE | 99 | PL B453 121 | R. Barate et al. | (ALEPH Collab.) |
| BARATE | 991 | PL B453 107 | R. Barate et al. | (ALEPH Collab.) |
| BARATE | 99L | PL B462 389 | R. Barate et al. | (ALEPH Collab.) |
| BARATE | 99M | PL B465 349 | R. Barate et al. | (ALEPH Collab.) |
| MOLNAR | 99 | PL B461 149 | P. Molnar, M. Grunewald |  |
| ABBOTT | 98N | PR D58 092003 | B. Abbott et al. | (D0 Collab.) |
| ABBOTT | 980 | PRL 803008 | B. Abbott et al. | (D0 Collab.) |
| ABBOTT | 98P | PR D58 012002 | B. Abbott et al. | (D0 Collab.) |
| ABE | 98H | PR D58 031101 | F. Abe et al. | (CDF Collab.) |
| ABE | 98P | PR D58 091101 | F. Abe et al. | (CDF Collab.) |
| ABREU | 98B | EPJ C2 581 | P. Abreu et al. | (DELPHI Collab.) |
| ABREU | 98C | PL B416 233 | P. Abreu et al. | (DELPHI Collab.) |
| ABREU | 98K | PL B423 194 | P. Abreu et al. | (DELPHI Collab.) |
| ABREU | 98N | PL B439 209 | P. Abreu et al. | (DELPHI Collab.) |
| ACCIARRI | 98N | PL B436 417 | M. Acciarri et al. | (L3 Collab.) |
| ACCIARRI | 98P | PL B436 437 | M. Acciarri et al. | (L3 Collab.) |
| ACKERSTAFF | 98D | EPJ C1 395 | K. Ackerstaff et al. | (OPAL Collab.) |
| BARATE | 98B | PL B422 384 | R. Barate et al. | (ALEPH Collab.) |
| BARATE | 98 Y | PL B422 369 | R. Barate et al. | (ALEPH Collab.) |
| BARATE,R | 98 | PL B445 239 | R. Barate et al. | (ALEPH Collab.) |
| ACCIARRI | 97 | PL B398 223 | M. Acciarri et al. | (L3 Collab.) |
| ACCIARRI | 97M | PL B407 419 | M. Acciarri et al. | (L3 Collab.) |
| ACCIARRI | 97 S | PL B413 176 | M. Acciarri et al. | (L3 Collab.) |
| BARATE | 97 | PL B401 347 | R. Barate et al. | (ALEPH Collab.) |
| BARATE | 97 S | PL B415 435 | R. Barate et al. | (ALEPH Collab.) |


| ABACHI | 96E | PRL 773309 | S. Abachi et al. | (D0 Collab.) |
| :---: | :---: | :---: | :---: | :---: |
| ACKERSTAFF | 96B | PL B389 416 | K. Ackerstaff et al. | (OPAL Collab.) |
| AID | 96D | PL B379 319 | S. Aid et al. | (H1 Collab.) |
| ABACHI | 95D | PRL 751456 | S. Abachi et al. | (D0 Collab.) |
| ABE | 95C | PRL 74341 | F. Abe et al. | (CDF Collab.) |
| ABE | 95G | PRL 741936 | F. Abe et al. | (CDF Collab.) |
| ABE | 95P | PRL 7511 | F. Abe et al. | (CDF Collab.) |
| Also | 95Q | PR D52 4784 | F. Abe et al. | (CDF Collab.) |
| ABE | 95W | PR D52 2624 | F. Abe et al. | (CDF Collab.) |
| Also | 94B | PRL 73220 | F. Abe et al. | (CDF Collab.) |
| ROSNER | 94 | PR D49 1363 | J.L. Rosner, M.P. Worah, T. Takeuchi | i (EFI, FNAL) |
| ABE | 92E | PRL 683398 | F. Abe et al. | (CDF Collab.) |
| ABE | 921 | PRL 6928 | F. Abe et al. | (CDF Collab.) |
| ALITTI | 92 | PL B276 365 | J. Alitti et al. | (UA2 Collab.) |
| ALITTI | 92B | PL B276 354 | J. Alitti et al. | (UA2 Collab.) |
| ALITTI | 92C | PL B277 194 | J. Alitti et al. | (UA2 Collab.) |
| ALITTI | 92D | PL B277 203 | J. Alitti et al. | (UA2 Collab.) |
| ALITTI | 92F | PL B280 137 | J. Alitti et al. | (UA2 Collab.) |
| SAMUEL | 92 | PL B280 124 | M.A. Samuel et al. | (OKSU, CARL) |
| ABE | 91C | PR D44 29 | F. Abe et al. | (CDF Collab.) |
| ALBAJAR | 91 | PL B253 503 | C. Albajar et al. | (UA1 Collab.) |
| ALITTI | 91C | ZPHY C52 209 | J. Alitti et al. | (UA2 Collab.) |
| SAMUEL | 91 | PRL 679 | M.A. Samuel et al. | (OKSU, CARL) |
| Also | 91C | PRL 672920 erratum |  |  |
| ABE | 90 | PRL 64152 | F. Abe et al. | (CDF Collab.) |
| Also | 91 C | PR D44 29 | F. Abe et al. | (CDF Collab.) |
| ABE | 90G | PRL 652243 | F. Abe et al. | (CDF Collab.) |
| Also | 91B | PR D43 2070 | F. Abe et al. | (CDF Collab.) |
| ALBAJAR | 90 | PL B241 283 | C. Albajar et al. | (UA1 Collab.) |
| ALITTI | 90B | PL B241 150 | J. Alitti et al. | (UA2 Collab.) |
| ALITTI | 90C | ZPHY C47 11 | J. Alitti et al. | (UA2 Collab.) |
| ABE | 891 | PRL 621005 | F. Abe et al. | (CDF Collab.) |
| ALBAJAR | 89 | ZPHY C44 15 | C. Albajar et al. | (UA1 Collab.) |
| BAUR | 88 | NP B308 127 | U. Baur, D. Zeppenfeld | (FSU, WISC) |
| GRIFOLS | 88 | IJMP A3 225 | J.A. Grifols, S. Peris, J. Sola | (BARC, DESY) |
| Also | 87 | PL B197 437 | J.A. Grifols, S. Peris, J. Sola | (BARC, DESY) |
| ALBAJAR | 87 | PL B185 233 | C. Albajar et al. | (UA1 Collab.) |
| ANSARI | 87 | PL B186 440 | R. Ansari et al. | (UA2 Collab.) |
| ANSARI | 87 C | PL B194 158 | R. Ansari et al. | (UA2 Collab.) |
| GROTCH | 87 | PR D36 2153 | H. Grotch, R.W. Robinett | (PSU) |
| HAGIWARA | 87 | NP B282 253 | K. Hagiwara et al. | (KEK, UCLA, FSU) |
| VANDERBIJ | 87 | PR D35 1088 | J.J. van der Bij | (FNAL) |
| APPEL | 86 | ZPHY C30 1 | J.A. Appel et al. | (UA2 Collab.) |
| ARNISON | 86 | PL 166B 484 | G.T.J. Arnison et al. | (UA1 Collab.) |
| ALTARELLI | 85B | ZPHY C27 617 | G. Altarelli, R.K. Ellis, G. Martinelli | (CERN+) |
| GRAU | 85 | PL 154B 283 | A. Grau, J.A. Grifols | (BARC) |
| SUZUKI | 85 | PL 153B 289 | M. Suzuki | (LBL) |
| ARNISON | 84D | PL 134B 469 | G.T.J. Arnison et al. | (UA1 Collab.) |
| HERZOG | 84 | PL 148B 355 | F. Herzog | (WISC) |
| Also | 84B | PL 155B 468 erratum | F. Herzog | (WISC) |
| ARNISON | 83 | PL 122B 103 | G.T.J. Arnison et al. | (UA1 Collab.) |
| BANNER | 83B | PL 122B 476 | M. Banner et al. | (UA2 Collab.) |

