## LIGHT QUARKS ( $u, d, s$ )

OMITTED FROM SUMMARY TABLE

## u-QUARK MASS

The $u$-, $d$-, and $s$-quark masses are estimates of so-called "current-quark masses," in a mass- independent subtraction scheme such as $\overline{\mathrm{MS}}$. The ratios $m_{u} / m_{d}$ and $m_{s} / m_{d}$ are extracted from pion and kaon masses using chiral symmetry. The estimates of $d$ and $u$ masses are not without controversy and remain under active investigation. Within the literature there are even suggestions that the $u$ quark could be essentially massless. The $s$-quark mass is estimated from $\operatorname{SU}(3)$ splittings in hadron masses.

Starting with this edition of the Review, we have normalized the $\overline{\mathrm{MS}}$ masses at a renormalization scale of $\mu=2 \mathrm{GeV}$. Results quoted in the literature at $\mu=1 \mathrm{GeV}$ have been rescaled by dividing by 1.35 .

## VALUE (MeV) <br> DOCUMENT ID TECN COMMENT

## 1.5 to 5 OUR EVALUATION

-     - We do not use the following data for averages, fits, limits, etc. - - -

| $3.9 \pm 1.1$ | 1 | JAMIN | 95 |
| :--- | :--- | :--- | :--- |
| $3.0 \pm 0.7$ | 2 NARISON | 95 CHEO | $\overline{\mathrm{MS}}$ tHEO scheme |
| $\overline{\mathrm{MS}}$ scheme |  |  |  |
|  | 3 CHOI | 92 B THEO |  |
| 4.3 | 4 BARDUCCI | 88 THEO |  |
| $3.8 \pm 1.1$ | 5 GASSER | 82 THEO |  |

1 JAMIN 95 uses QCD sum rules at next-to-leading order. We have rescaled $m_{u}(1 \mathrm{GeV})$ $=5.3 \pm 1.5$ to $\mu=2 \mathrm{GeV}$.
${ }^{2}$ For NARISON 95C, we have rescaled $m_{u}(1 \mathrm{GeV})=4 \pm 1$ to $\mu=2 \mathrm{GeV}$.
${ }^{3} \mathrm{CHOI} 92 \mathrm{~B}$ argues that $m_{u}=0$ is okay based on instanton contributions to the chiral coefficients. Disagrees with DONOGHUE 92 and DONOGHUE 92B.
${ }^{4}$ BARDUCCI 88 uses a calculation of the effective potential for $\bar{\psi} \psi$ in QCD, and estimates for $\Sigma\left(p^{2}\right)$. We have rescaled $m_{u}(1 \mathrm{GeV})=5.8$ to $\mu=2 \mathrm{GeV}$.
${ }^{5}$ GASSER 82 uses chiral perturbation theory for the mass ratios, and uses QCD sum rules to extract the absolute values. We have rescaled $m_{u}(1 \mathrm{GeV})=5.1 \pm 1.5$ to $\mu=2 \mathrm{GeV}$.

## d-QUARK MASS

See the comment for the $u$ quark above.
Starting with this edition of the Review, we have normalized the $\overline{\mathrm{MS}}$ masses at a renormalization scale of $\mu=2 \mathrm{GeV}$. Results quoted in the literature at $\mu=1 \mathrm{GeV}$ have been rescaled by dividing by 1.35 .
VALUE $(\mathrm{MeV})$ DOCUMENT ID $\quad$ TECN COMMENT

## 3 to 9 OUR EVALUATION

-     - We do not use the following data for averages, fits, limits, etc. • • •

| $7.0 \pm 1.1$ | 6 | JAMIN | 95 | THEO |
| :--- | :--- | :--- | :--- | :--- |
| $.4 \pm 0.7$ | 7 NARISON scheme | 95 THEO | $\overline{\mathrm{MS}}$ scheme |  |
|  | 8 ADAMI | 93 | THEO |  |
|  | 9 NEFKENS | 92 | THEO |  |
| 6.2 | 10 BARDUCCI | 88 | THEO |  |
|  | 11 DOMINGUEZ | 87 | THEO |  |
|  | 12 KREMER | 84 | THEO |  |
|  | 13 GASSER | 82 | THEO |  |

6 JAMIN 95 uses QCD sum rules at next-to-leading order. We have rescaled $m_{d}(1 \mathrm{GeV})$ $=9.4 \pm 1.5$ to $\mu=2 \mathrm{GeV}$.
${ }^{7}$ For NARISON 95C, we have rescaled $m_{d}(1 \mathrm{GeV})=10 \pm 1$ to $\mu=2 \mathrm{GeV}$.
${ }^{8}$ ADAMI 93 obtain $m_{d}-m_{u}=3 \pm 1 \mathrm{MeV}$ at $\mu=0.5 \mathrm{GeV}$ using isospin-violating effects in QCD sum rules.
${ }^{9}$ NEFKENS 92 results for $m_{d}-m_{u}$ are $3.1 \pm 0.4 \mathrm{MeV}$ from meson masses and $3.6 \pm 0.4$ MeV from baryon masses.
$1^{10}$ BARDUCCI 88 uses a calculation of the effective potential for $\bar{\psi} \psi$ in QCD, and estimates for $\Sigma\left(p^{2}\right)$. We have rescaled $m_{d}(1 \mathrm{GeV})=8.4$ to $\mu=2 \mathrm{GeV}$.
11 DOMINGUEZ 87 uses QCD sum rules to obtain $m_{u}+m_{d}=15.5 \pm 2.0 \mathrm{MeV}$ and $m_{d}-$ $m_{u}=6 \pm 1.5 \mathrm{MeV}$.
12 KREMER 84 obtain $m_{u}+m_{d}=21 \pm 2 \mathrm{MeV}$ at $Q^{2}=1 \mathrm{GeV}^{2}$ using SVZ values for quark condensates; they obtain $m_{u}+m_{d}=35 \pm 3 \mathrm{MeV}$ at $Q^{2}=1 \mathrm{GeV}^{2}$ using factorization values for quark condensates.
13 GASSER 82 uses chiral perturbation theory for the mass ratios, and uses QCD sum rules to extract the absolute values. We have rescaled $m_{d}(1 \mathrm{GeV})=8.9 \pm 2.6$ to $\mu=2 \mathrm{GeV}$.

$$
\bar{m}=\left(m_{u}+m_{d}\right) / 2
$$

See the comments for the $u$ quark above.
Starting with this edition of the Review, we have normalized the $\overline{\mathrm{MS}}$ masses at a renormalization scale of $\mu=2 \mathrm{GeV}$. Results quoted in the literature at $\mu=1 \mathrm{GeV}$ have been rescaled by dividing by 1.35 .


14 DOSCH 98 use sum rule determinations of the quark condensate and chiral perturbation theory to obtain $9.4 \leq\left(m_{u}+m_{d}\right)(1 \mathrm{GeV}) \leq 15.7 \mathrm{MeV}$. We have converted to result to $\mu=2 \mathrm{GeV}$.
15 EICKER 97 use lattice gauge computations with two dynamical light flavors.
16 GOUGH 97 use lattice gauge computations in the quenched approximation. Correcting for quenching gives $2.1<\bar{m}<3.5 \mathrm{MeV}$ at $\mu=2 \mathrm{GeV}$.
17 GUPTA 97 use Lattice Monte Carlo computations in the quenched approximation. The value for two light dynamic flavors at $\mu=2 \mathrm{GeV}$ is $2.7 \pm 0.3 \pm 0.3 \mathrm{MeV}$.
18 LELLOUCH 97 obtain lower bounds on quark masses using hadronic spectral functions.
19 BIJNENS 95 determines $m_{u}+m_{d}(1 \mathrm{GeV})=12 \pm 2.5 \mathrm{MeV}$ using finite energy sum rules. We have rescaled this to 2 GeV .

## $s$-QUARK MASS

See the comment for the $u$ quark above.
Starting with this edition of the Review, we have normalized the $\overline{\mathrm{MS}}$ masses at a renormalization scale of $\mu=2 \mathrm{GeV}$. Results quoted in the literature at $\mu=1 \mathrm{GeV}$ have been rescaled by dividing by 1.35 .
$\frac{\operatorname{VALUE}(\mathrm{MeV})}{\mathbf{6 0} \text { to } \mathbf{1 7 0} \text { OUR EVALUATION }}$ DOCUMENT ID TECN COMMENT

-     - We do not use the following data for averages, fits, limits, etc. - - -

| 148 | $\pm 48$ |  | 20 | CHETYRKIN | 98 |
| ---: | :--- | :--- | :--- | :--- | :--- |
| THEO | $\overline{\mathrm{MS}}$ scheme |  |  |  |  |
| 103 | $\pm 10$ | 21 CUCCHIERI | 98 | LATT | $\overline{\mathrm{MS}}$ scheme |
| 115 | $\pm 19$ | 22 DOMINGUEZ | 98 | THEO | $\overline{\mathrm{MS}}$ scheme |
| $>90$ | $\pm 9$ | 23 DOSCH | 98 | THEO | $\overline{\mathrm{MS}}$ scheme |
| $>30$ |  | 24 LEBED | 98 | THEO | $\overline{\mathrm{MS}}$ scheme |
| 84 | $\pm 80$ | 25 MALTMAN | 98 | THEO | $\overline{\mathrm{MS}}$ scheme |
| $<163$ | $\pm 81$ | 26 MALTMAN | 98 | THEO | $\overline{\mathrm{MS}}$ scheme |
| $152.4 \pm 14.1$ | 27 CHETYRKIN | 97 | THEO | $\overline{\mathrm{MS}}$ scheme |  |
| $\geq 89$ |  | 28 COLANGELO | 97 | THEO | $\overline{\mathrm{MS}}$ scheme |
| 140 | $\pm 20$ | 29 EICKER | 97 | LATT | $\overline{\mathrm{MS}}$ scheme |
| 95 | $\pm 16$ | 30 GOUGH | 97 | LATT | $\overline{\mathrm{MS}}$ scheme |
| 100 | $\pm 21$ | $\pm 10$ | 31 GUPTA | 97 | LATT |
| $\overline{\mathrm{MS}}$ scheme |  |  |  |  |  |
| $>37$ |  | 32 LELLOUCH | 97 | THEO | $\overline{\mathrm{MS}}$ scheme |
| 127 | $\pm 11$ | 33 CHETYRKIN | 95 | THEO | $\overline{\mathrm{MS}}$ scheme |
| 140 | $\pm 24$ | 34 JAMIN | 95 | THEO | $\overline{\mathrm{MS}}$ scheme |
| 146 | $\pm 22$ | 35 NARISON | $95 C$ | THEO | $\overline{\mathrm{MS}}$ scheme |
|  |  | 36 NEFKENS | 92 | THEO |  |
| 144 | $\pm 3$ | 37 DOMINGUEZ | 91 | THEO |  |
| 88 |  | 38 BARDUCCI | 88 | THEO |  |
|  |  | 39 KREMER | 84 | THEO |  |
| 130 | $\pm 41$ | 40 GASSER | 82 | THEO |  |

${ }^{20}$ CHETYRKIN 98 uses spectral moments of hadronic $\tau$ decays to determine $m_{s}(1 \mathrm{GeV})=200 \pm 70 \mathrm{MeV}$. We have rescaled the result to $\mu=2 \mathrm{GeV}$.
${ }^{21}$ CUCCHIERI 98 obtains the quark mass using a quenched lattice computation of the hadronic spectrum.
22 DOMINGUEZ 98 uses hadronic spectral function sum rules (to four loops, and including dimension six operators) to determine $m_{s}(1 \mathrm{GeV})<155 \pm 25 \mathrm{MeV}$. We have rescaled the result to $\mu=2 \mathrm{GeV}$.
23 DOSCH 98 use sum rule determinations of the quark condensate and chiral perturbation theory to obtain $m_{s}(1 \mathrm{GeV})>121 \pm 12 \mathrm{MeV}$. We have converted the result to $\mu=2 \mathrm{GeV}$.
${ }^{24}$ LEBED 98 obtain lower bounds of 41,90 , and 139 MeV for $m_{s}(1 \mathrm{GeV})$ using dispersion relations and chiral perturbation theory. The numbers assume the chiral perturbation theory form factor is accurate to $5 \%, 1 \%$, and $0.05 \%$, respectively. We have used the first number converted to $\mu=2$.
25 MALTMAN 98 uses $\tau$ decay sum rules to determine $m_{s}(1 \mathrm{GeV})=113 \pm 107 \mathrm{MeV}$. We have rescaled the result to $\mu=2 \mathrm{GeV}$.
${ }^{26}$ MALTMAN 98B uses spectral moments of hadronic $\tau$ decays to determine $m_{s}(1 \mathrm{GeV})<$ $220 \pm 110 \mathrm{MeV}$. We have rescaled the result to $\mu=2 \mathrm{GeV}$.
27 CHETYRKIN 97 obtains $205.5 \pm 19.1 \mathrm{MeV}$ at $\mu=1 \mathrm{GeV}$ from QCD sum rules including fourth-order QCD corrections. We have rescaled the result to 2 GeV .
${ }^{28}$ COLANGELO 97 is QCD sum rule computation. We have rescaled $m_{s}(1 \mathrm{GeV})>120$ to $\mu=2 \mathrm{GeV}$.
${ }^{29}$ EICKER 97 use lattice gauge computations with two dynamical light flavors.
30 GOUGH 97 use lattice gauge computations in the quenched approximation. Correcting for quenching gives $54<m_{s}<92 \mathrm{MeV}$ at $\mu=2 \mathrm{GeV}$.
31 GUPTA 97 use Lattice Monte Carlo computations in the quenched approximation. The value for two light dynamical flavors at $\mu=2 \mathrm{GeV}$ is $68 \pm 12 \pm 7 \mathrm{MeV}$.
${ }^{32}$ LELLOUCH 97 obtain lower bounds on quark masses using hadronic spectral functions.
${ }^{33}$ CHETYRKIN 95 uses QCD sum rules at next-to-leading order. We have rescaled $m_{s}(1 \mathrm{GeV})=171 \pm 15$ to $\mu=2 \mathrm{GeV}$.
34 JAMIN 95 uses QCD sum rules at next-to-leading order. We have rescaled $m_{s}(1 \mathrm{GeV})$ $=189 \pm 32$ to $\mu=2 \mathrm{GeV}$.
${ }^{35}$ For NARISON 95C, we have rescaled $m_{s}(1 \mathrm{GeV})=197 \pm 29$ to $\mu=2 \mathrm{GeV}$.
${ }^{36}$ NEFKENS 92 results for $m_{s}-\left(m_{u}+m_{d}\right) / 2$ are $111 \pm 10 \mathrm{MeV}$ from meson masses and $163 \pm 15 \mathrm{MeV}$ from baryon masses.
37 DOMINGUEZ 91 uses QCD sum rules with $\Lambda_{\mathrm{QCD}}=100-200 \mathrm{MeV}$ and the SVZ value for the gluon condensate. We have rescaled $m_{s}(1 \mathrm{GeV})=194 \pm 9$ to $\mu=2 \mathrm{GeV}$.
38 BARDUCCI 88 uses a calculation of the effective potential for $\bar{\psi} \psi$ in QCD, and estimates for $\Sigma\left(p^{2}\right)$. We have rescaled $m_{s}(1 \mathrm{GeV})=118$ to $\mu=2 \mathrm{GeV}$.
39 KREMER 84 obtain $m_{u}+m_{s}=245 \pm 10 \mathrm{MeV}$ at $Q^{2}=1 \mathrm{GeV}^{2}$ using SVZ values for quark condensates; they obtain $m_{u}+m_{s}=270 \pm 10 \mathrm{MeV}$ at $Q^{2}=1 \mathrm{GeV}^{2}$ using factorization values for quark condensates.
40 GASSER 82 uses chiral perturbation theory for the mass ratios, and uses QCD sum rules to extract the absolute values. We have rescaled $m_{s}(1 \mathrm{GeV})=175 \pm 55$ to $\mu=2 \mathrm{GeV}$.

## LIGHT QUARK MASS RATIOS

## $u / d$ MASS RATIO

VALUE DOCUMENT ID TECN COMMENT

## 0.2 to 0.7 OUR EVALUATION

-     - We do not use the following data for averages, fits, limits, etc.

| 0.44 | 41 GAO | 97 | THEO | $\overline{M S}$ scheme |
| :--- | :--- | :--- | :--- | :--- |
| $0.553 \pm 0.043$ | 42 LEUTWYLER | 96 | THEO | Compilation |
| $<0.3$ | 43 CHOI | 92 | THEO |  |
| 0.26 | 44 DONOGHUE | 92 | THEO |  |
| $0.30 \pm 0.07$ | 45 DONOGHUE | 92 B | THEO |  |
| 0.66 | 46 GERARD | 90 | THEO |  |
| 0.4 | to 0.65 | 47 LEUTWYLER | 90 THEO |  |
| 0.05 to 0.78 | 48 MALTMAN | 90 THEO |  |  |
| 0.0 to 0.56 | 49 CHOI | 89 B | THEO |  |
| 0.0 | to 0.8 | 50 KAPLAN | 86 THEO |  |
| $0.57 \pm 0.04$ | 51 GASSER | 82 | THEO |  |
| $0.38 \pm 0.13$ | 52 LANGACKER | 79 | THEO |  |
| $0.47 \pm 0.11$ | 53 LANGACKER | 79 B | THEO |  |
| 0.56 | 54 WEINBERG | 77 | THEO |  |

41 GAO 97 uses electromagnetic mass splittings of light mesons.
${ }^{42}$ LEUTWYLER 96 uses a combined fit to $\eta \rightarrow 3 \pi$ and $\psi^{\prime} \rightarrow J / \psi(\pi, \eta)$ decay rates, and the electromagnetic mass differences of the $\pi$ and $K$.
${ }^{43} \mathrm{CHOI} 92$ result obtained from the decays $\psi(2 S) \rightarrow J / \psi(1 S) \pi$ and $\psi(2 S) \rightarrow J / \psi(1 S) \eta$, and a dilute instanton gas estimate of some unknown matrix elements.
44 DONOGHUE 92 result is from a combined analysis of meson masses, $\eta \rightarrow 3 \pi$ using second-order chiral perturbation theory including nonanalytic terms, and $(\psi(2 S) \rightarrow$ $J / \psi(1 S) \pi) /(\psi(2 S) \rightarrow J / \psi(1 S) \eta)$.

45 DONOGHUE 92B computes quark mass ratios using $(\psi(2 S) \rightarrow J / \psi(1 S) \pi) /(\psi(2 S) \rightarrow$ $J / \psi(1 S) \eta$ ), and an estimate of $L_{14}$ using Weinberg sum rules.
46 GERARD 90 uses large $N$ and $\eta-\eta^{\prime}$ mixing.
47 LEUTWYLER 90B determines quark mass ratios using second-order chiral perturbation theory for the meson and baryon masses, including nonanalytic corrections. Also uses Weinberg sum rules to determine $L_{7}$.
48 MALTMAN 90 uses second-order chiral perturbation theory including nonanalytic terms for the meson masses. Uses a criterion of "maximum reasonableness" that certain coefficients which are expected to be of order one are $\leq 3$.
${ }^{49} \mathrm{CHOI} 89$ uses second-order chiral perturbation theory and a dilute instanton gas estimate of second-order coefficients in the chiral lagrangian.
50 KAPLAN 86 uses second-order chiral perturbation theory including nonanalytic terms for the meson masses. Assumes that less than $30 \%$ of the mass squared of the pion is due to second-order corrections.
51 GASSER 82 uses chiral perturbation theory for the meson and baryon masses.
52 LANGACKER 79 result is from a fit to the meson and baryon mass spectrum, and the decay $\eta \rightarrow 3 \pi$. The electromagnetic contribution is taken from Socolow rather than from Dashen's formula.
53 LANGACKER 79B result uses LANGACKER 79 and also $\rho-\omega$ mixing.
54 WEINBERG 77 uses lowest-order chiral perturbation theory for the meson and baryon masses and Dashen's formula for the electromagnetic mass differences.

## $s / d$ MASS RATIO

VALUE DOCUMENT ID TECN COMMENT

## 17 to 25 OUR EVALUATION

-     -         - We do not use the following data for averages, fits, limits, etc. • • •

| 20.0 | 55 GAO | 97 | THEO | $\overline{M S}$ scheme |
| :--- | :--- | :--- | :--- | :--- |
| $18.9 \pm 0.8$ | 56 LEUTWYLER | 96 | THEO | Compilation |
| 21 | 57 DONOGHUE | 92 | THEO |  |
| 18 | 58 GERARD | 90 | THEO |  |
| 18 to 23 | 59 LEUTWYLER | 90 THEO |  |  |
| 15 to 26 | 60 KAPLAN | 86 | THEO |  |
| $19.6 \pm 1.5$ | 61 GASSER | 82 | THEO |  |
| $22 \pm 5$ | 62 LANGACKER | 79 | THEO |  |
| $24 \pm 4$ | 63 LANGACKER | 79 THEO |  |  |
| 20 | 64 WEINBERG | 77 | THEO |  |

55 GAO 97 uses electromagnetic mass splittings of light mesons.
56 LEUTWYLER 96 uses a combined fit to $\eta \rightarrow 3 \pi$ and $\psi^{\prime} \rightarrow J / \psi(\pi, \eta)$ decay rates, and the electromagnetic mass differences of the $\pi$ and $K$.
57 DONOGHUE 92 result is from a combined analysis of meson masses, $\eta \rightarrow 3 \pi$ using second-order chiral perturbation theory including nonanalytic terms, and $(\psi(2 S) \rightarrow$ $J / \psi(1 S) \pi) /(\psi(2 S) \rightarrow J / \psi(1 S) \eta)$.
58 GERARD 90 uses large $N$ and $\eta-\eta^{\prime}$ mixing.
59 LEUTWYLER 90B determines quark mass ratios using second-order chiral perturbation theory for the meson and baryon masses, including nonanalytic corrections. Also uses Weinberg sum rules to determine $L_{7}$.
60 KAPLAN 86 uses second-order chiral perturbation theory including nonanalytic terms for the meson masses. Assumes that less than $30 \%$ of the mass squared of the pion is due to second-order corrections.
61 GASSER 82 uses chiral perturbation theory for the meson and baryon masses.
62 LANGACKER 79 result is from a fit to the meson and baryon mass spectrum, and the decay $\eta \rightarrow 3 \pi$. The electromagnetic contribution is taken from Socolow rather than from Dashen's formula.
63 LANGACKER 79B result uses LANGACKER 79 and also $\rho-\omega$ mixing.
64 WEINBERG 77 uses lowest-order chiral perturbation theory for the meson and baryon masses and Dashen's formula for the electromagnetic mass differences.

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\(\left(m_{s}-m\right) /\left(m_{d}-m_{u}\right)\) MASS RATIO
    \(\bar{m} \equiv\left(m_{u}+m_{d}\right) / 2\)
VALUE DOCUMENTID TECN
34 to 51 OUR EVALUATION
- - We do not use the following data for averages, fits, limits, etc. • • -
\begin{tabular}{llll}
\(36 \pm 5\) & 66 NEFKENS & 92 & THEO \\
\(45 \pm 3\) & 67 NEFKENS & 92 & THEO \\
\(38 \pm 9\) & 68 AMETLLER & 84 & THEO \\
\(43.5 \pm 2.2\) & GASSER & 82 & THEO \\
34 to 51 & GASSER & 81 & THEO \\
\(48 \pm 7\) & MINKOWSKI & 80 & THEO
\end{tabular}
    \({ }^{65}\) ANISOVICH 96 find \(Q=22.7 \pm 0.8\) with \(Q^{2} \equiv\left(m_{s}^{2}-m^{2}\right) /\left(m_{d}^{2}-m_{s}^{2}\right)\) from \(\eta \rightarrow\)
        \(\pi^{+} \pi^{-} \pi^{0}\) decay using dispersion relations and chiral perturbation theory.
    66 NEFKENS 92 result is from an analysis of meson masses, mixing, and decay.
    67 NEFKENS 92 result is from an analysis of of baryon masses.
    68 AMETLLER 84 uses \(\eta \rightarrow \pi^{+} \pi^{-} \pi^{0}\) and \(\rho\) dominance.
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## LIGHT QUARKS ( $u, d, s$ ) REFERENCES

| CHETYRKIN | 98 | NP B533 473 |
| :---: | :---: | :---: |
| CUCCHIERI | 98 | PL B422 212 |
| DOMINGUEZ | 98 | PL B425 193 |
| DOSCH | 98 | PL B417 173 |
| LEBED | 98 | PL B430 341 |
| MALTMAN | 98 | PL B428 179 |
| MALTMAN | 98B | PR D58 093015 |
| CHETYRKIN | 97 | PL B404 337 |
| COLANGELO | 97 | PL B408 340 |
| EICKER | 97 | PL B407 290 |
| GAO | 97 | PR D56 4115 |
| GOUGH | 97 | PRL 791622 |
| GUPTA | 97 | PR D55 7203 |
| LELLOUCH | 97 | PL B414 195 |
| ANISOVICH | 96 | PL B375 335 |
| LEUTWYLER | 96 | PL B378 313 |
| BIJNENS | 95 | PL B348 226 |
| CHETYRKIN | 95 | PR D51 5090 |
| JAMIN | 95 | ZPHY C66 633 |
| NARISON | 95 C | PL B358 113 |
| ADAMI | 93 | PR D48 2304 |
| CHOI | 92 | PL B292 159 |
| CHOI | 92B | NP B383 58 |
| DONOGHUE | 92 | PRL 693444 |
| DONOGHUE | 92B | PR D45 892 |
| NEFKENS | 92 | CNPP 20221 |
| DOMINGUEZ | 91 | PL B253 241 |
| GERARD | 90 | MPL A5 391 |
| LEUTWYLER | 90B | NP B337 108 |
| MALTMAN | 90 | PL B234 158 |
| CHOI | 89 | PRL 62849 |
| CHOI | 89B | PR D40 890 |
| BARDUCCI | 88 | PR D38 238 |
| Also | 87 | PL B193 305 |


| A. Chucchieri+ <br> C.A. Dominguez, L. Pirovano, K. Schilcher |  |
| :---: | :---: |
| H.G. Dosch, S. Narison |  |
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| A.V. Anisovich, H. Leutwyler |  |
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| +Dominguez, Pirjol, Schilche | (INRM, CAPE, MANZ) |
| +Munz | (HEIDT, MUNT) |
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| +van Gend, Paver | (CAPE, TRST, INFN) |
|  | (MPIM) |
|  | (BERN) |
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| +Casalbuoni, De Curtis+ | (FIRZ, INFN, LECE, GEVA) |
| Barducci, Casalbuoni+ | (FIRZ, INFN, LECE, GEVA) |


| DOMINGUEZ | 87 | ANP 174 372 | +de Rafael | (ICTP, MARS, WIEN) |
| :--- | :--- | :--- | :--- | ---: |
| KAPLAN | 86 | PRL 56 2004 | +Manohar | (HARV) |
| AMETLLER | 84 | PR D30 674 | +Ayala, Bramon | (BARC) |
| KREMER | 84 | PL 143B 476 | +Papadopoulos, Schilcher | (MANZ) |
| GASSER | 82 | PRPL 8777 | +Leutwyler | (BERN) |
| GASSER | 81 | ANP 136 62 |  | (BERN) |
| MINKOWSKI | 80 | NP B164 25 | +Zepeda | (BERN) |
| LANGACKER | 79 | PR D19 2070 | +Pagels | (DESY, PRIN) |
| LANGACKER | $79 B$ | PR D20 2983 |  | (PENN) |
| WEINBERG | 77 | ANYAS 38185 |  | (HARV) |
|  |  |  |  |  |

