## THE $\eta(1440), f_{1}(1420)$, AND $f_{1}(1510)$

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The first observation of $\eta(1440)$ was made in $p \bar{p}$ annihilation at rest into $\eta(1440) \pi^{+} \pi^{-}, \eta(1440) \rightarrow K \bar{K} \pi$ (BAILLON 67 ). This state was reported to decay through $a_{0}(980) \pi$ and $K^{*}(892) \bar{K}$ with roughly equal contributions. The $\eta(1440)$ has also been observed in radiative $J / \psi(1 S)$ decay to $K \bar{K} \pi$ (SCHARRE 80, EDWARDS 82E, AUGUSTIN 90).

The $f_{1}(1420)$, decaying to $K^{*} \bar{K}$ was reported in $\pi^{-} p$ reactions at $4 \mathrm{GeV} / c$ (DIONISI 80). However, later analyses found that the $1400-1500 \mathrm{MeV}$ region is far more complex. In $\pi^{-} p$ experiments (CHUNG 85, REEVES 86, BIRMAN 88) reported $0^{-+}$with a dominant $a_{0}(980) \pi$ contribution to $K \bar{K} \pi$. The $\pi^{-} p$ data of RATH 89 at $21 \mathrm{GeV} / c$ suggest the presence of two pseudoscalars decaying to $K \bar{K} \pi$, one around 1410 MeV decaying through $a_{0}(980) \pi$ and the other around 1470 MeV , decaying to $K^{*} \bar{K}$. A reanalysis of the MARK III data in radiative $J / \psi(1 S)$ decay to $K \bar{K} \pi$ (BAI 90C) also claims the existence of two pseudoscalars in the $1400-1500 \mathrm{MeV}$ range, the lower mass state decaying through $a_{0}(980) \pi$ and the higher mass state decaying via $K^{*} \bar{K}$. In addition, $f_{1}(1420)$ is observed to decay into $K^{*} \bar{K}$.

In $\pi^{-} p \rightarrow \eta \pi \pi n$ charge-exchange reactions at $8-9 \mathrm{GeV} / c$ the $\eta \pi \pi$ mass spectrum is dominated by $\eta(1440)$ and $\eta(1295)$ (ANDO 86, FUKUI 91C) and at 100 GeV ALDE 97B report $\eta(1295)$ and $\eta(1440)$ decaying to $\eta \pi^{0} \pi^{0}$ with a weak $f_{1}(1285)$ and no evidence for $f_{1}(1420)$.

An experiment in $\bar{p} p$ annihilation at rest into $K \bar{K} 3 \pi$ (BERTIN 95) reports two pseudoscalars with decay properties similar to BAI 90C, although the lower state shows, apart from $a_{0}(980) \pi$, a large contribution from the direct decay $\eta(1440) \rightarrow K \bar{K} \pi$. We note that the data from AUGUSTIN 92 also suggest two states but their intermediate states, $a_{0}(980) \pi$ and $K^{*} \bar{K}$, are reversed relative to BAI 90C.

In $J / \psi(1 S)$ radiative decay $\eta(1440)$ decays to $K \bar{K} \pi$ through $a_{0}(980) \pi$ and hence a signal is also expected in the $\eta \pi \pi$ mass spectrum. This has indeed been observed by MARK III in
$\eta \pi^{+} \pi^{-}$(BOLTON 92B) which report a mass of 1400 MeV , in line with the existence of a low mass pseudoscalar in the $\eta(1440)$ structure, decaying to $a_{0}(980) \pi$. This state is also observed in $\bar{p} p$ annihilation at rest into $\eta \pi^{+} \pi^{-} \pi^{0} \pi^{0}$ where it decays to $\eta \pi \pi$ (AMSLER 95F). The intermediate $a_{0}(980) \pi$ accounts for roughly half of the $\eta \pi \pi$ rate, in accord with MARK III (BOLTON 92B) and DM2 (AUGUSTIN 90). However, ALDE 97B reports only a very small contribution of $a_{0}(980) \pi$.

One of these two pseudoscalars could be the first radial excitation of the $\eta^{\prime}$, with $\eta(1295)$ the first radial of the $\eta$. Ideal mixing suggested by the $\eta(1295)$ and $\pi(1300)$ mass degeneracy would then imply that the second isoscalar in the nonet is mainly $s \bar{s}$ and hence couples to $K^{*} \bar{K}$, in accord with observations for the upper $\eta(1440)$ state. This scheme then favors an exotic interpretation of the lower state, perhaps gluonium mixed with $q \bar{q}$ (CLOSE 97B) or a bound state of gluinos (FARRAR 96). The gluonium interpretation is, however, not favoured by lattice gauge theories, which predict the $0^{-+}$state above 2 GeV (BALI 93).

Axial $\left(1^{++}\right)$mesons are not observed in $\bar{p} p$ annihilation at rest in liquid hydrogen which proceeds dominantly through $S$-wave annihilation. However, in gaseous hydrogen $P$-wave annihilation is enhanced and, indeed, BERTIN 97 report $f_{1}(1420)$ decaying to $K^{*} \bar{K}$ in gaseous hydrogen, while confirming their earlier evidence for two pseudoscalars (BERTIN 95).

In $\gamma \gamma$ fusion from $e^{+} e^{-}$annihilations, a signal around 1420 MeV is seen in single-tag events (GIDAL 87B, AIHARA 88B, BEHREND 89, HILL 89) where one of the two photons is off-shell. However, it is totally absent in the untagged events where both photons are real. This points to a spin 1 object which is not produced by two real (massless) photons (YangLandau theorem). The $2 \gamma$ decays also implies $C=+1$. For the parity, AIHARA 88C and BEHREND 89 both find angular distributions with positive parity preferred, but negative parity cannot be excluded.

The $f_{1}(1420)$ is definitively observed in $K \bar{K} \pi$ in $p p$ central production at 300 and 450 GeV , together with $f_{1}(1285)$. The latter decays via $a_{0}(980) \pi$ and the former only via $K^{*} \bar{K}$, while
$\eta(1440)$ is absent (ARMSTRONG 89, BARBERIS 97C). The $K_{S} K_{S} \pi^{0}$ decay mode of $f_{1}(1420)$ establishes unambiguously that $C=+1$. On the other hand, there is no evidence for any state decaying to $\eta \pi \pi$ around 1400 MeV and hence the $\eta \pi \pi$ mode of $f_{1}(1420)$ is suppressed (ARMSTRONG 91B).

We now turn to the experimental evidence for $f_{1}(1510)$. Two states, $f_{1}(1420)$ and $f_{1}(1510)$, decaying to $K^{*} \bar{K}$, compete for the $s \bar{s}$ assignment in the $1^{++}$nonet. The $f_{1}(1510)$ was seen in $K^{-} p \rightarrow \Lambda K \bar{K} \pi$ at $4 \mathrm{GeV} / c$ (GAVILLET 82) and at 11 $\mathrm{GeV} / c$ (ASTON 88C). Evidence is also reported in $\pi^{-} p$ at $8 \mathrm{GeV} / c$, based on the phase motion of the $1^{++} K^{*} \bar{K}$ wave (BIRMAN 88).

The absence of $f_{1}(1420)$ in $K^{-} p$ (ASTON 88C) argues against $f_{1}(1420)$ being the $s \bar{s}$ member of the $1^{++}$nonet. However, $f_{1}(1420)$ has been reported in $K^{-} p$ but not in $\pi^{-} p$ (BITYUKOV 84) while two experiments do not observe $f_{1}(1510)$ in $K^{-} p$ (BITYUKOV 84, KING 91). It is also not seen in radiative $J / \psi(1 S)$ decay (BAI 90C, AUGUSTIN 92), central collisions (BARBERIS 97C), nor in $\gamma \gamma$ collisions (AIHARA 88C), although and surprisingly for an $s \bar{s}$ state, a signal is reported in $4 \pi$ decays (BAUER 93B). These facts led to the conclusion that $f_{1}(1510)$ is not well established and that its assignment as $s \bar{s}$ member of the $1^{++}$nonet is premature (CLOSE 97D). The Particle Data Group agrees and has removed this state from the Summary Table. Assigning instead $f_{1}(1420)$ to the $1^{++}$nonet one finds a nonet mixing angle of $\sim$ $50^{\circ}$ (CLOSE 97D). This is derived from the mass formula and from $f_{1}(1285)$ radiative decays to $\phi \gamma$ (BITYUKOV 88) and $\rho \gamma$ (AMELIN 95).

Arguments favoring $f_{1}(1420)$ being a hybrid $q \bar{q} g$ meson or a four-quark state are put forward by ISHIDA 89 and by CALDWELL 90, respectively, while LONGACRE 90 argues that this particle is a molecular state formed by the $\pi$ orbiting in a $P$-wave around an $S$-wave $K \bar{K}$ state.

Summarizing, there is strong evidence for $f_{1}(1420)$, mostly produced in central collisions and decaying to $K^{*} \bar{K}$, and for $\eta(1440)$ mostly produced in radiative $J / \psi(1 S)$ decay and $\bar{p} p$ annihilation at rest, decaying to $K^{*} \bar{K}$ and $a_{0}(980) \pi$. Confusion
remains as to which states are observed in $\pi^{-} p$ interactions. The $f_{1}(1510)$ is not well established. Furthermore, there are experimental indications for the presence of two pseudoscalars in the $\eta(1440)$ structure. Accordingly, the Particle Data Group has split the $K \bar{K} \pi$ entry for $\eta(1440)$ into $a_{0}(980) \pi$ and $K^{*} \bar{K}$.

